

Strongly-interacting mirror fermions at the LHC

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XIIth Quark Confinement and the Hadron Spectrum

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Motivation

- Explore the inner works of **mass** generation
- Understand the Physics behind the M_{Planck}/M_{weak} **hierarchy**
- Results from the **LHC** might provide indirect clues for **Planck-energy** Physics

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- 1 Motivation
- 2 Inductive approach
 - Hierarchy stabilization and extrapolation to M_{Planck}
 - Mirror fermions: a compelling completion of the SM
- 3 Mirror-meson phenomenology
 - *A déjà* – *vu* reminiscent of QCD
 - Mirror-meson decays at the LHC
- 4 Conclusions

Some natural solutions

- Large extra dimensions
Stabilization via **size** of extra dimensions
- Known particles have spin-zero partners (SUSY)
Stabilization via **space-time symmetry**
Interaction: **weak**
- Known particles have mirror partners (Katoptrons)
Stabilization via **gauge symmetry**
Interaction: **strong**

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A field fertilized by great old ideas

- T.D. Lee and C.N. Yang (1956): **mirror** fermions
- J.C. Pati and A. Salam (1973): Coupling **unification**
- S. Weinberg (1976), L. Susskind (1979):
Dynamical BEH mechanism
Universe: a "superconductor", Composite Higgs $\sim \langle \bar{\Psi}\Psi \rangle$
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An (almost) **left-right symmetric** Lagrangian

$$\mathcal{L} = \mathcal{L}_{YM} + \mathcal{L}_{int}$$

$$\mathcal{L}_{YM} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu} - \frac{1}{4} G_{\mu\nu}^e G^{e\mu\nu} - \frac{1}{4} G_{\mu\nu}^K G^{K\mu\nu}$$

$$\mathcal{L}_{int} = i \sum_{j,k} \left[(\bar{\Psi}_u^{jk}, \bar{\Psi}_d^{jk}) \gamma_\mu \mathcal{D}_k^\mu \begin{pmatrix} \Psi_u^{jk} \\ \Psi_d^{jk} \end{pmatrix} + (\bar{\tilde{\Psi}}_u^{jk}, \bar{\tilde{\Psi}}_d^{jk}) \gamma_\mu \hat{\mathcal{D}}_k^\mu \begin{pmatrix} \hat{\Psi}_u^{jk} \\ \hat{\Psi}_d^{jk} \end{pmatrix} \right]$$

$$\Psi_u^{jk} = (N_L^j, U_L^j, N_R^K \delta^{3j}, U_R^K \delta^{3j})$$

$$\Psi_d^{jk} = (E_L^j, D_L^j, E_R^K \delta^{3j}, D_R^K \delta^{3j})$$

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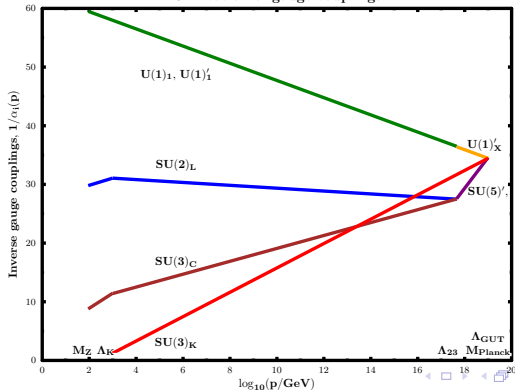
$$\begin{aligned} \Psi_u^{j|k} &= (N_L^j, U_L^j, N_R^K \delta^{3j}, U_R^K \delta^{3j}) \\ \Psi_d^{j|k} &= (E_L^j, D_L^j, E_R^K \delta^{3j}, D_R^K \delta^{3j}) \\ \hat{\Psi}_u^{j|k} &= (N_R^j, U_R^j, N_L^K \delta^{3j}, U_L^K \delta^{3j}) \\ \hat{\Psi}_d^{j|k} &= (E_R^j, D_R^j, E_L^K \delta^{3j}, D_L^K \delta^{3j}), \end{aligned} \quad k = 1, \dots, 4$$

Gauged mirror generation symmetry: **strong** at 1 TeV
 $\sim M_{\text{Planck}} \exp(-1/\alpha_{\text{GUT}})$: G.T., *EJTP* **10** (2013) 135.

$$E_8 \times E'_8(M_{\text{Planck}}) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_1 \times SU(3)_K(1 \text{ TeV}) \rightarrow$$

$$\rightarrow \text{Standard Model}$$

Unification of gauge couplings



The (parity-odd) spectrum

- **Color singlets:**

$\pi^{K 2 0}$ and $\pi^{K 2 \pm}$, "eaten" by Z^0 , W^\pm

$\pi^{K 1 0}$, $\pi^{K 1 \pm}$, $\pi^{K 1 0'}$, η^K (spin – 0)
 $\rho^{K 1,2 0}$, $\rho^{K 1,2 \pm}$, $\rho^{K 0'}$, ω^K (spin – 1)

- **Color triplets (leptoquarks):**

$\pi_3^{K 1,2,2',5}$, $\bar{\pi}_3^{K 1,2,2',5}$ (spin – 0)
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- **Color octets:**

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- SM-fermion mass generation:
consecutive breaking of the
 $SU(3)_K$ katoptron-generation symmetry \rightarrow

Doubling of the mirror-meson spectrum \rightarrow
heavy (group "B") and light (group "A") mirror mesons:

$$r = M_B/M_A \sim \exp\left(3(C_2(SU(3)_K) - C_2(SU(2)_K))\right) \sim 5.75$$

G.T., *Mod. Phys. Lett. A* **16** (2001) 53.

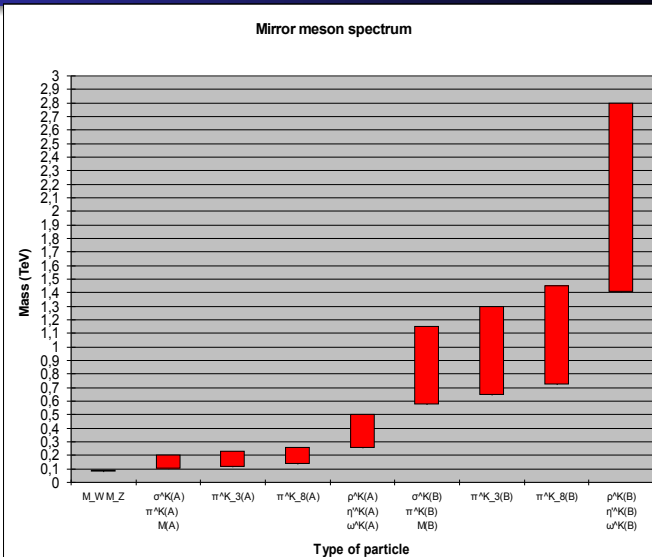
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Group A & B mirror-meson mass spectra



Dynamical Masses and the **S** parameter

$$\langle \bar{\Psi}_{u,d}^{3|3,4} \hat{\Psi}_{u,d}^{3|3,4} + \text{h.c.} \rangle \approx -\frac{N_i}{4\pi^2} \int dp^2 M_i(p^2)$$

$$m_f = \begin{pmatrix} m_{SM} & m \\ m & M \end{pmatrix}, \text{ with } M \equiv \begin{pmatrix} M_A & 0 \\ 0 & M_B \end{pmatrix} \text{ and } m \equiv \begin{pmatrix} m_{AA} & m_{AB} \\ m_{AB} & m_{BB} \end{pmatrix}$$

$$S = S_A + S_B = 4\pi \Sigma_A \left(\frac{F_{\rho_A^K}^2}{M_{\rho_A^K}^2} - \frac{F_{a_A^K}^2}{M_{a_A^K}^2} \right) + 4\pi \Sigma_B \left(\frac{F_{\rho_B^K}^2}{M_{\rho_B^K}^2} - \frac{F_{a_B^K}^2}{M_{a_B^K}^2} \right)$$

$$S_B \approx 4\pi(v/M_{\rho_B^K})^2 \lesssim 0.122 \text{ for } M_{\rho_B^K} > 2.5 \text{ TeV}$$

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Deviations from SM due to **radiative corrections**

Importance of heavy fermions:

- top- and bottom-quark **left-right** asymmetries
- V_{tb} not assuming 3×3 CKM matrix unitarity
- lepton **universality** (B-meson decays?)

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Most promising processes

- $gg \longrightarrow \pi_{8(A)B}^{K 0'} \longrightarrow (\bar{b}b) \bar{t}t$
- $gg \longrightarrow \sigma_{(A)B}^{K 0} \longrightarrow (\bar{b}b) \bar{t}t, \gamma\gamma$
- **Acollinear** $\bar{f}f'$ jets:

$$gg \longrightarrow \text{direct, } g, \rho_{8B}^{K 0'} \longrightarrow \pi_{8B}^{K +} \pi_{8B}^{K -} \longrightarrow \bar{b}t + \bar{t}b$$

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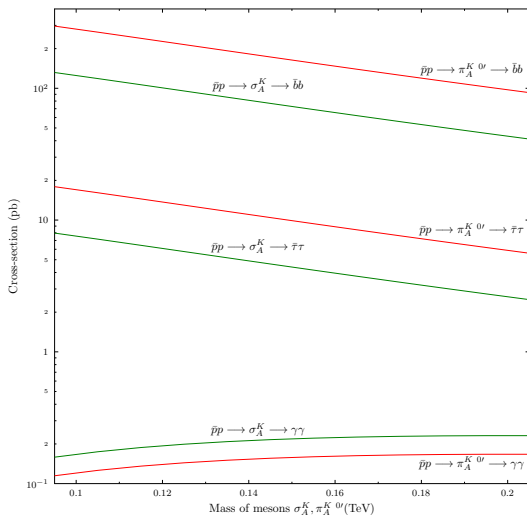
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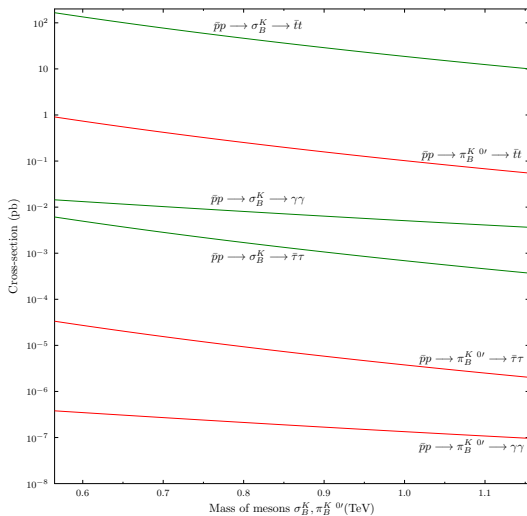
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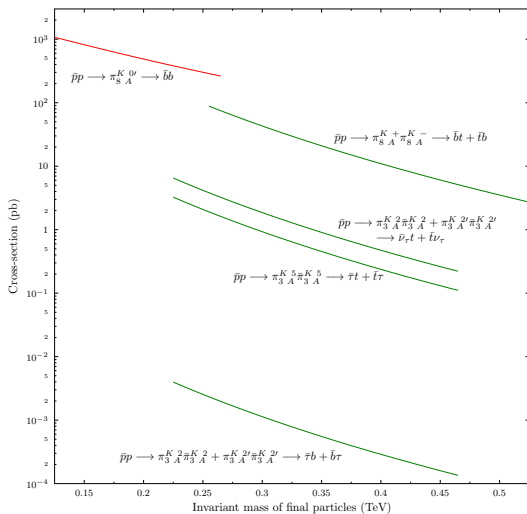
Color-singlet A-mirror-meson processes



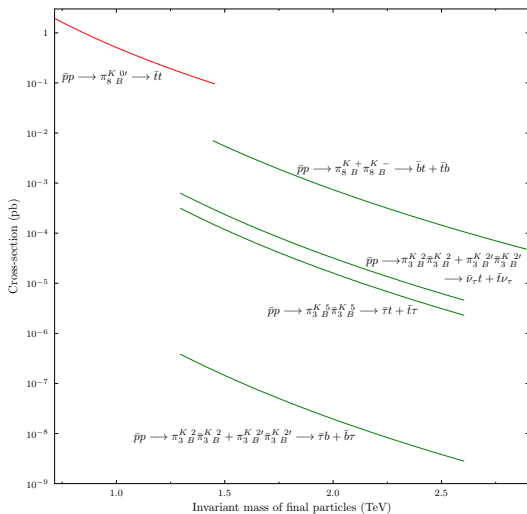
Color-singlet B-mirror-meson processes



Color-octet & color-triplet A-mirror-meson decays



Color-octet & color-triplet B-mirror-meson decays



Conclusions

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- Rich **mirror-meson** phenomenology → Holistic approach enabling **differentiation** of competing models predicting similar signals
- Need for a **3-4 TeV leptonic** collider due to the large QCD background at the LHC
- Quantum-gravity implications:
 - **space-time discreteness** → connected to **mirror** fermions
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