

Gauge Engineering and Propagators

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Thessaloniki
Greece

The logo for UNI GRAZ, featuring a yellow square on the right and a white square on the left, with the text 'UNI' above 'GRAZ' in black.

**UNI
GRAZ**

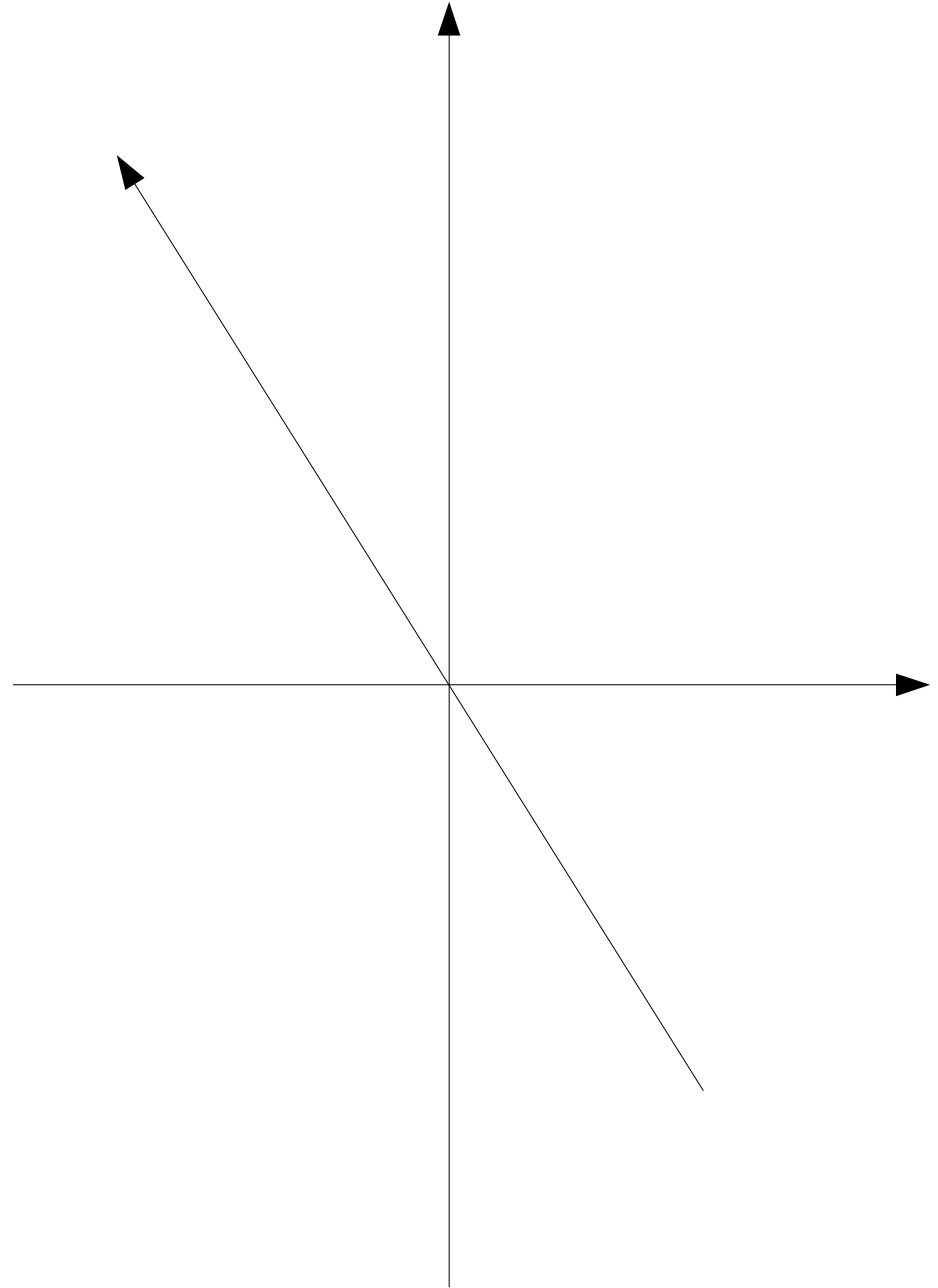


NAWI Graz
Natural Sciences

Gauge-fixing is engineering

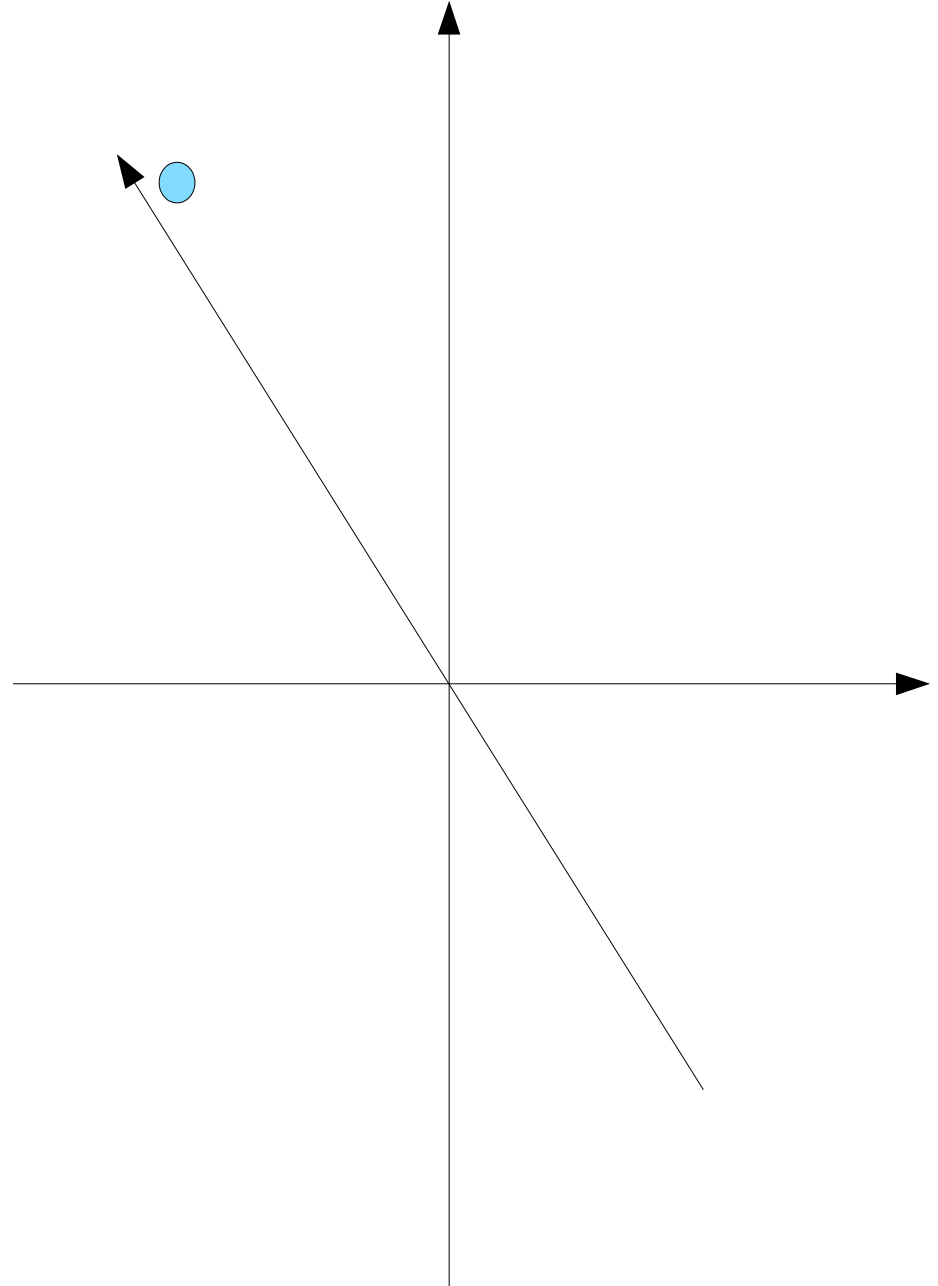
- All information of a theory is encoded in gauge-, renormalization-point, and renormalization-scheme-invariant quantities
 - I.e. measurable quantities, observables
- Other quantities helpful in determining these
 - Gauge-fixed quantities are helpful tools
- Engineering useful gauges is like choosing an appropriate coordinate system

Configuration space (artist's view)



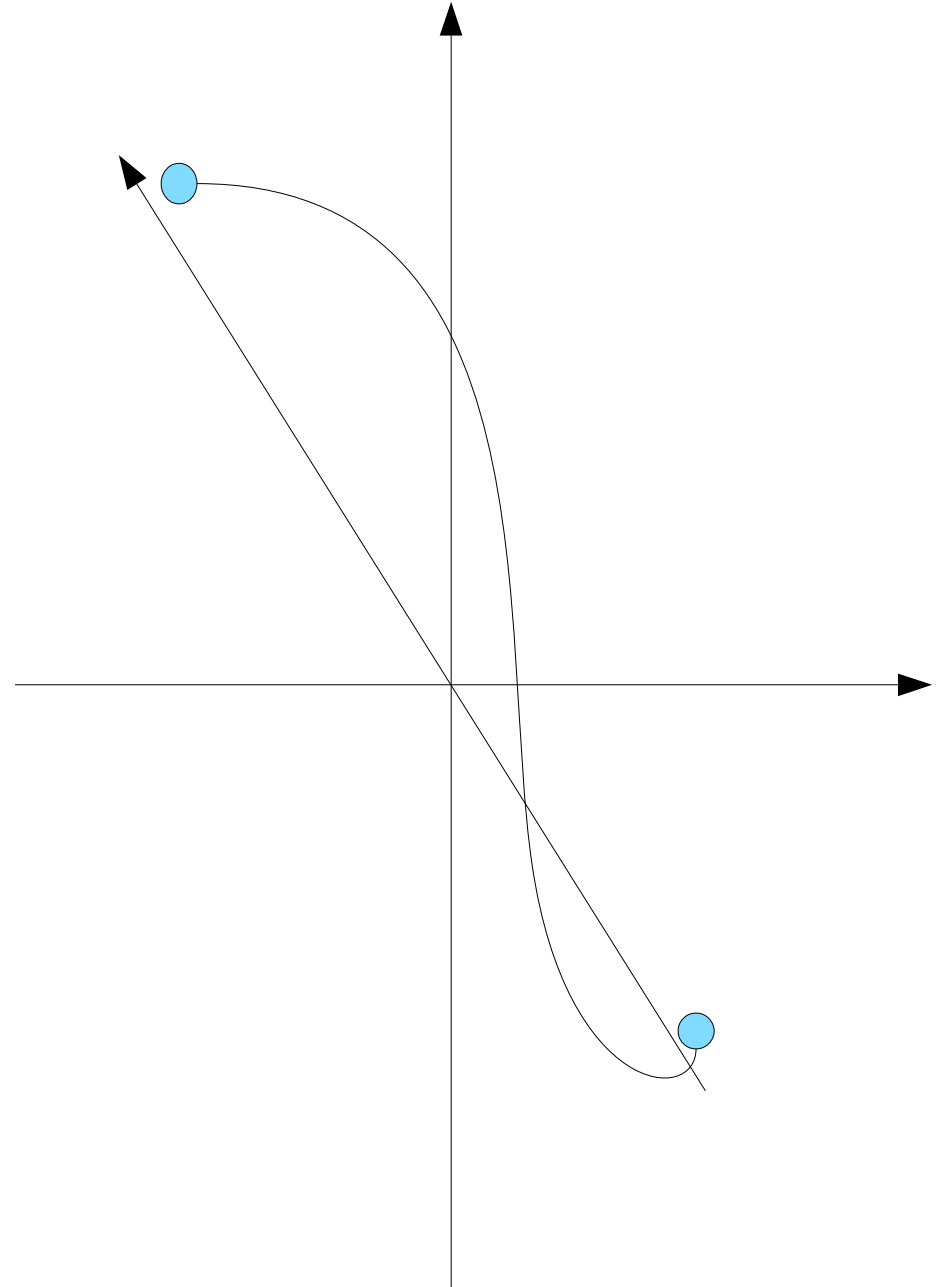
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- Gauge fields not unique



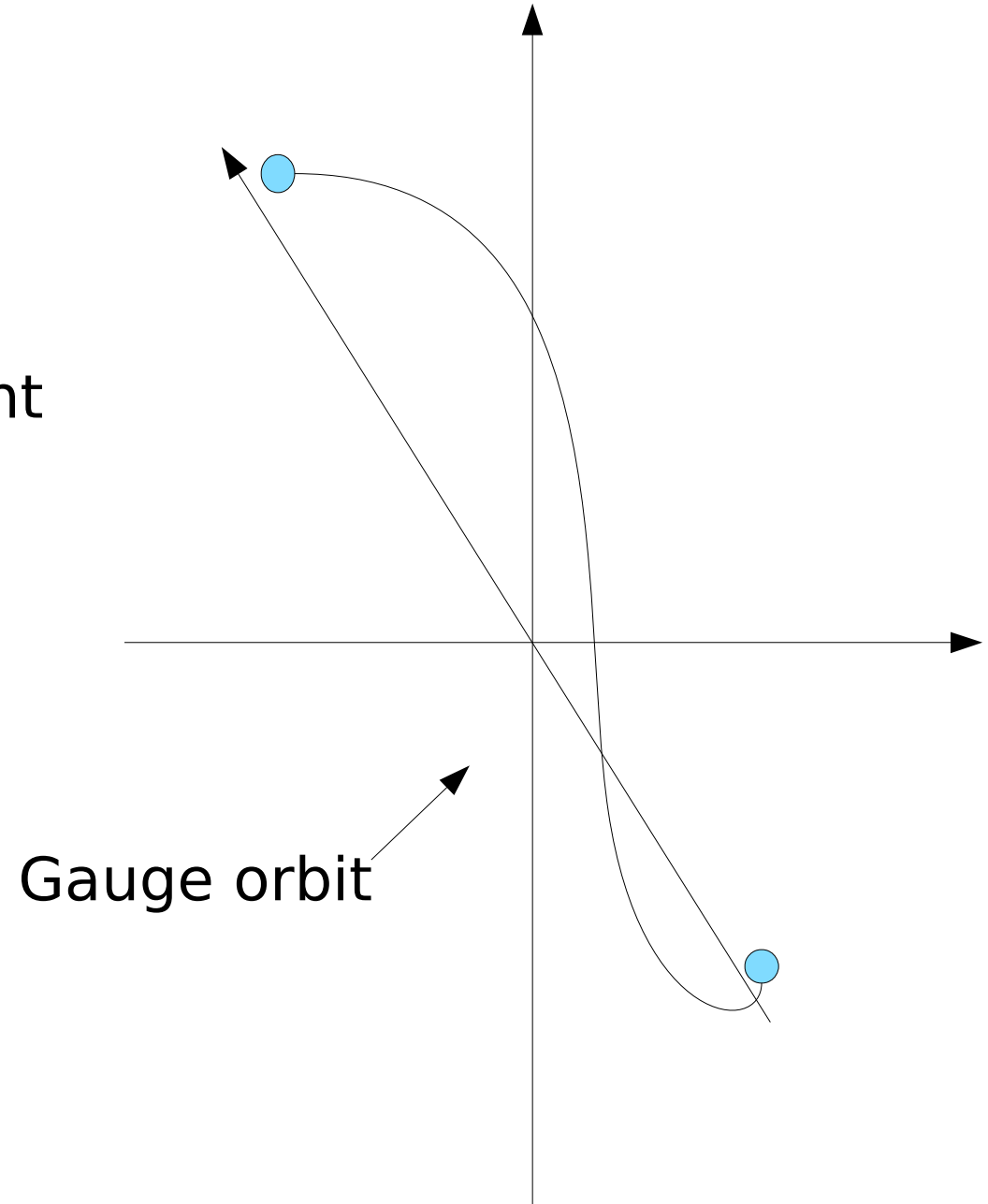
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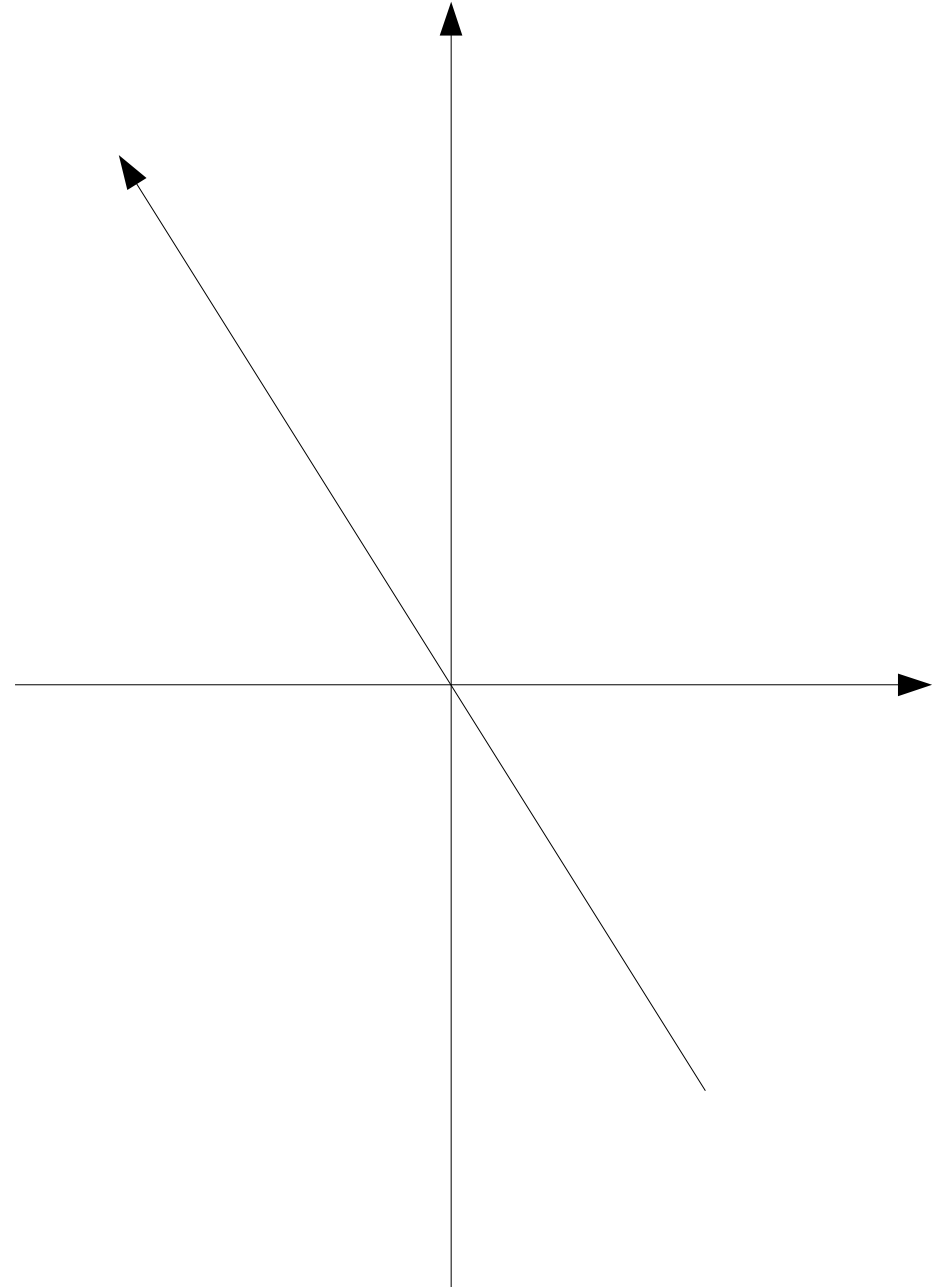
Configuration space (artist's view)

- Gauge fields not unique
 - Gauge transformation does not change physics
 - Set of all gauge-equivalent copies is a gauge orbit



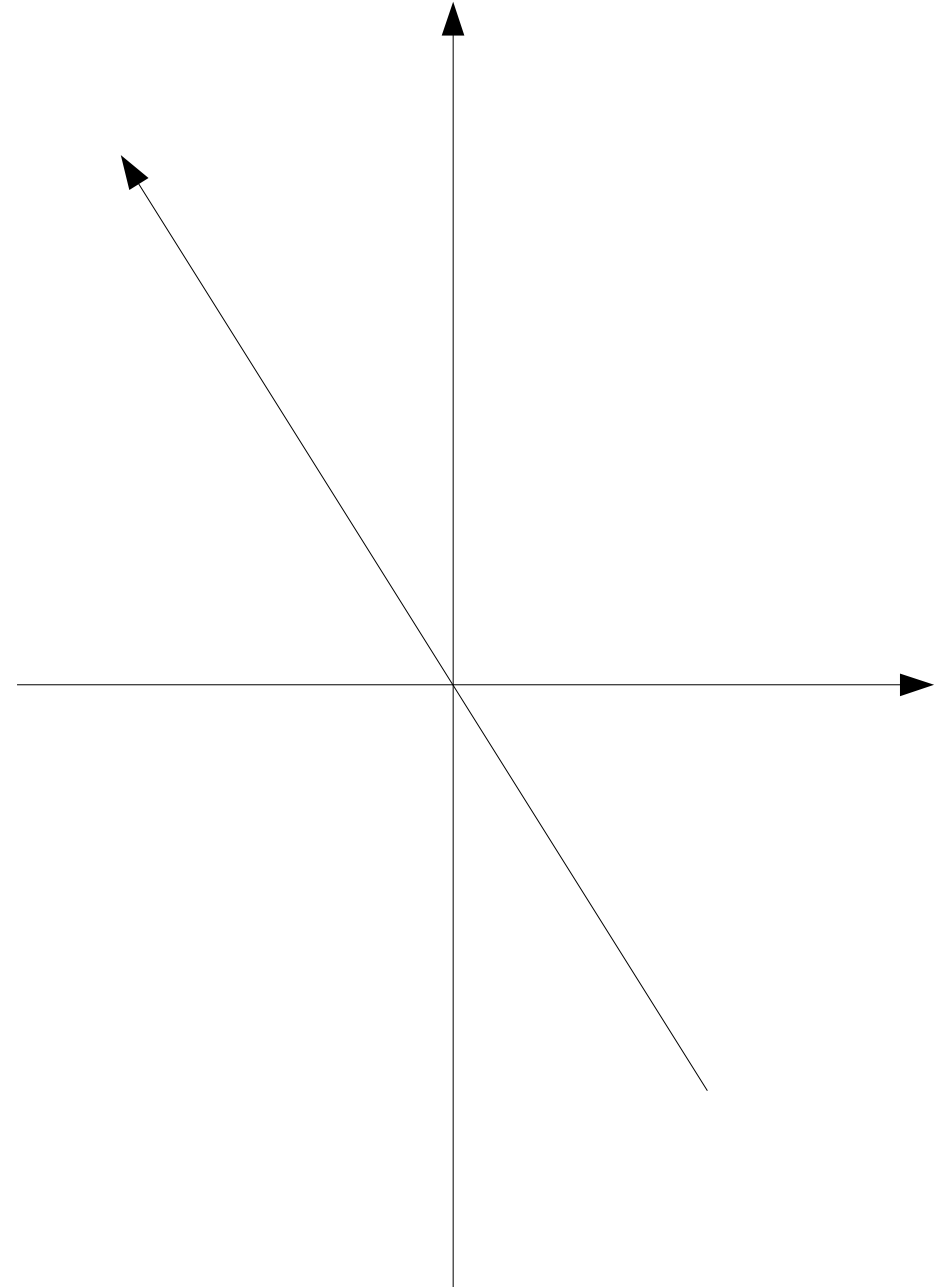
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- Fields and most correlation functions change under gauge transformations
- Requires a choice of gauge/coordinate system



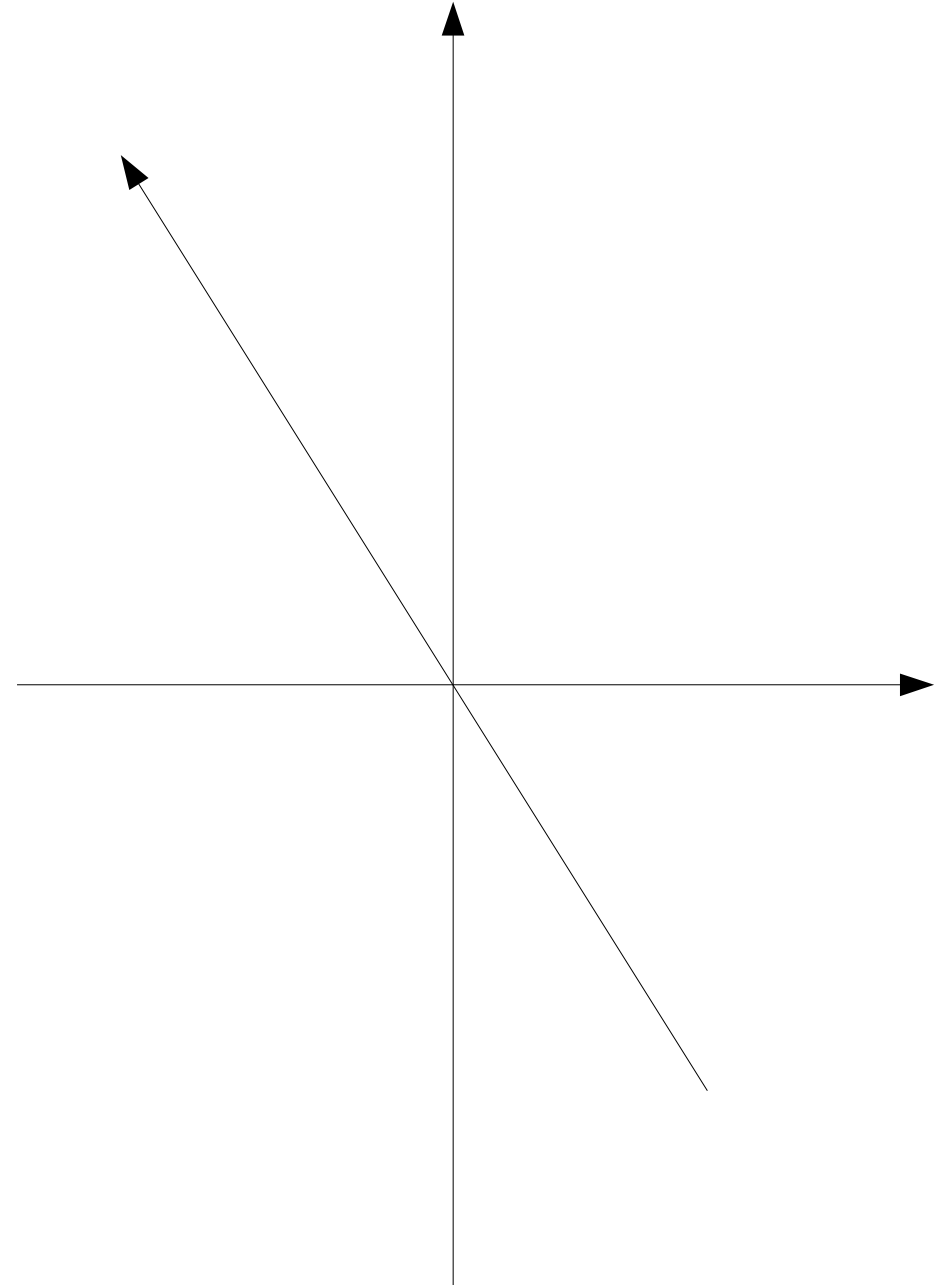
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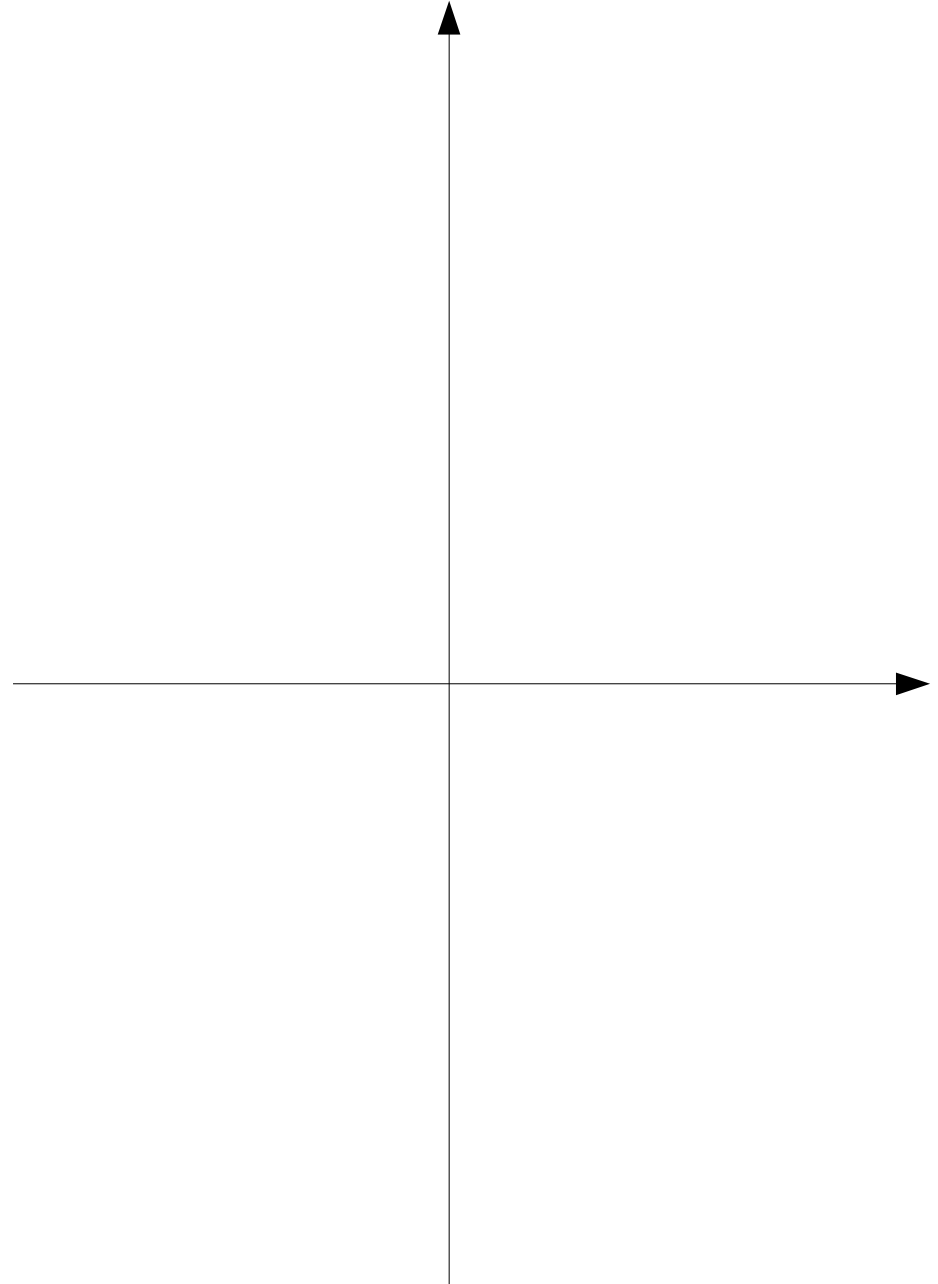
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- Local condition in perturbation theory
 - Here: Landau gauge $\partial^\mu A_\mu^a = 0$
 - Reduces configuration space to a hypersurface



(Perturbative) Landau gauge

- Condition can be implemented using auxiliary fields, the ghost fields

- Lagrangian:
$$L = -\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a - \bar{c}^a \partial_\mu D_\mu^{ab} c^b$$

$$D_\mu^{ab} = \delta^{ab} \partial_\mu - gf^{abc} A_\mu^c$$

$$\lim_{\zeta \rightarrow 0} \int DA c \bar{c} \exp\left(\int dx \left(-\frac{1}{4} F^2 - \bar{c} \partial D c - \frac{1}{2\zeta} (\partial A)^2\right)\right)$$

- Degrees of freedom: Gluons A_μ^a Ghosts \bar{c}^a, c^a
 - Not physical objects
 - Pure mathematical convenience

Unique gauge-fixing

[Sobreiro & Sorella, 2005]

- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
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- Local gauge condition
 - Landau gauge: $\partial_\mu A_\mu^a = 0$
- Sufficient for perturbation theory
- Insufficient beyond perturbation theory
 - There are gauge-equivalent configurations:
Gribov copies [Gribov NPB 78]
- No local gauge conditions known to select a unique gauge copy: Gribov-Singer ambiguity

[Singer CMP 78]

Residual freedom

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- Method: Lattice

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- Two (renormalized and integrated) 2-point functions
 - Momentum-integrated gluon propagator $F \sim -\int dp D_{\mu\mu}^{aa}(p)$
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- Why these two?

First horizon gauges

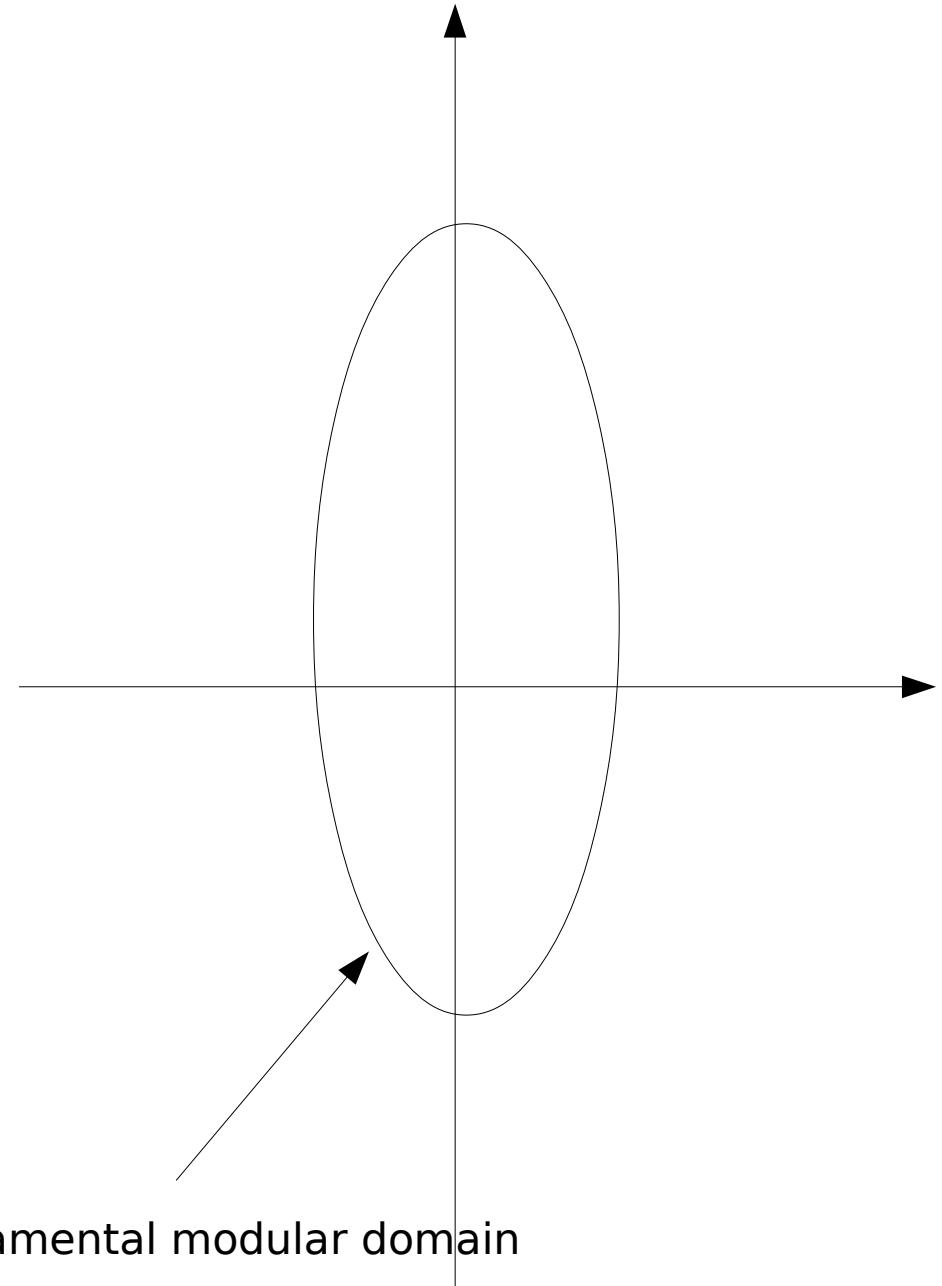
[Gribov NPA 78, Zwanziger 93...03]

- A global minimum of

$$F \sim - \int d^d x A_\mu^a(x) A_\mu^a(x)$$

defines the fundamental modular domain

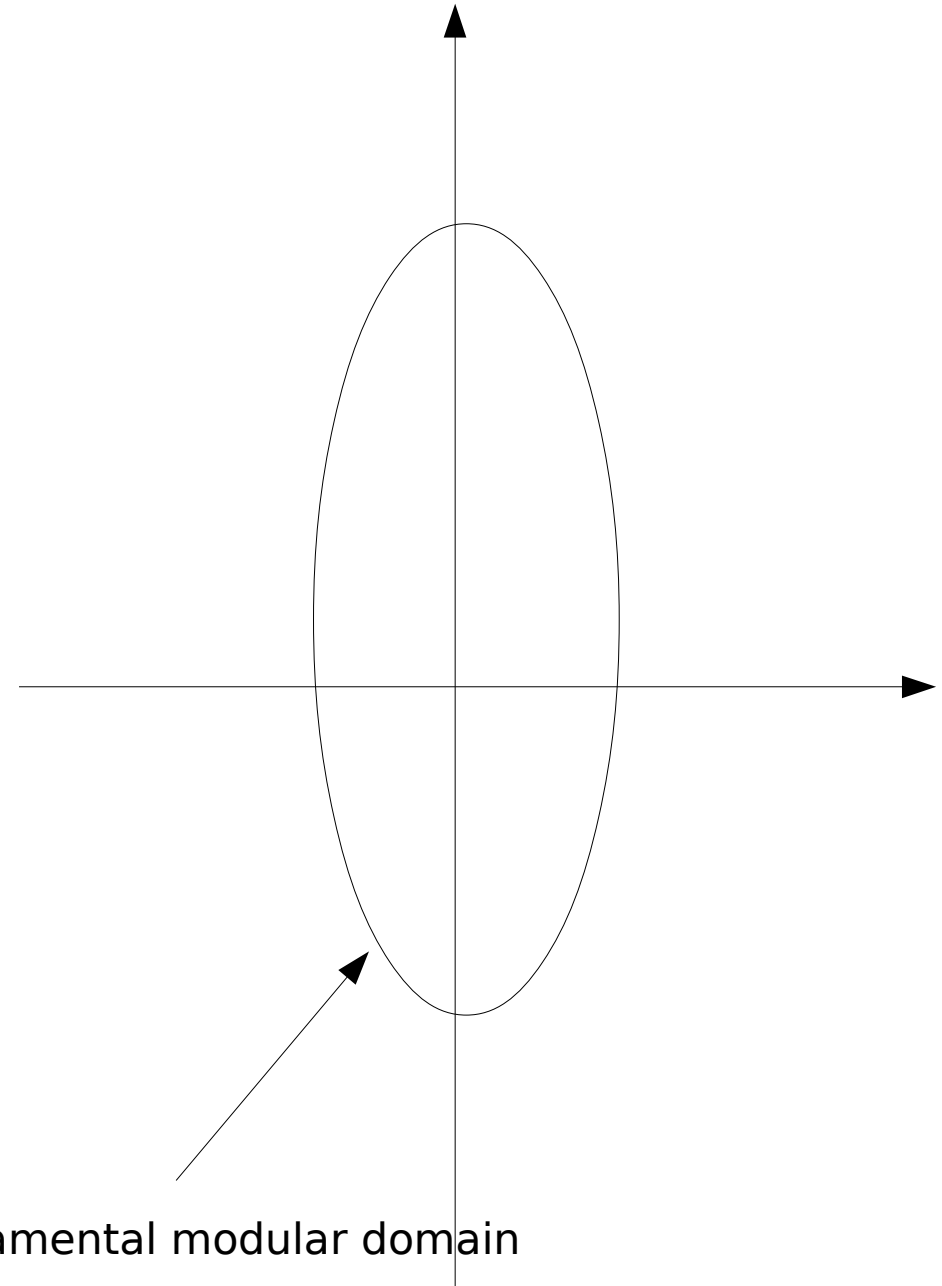
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First horizon gauges

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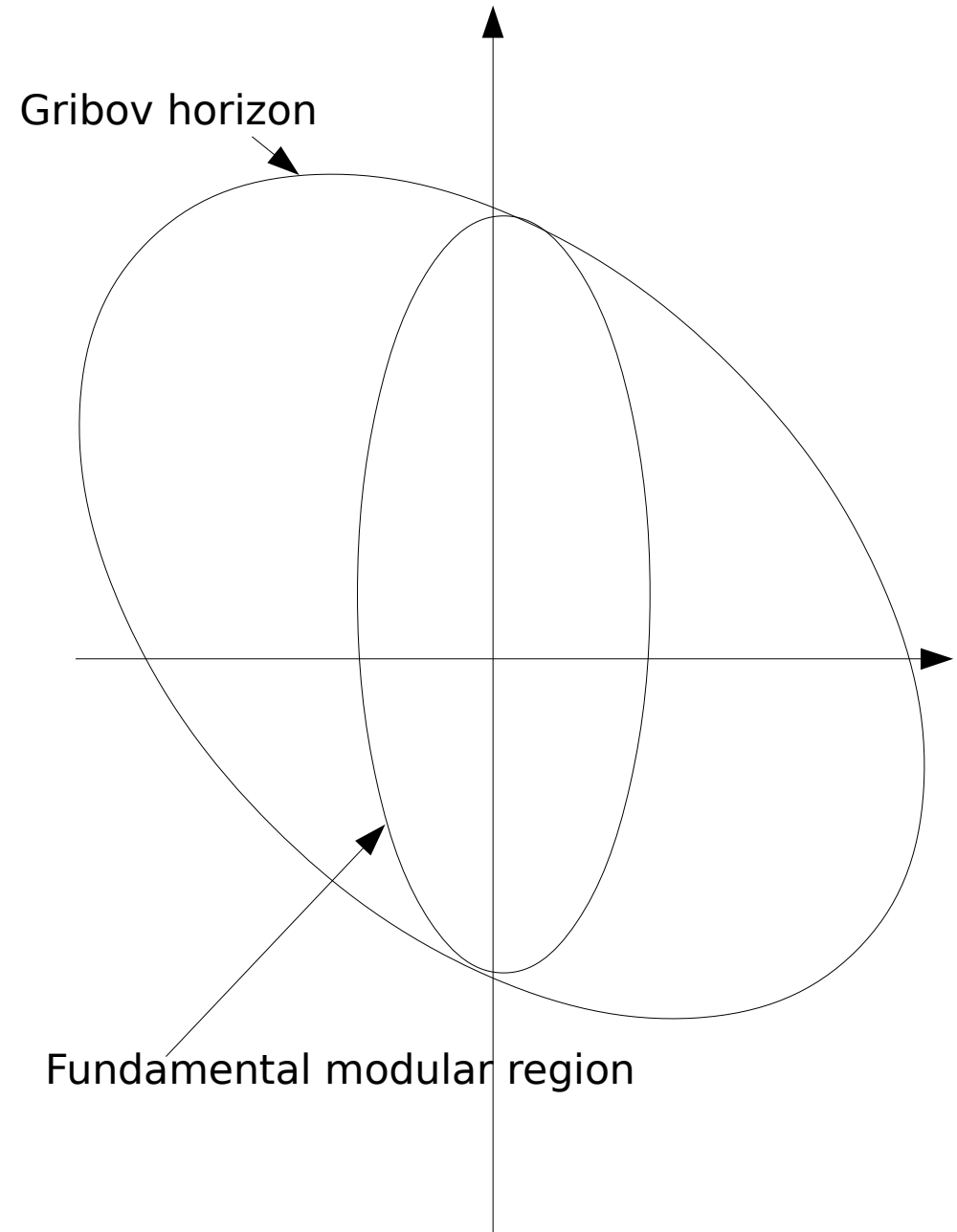
- A global minimum of
$$F \sim - \int d^d x A_\mu^a(x) A_\mu^a(x)$$
defines the fundamental modular domain
- This region is bounded and convex
- Singles out exactly one gauge copy
 - Defines absolute Landau gauge
 - Boundary subtle



First horizon gauges

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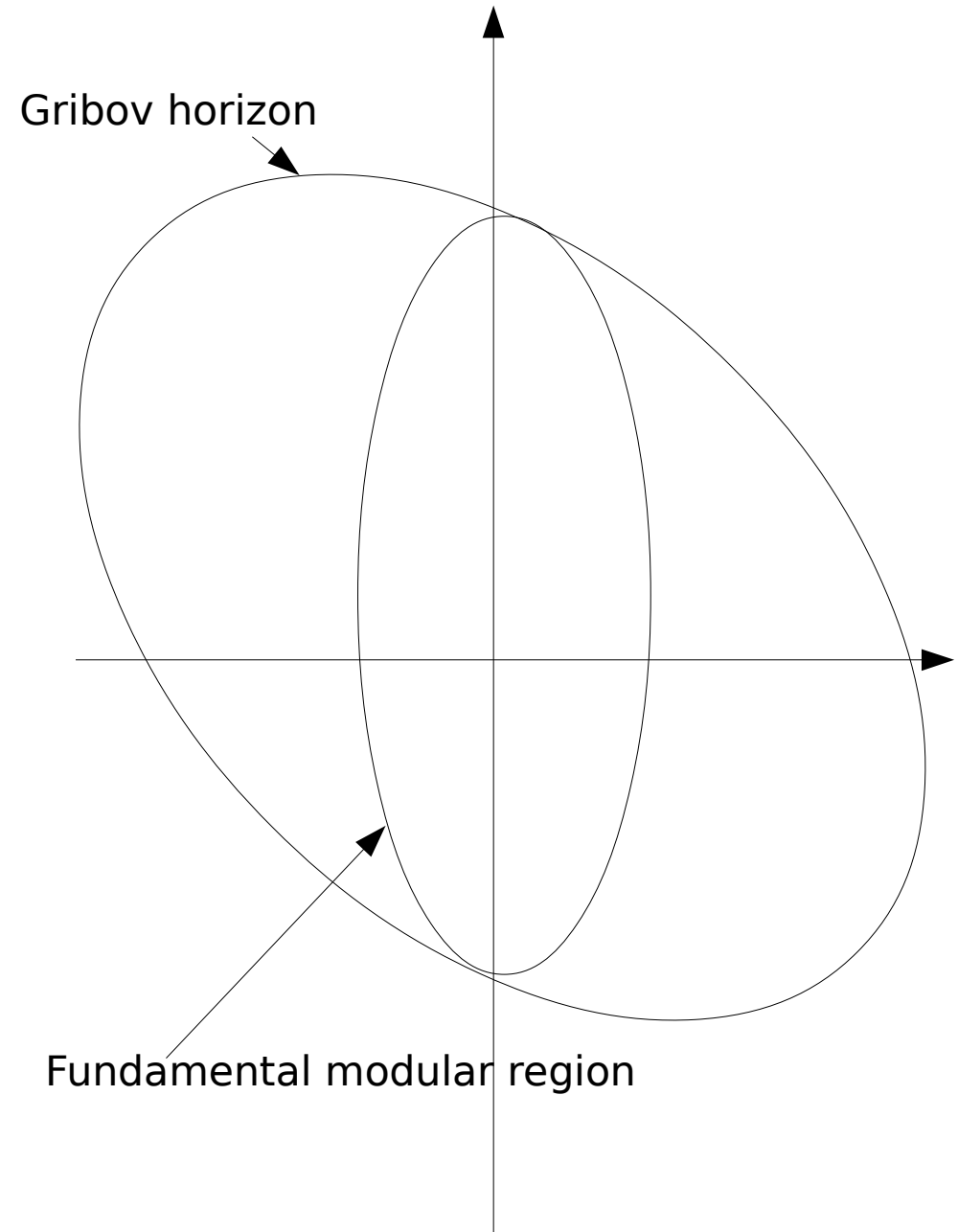
- Gribov horizon encloses all field configurations with non-negative Faddeev-Popov operator $-\partial_\mu D_\mu$
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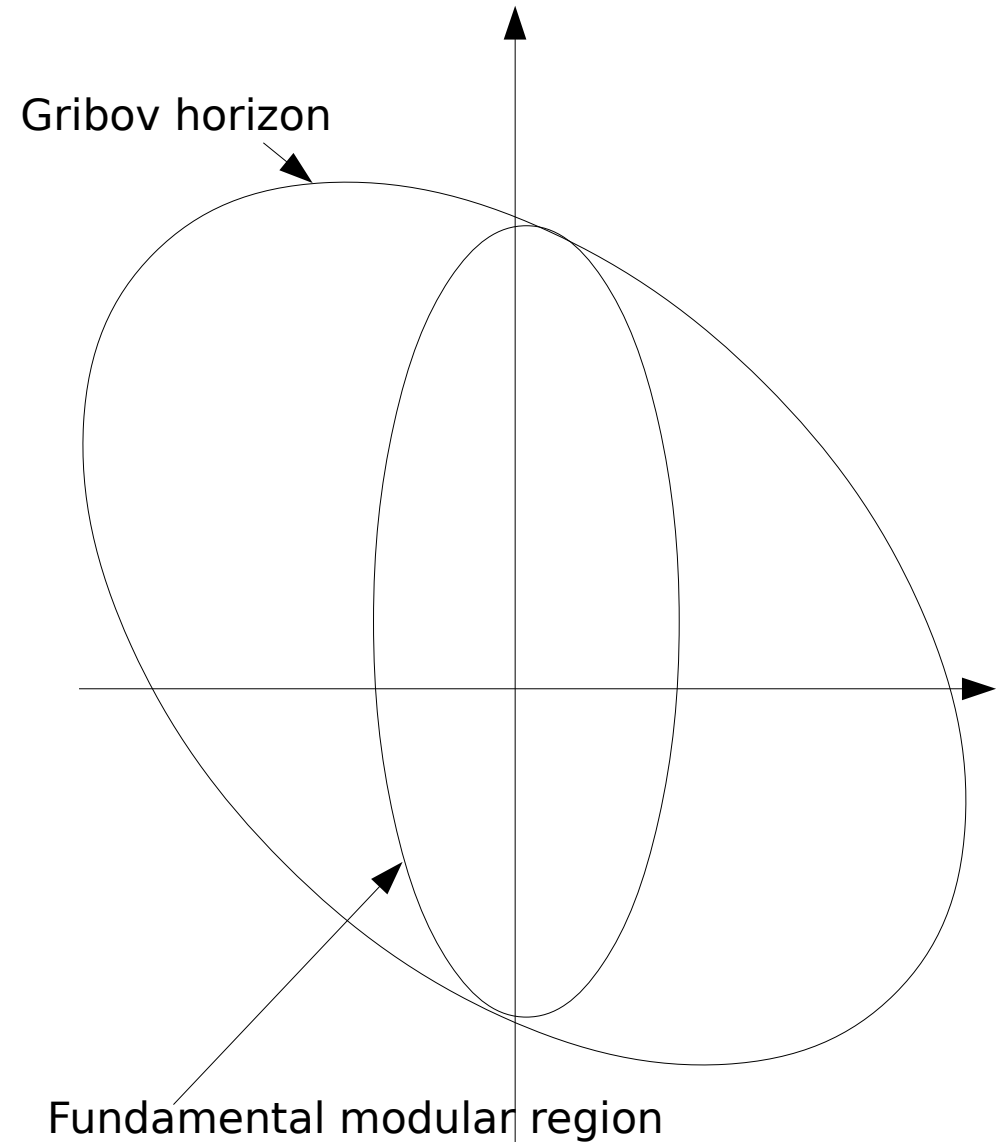


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- All gauge orbits pass through this region
 - Many Gribov copies
- Only this region:

$$\lim_{\zeta \rightarrow 0} \int DA c \bar{c} \exp \left(\int dx \left(-\frac{1}{4} F^2 - \bar{c} \partial D c - \frac{1}{2\zeta} (\partial A)^2 \right) \right) \Theta(-\partial D)$$



Identifying Gribov copies

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- Identifying Gribov copies non-trivial task
 - No construction to generate Gribov copies
 - Practical solution: Generate for each residual gauge orbit many random copies and select different ones [Cucchieri NPB 97]

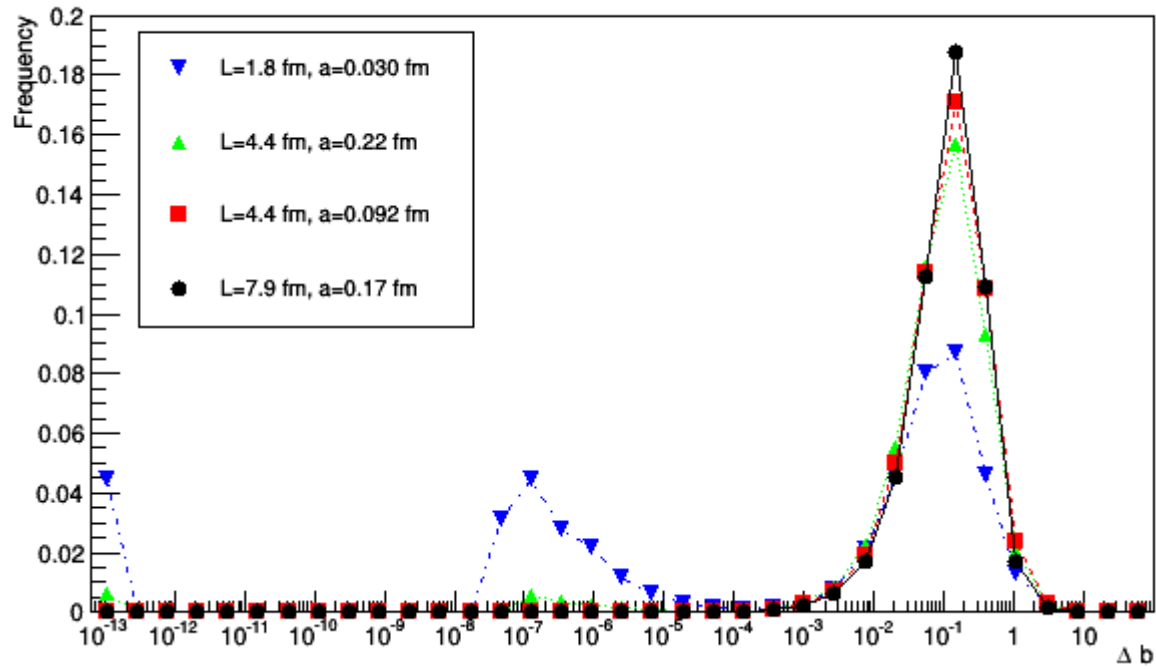
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 - Practical solution: Generate for each residual gauge orbit many random copies and select different ones [Cucchieri NPB 97]
 - When are two copies different? [Maas PRD 15]
 - Connected by non-trivial transformation
 - Practical: Sufficiently different value of F and B
 - Distinguishing power of B much better
 - Other dismissed as numerical/lattice artifacts

Separating Gribov copies

[Maas PRD 15]

Distribution of Δb in 3 dimensions

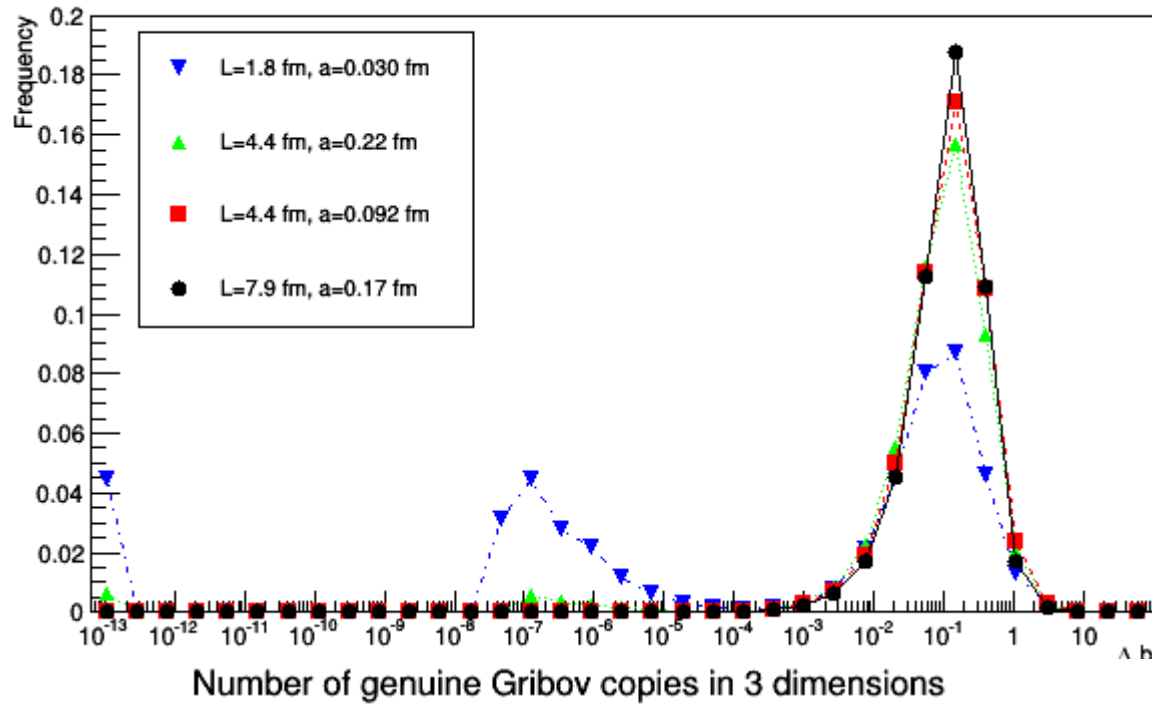


- Distinction better for larger volume and dimension

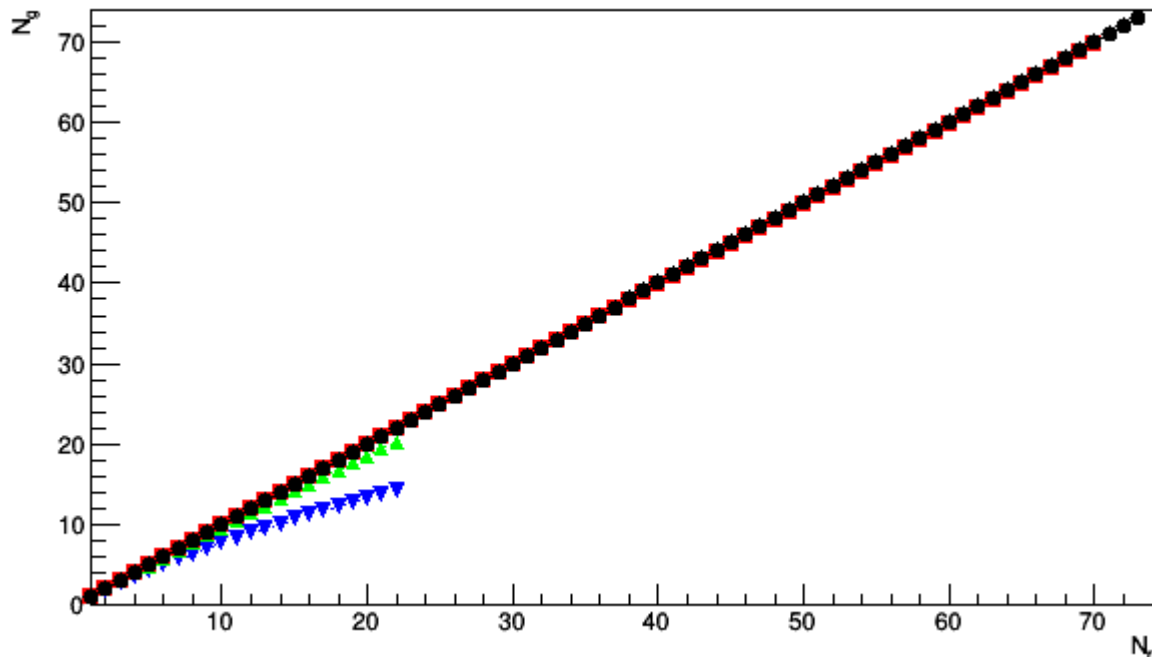
Separating Gribov copies

[Maas PRD 15]

Distribution of Δb in 3 dimensions



- Distinction better for larger volume and dimension
- Number of identified Gribov copies rises quickly
 - Linear on large volumes
 - No extrapolation to infinite-volume limit possible



Structure of the residual gauge orbit

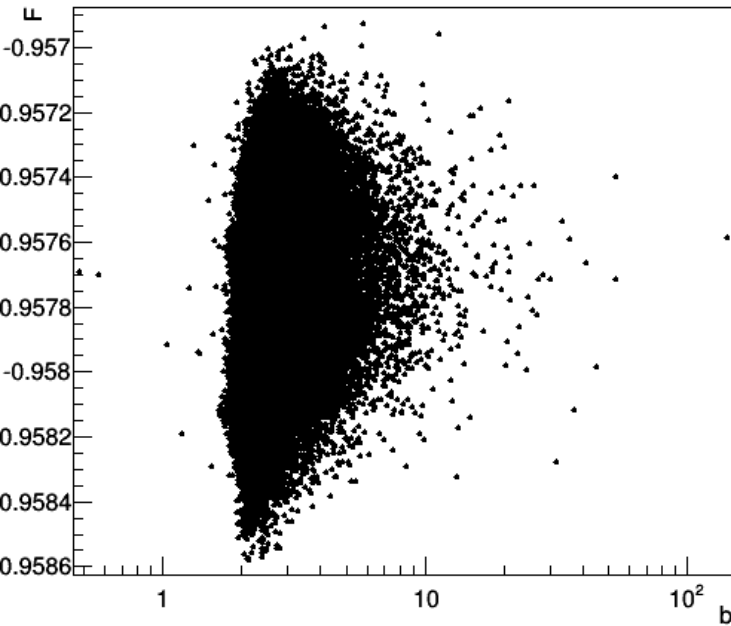
[Maas PRD 15]

- Projection of first Gribov region to B and F

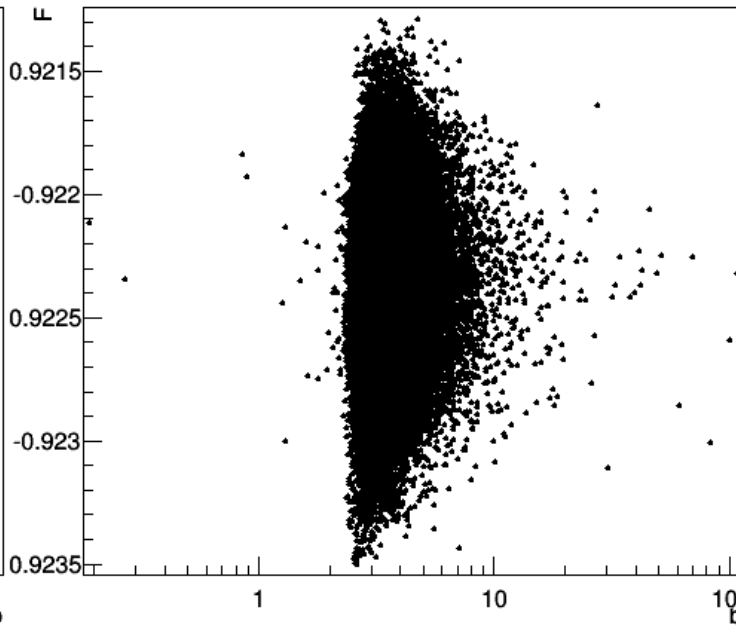
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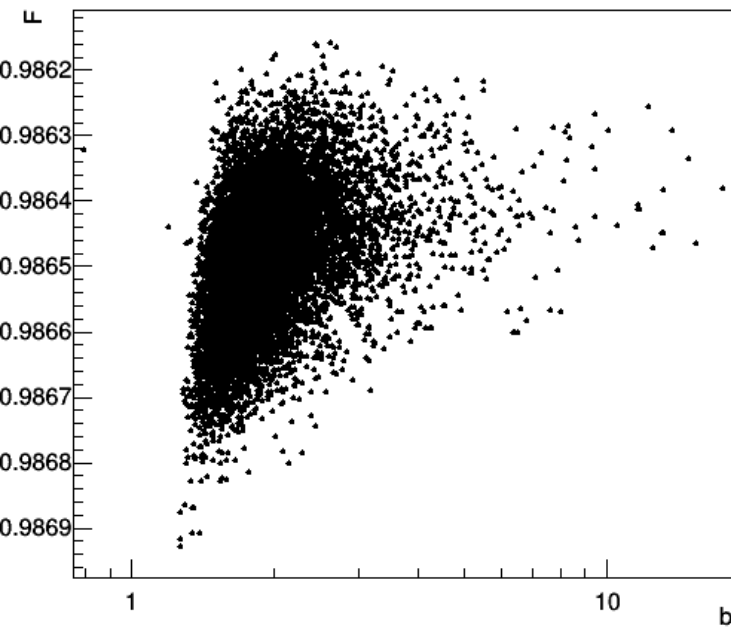
Gribov region for $a=0.092$ fm and $V=(4.4 \text{ fm})^3$



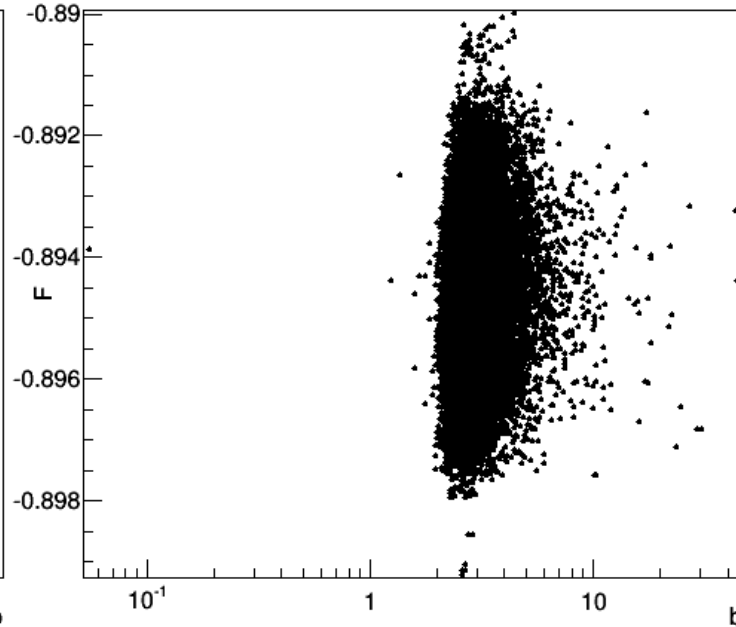
Gribov region for $a=0.17$ fm and $V=(7.9 \text{ fm})^3$



Gribov region for $a=0.030$ fm and $V=(1.8 \text{ fm})^3$



Gribov region for $a=0.22$ fm and $V=(4.4 \text{ fm})^3$

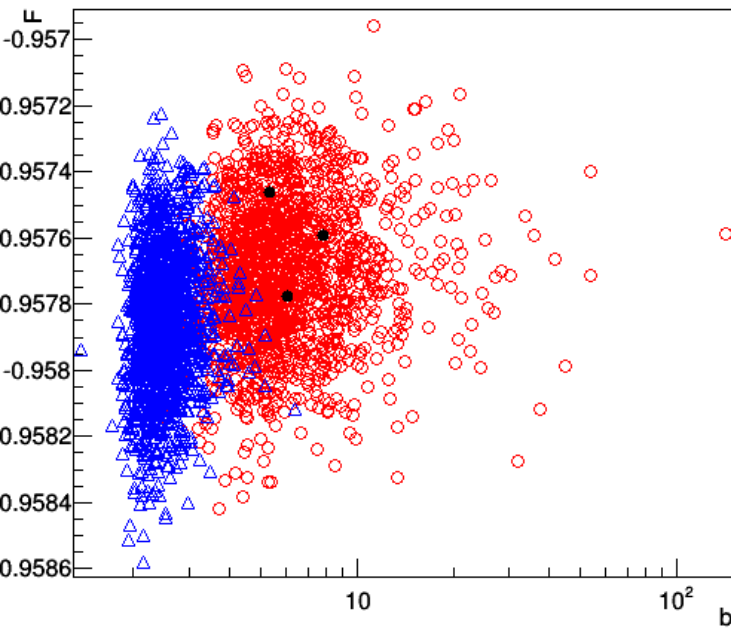


- Projection of first Gribov region to B and F
- Independent coordinates at sufficiently large volumes and fine lattices

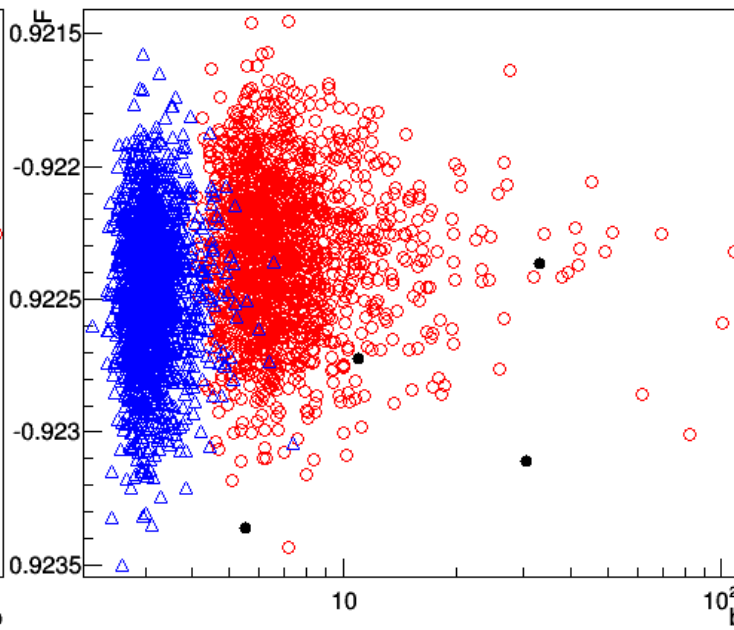
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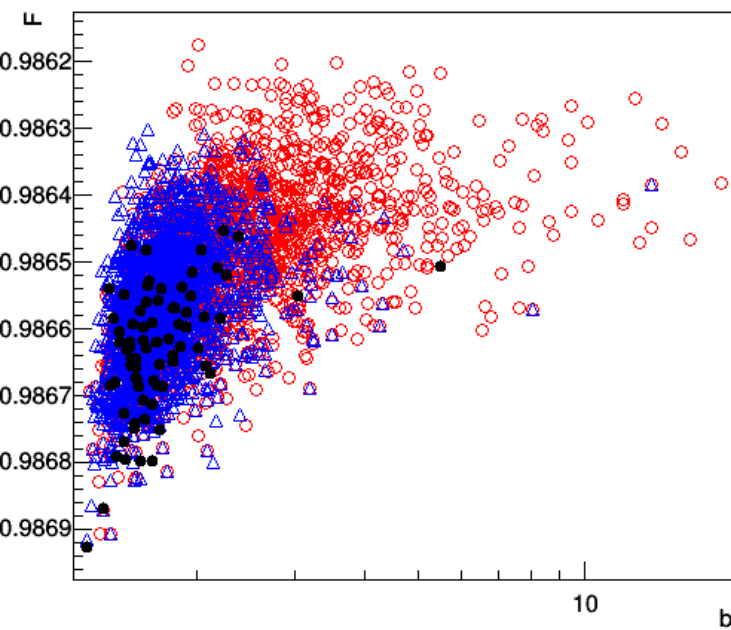
Gribov horizon and FMR for $a=0.092$ fm and $V=(4.4 \text{ fm})^3$



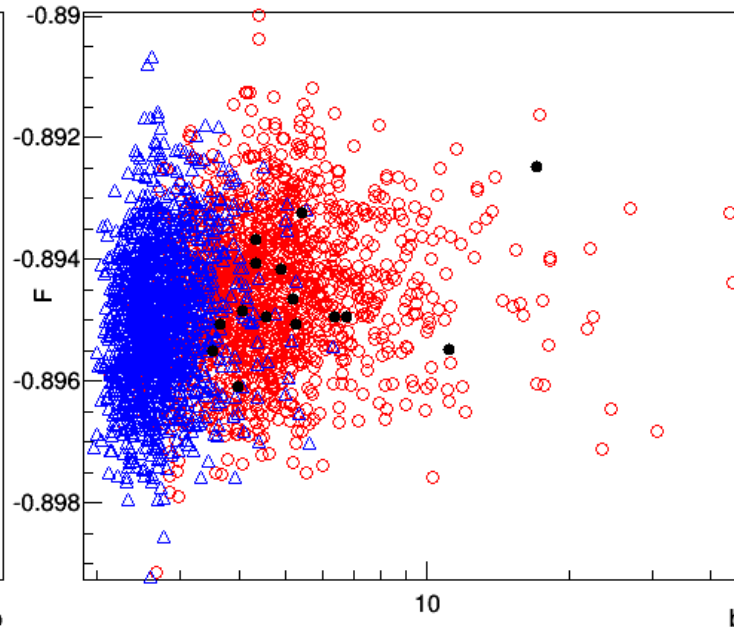
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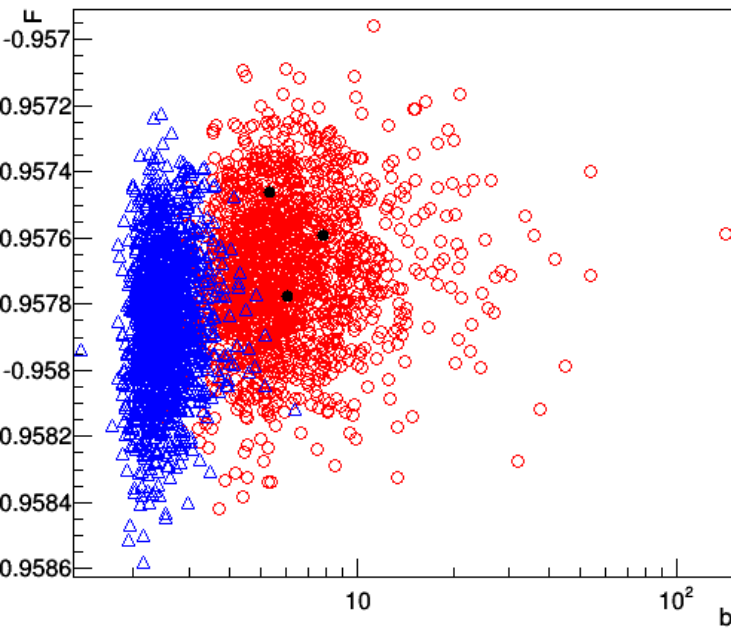
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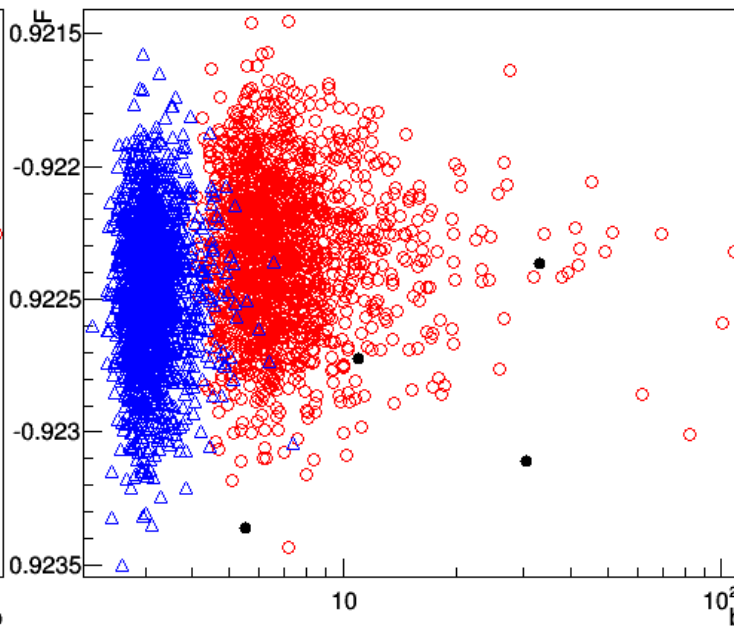
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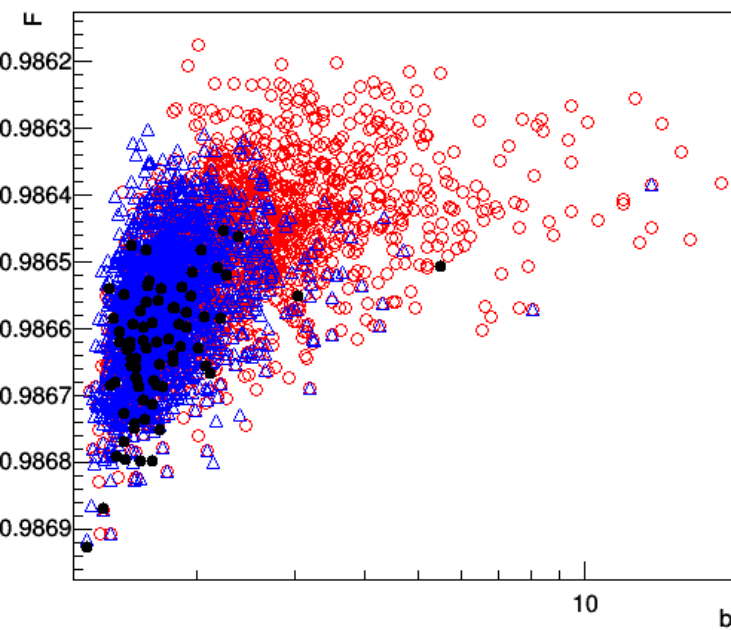
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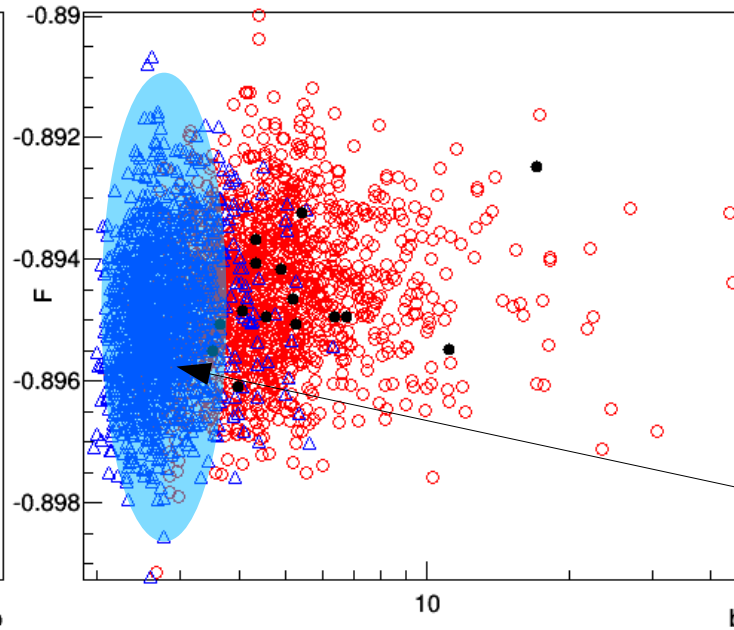
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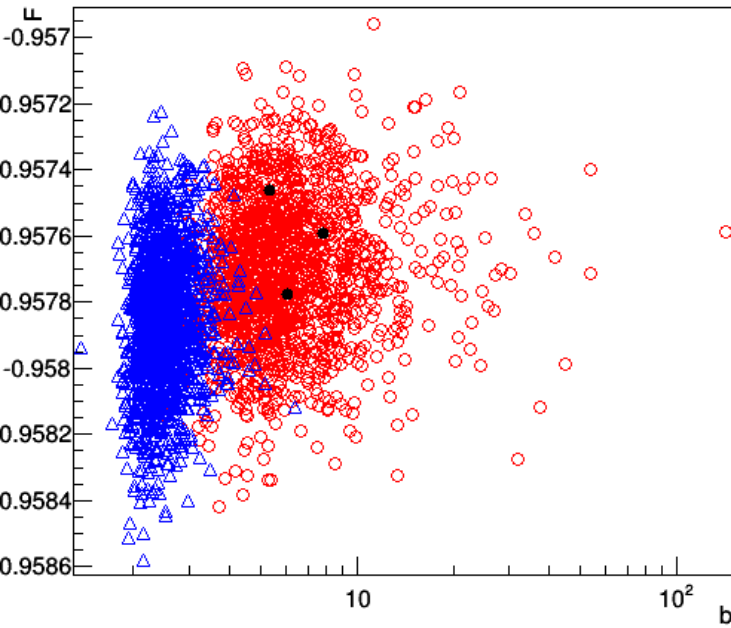


FMR

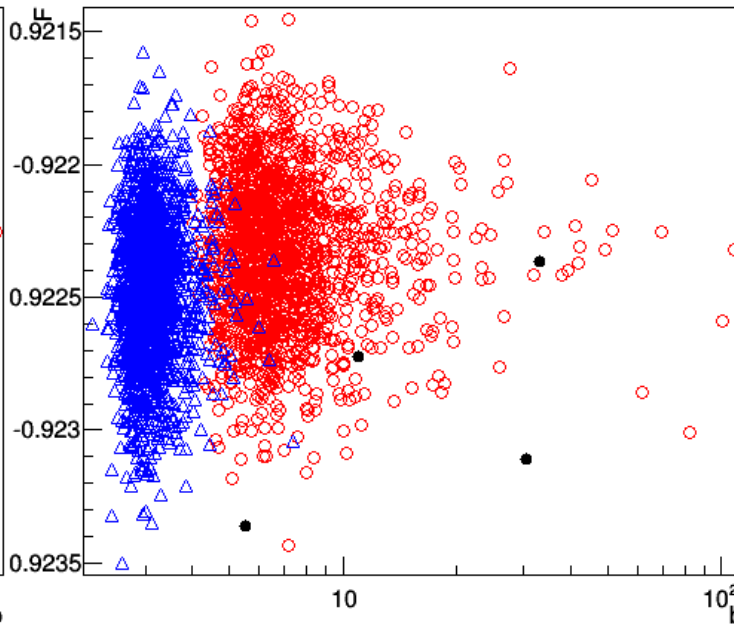
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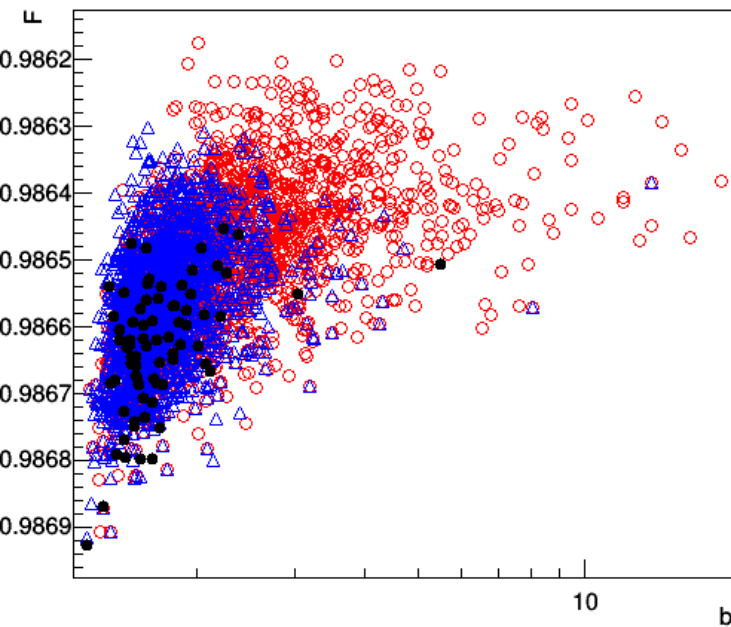
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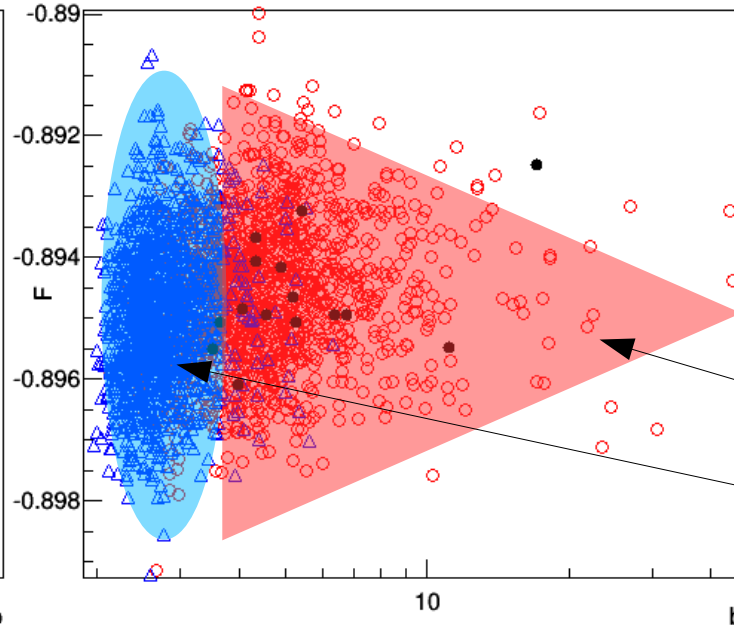
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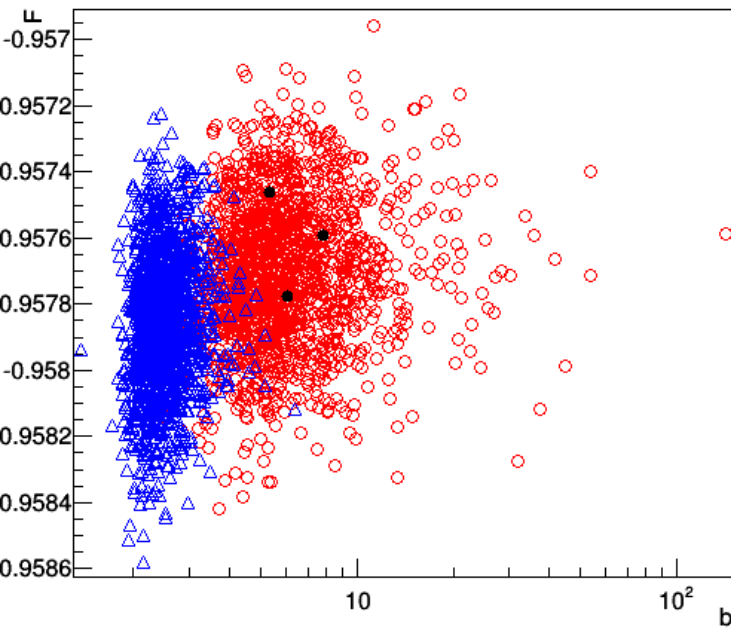
Horizon

FMR

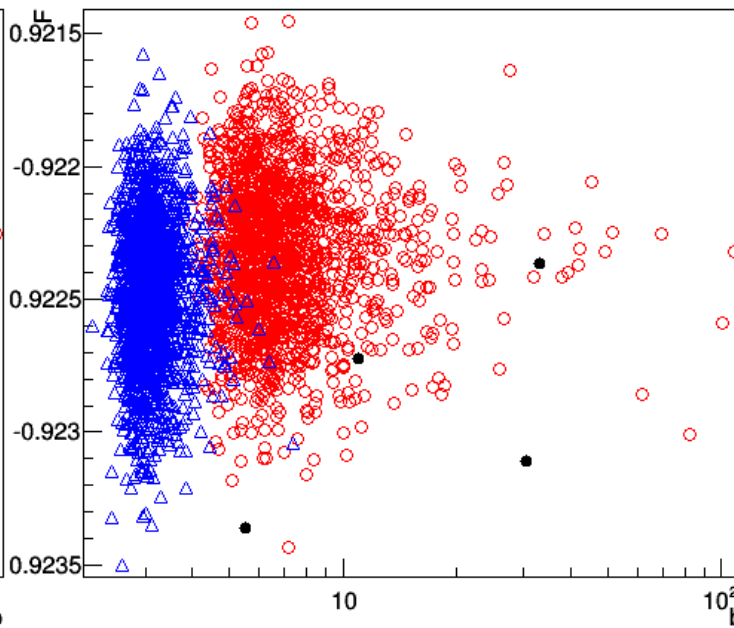
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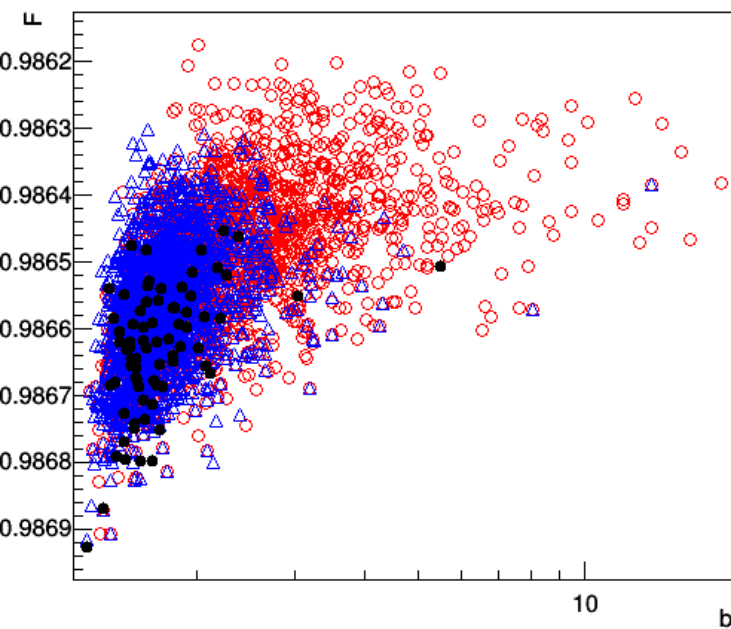
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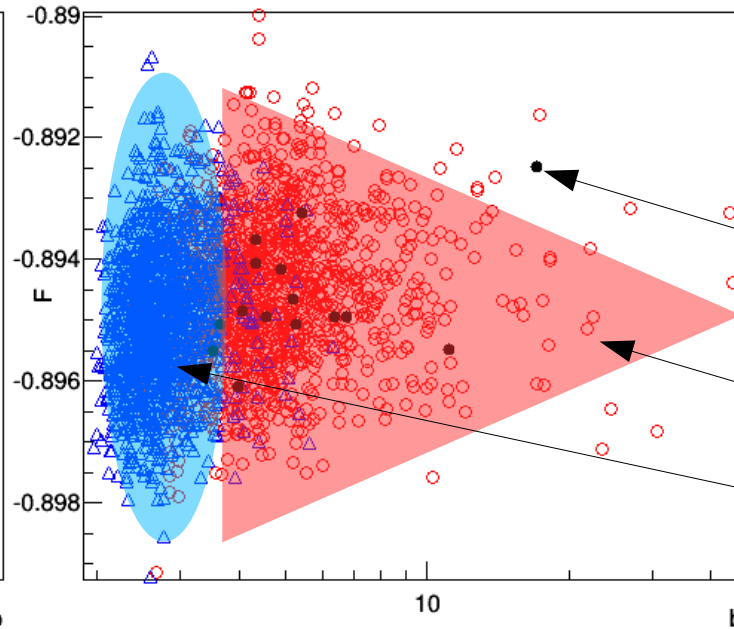
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Common
boundary
Horizon
FMR

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 - Maximum value: **$\max B$**

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- Average among still equal copies

Constructing a gauge

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 - Like in perturbative covariant gauges

Constructing a gauge

- Average over Gribov copies
 - Like in perturbative covariant gauges
- Constraints by Lagrange multiplier/gauge parameter

$$\lim_{\xi \rightarrow 0} \int DAc \bar{c} \exp\left(\int dx \left(-\frac{1}{4} F^2 - \bar{c} \partial D c - \frac{1}{2\xi} (\partial A)^2\right)\right) \times \\ \times \Theta(-\partial D) \exp\left(N + \frac{\lambda}{V} \int dx \partial_\mu^x \bar{c}(x) \int dy \partial_\mu^y c(y) + \frac{\omega}{V} \int dx A^2\right)$$

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- **Absolute Landau gauge** $\lambda = 0 \wedge \omega = -\infty$

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$$\lim_{\xi \rightarrow 0} \int DAc \bar{c} \exp\left(\int dx \left(-\frac{1}{4} F^2 - \bar{c} \partial D c - \frac{1}{2\xi} (\partial A)^2\right)\right) \times \\ \times \Theta(-\partial D) \exp\left(N + \frac{\lambda}{V} \int dx \partial_\mu^x \bar{c}(x) \int dy \partial_\mu^y c(y) + \frac{\omega}{V} \int dx A^2\right)$$

- Normalization N if Gribov copies/orbit differ
- Minimal Landau gauge $\lambda = \omega = 0$
- **Absolute Landau gauge** $\lambda = 0 \wedge \omega = -\infty$
- **Inverse Landau gauge** $\lambda = 0 \wedge \omega = \infty$
- **maxB gauge** $\lambda = -\infty \wedge \omega = 0$

Constructing a gauge

- Average over Gribov copies
 - Like in perturbative covariant gauges
- Constraints by Lagrange multiplier/gauge parameter

$$\lim_{\xi \rightarrow 0} \int DAc \bar{c} \exp\left(\int dx \left(-\frac{1}{4} F^2 - \bar{c} \partial D c - \frac{1}{2\xi} (\partial A)^2\right)\right) \times \\ \times \Theta(-\partial D) \exp\left(N + \frac{\lambda}{V} \int dx \partial_\mu^x \bar{c}(x) \int dy \partial_\mu^y c(y) + \frac{\omega}{V} \int dx A^2\right)$$

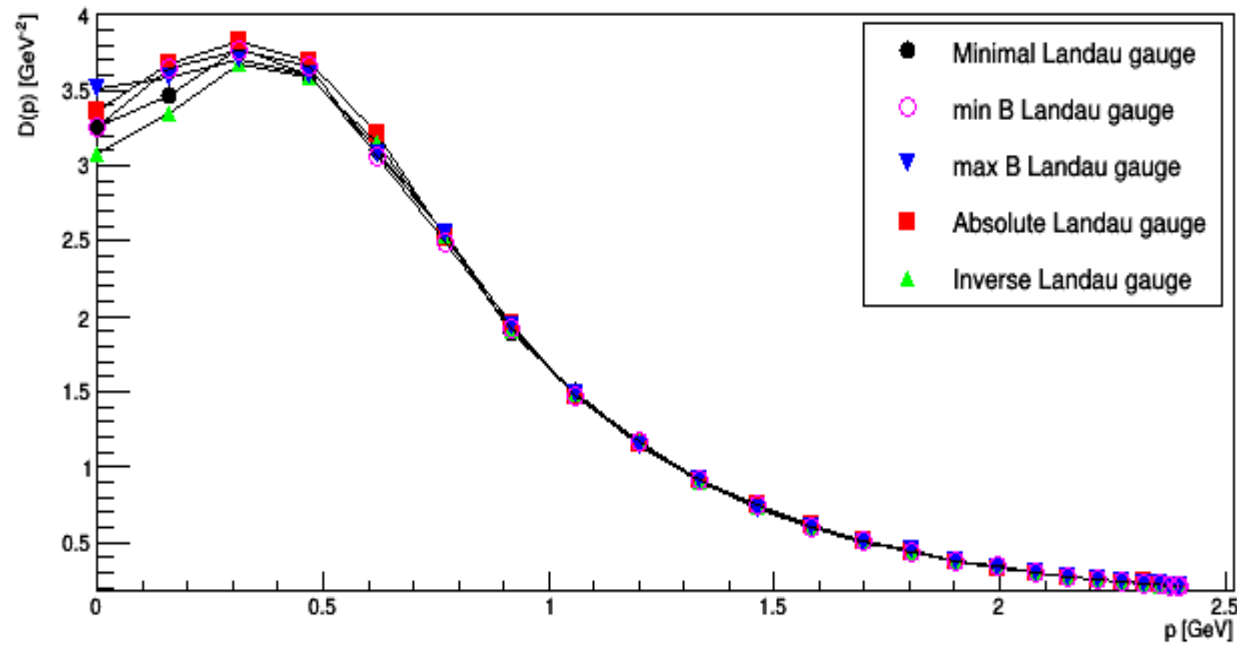
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Gluon propagator

[Maas PoS 12, unpublished]

[3d $V=(7.9 \text{ fm})^3$, $a=0.17 \text{ fm}$]

Extreme gluon propagators

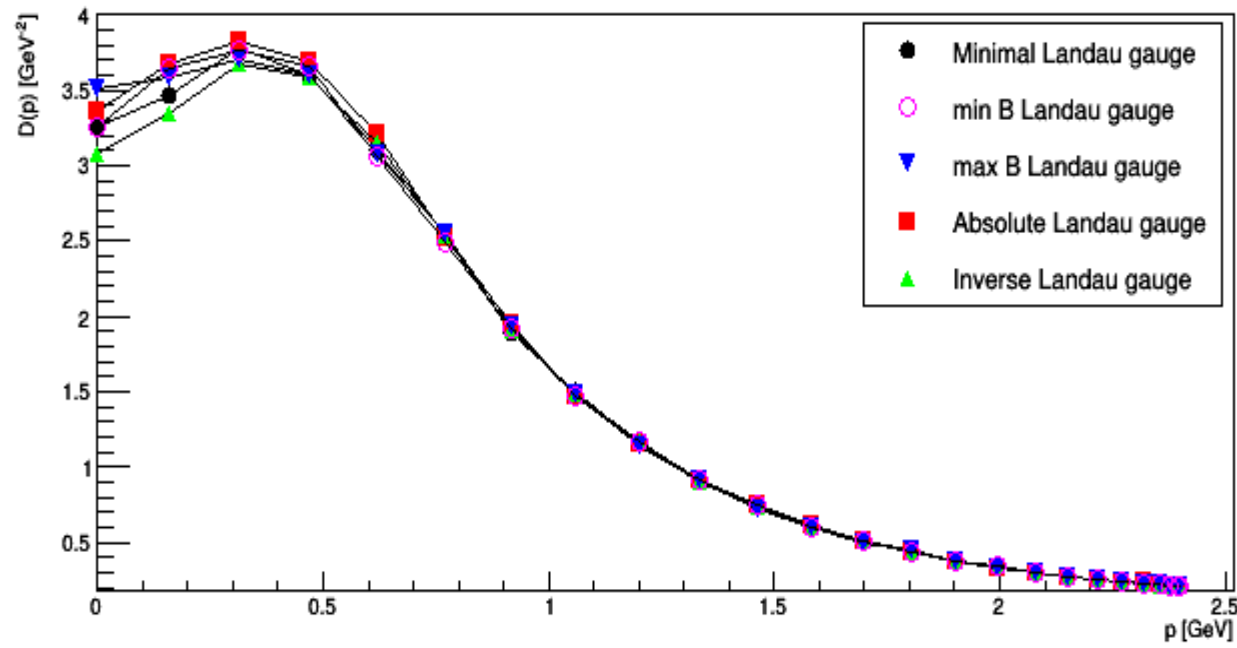


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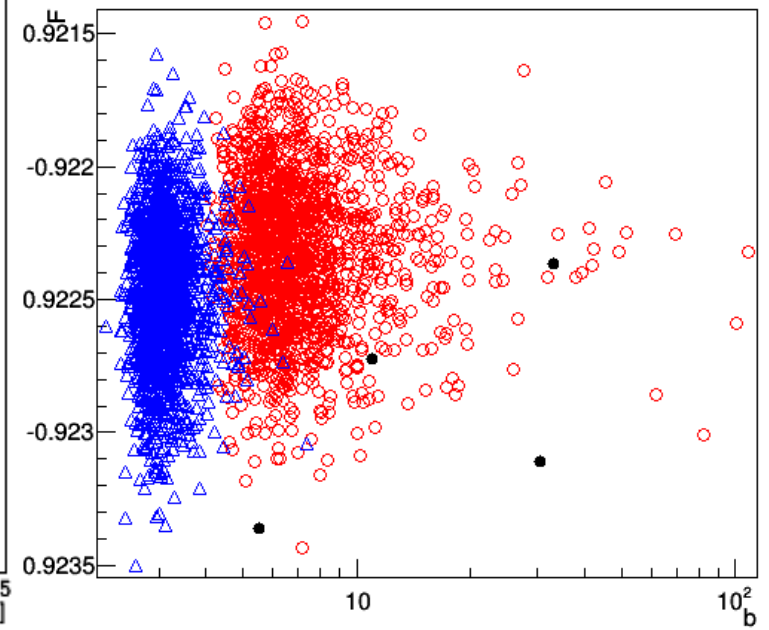
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Extreme gluon propagators



Gribov horizon and FMR for $a=0.18 \text{ fm}$ and $V=(7.9 \text{ fm})^3$

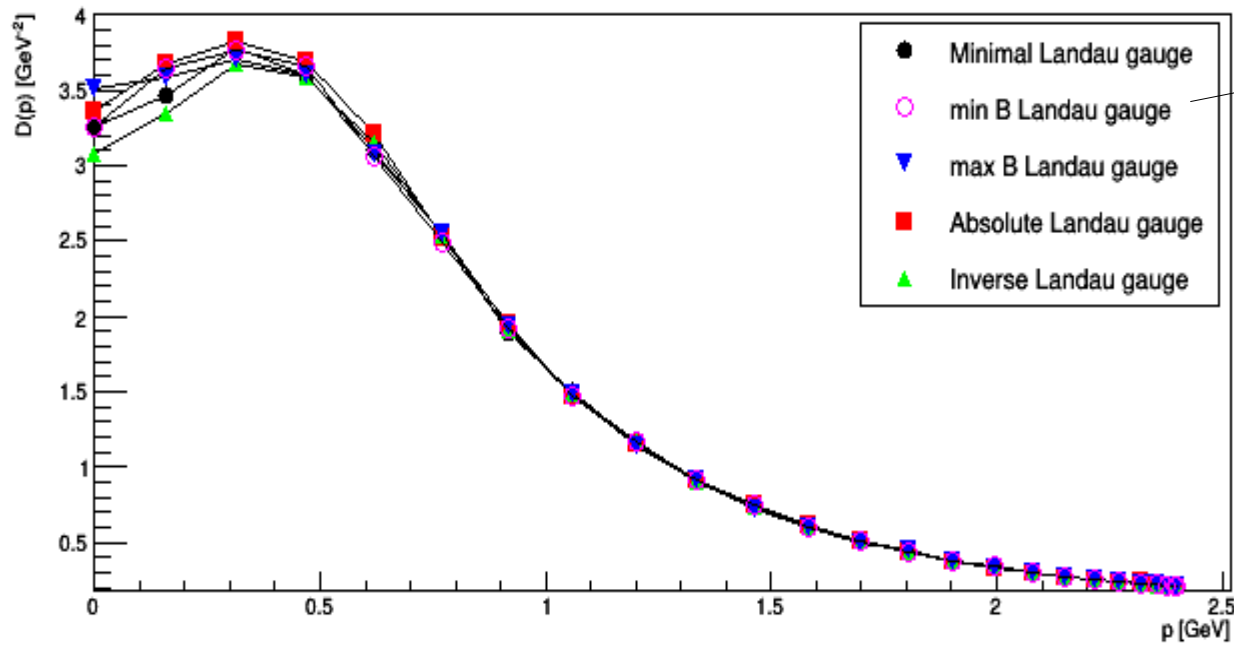


Gluon propagator

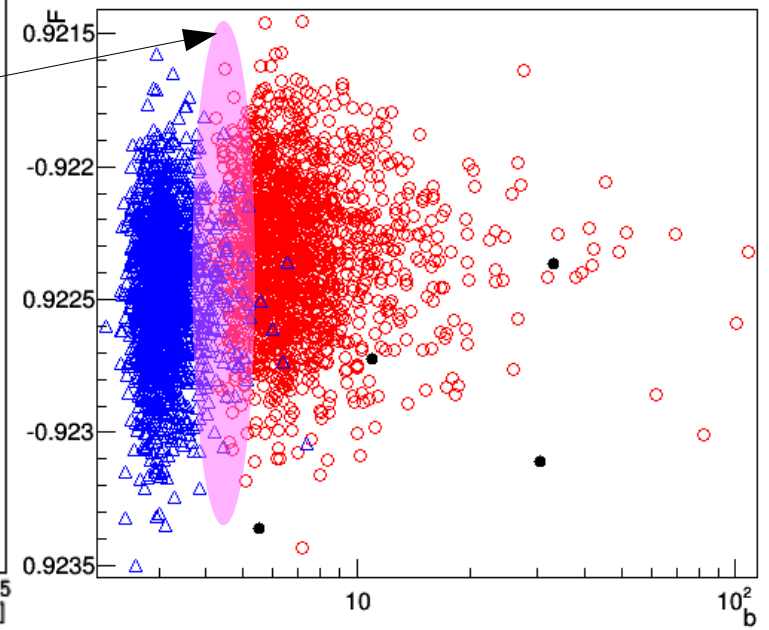
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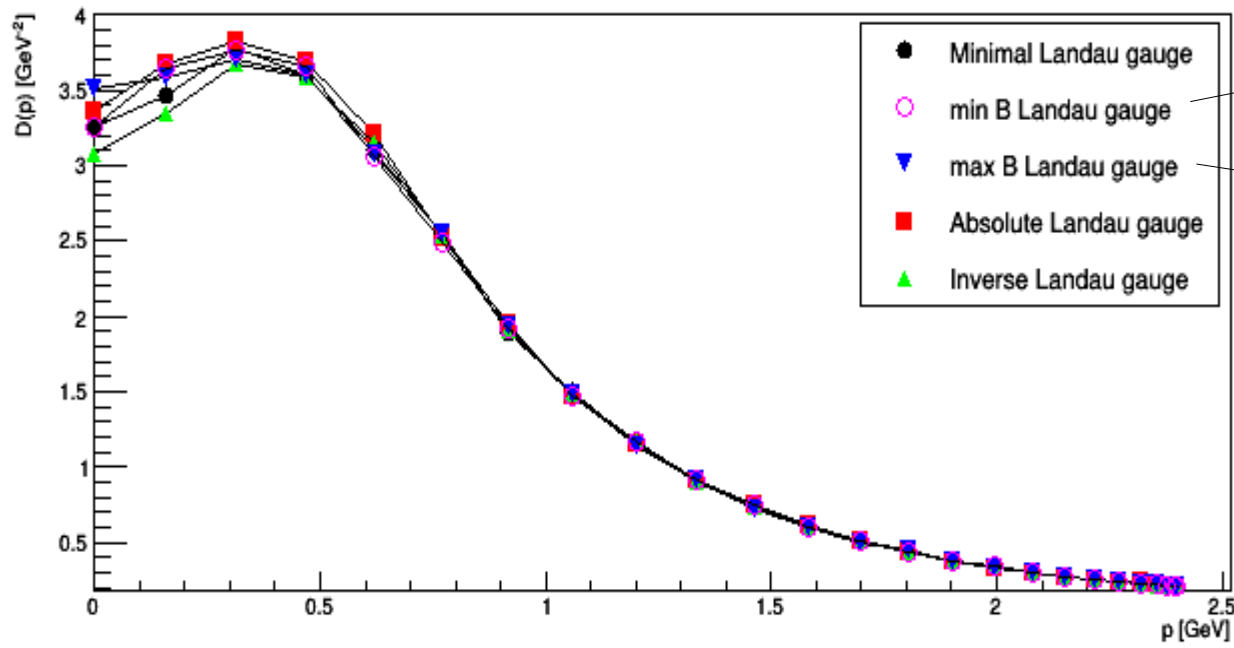


Gluon propagator

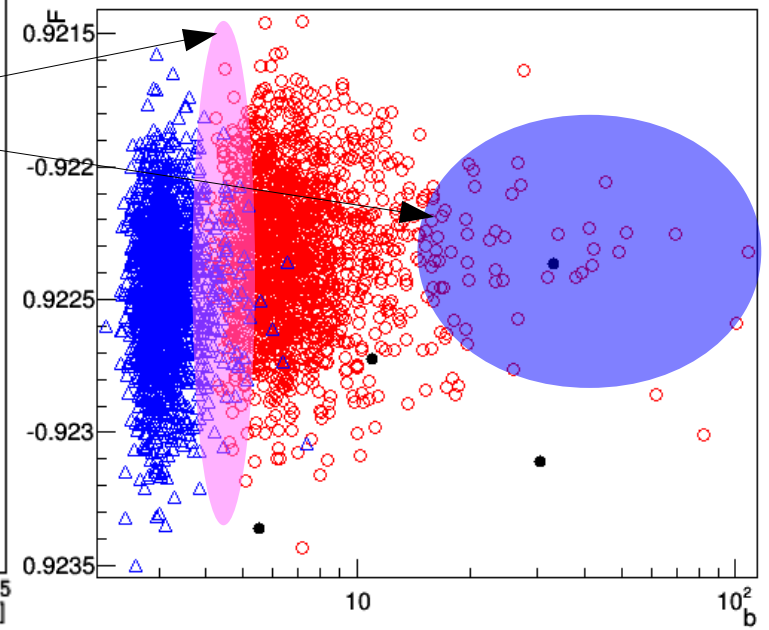
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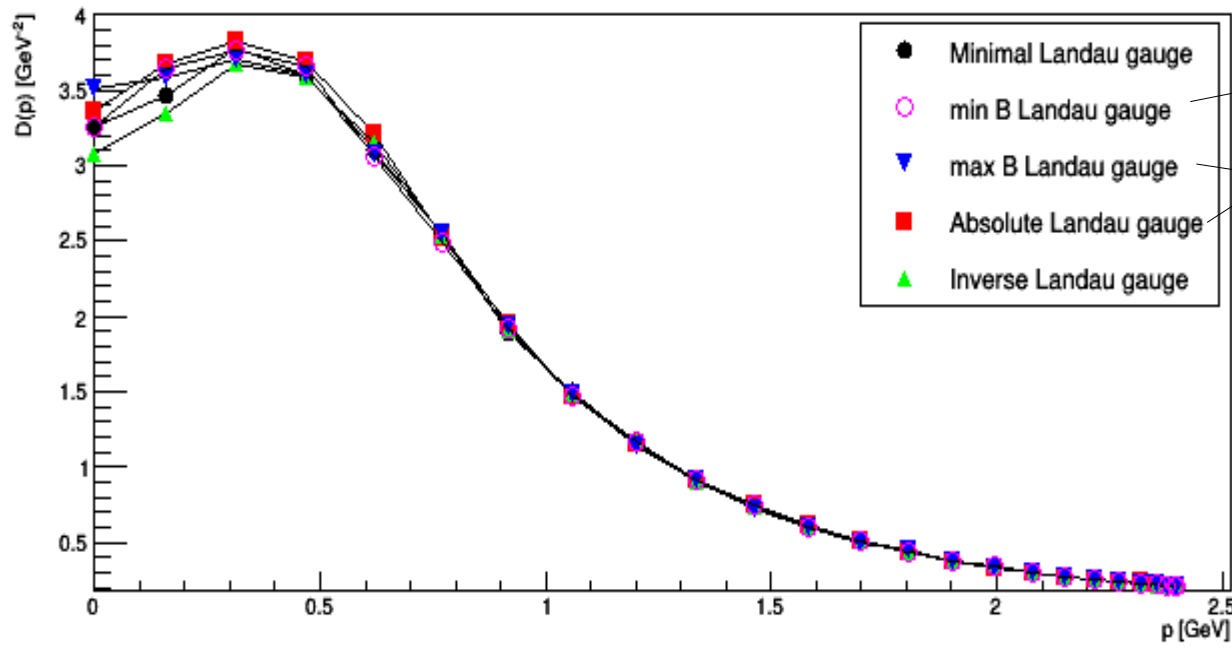


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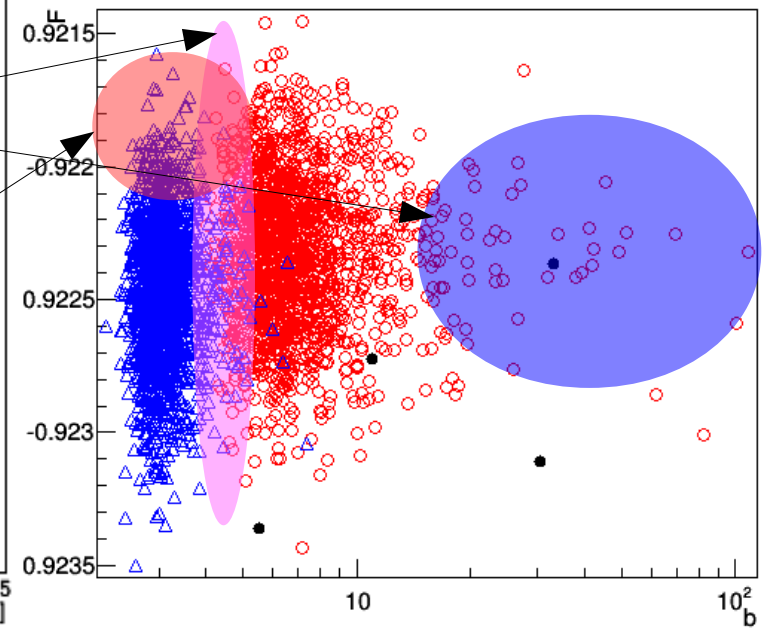
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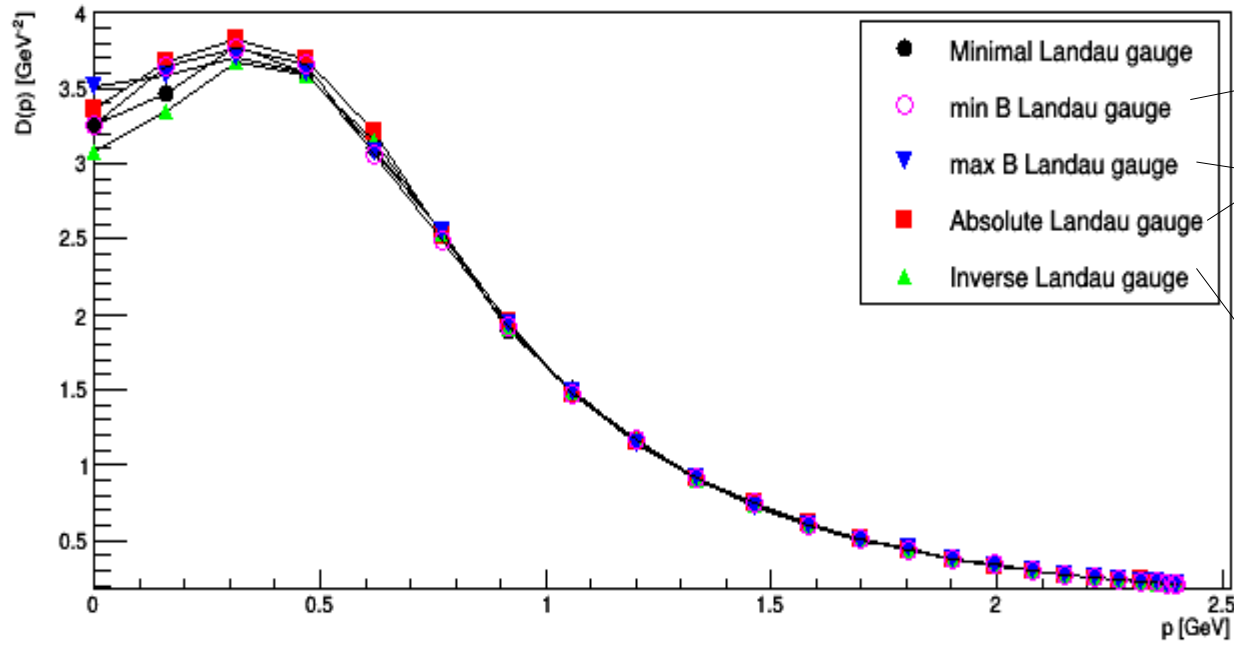


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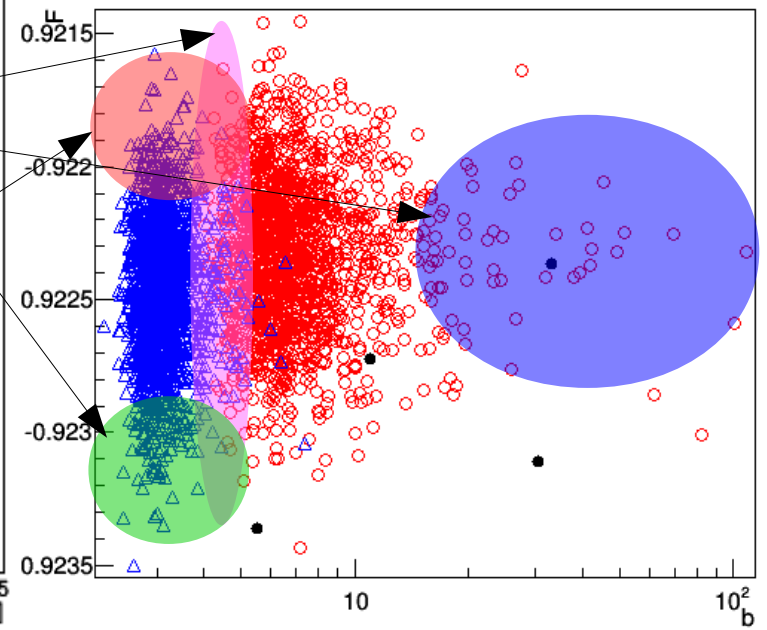
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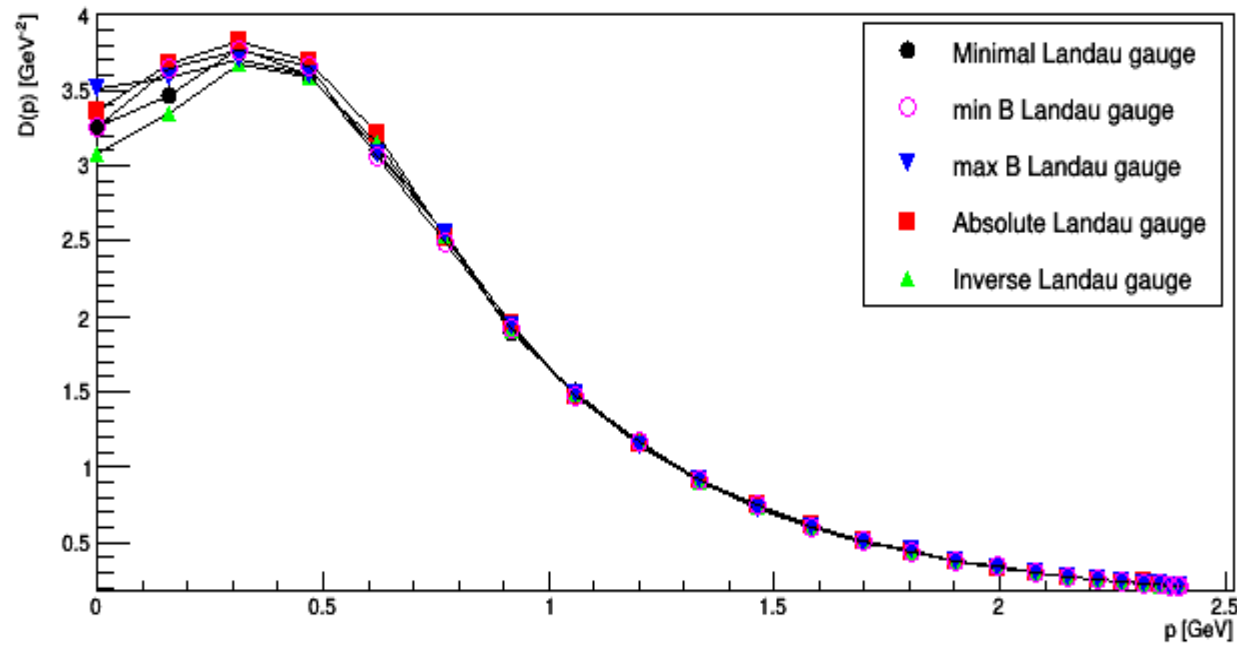


Gluon propagator

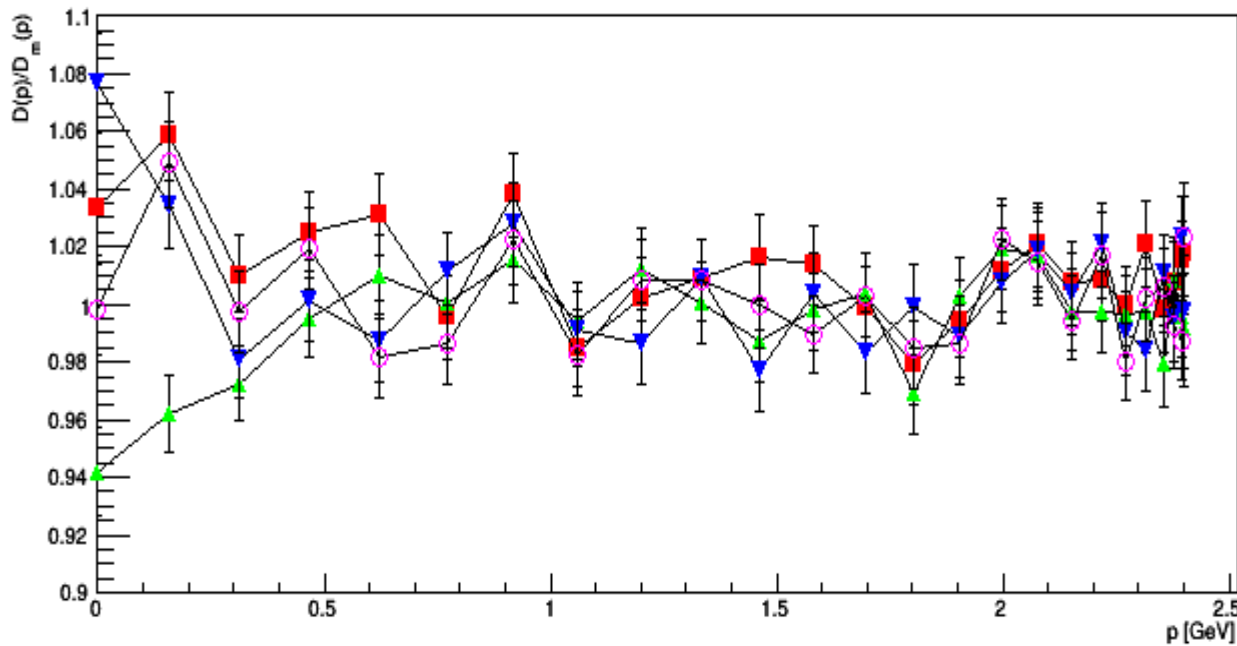
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Ratios to minimal Landau gauge

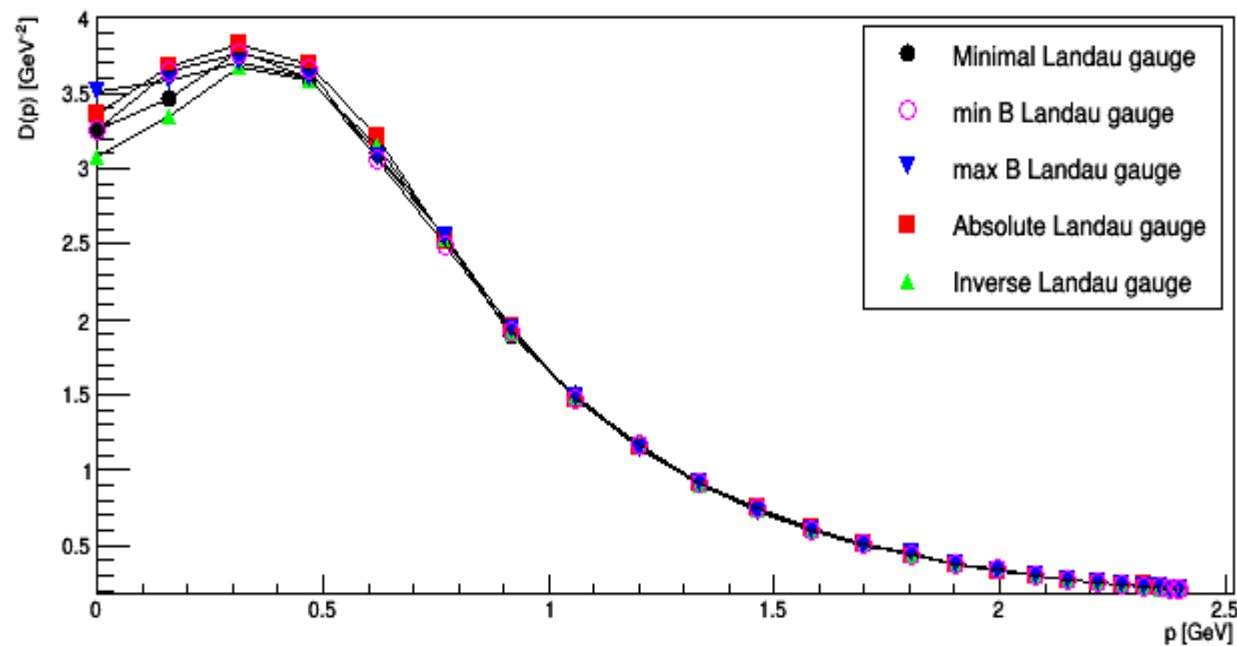


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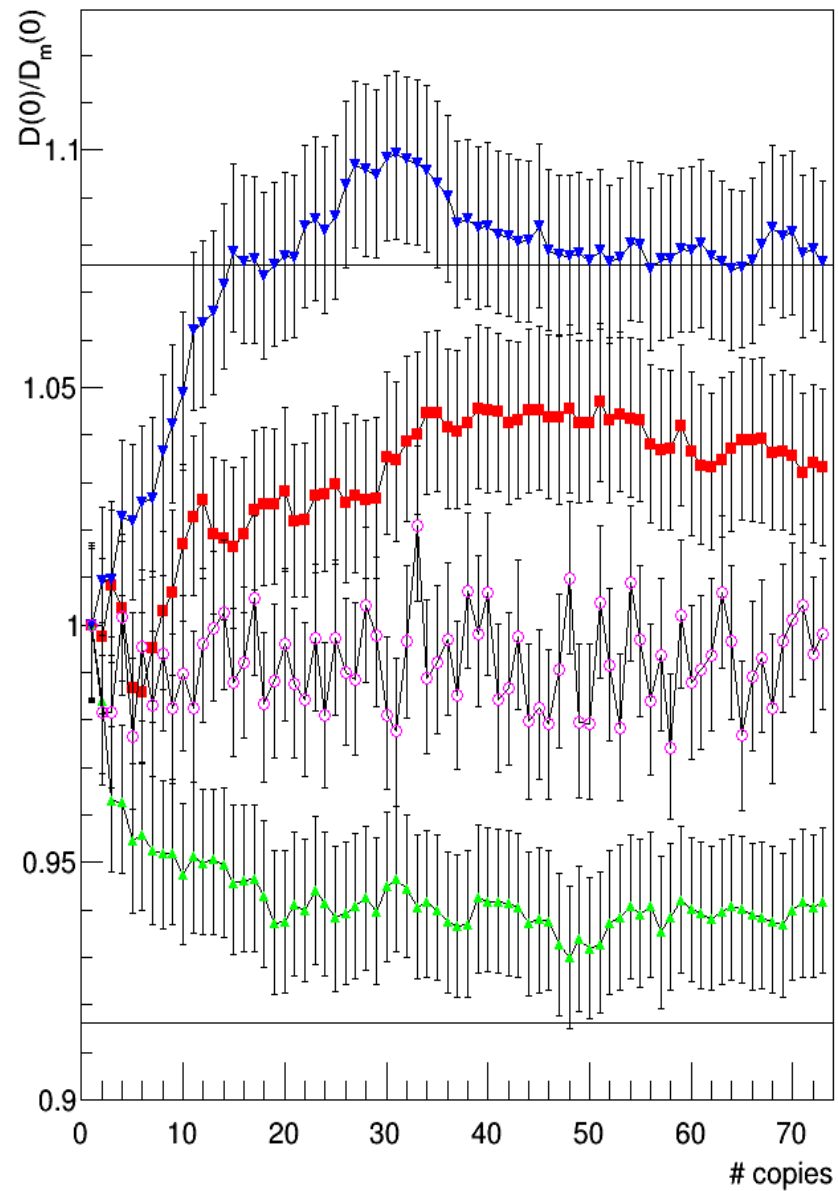
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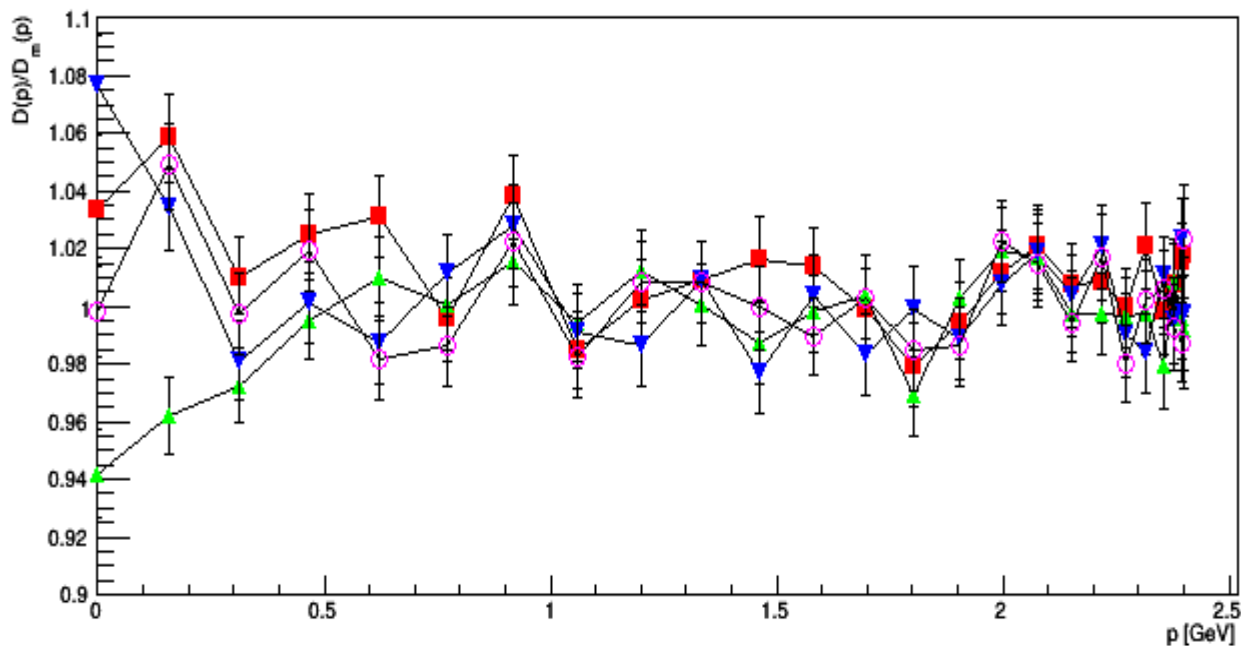
Extreme gluon propagators



Evolution with # of copies



Ratios to minimal Landau gauge

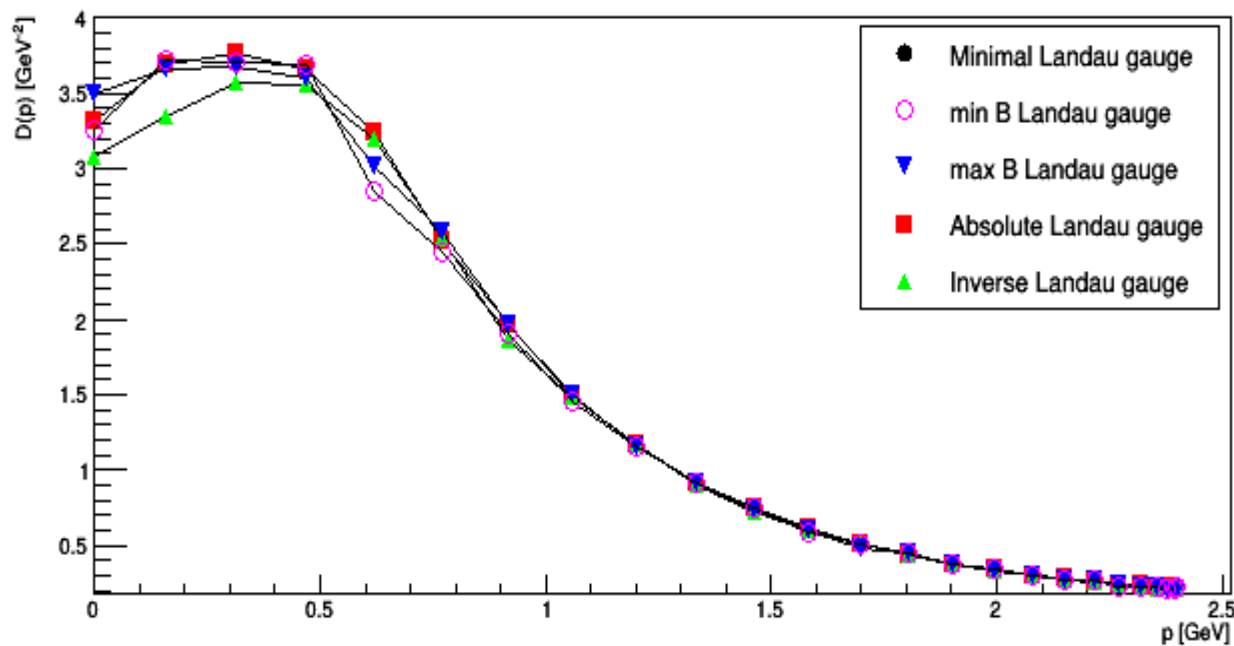


Glue propagator - extrapolated

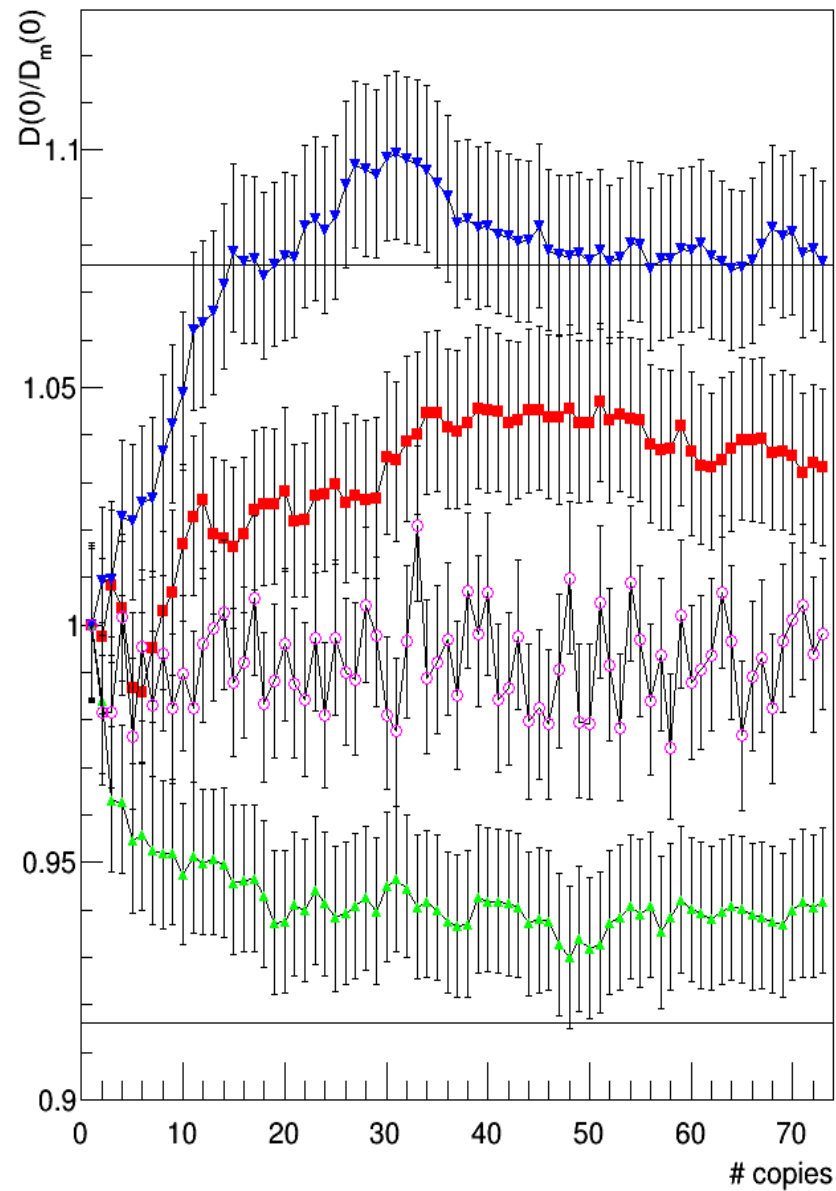
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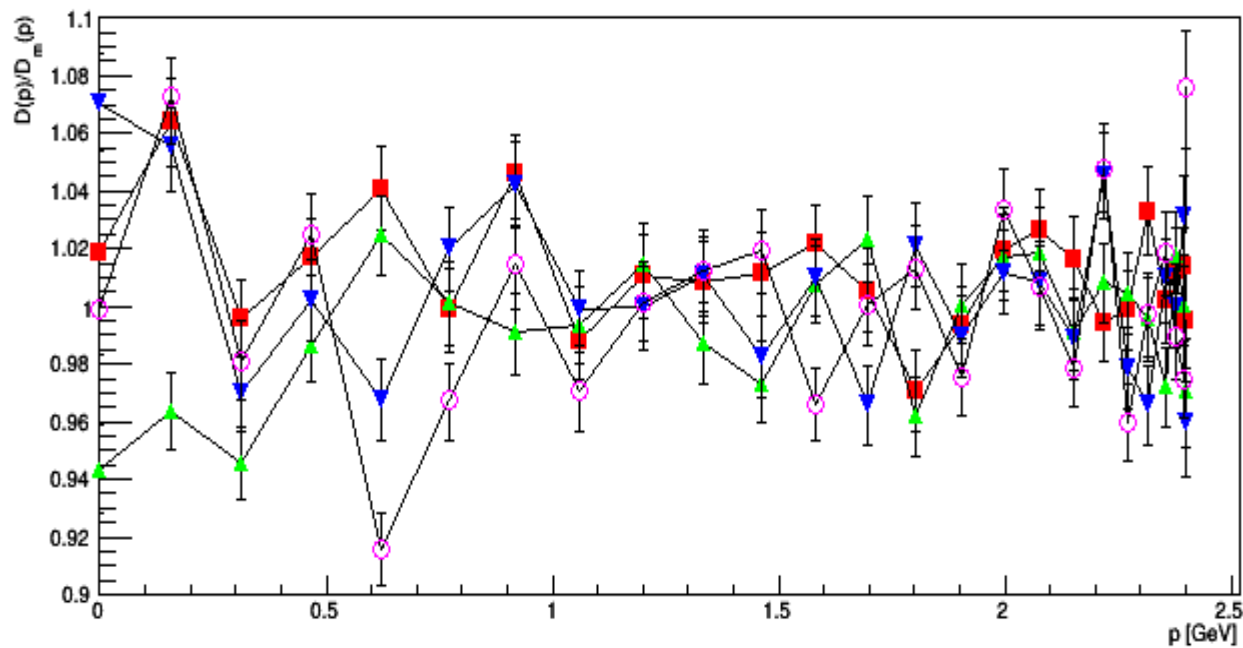
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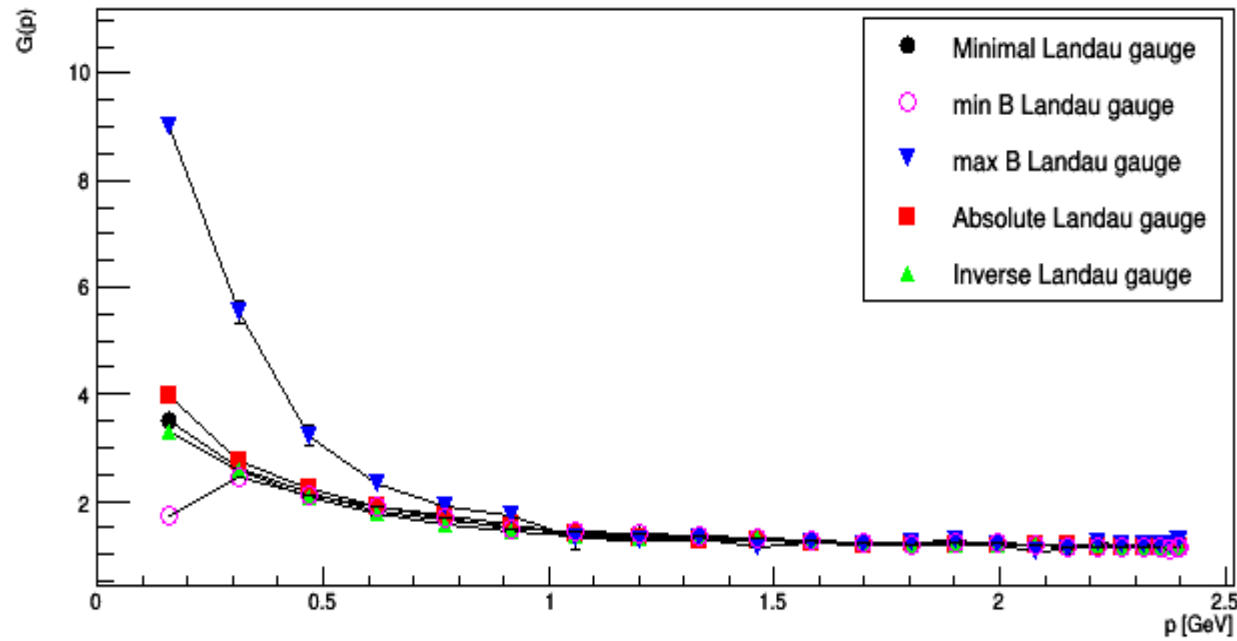


Ghost propagator

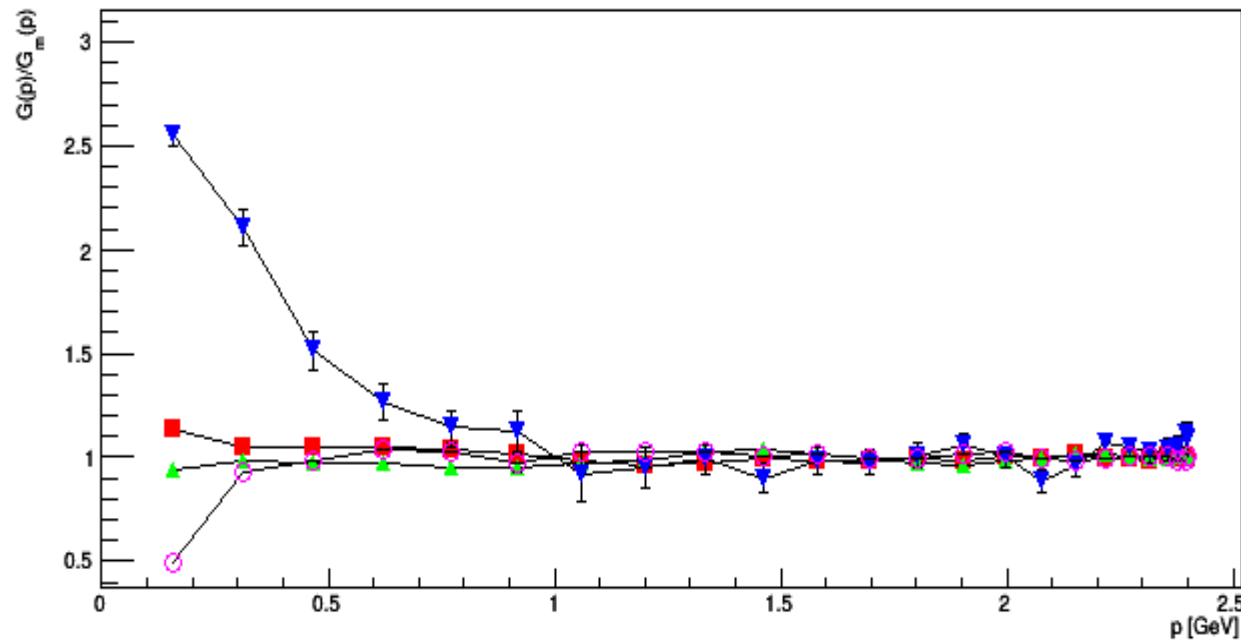
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Extreme ghost dressing function



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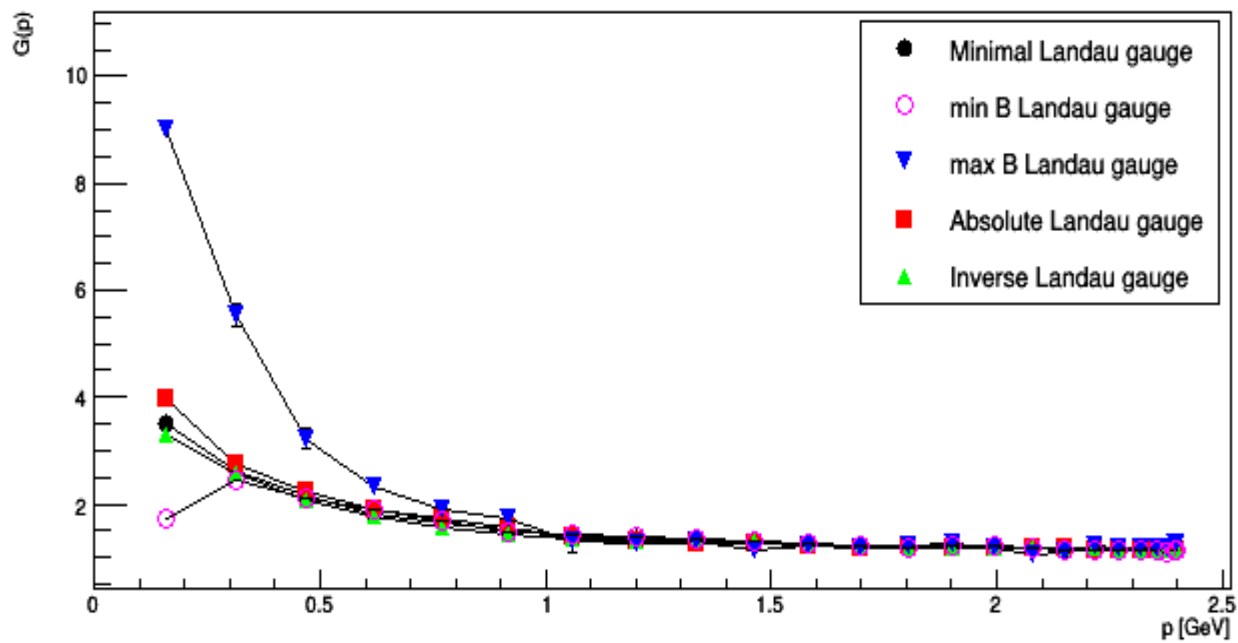


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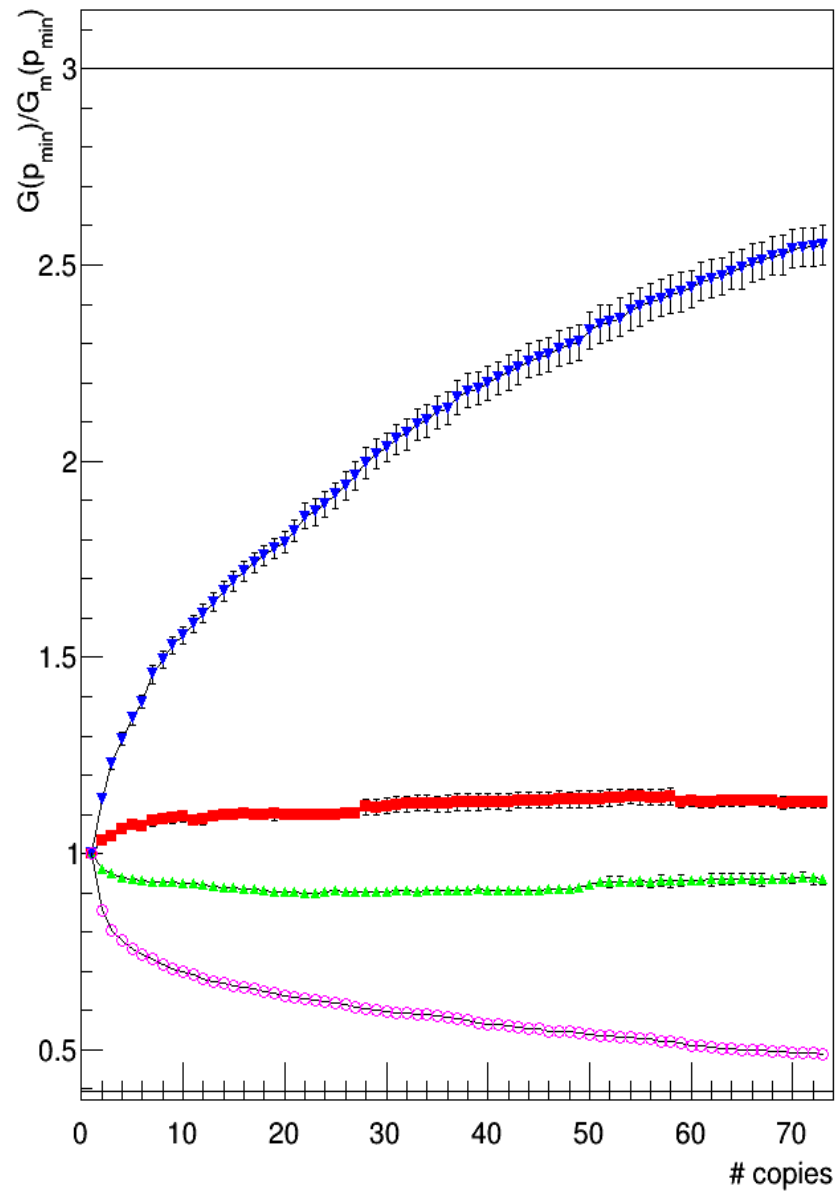
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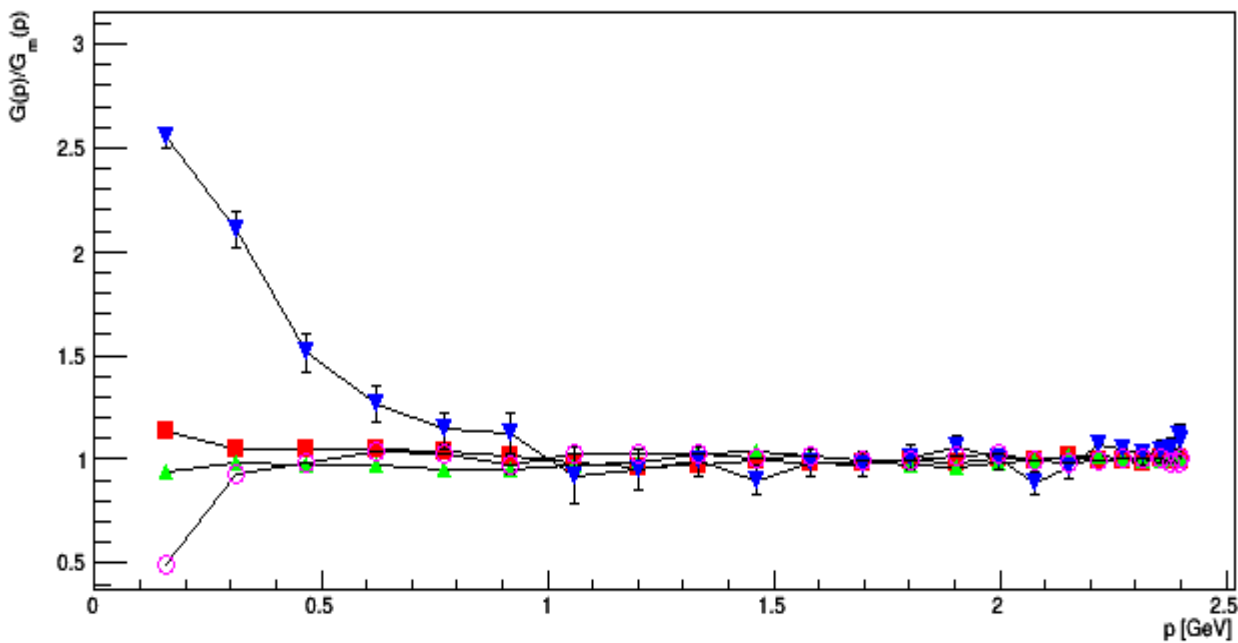
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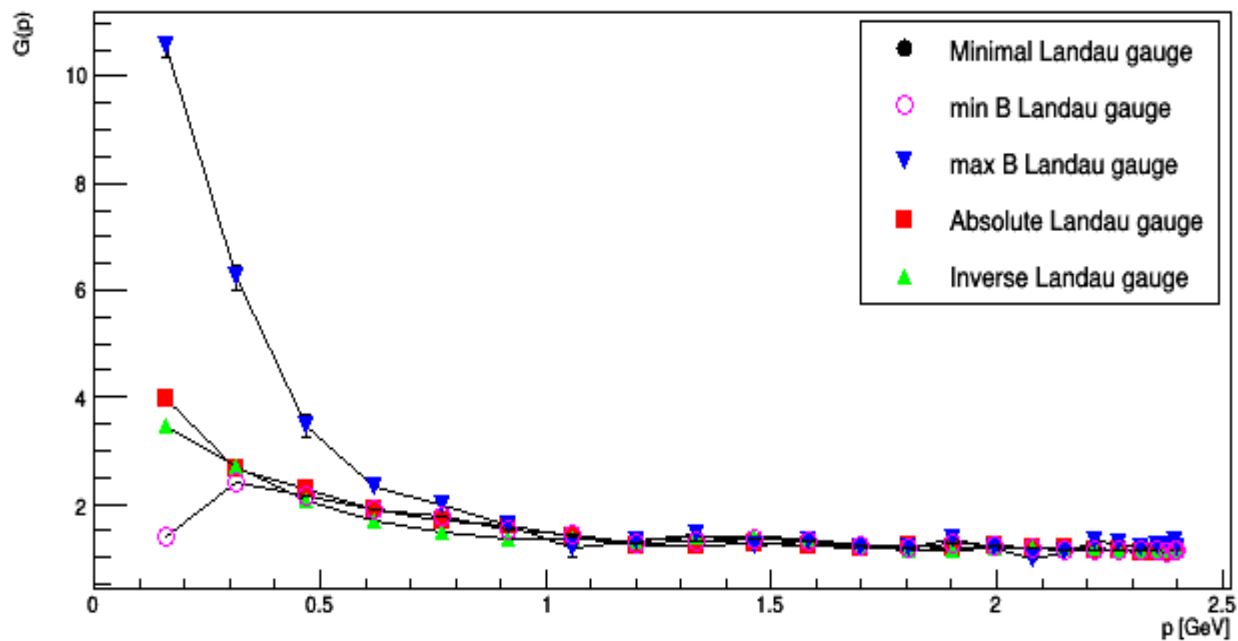


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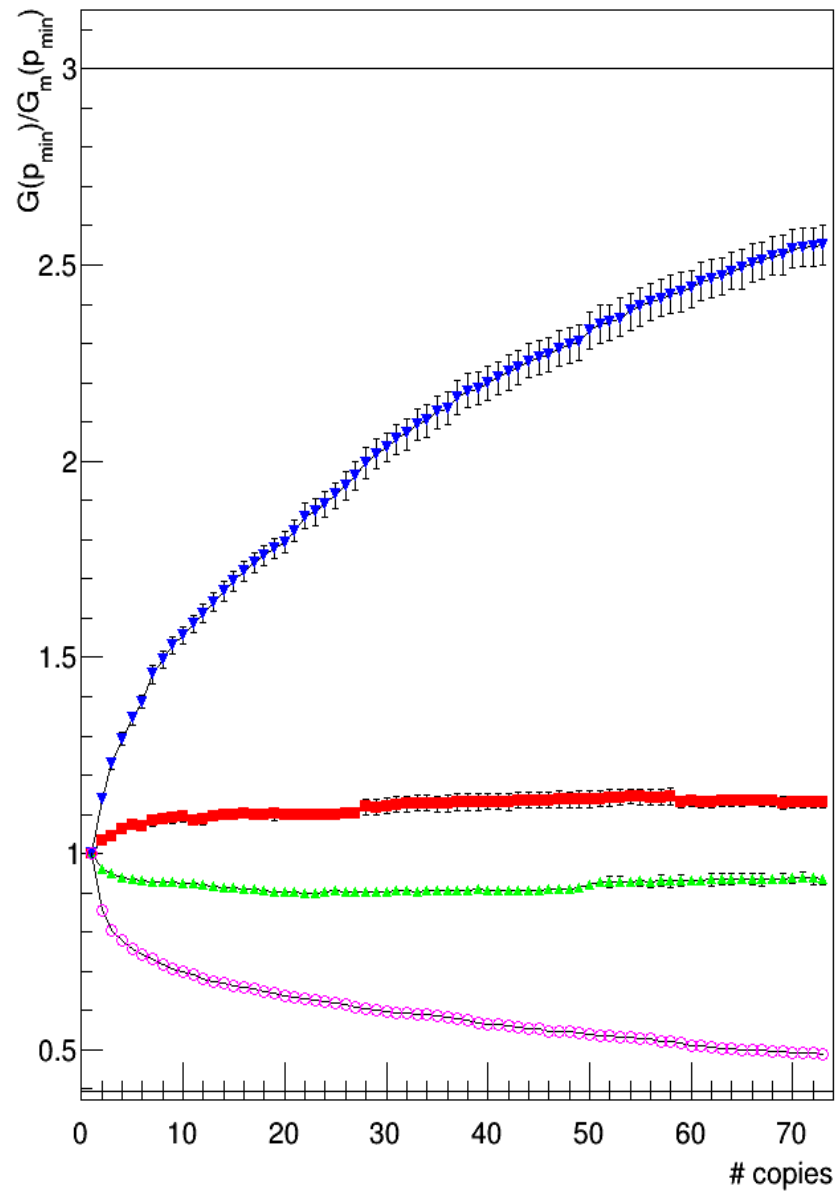
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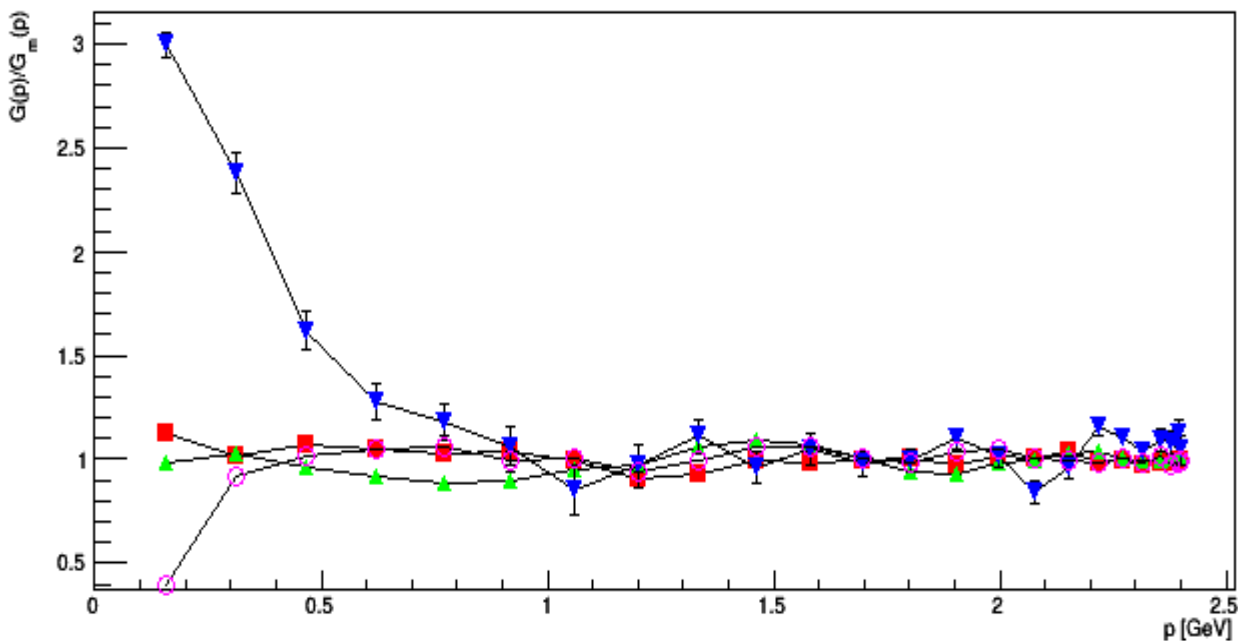
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Evolution with # of copies



Ratios to minimal Landau gauge

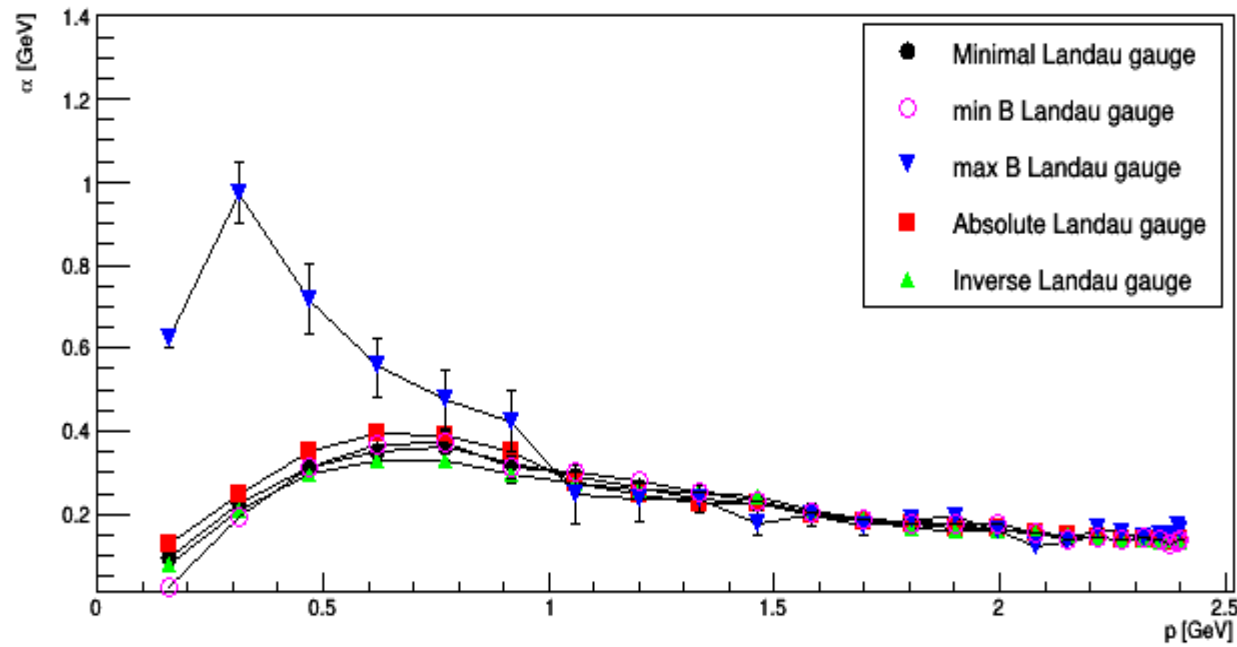


Running coupling

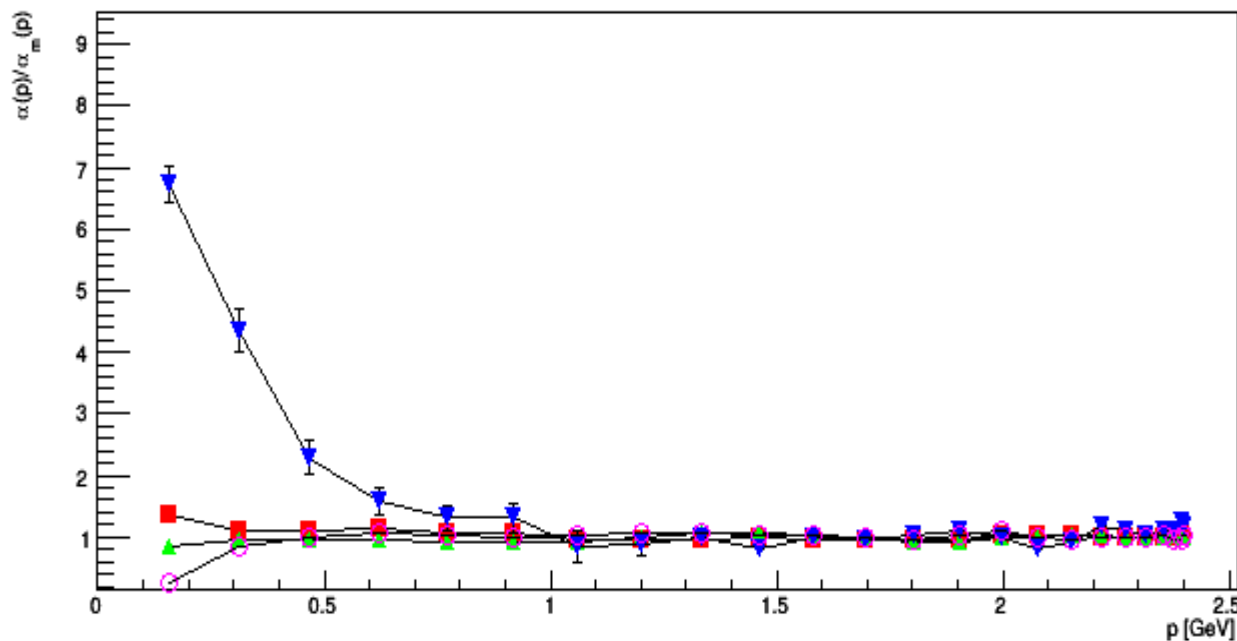
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Extreme α



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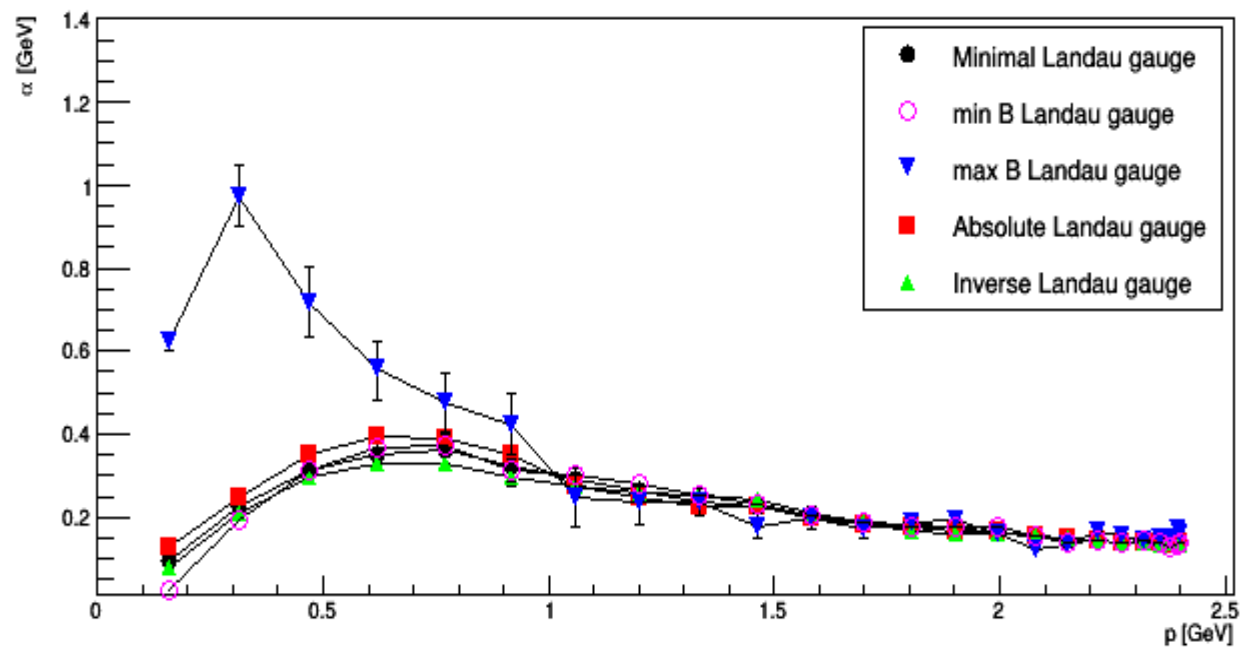


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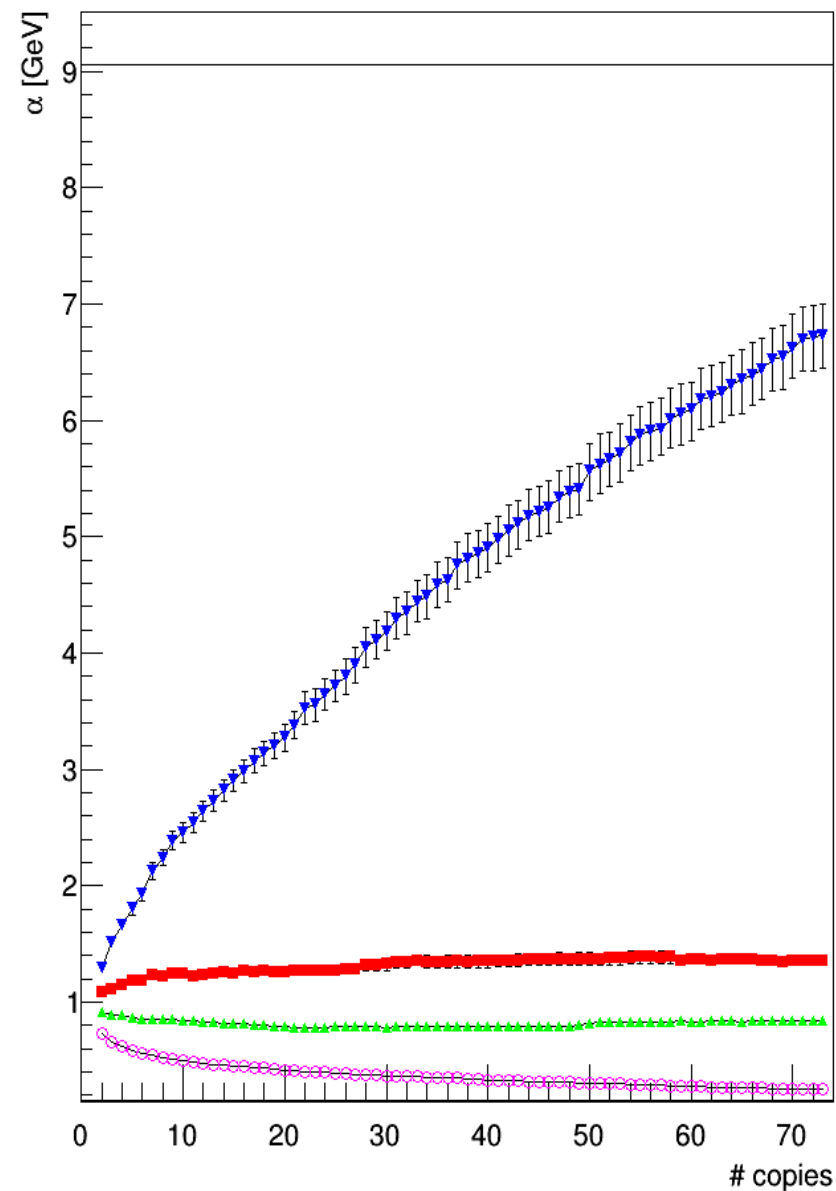
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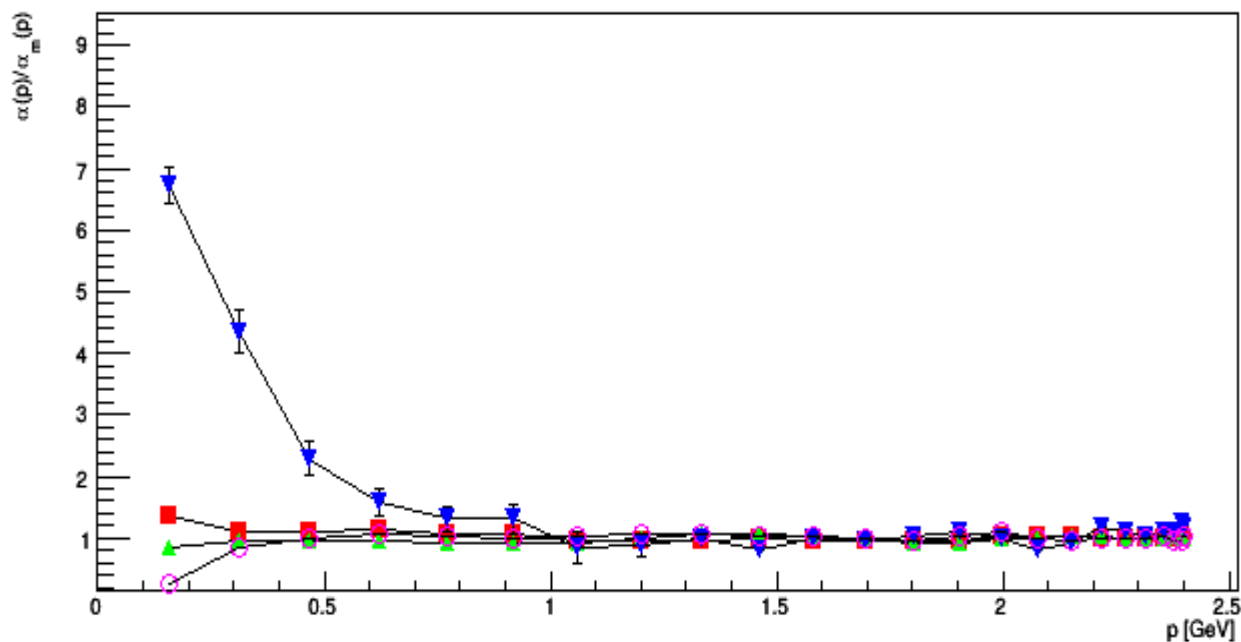
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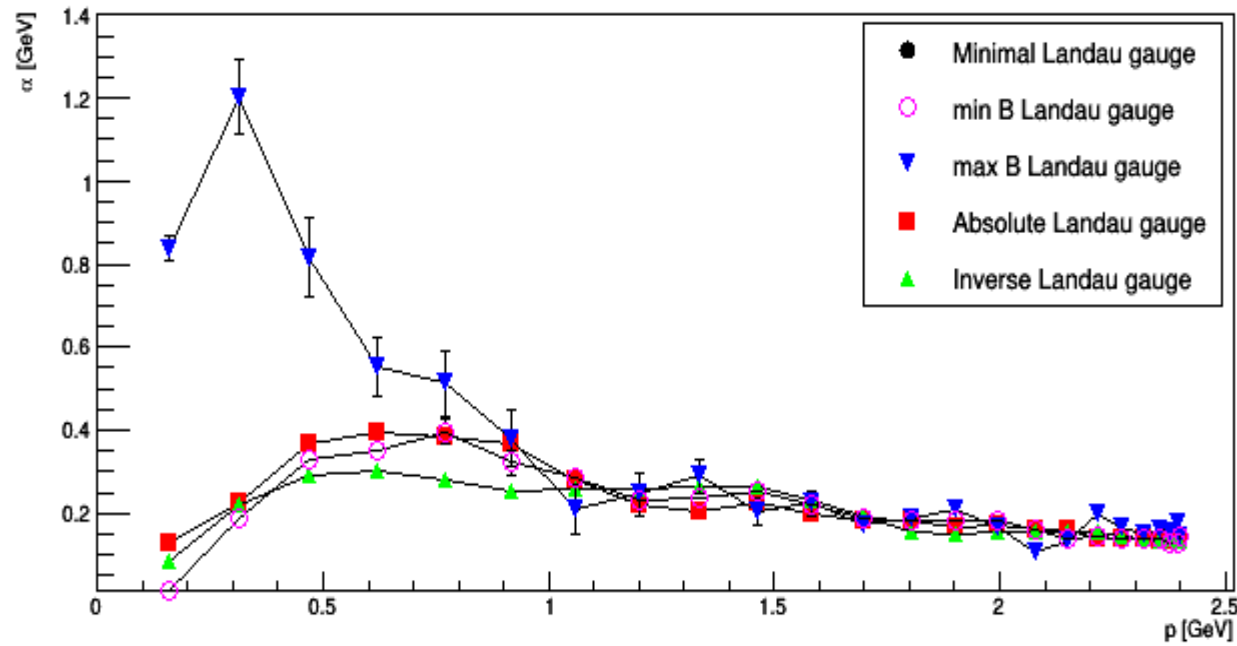
Evolution with # of copies



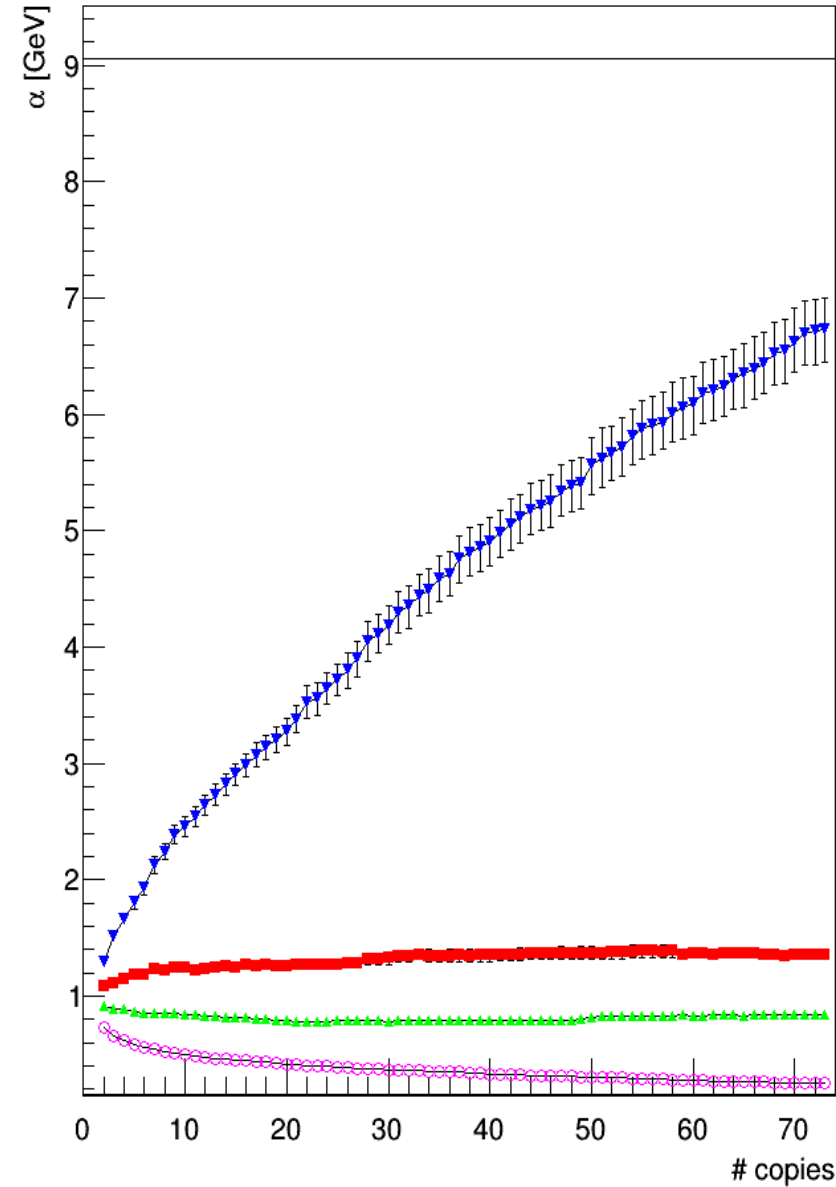
Ratio to minimal Landau gauge



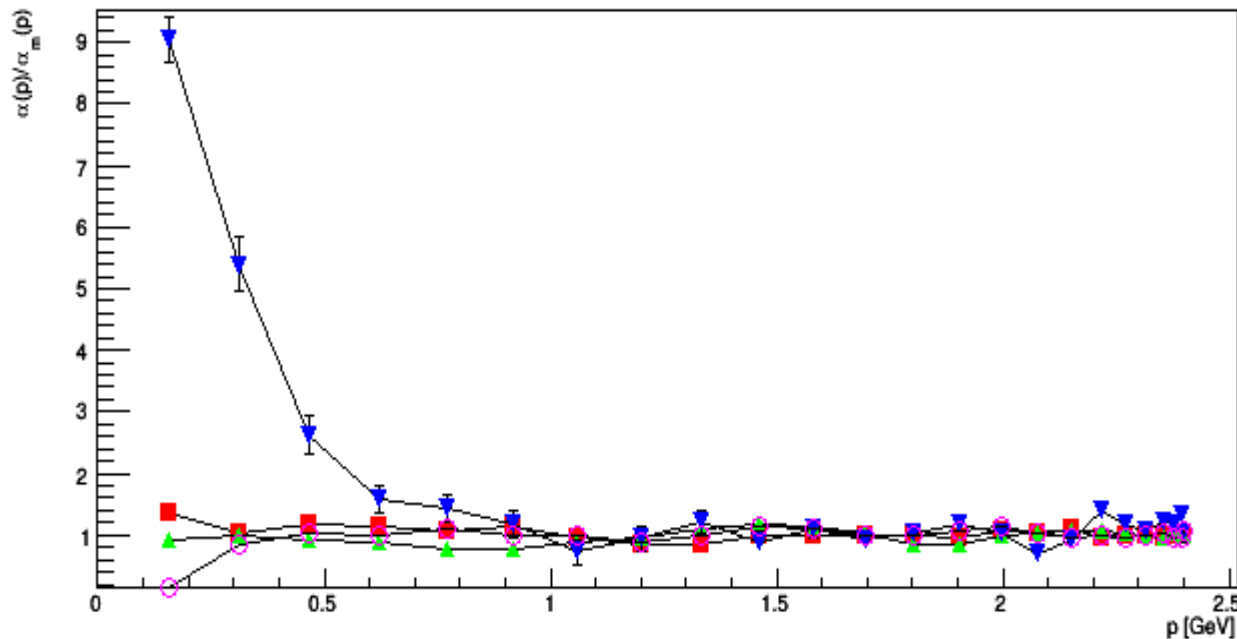
Running coupling - extrapolated

[3d V=(7.9 fm)³, a=0.17 fm]Extreme α 

Evolution with # of copies



Ratio to minimal Landau gauge



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 - Superconformal theories [Capri et al. PLB 14, Maas & Zitz EPJC 16]
 - No lattice yet, claim of difference/absence

Summary

- The structure of the residual gauge orbit is complicated due to the Gribov-Singer ambiguity
- Number of gauge copies very large
 - Multiple coordinates required for distinction
- Requires additional gauge conditions
- Influence on correlation functions substantial
 - Especially ghost and ghost-dependent quantities
 - Requires extrapolation in Gribov copies
- Not yet resolved
 - Especially in the continuum

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