

# XIIth Confinement and the Hadron Spectrum,

*Θεσσαλονικη.*

**Section E; QCD and New Physics**

**August 30, 2016**

Recent progress on **intrinsic charm**

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# motivation

- though “exotic,” **IC** is still **SM physics**, and potentially crucial to **HEP phenomenology**

→ through constraints to **PDF sets** and models,

→ and **background** processes in collider physics

*(see talk by W. Melnitchouk)*

→ but also studies of hadronic bound-state structure!

- knowledge of charm distributions may influence dynamics in **heavy-ion** physics

*(see talks by R. Vogt/J. P. Lansberg)*

...e.g., through effects in  $J/\psi$  **production** and decay

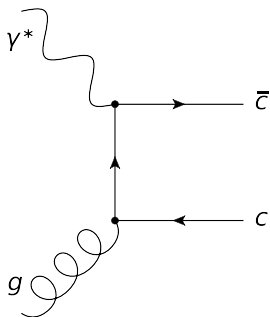
- relevance to searches for **new physics??**

...i.e., through enhanced BSM cross sections...

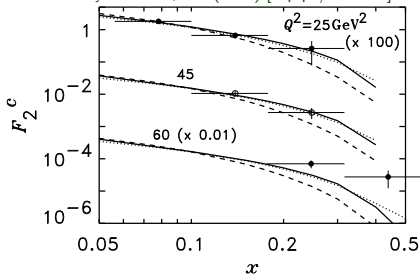
$$\rightarrow \text{DM cross sections} \sim \left| m_q \langle p | \bar{q}q | p \rangle \right|^2$$

charm in *perturbative* QCD (pQCD)

$$\bullet c(x, Q^2 \leq m_c^2) = \bar{c}(x, Q^2 \leq m_c^2) = 0$$



F. M. Steffens, W. Melnitchouk and A. W. Thomas,  
Eur. Phys. J. C **11**, 673 (1999) [hep-ph/9903441].

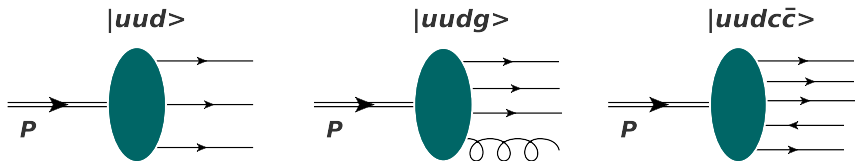


- *intermediate*  $Q^2$ :

$$F_{2, \text{PGF}}^c(x, Q^2) = \frac{\alpha_s(\mu^2)}{9\pi} \int_x^{z'} \frac{dz'}{z} C^{\text{PGF}}(z, Q^2, m_c^2) \cdot xg\left(\frac{x}{z}, \mu^2\right)$$

- *high*  $Q^2$ :

massless **DGLAP** (i.e., *variable flavor-number schemes*)

simplest *nonperturbative* model calculations

→ original models possessed *scalar* vertices...

- Brodsky et al. (1980):

$$P(p \rightarrow uudc\bar{c}) \sim \left[ M^2 - \sum_{i=1}^5 \frac{k_{\perp i}^2 + m_i^2}{x_i} \right]^{-2}$$

→ produces *intrinsic* PDF,  $c^{\text{IC}}(x) = \bar{c}^{\text{IC}}(x)$

- Blümlein (2015):

$$\tau_{\text{life}} = \frac{1}{\sum_i E_i - \bar{E}} = \frac{2P}{\left( \sum_{i=1}^5 \frac{k_{\perp i}^2 + m_i^2}{x_i} - M^2 \right) \Big|_{\sum_j x_j = 1}} \quad \text{vs.} \quad \tau_{\text{int}} = \frac{1}{q_0}$$

→ comparison constrains  $x - Q^2$  space over which IC is observable



# amplitudes from hadronic **EFT**

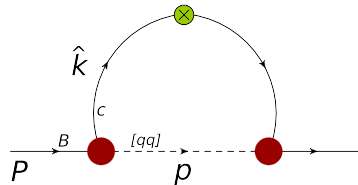
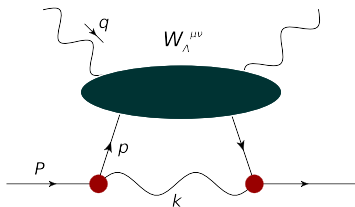
- e.g., for the **dominant** contribution to  $c(x)$ , i.e.,  $\Lambda_c D^*$ :

$$c(x) = \int_x^1 \frac{d\bar{y}}{\bar{y}} f_{\Lambda D^*}(\bar{y}) \cdot c_{\Lambda}\left(\frac{x}{\bar{y}}\right):$$

$$\mathcal{L}_{D^* \Lambda N} = g \bar{\psi}_N \gamma_{\mu} \psi_{\Lambda} \theta_{D^*}^{\mu} + \frac{f}{4M} \bar{\psi}_N \sigma_{\mu\nu} \psi_{\Lambda} F_{D^*}^{\mu\nu} + \text{h.c.}$$

$$\mathcal{L}_{c[qq]\Lambda} = g \bar{\psi}_{\Lambda} \psi_c \phi_{[qq]} + \text{h.c.}$$

quark model  $\rightarrow$  had.  $g, f$

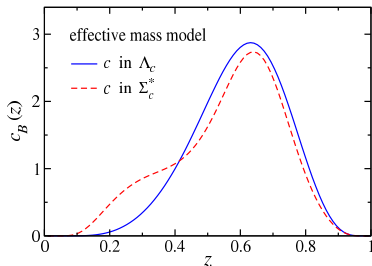
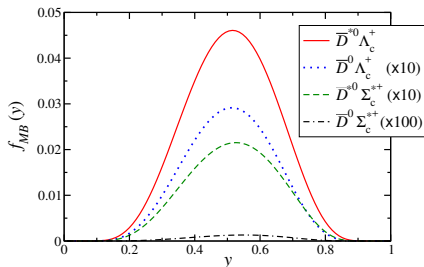


$\rightarrow$  evaluate forward-moving **TOPT** diagrams

hadron/parton **distributions**

$$f_{BD^*}(\bar{y}) = T_B \frac{1}{16\pi^2} \int dk_{\perp}^2 \frac{|\mathcal{F}(s_{BM})|^2}{(s_{BM} - M^2)^2} \frac{1}{\bar{y}(1-\bar{y})} \\ \times \left[ g^2 G_v(\bar{y}, k_{\perp}^2) + \frac{gf}{M} G_{vt}(\bar{y}, k_{\perp}^2) + \frac{f^2}{M^2} G_t(\bar{y}, k_{\perp}^2) \right]$$

$$c_B(z) = N_B \frac{1}{16\pi^2} \int d\hat{k}_{\perp}^2 \frac{1}{z^2(1-z)} \frac{|\mathcal{F}(\hat{s})|^2}{(\hat{s} - M_B^2)^2} \left[ \hat{k}_{\perp}^2 + (m_c + zM_B)^2 \right]$$



→ **model dependence** mainly from  $\mathcal{F}(s)$ ,

$$s(\bar{y}, k_{\perp}^2) = (M_{\Lambda}^2 + k_{\perp}^2)/\bar{y} + (m_D^2 + k_{\perp}^2)/(1-\bar{y})$$

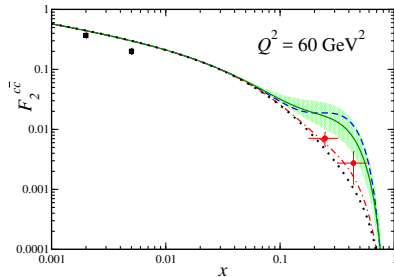
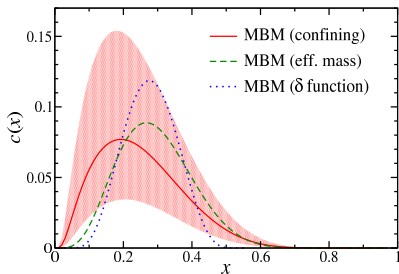
# charm in the nucleon

- tune **universal** cutoff  $\Lambda = \hat{\Lambda}$  to fit **ISR**  $pp \rightarrow \Lambda_c X$  collider data

**multiplicities, momentum sum:**

$$\langle n \rangle_{MB}^{(\text{charm})} = 2.40\% \begin{matrix} +2.47 \\ -1.36 \end{matrix};$$

$$P_c := \langle x \rangle_{IC} = 1.34\% \begin{matrix} +1.35 \\ -0.75 \end{matrix}$$



$$F_2^{c\bar{c}}(x, Q^2) = \frac{4x}{9} [c(x, Q^2) + \bar{c}(x, Q^2)]$$

→ evolve to **EMC** scale,  $Q^2 = 60 \text{ GeV}^2$

low- $x$  H1/ZEUS data check *massless* **DGLAP** evolution



constraints from **global** fits...

P. Jimenez-Delgado, TH, J. T. Londergan and W. Melnitchouk; PRL 114, no. 8, 082002 (2015).

26 sets:

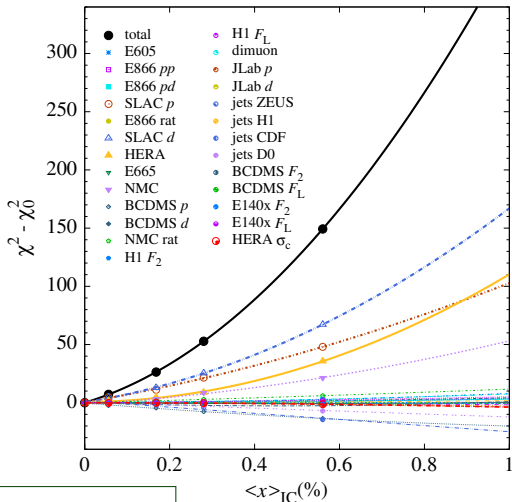
$$N_{dat} = 4296$$

$$Q^2 \geq 1 \text{ GeV}^2$$

$$W^2 \geq 3.5 \text{ GeV}^2$$



\*\* HTs, TMCs,  
smearing...



• constrain:

$$\langle x \rangle_{IC} = \int_0^1 dx x \cdot [c + \bar{c}](x)$$

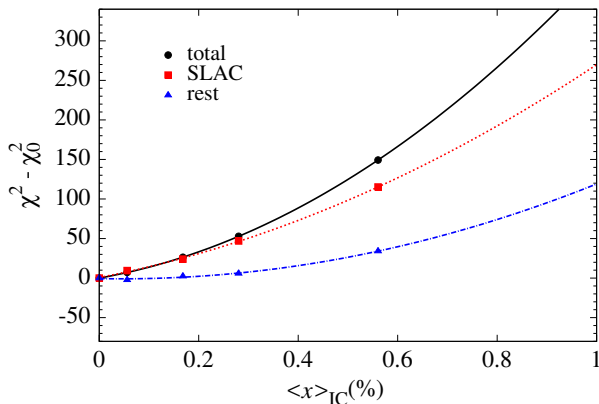
... 'total IC momentum'

...without EMC  $F_2^{c\bar{c}}$  ...SLAC *ep, ed* data!

$$\langle Q^2 \rangle \sim 15 \text{ GeV}^2$$

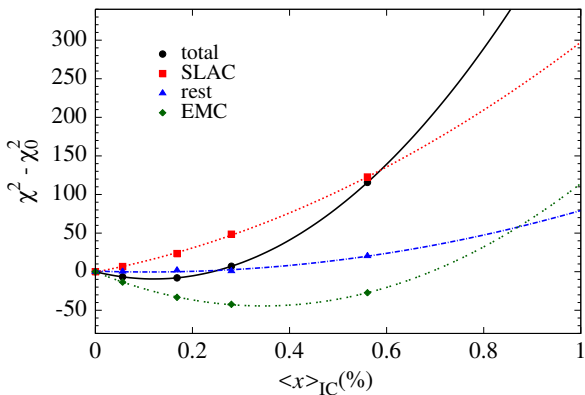
$$0.06 \leq x \leq 0.9$$

$$(\chi^2/N_{dat} \sim 1.25)$$

'SLAC + REST'  $\Rightarrow$   $\langle x \rangle_{IC} < 0.1\%$ ; at  $5\sigma$  !'REST' only  $\Rightarrow$   $\langle x \rangle_{IC} < 0.1\%$ ; at  $1\sigma$ cf.,  $\langle x \rangle_{IC} \sim 2 - 3\%$ 

e.g., [S. Dulat et al., Phys. Rev. D 89, 073004 (2014).]

N.B.: different tolerances:  $\Delta\chi^2 = 1$  vs.  $\Delta\chi_{CT}^2 = 100$

...and constrained by **EMC**

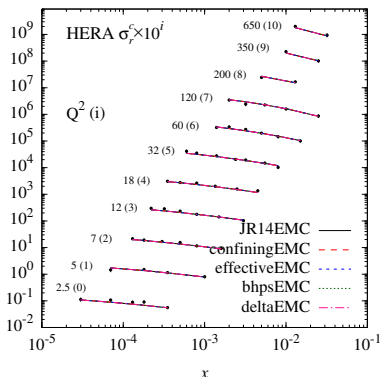
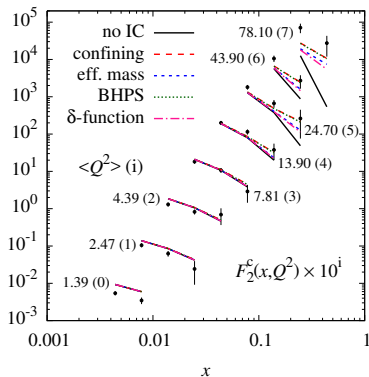
EMC alone:  $\langle x \rangle_{IC} = 0.3 - 0.4\%$

+ **SLAC**/**'REST'**:  $\langle x \rangle_{IC} = 0.13 \pm 0.04\%$

...but  $F_2^{c\bar{c}}$  poorly fit —  $\chi^2 \sim 4.3$  per datum!

# data comparisons:

...full fits, constrained by EMC  $F_2^{c\bar{c}}$  measurements:



- EMC: low- $x$ /low- $Q^2$  tension with HERA  $\sigma_r^c$
- $\frac{\tau_{life}}{\tau_{int}} = 5 \rightarrow$  for  $Q^2 = 170 \text{ GeV}^2$ , EMC sensitive to IC at  $x \lesssim 0.01$

$\rightarrow$  more  $F_2^{c\bar{c}}$  data are needed!

# new/ongoing global analyses

- **NNPDF3**: not anchored to specific parametrizations/models

see: Ball *et al.* 1605.06515

- *included* EMC:  $\langle x \rangle_{\text{IC}} = 0.7 \pm 0.3\%$  at  $Q \sim 1.5$  GeV
  - drove a very **hard**  $c(x) = \bar{c}(x)$  distribution
    - peaked at  $x \sim 0.5$
    - AND, required a **negative** IC component to describe EMC  $F_2^{c\bar{c}}$ !

- complementary analyses for possible intrinsic **bottom**

see: Lyonnet *et al.* JHEP07 (2015) 141.

→ would be negligible based on the analysis presented here...

# future **experimental** prospects?

- jet hadroproduction:  $pp \rightarrow (Zc) + X$  at **LHCb**

*e.g.*, Boettcher, Ilten, Williams, PRD**93**, 074008 (2016).

→ a “direct” measure in the forward region,  $2 < \eta < 5$   
 ... sensitive to  $c(x)$ ,  $x \sim 1$  for *one* colliding proton

→ can discriminate  $\langle x \rangle_{\text{IC}} \gtrsim 0.3\%$  (“valencelike”), 1% (“sealike”)

- **prompt atmospheric neutrinos?**

see: Laha & Brodsky, 1607.08240 (2016).

→ IceCube  $\nu$  spectra may constrain IC normalization

- possible impact upon **hidden charm pentaquark**,  $P_c^+$ ?

*e.g.*, Schmidt & Siddikov, PRD**93**, 094005 (2016).

- **AFTER@LHC?** ... fixed-target  $pp$  at  $\sqrt{s} = 115$  GeV

Brodsky *et al.* Adv. High Energy Phys. **2015**, 231547 (2015).

# outlook

- have **model framework** for IC; contacts the  $SU(4)$  spectrum
  - ... when constrained by  $pp \rightarrow \Lambda_c X$  data, *overpredicts*  $F_2^{c\bar{c}}$ 
    - generates “smoking gun”  $c(x) \neq \bar{c}(x)$  signal
- modeling necessitated a **QCD global analysis**:
  - severely limits  $\langle x \rangle_{IC} < 0.1\%$ ,  $5\sigma$  (without EMC)
  - with EMC,  $\langle x \rangle_{IC} = 0.13 \pm 0.04\%$
  - improved measurements at large  $x$  would be definitive (e.g., by **fixed target EICs**!)
- modeling/numerical analyses have reached an advanced stage
  - more **experimental information** is required, and diverse channels are available

... thank you ...