

Inverse Magnetic Catalysis in Holographic Models of QCD

based on
JHEP 1505:121,2015 (1501.03262) KM

Kiminad Mamo

University of Illinois at Chicago

XIIth Quark Confinement at Thessaloniki (Greece), August 29, 2016

Outline

1 Introduction

Outline

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD

Outline

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$

Outline

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$
- 4 Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$
- 4 Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$
- 4 Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$

- 1 Introduction
- 2 Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD
 - for $B = 0$
 - for $B \neq 0$
- 3 Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$
- 4 Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$
 - for $B = 0$
 - for $B \neq 0$
- 5 Summary and conclusion

Introduction

Introduction

- strong magnetic field B is produced

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$,

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$,

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$, and in the interior of dense neutron stars $eB \sim 4 \text{ MeV}^2$

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$, and in the interior of dense neutron stars $eB \sim 4 \text{ MeV}^2$
- deconfinement and chiral symmetry breaking transitions in the presence of B from lattice QCD simulation:

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$, and in the interior of dense neutron stars $eB \sim 4 \text{ MeV}^2$
- deconfinement and chiral symmetry breaking transitions in the presence of B from lattice QCD simulation:
 - * the critical temperatures of both the deconfinement T_{dec} and chiral symmetry breaking T_χ transitions decrease with increasing B

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$, and in the interior of dense neutron stars $eB \sim 4 \text{ MeV}^2$
- deconfinement and chiral symmetry breaking transitions in the presence of B from lattice QCD simulation:
 - * the critical temperatures of both the deconfinement T_{dec} and chiral symmetry breaking T_χ transitions decrease with increasing B (*inverse magnetic catalysis*)

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$, and in the interior of dense neutron stars $eB \sim 4 \text{ MeV}^2$
- deconfinement and chiral symmetry breaking transitions in the presence of B from lattice QCD simulation:
 - * the critical temperatures of both the deconfinement T_{dec} and chiral symmetry breaking T_χ transitions decrease with increasing B (*inverse magnetic catalysis*)
 - * $T_{dec}(B) \simeq$

Introduction

- strong magnetic field B is produced in ultrarelativistic heavy-ion collision experiments at RHIC $eB \sim 0.01 \text{ GeV}^2$ and LHC $eB \sim 0.25 \text{ GeV}^2$, during the electroweak phase transition of the early Universe $eB \sim 4 \text{ GeV}^2$, and in the interior of dense neutron stars $eB \sim 4 \text{ MeV}^2$
- deconfinement and chiral symmetry breaking transitions in the presence of B from lattice QCD simulation:
 - * the critical temperatures of both the deconfinement T_{dec} and chiral symmetry breaking T_χ transitions decrease with increasing B (*inverse magnetic catalysis*)
 - * $T_{dec}(B) \simeq T_\chi(B)$

Hard-wall AdS/QCD

Hard-wall AdS/QCD

- in the hard-wall AdS/QCD:

Hard-wall AdS/QCD

- in the hard-wall AdS/QCD:

Hard-wall AdS/QCD

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon

Hard-wall AdS/QCD

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime)

Hard-wall AdS/QCD

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$ where $r = \infty$ is the boundary of the thermal- AdS_5 spacetime

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$ where $r = \infty$ is the boundary of the thermal- AdS_5 spacetime
 - * the confining IR physics is imposed by boundary conditions at the IR cut-off $r = r_0$

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$ where $r = \infty$ is the boundary of the thermal- AdS_5 spacetime
 - * the confining IR physics is imposed by boundary conditions at the IR cut-off $r = r_0$
 - * the scalar glueball, spin-2 glueball, and meson spectrum is determined by

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$ where $r = \infty$ is the boundary of the thermal- AdS_5 spacetime
 - * the confining IR physics is imposed by boundary conditions at the IR cut-off $r = r_0$
 - * the scalar glueball, spin-2 glueball, and meson spectrum is determined by solving for the eigenmodes of a 5d scalar, gravitational, and gauge field fluctuations, respectively, living on the thermal- AdS_5 spacetime with an IR cut-off at $r = r_0$

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$ where $r = \infty$ is the boundary of the thermal- AdS_5 spacetime
 - * the confining IR physics is imposed by boundary conditions at the IR cut-off $r = r_0$
 - * the scalar glueball, spin-2 glueball, and meson spectrum is determined by solving for the eigenmodes of a 5d scalar, gravitational, and gauge field fluctuations, respectively, living on the thermal- AdS_5 spacetime with an IR cut-off at $r = r_0$
 - * the eigenvalues (which are the masses of the glueballs and mesons)

- in the hard-wall AdS/QCD:
 - * the bulk gravitational background consists of a slice of AdS_5 spacetime without a horizon (thermal- AdS_5 spacetime) but with an IR cut-off at $r = r_0$ where $r = \infty$ is the boundary of the thermal- AdS_5 spacetime
 - * the confining IR physics is imposed by boundary conditions at the IR cut-off $r = r_0$
 - * the scalar glueball, spin-2 glueball, and meson spectrum is determined by solving for the eigenmodes of a 5d scalar, gravitational, and gauge field fluctuations, respectively, living on the thermal- AdS_5 spacetime with an IR cut-off at $r = r_0$
 - * the eigenvalues (which are the masses of the glueballs and mesons) are given by $m_n^2 \sim r_0^2 n^2$ for large n

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} +$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + \right.$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + \frac{12}{L^2} \right) ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bdy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + \frac{12}{L^2} \right) ,$$

and the boundary action S_{bdy} is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + \frac{12}{L^2} \right) ,$$

and the boundary action S_{bndy} is

$$S_{\text{bndy}} = \frac{1}{8\pi G_5} \int d^4x \sqrt{-\gamma} \left(K - \right.$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + \frac{12}{L^2} \right) ,$$

and the boundary action S_{bndy} is

$$S_{\text{bndy}} = \frac{1}{8\pi G_5} \int d^4x \sqrt{-\gamma} \left(K - \frac{3}{L} \right) \Big|_{r=r_\Lambda}$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{2\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{2\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} ,$$

and, the on-shell Euclidean boundary action S_{bndy}^E is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{2\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} ,$$

and, the on-shell Euclidean boundary action S_{bndy}^E is

$$S_{\text{bndy}}^E = -\frac{V_3}{8\pi G_5} \int_0^\beta dt_E \sqrt{\gamma} (K - 3)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- using $R = -20$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{2\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} ,$$

and, the on-shell Euclidean boundary action S_{bndy}^E is

$$S_{\text{bndy}}^E = -\frac{V_3}{8\pi G_5} \int_0^\beta dt_E \sqrt{\gamma} (K - 3)$$

- the free energy $F = \frac{S_E}{\beta}$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- AdS_5 black hole

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- AdS_5 black hole

$$ds_{\text{bh}}^2 = r^2 (f(r)dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{f(r)r^2},$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- AdS_5 black hole

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4},$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- AdS_5 black hole

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4},$$

$$T = \frac{1}{\beta} = \frac{r_h}{\pi}$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- AdS_5 black hole

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4},$$

$$T = \frac{1}{\beta} = \frac{r_h}{\pi}$$

- the free energy of the AdS_5 black hole

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- AdS_5 black hole

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4},$$

$$T = \frac{1}{\beta} = \frac{r_h}{\pi}$$

- the free energy of the AdS_5 black hole

$$F_{\text{bh}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} \right)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- thermal- AdS_5

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- thermal-AdS₅

$$ds_{\text{th}}^2 = r^2 (dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{r^2},$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- thermal-AdS₅

$$ds_{\text{th}}^2 = r^2 (dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{r^2},$$
$$T = \frac{1}{\beta}$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- thermal-AdS₅

$$ds_{\text{th}}^2 = r^2 (dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{r^2},$$
$$T = \frac{1}{\beta}$$

- the free energy of the thermal-AdS₅

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- thermal-AdS₅

$$ds_{\text{th}}^2 = r^2 (dt_E^2 + dx^2 + dy^2 + dz^2) + \frac{dr^2}{r^2},$$
$$T = \frac{1}{\beta}$$

- the free energy of the thermal-AdS₅

$$F_{\text{th}} = -\frac{V_3}{8\pi G_5} (r_0^4)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the difference between the free energies

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the difference between the free energies

$$\Delta F = F_{\text{bh}} - F_{\text{th}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - r_0^4 \right)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the difference between the free energies

$$\Delta F = F_{\text{bh}} - F_{\text{th}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - r_0^4 \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the difference between the free energies

$$\Delta F = F_{\text{bh}} - F_{\text{th}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - r_0^4 \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(\frac{r_{hc}^4}{2} - r_0^4 \right) = 0,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B = 0$

- the difference between the free energies

$$\Delta F = F_{\text{bh}} - F_{\text{th}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - r_0^4 \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(\frac{r_{hc}^4}{2} - r_0^4 \right) = 0,$$

$$T_c = T(r_h = r_{hc}) = \frac{2^{1/4}}{\pi} r_0 = T_{dec}$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R - F^{MN} F_{MN} + \frac{12}{L^2} \right) ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bdy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R - F^{MN} F_{MN} + \frac{12}{L^2} \right) ,$$

and the boundary action S_{bdy} is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bdy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R - F^{MN} F_{MN} + \frac{12}{L^2} \right) ,$$

and the boundary action S_{bdy} is

$$S_{\text{bdy}} = \frac{1}{8\pi G_5} \int d^4x \sqrt{-\gamma} \left(K - \frac{3}{L} + \right.$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the action

$$S = S_{\text{bulk}} + S_{\text{bndy}} ,$$

where the bulk action S_{bulk} is

$$S_{\text{bulk}} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R - F^{MN} F_{MN} + \frac{12}{L^2} \right) ,$$

and the boundary action S_{bndy} is

$$S_{\text{bndy}} = \frac{1}{8\pi G_5} \int d^4x \sqrt{-\gamma} \left(K - \frac{3}{L} + \frac{L}{2} \left(\ln \frac{r}{L} \right) F^{\mu\nu} F_{\mu\nu} \right) \Big|_{r=r_\Lambda}$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- using $R = -20 + \frac{2}{3}B^2 g^{xx} g^{yy}$, the on-shell Euclidean action S_E

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- using $R = -20 + \frac{2}{3}B^2 g^{xx} g^{yy}$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- using $R = -20 + \frac{2}{3}B^2 g^{xx} g^{yy}$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- using $R = -20 + \frac{2}{3}B^2 g^{xx} g^{yy}$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{8\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} \left(4 + \frac{2}{3} B^2 g^{xx} g^{yy} \right) ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- using $R = -20 + \frac{2}{3}B^2 g^{xx} g^{yy}$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bdy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{8\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} \left(4 + \frac{2}{3} B^2 g^{xx} g^{yy} \right) ,$$

and, the on-shell Euclidean boundary action S_{bdy}^E is

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- using $R = -20 + \frac{2}{3}B^2 g^{xx} g^{yy}$, the on-shell Euclidean action S_E

$$S_E = S_{\text{bulk}}^E + S_{\text{bndy}}^E ,$$

where the on-shell Euclidean bulk action S_{bulk}^E is

$$S_{\text{bulk}}^E = \frac{V_3}{8\pi G_5} \int_0^\beta dt_E \int_{r'}^{r_\Lambda} dr \sqrt{g} \left(4 + \frac{2}{3} B^2 g^{xx} g^{yy} \right) ,$$

and, the on-shell Euclidean boundary action S_{bndy}^E is

$$S_{\text{bndy}}^E = -\frac{V_3}{8\pi G_5} \int_0^\beta dt_E \sqrt{\gamma} \left(K - 3 + B^2 g^{xx} g^{yy} \ln r_\Lambda \right)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r)dt_E^2 + q(r)dz^2 + h(r)(dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2} ,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r)dt_E^2 + q(r)dz^2 + h(r)(dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4} - \frac{2}{3}B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4} - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4} - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4} - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$T = \frac{1}{\beta} = \frac{r_h}{2\pi} \left(1 + \frac{M}{r_h^4} - \frac{2}{3} B^2 \left(\frac{1}{2r_h^4} - \frac{\ln r_h}{r_h^4} \right) \right) + \mathcal{O}(B^4)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r)dt_E^2 + q(r)dz^2 + h(r)(dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4} - \frac{2}{3}B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3}B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3}B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$T = \frac{1}{\beta} = \frac{r_h}{2\pi} \left(1 + \frac{M}{r_h^4} - \frac{2}{3}B^2 \left(\frac{1}{2r_h^4} - \frac{\ln r_h}{r_h^4} \right) \right) + \mathcal{O}(B^4)$$

- the free energy of the AdS_5 black hole

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- AdS_5 black hole [D'Hoker and Kraus (2009)]

$$ds_{\text{bh}}^2 = r^2 (f(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f(r)r^2},$$

$$f(r) = 1 - \frac{M}{r^4} - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$T = \frac{1}{\beta} = \frac{r_h}{2\pi} \left(1 + \frac{M}{r_h^4} - \frac{2}{3} B^2 \left(\frac{1}{2r_h^4} - \frac{\ln r_h}{r_h^4} \right) \right) + \mathcal{O}(B^4)$$

- the free energy of the AdS_5 black hole

$$F_{\text{bh}} = -\frac{V_3}{8\pi G_5} \left(r_h^4 - \frac{1}{2} M + \frac{2}{3} B^2 \ln r_h - \frac{1}{3} B^2 \right) + \mathcal{O}(B^4)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- thermal- AdS_5

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- thermal-AdS₅

$$ds_{\text{thermal}}^2 = r^2 (f_0(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f_0(r)r^2},$$

$$f_0(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- thermal-AdS₅

$$ds_{\text{thermal}}^2 = r^2 (f_0(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f_0(r)r^2},$$

$$f_0(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$T = \frac{1}{\beta}$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- thermal-AdS₅

$$ds_{\text{thermal}}^2 = r^2 (f_0(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f_0(r)r^2},$$

$$f_0(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$T = \frac{1}{\beta}$$

- the free energy of the thermal-AdS₅

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- thermal-AdS₅

$$ds_{\text{thermal}}^2 = r^2 (f_0(r) dt_E^2 + q(r) dz^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f_0(r)r^2},$$

$$f_0(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$T = \frac{1}{\beta}$$

- the free energy of the thermal-AdS₅

$$F_{\text{th}} = -\frac{V_3}{8\pi G_5} \left(r_0^4 + \frac{2}{3} B^2 \ln r_0 - \frac{1}{3} B^2 \right) + \mathcal{O}(B^4)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_h}{r_0}\right) \right) + \mathcal{O}(B^4)$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_h}{r_0}\right) \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_h}{r_0}\right) \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(r_{hc}^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_{hc}}{r_0}\right) \right) + \mathcal{O}(B^4) = 0,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_h}{r_0}\right) \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(r_{hc}^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_{hc}}{r_0}\right) \right) + \mathcal{O}(B^4) = 0,$$

- the constraint equation

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_h}{r_0}\right) \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(r_{hc}^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_{hc}}{r_0}\right) \right) + \mathcal{O}(B^4) = 0,$$

- the constraint equation

$$r_{hc}^4 + \frac{2}{3}B^2 \ln\left(\frac{r_{hc}}{r_0}\right) - 2r_0^4 + \mathcal{O}(B^4) = 0,$$

Hawking-Page or confinement/deconfinement transition in hard-wall AdS/QCD for $B \neq 0$

- the difference between the free energies

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_h}{r_0}\right) \right) + \mathcal{O}(B^4)$$

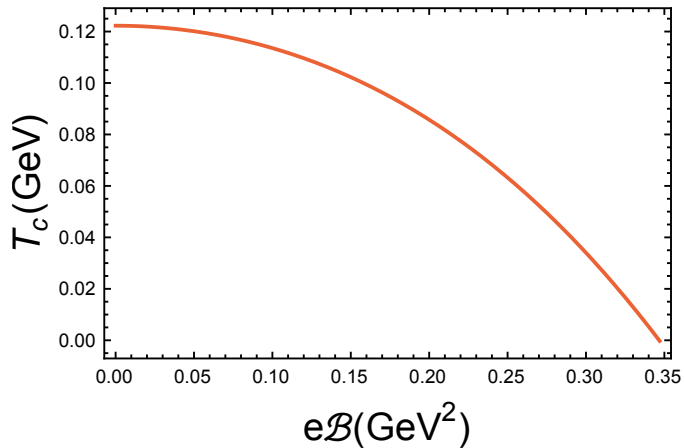
- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(r_{hc}^4 - r_0^4 - \frac{1}{2}M + \frac{2}{3}B^2 \ln\left(\frac{r_{hc}}{r_0}\right) \right) + \mathcal{O}(B^4) = 0,$$

- the constraint equation

$$r_{hc}^4 + \frac{2}{3}B^2 \ln\left(\frac{r_{hc}}{r_0}\right) - 2r_0^4 + \mathcal{O}(B^4) = 0,$$

$$T_c = T(r_h = r_{hc}) = T_{dec}$$



$\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

$\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

$\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- the fermions obey antiperiodic boundary conditions around the circle S^1_ℓ , and they acquire a tree-level mass $\sim \frac{1}{\ell}$

$\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- the fermions obey antiperiodic boundary conditions around the circle S^1_ℓ , and they acquire a tree-level mass $\sim \frac{1}{\ell}$
- scalars obey periodic boundary conditions around the circle S^1_ℓ , hence they are massless at tree-level and acquire masses only at the quantum level through their couplings to the fermions

$\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- the fermions obey antiperiodic boundary conditions around the circle S^1_ℓ , and they acquire a tree-level mass $\sim \frac{1}{\ell}$
- scalars obey periodic boundary conditions around the circle S^1_ℓ , hence they are massless at tree-level and acquire masses only at the quantum level through their couplings to the fermions
- the gluons do not acquire masses

$\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- the fermions obey antiperiodic boundary conditions around the circle S^1_ℓ , and they acquire a tree-level mass $\sim \frac{1}{\ell}$
- scalars obey periodic boundary conditions around the circle S^1_ℓ , hence they are massless at tree-level and acquire masses only at the quantum level through their couplings to the fermions
- the gluons do not acquire masses
- therefore, at low-energy, $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ reduce to pure 3D Yang-Mills theory

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

$$ds_{\text{soliton}}^2 = r^2 (f_s(r) dz'^2 + dt_E'^2 + dx^2 + dy^2) + \frac{dr^2}{f_s(r)r^2} ,$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

$$ds_{\text{soliton}}^2 = r^2 (f_s(r) dz'^2 + dt_E'^2 + dx^2 + dy^2) + \frac{dr^2}{f_s(r)r^2},$$

$$f_s(r) = 1 - \frac{r_0^4}{r^4},$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

$$ds_{\text{soliton}}^2 = r^2 (f_s(r) dz'^2 + dt_E'^2 + dx^2 + dy^2) + \frac{dr^2}{f_s(r)r^2},$$

$$f_s(r) = 1 - \frac{r_0^4}{r^4},$$

$$\frac{1}{\ell} = \frac{r_0}{\pi},$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

$$ds_{\text{soliton}}^2 = r^2 (f_s(r) dz'^2 + dt_E'^2 + dx^2 + dy^2) + \frac{dr^2}{f_s(r)r^2},$$

$$f_s(r) = 1 - \frac{r_0^4}{r^4},$$

$$\frac{1}{\ell} = \frac{r_0}{\pi},$$

$$T = \frac{1}{\beta}$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = F_{\text{bh}} - F_{\text{soliton}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - \frac{r_0^4}{2} \right)$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = F_{\text{bh}} - F_{\text{soliton}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - \frac{r_0^4}{2} \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = F_{\text{bh}} - F_{\text{soliton}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - \frac{r_0^4}{2} \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(\frac{r_{hc}^4}{2} - \frac{r_0^4}{2} \right) = 0,$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = F_{\text{bh}} - F_{\text{soliton}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - \frac{r_0^4}{2} \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(\frac{r_{hc}^4}{2} - \frac{r_0^4}{2} \right) = 0,$$

$$T_c = T(r_h = r_{hc}) =$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = F_{\text{bh}} - F_{\text{soliton}} = -\frac{V_3}{8\pi G_5} \left(\frac{r_h^4}{2} - \frac{r_0^4}{2} \right)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(\frac{r_{hc}^4}{2} - \frac{r_0^4}{2} \right) = 0,$$

$$T_c = T(r_h = r_{hc}) = \frac{r_0}{\pi} = \frac{1}{\ell} = T_{dec}$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the AdS_5 -soliton solution can be determined from the black hole solution by "double Wick rotation" $t = iz'$ and $z = it'$

$$ds_{\text{soliton}}^2 = r^2 (f_s(r) dz'^2 - q(r) dt'^2 + h(r) (dx^2 + dy^2)) + \frac{dr^2}{f_s(r)r^2},$$

$$f_s(r) = 1 - \frac{M}{r^4} - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$q(r) = 1 - \frac{2}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$h(r) = 1 + \frac{1}{3} B^2 \frac{\ln r}{r^4} + \mathcal{O}(B^4),$$

$$\frac{1}{\ell} = \frac{r_0}{2\pi} \left(1 + \frac{M}{r_0^4} + \frac{2}{3} B^2 \left(\frac{\ln r_0}{r_0^4} - \frac{1}{2r_0^4} \right) \right) + \mathcal{O}(B^4).$$

$$T = \frac{1}{\beta}$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 + \frac{2}{3} B^2 \ln \frac{r_h}{r_0} \right) + \mathcal{O}(B^4)$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 + \frac{2}{3} B^2 \ln \frac{r_h}{r_0} \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 + \frac{2}{3} B^2 \ln \frac{r_h}{r_0} \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(r_{hc}^4 - r_0^4 + \frac{2}{3} B^2 \ln \frac{r_{hc}}{r_0} \right) + \mathcal{O}(B^4) = 0,$$

Hawking-Page or confinement/deconfinement transition in $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the free energy of the AdS_5 -soliton can be found by just replacing $r_h \leftrightarrow r_0$ in the free energy of the AdS_5 black hole
- therefore,

$$\Delta F = -\frac{V_3}{8\pi G_5} \left(r_h^4 - r_0^4 + \frac{2}{3} B^2 \ln \frac{r_h}{r_0} \right) + \mathcal{O}(B^4)$$

- the critical temperature for the Hawking-Page or confinement/deconfinement transition

$$\Delta F(r_h = r_{hc}) = -\frac{V_3}{8\pi G_5} \left(r_{hc}^4 - r_0^4 + \frac{2}{3} B^2 \ln \frac{r_{hc}}{r_0} \right) + \mathcal{O}(B^4) = 0,$$

$$T_c = T(r_h = r_{hc}) = T_c^0 \left(1 - \left(\frac{B}{B_c} \right)^2 \right) + \mathcal{O}(B^4) = \frac{1}{\ell}$$

Flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

Flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- we add flavor to $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ by using N_f D7 backreacting branes

Flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$

- we add flavor to $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ by using N_f D7 backreacting branes
- Hawking-Page transition is between AdS_5 black hole and AdS_5 -soliton geometries, now including the backreaction of the D7 branes

Hawking-Page or confinement/deconfinement transition in
flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

$$T = \frac{r_h}{\pi} \left(1 - \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

$$T = \frac{r_h}{\pi} \left(1 - \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where λ_h is the value of the 't Hooft coupling fixed at the horizon r_h , that is, $\lambda_h = 4\pi g_s e^{\phi_h} N_c$

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

$$T = \frac{r_h}{\pi} \left(1 - \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where λ_h is the value of the 't Hooft coupling fixed at the horizon r_h , that is, $\lambda_h = 4\pi g_s e^{\phi_h} N_c$

- the critical temperature for Hawking-Page or confinement/deconfinement transition

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

$$T = \frac{r_h}{\pi} \left(1 - \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where λ_h is the value of the 't Hooft coupling fixed at the horizon r_h , that is, $\lambda_h = 4\pi g_s e^{\phi_h} N_c$

- the critical temperature for Hawking-Page or confinement/deconfinement transition

$$T_c = T(r_h = r_0) = T_c^0 \left(1 - \frac{\lambda_0}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

$$T = \frac{r_h}{\pi} \left(1 - \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where λ_h is the value of the 't Hooft coupling fixed at the horizon r_h , that is, $\lambda_h = 4\pi g_s e^{\phi_h} N_c$

- the critical temperature for Hawking-Page or confinement/deconfinement transition

$$T_c = T(r_h = r_0) = T_c^0 \left(1 - \frac{\lambda_0}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where $T_c^0 = T_c(N_f = 0) = \frac{r_0}{\pi}$

Hawking-Page or confinement/deconfinement transition in flavored $\mathcal{N} = 4$ SYM on $S^1_\beta \times R^2 \times S^1_\ell$ for $B = 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes is [Bigazzi et al. (2009)]

$$T = \frac{r_h}{\pi} \left(1 - \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where λ_h is the value of the 't Hooft coupling fixed at the horizon r_h , that is, $\lambda_h = 4\pi g_s e^{\phi_h} N_c$

- the critical temperature for Hawking-Page or confinement/deconfinement transition

$$T_c = T(r_h = r_0) = T_c^0 \left(1 - \frac{\lambda_0}{64\pi^2} \frac{N_f}{N_c} \right) + \mathcal{O}((N_f/N_c)^2),$$

where $T_c^0 = T_c(N_f = 0) = \frac{r_0}{\pi}$

- consistent with lattice QCD results

Hawking-Page or confinement/deconfinement transition in $D3/D7$ branes on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

Hawking-Page or confinement/deconfinement transition in $D3/D7$ branes on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes with magnetic field B is [Ammon et al. (2012)]

Hawking-Page or confinement/deconfinement transition in $D3/D7$ branes on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes with magnetic field B is [Ammon et al. (2012)]

$$T = \frac{r_h}{\pi} \left(1 + \frac{\lambda_0}{64\pi^2} \frac{N_f}{N_c} \left(1 - 2\sqrt{1 + \frac{B^2}{r_h^4}} \right) \right) + \mathcal{O}((N_f/N_c)^2)$$

Hawking-Page or confinement/deconfinement transition in $D3/D7$ branes on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes with magnetic field B is [Ammon et al. (2012)]

$$T = \frac{r_h}{\pi} \left(1 + \frac{\lambda_0}{64\pi^2} \frac{N_f}{N_c} \left(1 - 2\sqrt{1 + \frac{B^2}{r_h^4}} \right) \right) + \mathcal{O}((N_f/N_c)^2)$$

- the critical temperature for Hawking-Page or confinement/deconfinement transition is at $r_h = r_0$

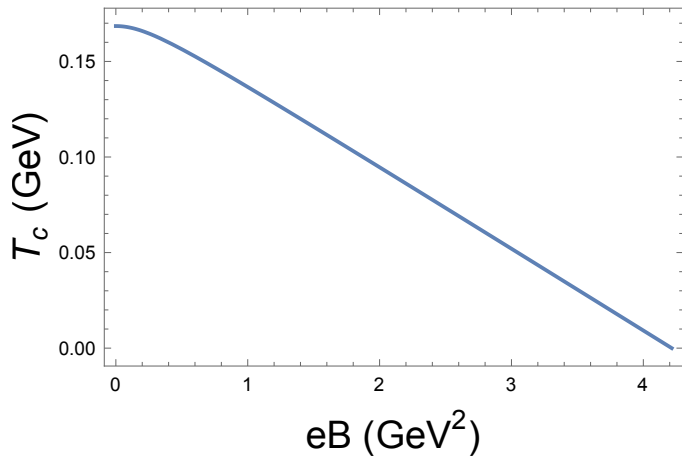
Hawking-Page or confinement/deconfinement transition in $D3/D7$ branes on $S^1_\beta \times R^2 \times S^1_\ell$ for $B \neq 0$

- the Hawking temperature of the AdS_5 black hole after including the backreaction of the N_f D7 branes with magnetic field B is [Ammon et al. (2012)]

$$T = \frac{r_h}{\pi} \left(1 + \frac{\lambda_0}{64\pi^2} \frac{N_f}{N_c} \left(1 - 2\sqrt{1 + \frac{B^2}{r_h^4}} \right) \right) + \mathcal{O}((N_f/N_c)^2)$$

- the critical temperature for Hawking-Page or confinement/deconfinement transition is at $r_h = r_0$

$$T_c = T(r_0) = T_c^0 \left(1 + \frac{\lambda_h}{64\pi^2} \frac{N_f}{N_c} \left(1 - 2\sqrt{1 + \frac{B^2}{r_0^4}} \right) \right) + \mathcal{O}((N_f/N_c)^2)$$



Summary and conclusion

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:
 1. is there *inverse magnetic catalysis* for the chiral transition at $T = T_\chi$ in holographic models of QCD?

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:
 1. is there *inverse magnetic catalysis* for the chiral transition at $T = T_\chi$ in holographic models of QCD? 1511.04042 (Dudal et al.)

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:
 1. is there *inverse magnetic catalysis* for the chiral transition at $T = T_\chi$ in holographic models of QCD? 1511.04042 (Dudal et al.) Phys.Lett. B758 (2016) (Fang)

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:
 1. is there *inverse magnetic catalysis* for the chiral transition at $T = T_\chi$ in holographic models of QCD? 1511.04042 (Dudal et al.) Phys.Lett. B758 (2016) (Fang)
 2. is $T_{dec}(B) = T_\chi(B)$ in holographic models of QCD?

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:
 1. is there *inverse magnetic catalysis* for the chiral transition at $T = T_\chi$ in holographic models of QCD? 1511.04042 (Dudal et al.) Phys.Lett. B758 (2016) (Fang)
 2. is $T_{dec}(B) = T_\chi(B)$ in holographic models of QCD? 1604.07197 (Li et al.)

Summary and conclusion

- I have studied the effect of an external magnetic field B in the gluon sector by including the backreaction of B in $\mathcal{N}=4$ SYM, hard-wall, and D3/D7 models of QCD
- I have shown that the critical temperature for the confinement/deconfinement phase transition T_{dec} decreases with increasing B (*inverse magnetic catalysis*)
- the result is consistent with the recent lattice QCD finding of *inverse magnetic catalysis* for the deconfinement transition
- however, there remains two questions which are not addressed by this work:
 1. is there *inverse magnetic catalysis* for the chiral transition at $T = T_\chi$ in holographic models of QCD? 1511.04042 (Dudal et al.) Phys.Lett. B758 (2016) (Fang)
 2. is $T_{dec}(B) = T_\chi(B)$ in holographic models of QCD? 1604.07197 (Li et al.)
- including the backreaction of B plays a key role in answering many questions related to QCD vacuum in magnetic field

Thank You!