

Scheme variations of the QCD coupling



Strong coupling

Scale evolution

The coupling α_s

Adler function

Tau decay

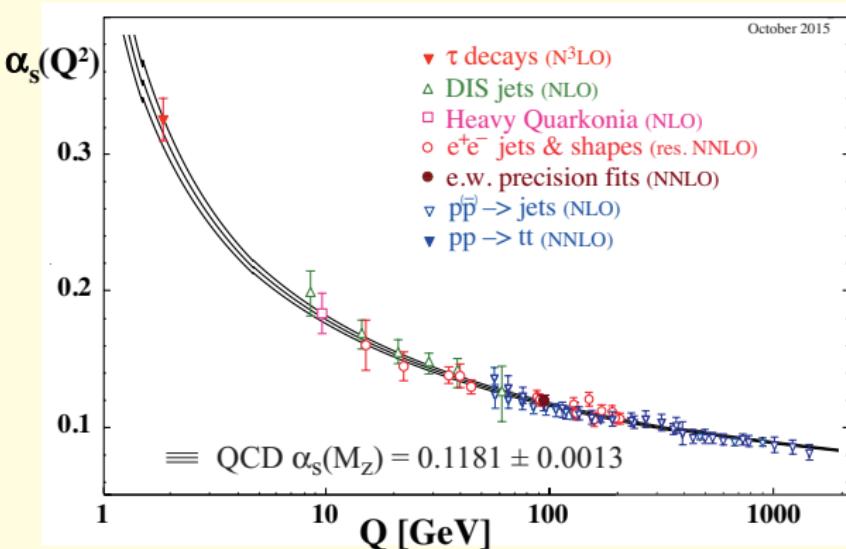
Higgs decay

Summary

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Strong coupling

PDG 2015



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Scale evolution

Scale evolution of α_s is given by the β -function:

$$-Q \frac{da_Q}{dQ} \equiv \beta(a_Q) = \beta_1 a_Q^2 + \beta_2 a_Q^3 + \beta_3 a_Q^4 + \dots$$

with $a_Q = \alpha_s / \pi$.

The scale invariant parameter Λ can be defined by:

$$\Lambda \equiv Q e^{-\frac{1}{\beta_1 a_Q}} [a_Q]^{-\frac{\beta_2}{\beta_1^2}} \exp \left\{ \int_0^{a_Q} \frac{da}{\tilde{\beta}(a)} \right\},$$

where

$$\frac{1}{\tilde{\beta}(a)} \equiv \frac{1}{\beta(a)} - \frac{1}{\beta_1 a^2} + \frac{\beta_2}{\beta_1^2 a}$$

is free of singularities at $a \rightarrow 0$.

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However, Λ depends on the renormalisation scheme.

$$a' \equiv a + c_1 a^2 + c_2 a^3 + c_3 a^4 + \dots$$

Then, Λ transforms as:

(Celmaster, Gonsalves 1979)

$$\Lambda' = \Lambda e^{c_1/\beta_1}.$$

This suggests to define a new coupling \hat{a}_Q :

$$\begin{aligned}\frac{1}{\hat{a}_Q} + \frac{\beta_2}{\beta_1} \ln \hat{a}_Q &\equiv \beta_1 \left(\ln \frac{Q}{\Lambda} + \frac{C}{2} \right) \\ &= \frac{1}{a_Q} + \frac{\beta_1}{2} C + \frac{\beta_2}{\beta_1} \ln a_Q - \beta_1 \int_0^{a_Q} \frac{da}{\tilde{\beta}(a)}\end{aligned}$$

(Boito, MJ, Miravitllas 2016)



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The β -function of \hat{a}_Q takes the simple form:

$$-Q \frac{d\hat{a}_Q}{dQ} \equiv \hat{\beta}(\hat{a}_Q) = \frac{\beta_1 \hat{a}_Q^2}{\left(1 - \frac{\beta_2}{\beta_1} \hat{a}_Q\right)}.$$

Perturbatively, \hat{a} and a are related by:

$$\begin{aligned}\hat{a}(a) &= a - \frac{9}{4} Ca^2 - \left(\frac{3397}{2592} + 4C - \frac{81}{16} C^2\right)a^3 \\ &- \left(\frac{741103}{186624} + \frac{233}{192} C - \frac{45}{2} C^2 + \frac{729}{64} C^3 + \frac{445}{144} \zeta_3\right)a^4 \\ &- \left(\frac{727240925}{80621568} - \frac{869039}{41472} C - \frac{26673}{512} C^2 + \frac{351}{4} C^3 - \frac{6561}{256} C^4\right. \\ &\left.- \frac{445}{32} \zeta_3 C + \frac{10375693}{373248} \zeta_3 - \frac{1335}{256} \zeta_4 - \frac{534385}{20736} \zeta_5\right)a^5 + \mathcal{O}(a^6)\end{aligned}$$

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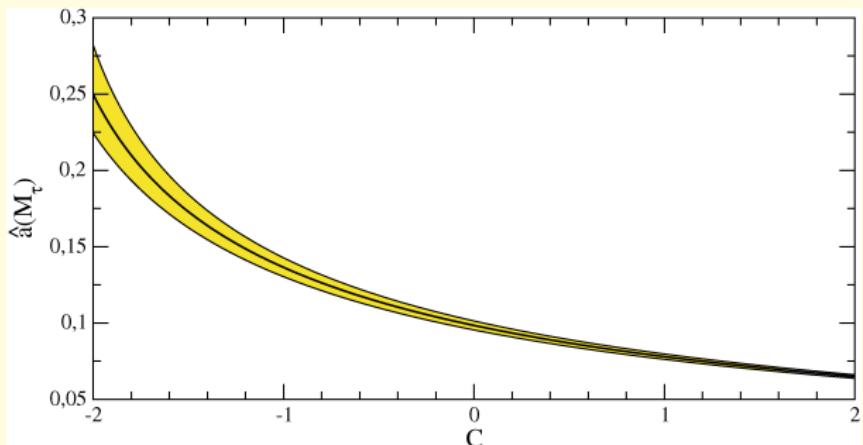
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The coupling \hat{a} as a function of C for $\alpha_s(M_\tau) = 0.316(10)$.

Adler function

$$4\pi^2 D(a_Q) - 1 \equiv \hat{D}(a_Q) = \sum_{n=1}^{\infty} c_{n,1} a_Q^n$$
$$= a_Q + 1.640 a_Q^2 + 6.371 a_Q^3 + 49.08 a_Q^4 + \dots$$

Expressed in terms of the coupling \hat{a} :

$$\begin{aligned}\hat{D}(\hat{a}_Q) &= \hat{a}_Q + (1.640 + 2.25C) \hat{a}_Q^2 \\ &+ (7.682 + 11.38C + 5.063C^2) \hat{a}_Q^3 \\ &+ (61.06 + 72.08C + 47.40C^2 + 11.39C^3) \hat{a}_Q^4 + \dots\end{aligned}$$



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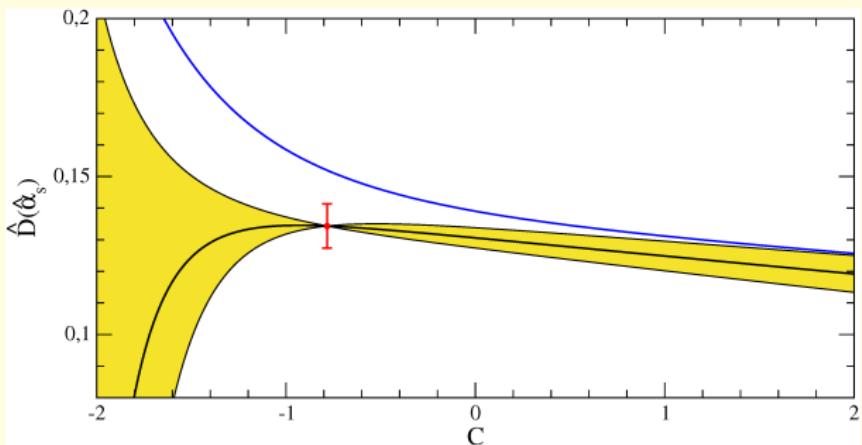
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$$\hat{D}(\hat{\alpha}_{M_\tau}, C = -0.78) = 0.1343 \pm 0.0070 \pm 0.0067$$

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The hadronic τ decay rate takes the form:

$$R_\tau = 3 S_{\text{EW}}(|V_{ud}|^2 + |V_{us}|^2)(1 + \delta^{(0)} + \dots)$$

The perturbative part in fixed-order PT reads:

$$\delta_{\text{FO}}^{(0)}(a_Q) = a_Q + 5.202 a_Q^2 + 26.37 a_Q^3 + 127.1 a_Q^4 + \dots$$

Expressed in terms of the coupling \hat{a} :

$$\begin{aligned}\delta_{\text{FO}}^{(0)}(\hat{a}_Q) &= \hat{a}_Q + (5.202 + 2.25C)\hat{a}_Q^2 \\ &+ (27.68 + 27.41C + 5.063C^2)\hat{a}_Q^3 \\ &+ (148.4 + 235.5C + 101.5C^2 + 11.39C^3)\hat{a}_Q^4 + \dots\end{aligned}$$



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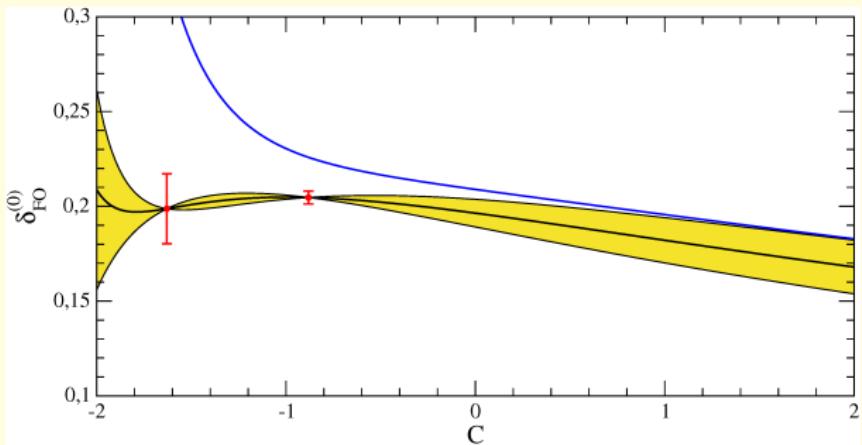
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$$\delta_{FO}^{(0)}(\hat{a}_{M_\tau}, C = -0.88) = 0.2047 \pm 0.0034 \pm 0.0133$$



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The **decay rate** of H into $\bar{q}q$ is given by:

$$\Gamma(H \rightarrow q\bar{q}) = \frac{\sqrt{2} G_F}{M_H} \text{Im}\Psi(M_H^2 + i0) \equiv \frac{N_c G_F M_H}{4\sqrt{2}\pi} \hat{m}_q^2 \hat{R}(\alpha_s(M_H))$$

\hat{R} is **only** a **function** of the **coupling**: (MJ, Miravitllas 2016)

$$\begin{aligned} \hat{R}(\hat{\alpha}_s) &= [\hat{\alpha}_s(Q)]^{24/23} \left\{ 1 + (8.0176 + 2C) \hat{a}_Q \right. \\ &\quad + (46.732 + 33.924C + 3.9167C^2) \hat{a}_Q^2 \\ &\quad + (141.19 + 315.38C + 103.88C^2 + 7.6157C^3) \hat{a}_Q^3 \\ &\quad - (524.03 - 1491.9C - 1353.1C^2 - 277.97C^3 \\ &\quad \left. - 14.756C^4) \hat{a}_Q^4 + \dots \right\} \end{aligned}$$



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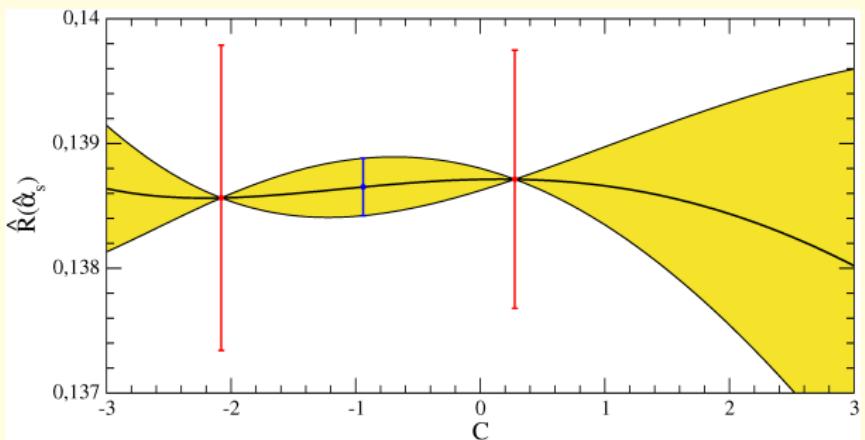
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$$\hat{R}(C = -0.94) = 0.1387 \pm 0.0002 \pm 0.0020 = 0.1387 \pm 0.0020$$

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Summary

- The scheme dependence of the novel coupling \hat{a} can be parameterised by a single parameter C .
- Its corresponding β -function turns out to be manifestly scheme invariant.
- The coupling \hat{a} allows to easily study scheme dependencies, and to optimise the perturbative expansion.
- This appears to be particularly useful in observables that contain global multiplicative factors of the coupling.



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Thank You!