

Parity doubling of nucleons, Δ and Ω baryons across the deconfinement phase transition

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[1607.05082v1]
[PRD 92 (2015) 014503] [1502.03603v2]

Motivation

$m_q = 0 \Rightarrow$ **chiral symmetry** of QCD action

$$\psi' = e^{i\alpha\gamma_5 T_i} \psi, \quad \bar{\psi}' = \bar{\psi} e^{i\alpha\gamma_5 T_i}$$

T_i generators of $SU(N_f)$, $i = 1, \dots, N_f^2 - 1$

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Positive and negative parity baryonic correlators (zero momentum)

$$C_{\pm}(\tau) = \int d\mathbf{x} \langle \text{tr} O(\mathbf{x}, \tau) P_{\pm} \bar{O}(\mathbf{0}, 0) \rangle, \quad P_{\pm} = \frac{1}{2}(\mathbb{1} \pm \gamma_4)$$

For nucleon $O(\mathbf{x}, \tau) = \epsilon_{abc} u_a(\mathbf{x}, \tau) \left(u_b^T(\mathbf{x}, \tau) C \gamma_5 d_c(\mathbf{x}, \tau) \right)$

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$$C_{\pm}(\tau) \approx A_{\pm} e^{-M_{\pm}\tau} + A_{\mp} e^{-M_{\mp}(a_{\tau}N_{\tau}-\tau)}$$

Chiral symmetry $\Rightarrow C_+ = -C_- \Rightarrow M_+ = M_-$

Motivation

In Nature ($T = 0$) $M_{N^*} - M_N \approx 600 \text{ MeV} \gg m_{u,d} \approx 5 \text{ MeV}$

- Explicit chiral symmetry breaking ($m_{u,d} \neq 0$) is not enough to account for this big difference

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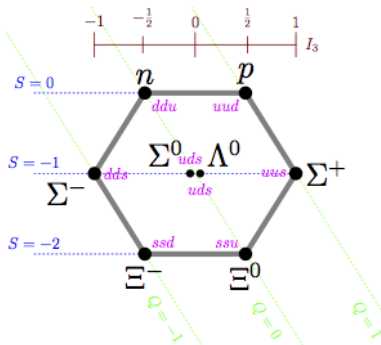
- What happens at high temperature?
- Parity restoration above T_c for N and Δ baryons

Even if chiral symmetry is slightly explicitly broken by $m_{u,d}$ and lattice artefacts
Wilson fermions \rightarrow No chiral symmetry at short distances

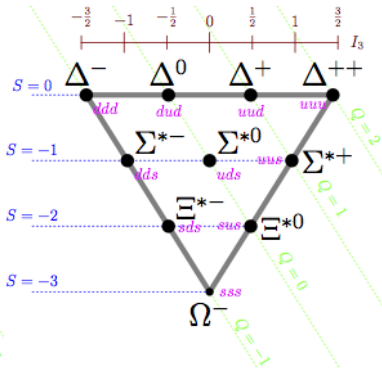
- Signal of parity restoration for Ω around T_c
Chiral symmetry is strongly explicitly broken by $m_s \approx 100 \text{ MeV}$

Name	N	Δ	Λ	Σ	Ξ	Ω
Isospin	$1/2$	$3/2$	0	1	$1/2$	0
Strangeness		0	-1	-1	-2	-3
Number of s-quarks		0	1	1	2	3

Spin 1/2 octet



Spin 3/2 decuplet



Lattice setup

FASTSUM ensembles and tuning by HadSpec collaboration

- $N_f = 2 + 1$ non-perturbatively improved Wilson fermions
- Anisotropic lattice: $a_s/a_\tau = 3.5$, $a_\tau^{-1} \approx 5.6$ GeV
→ Important for constructing spectral functions
- $T = \frac{1}{a_\tau N_\tau}$ varies by changing N_τ from 128 to 16
- Large volume of the box $\sim (3 \text{ fm})^3$, $N_s = 24$
- Degenerate u and d quarks, heavier than physical ones
($m_\pi = 384(4)$ MeV, $m_\pi/m_\rho = 0.466(3)$)
- Physical strange quark mass
- Gaussian smearing on both source and sink to enhance ground state signal

R factor for measuring parity doubling

$$R(\tau) \equiv \frac{C(\tau) - C(1/T - \tau)}{C(\tau) + C(1/T - \tau)}$$

- No parity doubling and $M_- \gg M_+ \Rightarrow R(\tau) = 1$,
 $0 \leq \tau < 1/(2T)$
- Parity doubling $\Rightarrow R(\tau) = 0$
- Note that $R(1/T - \tau) = -R(\tau)$ and $R(1/(2T)) = 0$

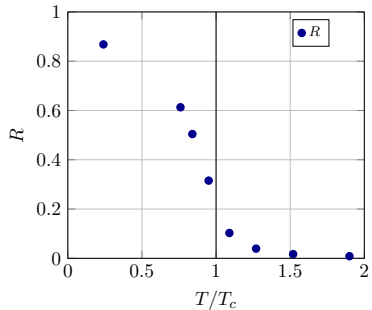
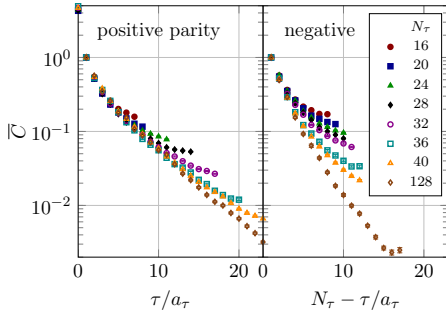
We consider the weighted average

[Datta, Mathur et al. (2013)]

$$R \equiv \frac{\sum_{n=1}^{N_\tau/2-1} R(\tau_n)/\sigma^2(\tau_n)}{\sum_{n=1}^{N_\tau/2-1} 1/\sigma^2(\tau_n)}$$

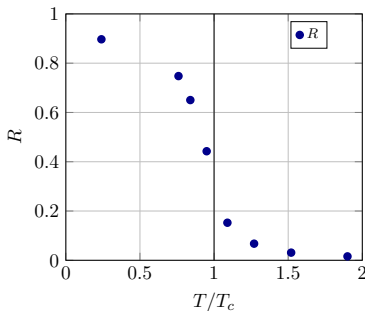
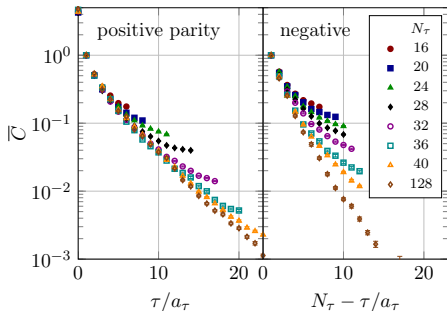
Technical note: Smearing essential to have a clear ground state

Nucleon (spin 1/2)



- Nucleon ground state largely independent of temperature
- Negative parity partner much more sensitive to temperature
- Strong signal of parity restoration around T_c

Δ -baryon (spin 3/2)



- Δ baryon ground state largely independent of temperature
- Negative parity partner much more sensitive to temperature
- Strong signal of parity restoration around deconfinement transition (tied to restoration of chiral symmetry)

Spectral functions

For baryons:

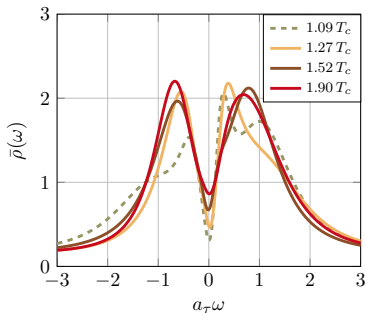
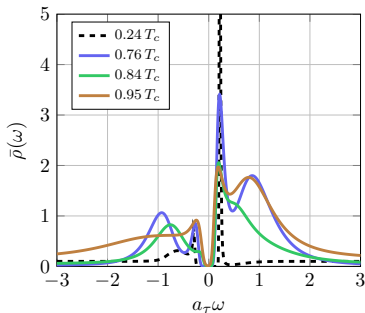
$$C_{\pm}(\tau, \mathbf{p}) = \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} \rho_{\pm}(\omega, \mathbf{p}) \frac{e^{-\omega\tau}}{1 + e^{-\omega/T}}, \quad \rho_{\pm} = \text{tr}[P_{\pm}\rho]$$

- *Ill-posed problem*: To extract $\sim 10^3$ points for $\rho_{\pm}(\omega, \mathbf{p} = \mathbf{0})$ given ~ 50 noisy data for $C_{\pm}(\tau, \mathbf{p} = \mathbf{0})$
- The **Maximum Entropy Method** (MEM) is an unbiased method to get a unique solution for ρ_{\pm} [Asakawa et al. hep-lat/0011040v2]

Important property for MEM: $\rho_{+}(\omega, \mathbf{p}) \geq 0 \quad \forall \omega, \mathbf{p}$
($\rho_{\pm}(-\omega, -\mathbf{p}) = -\rho_{\mp}(\omega, \mathbf{p})$)

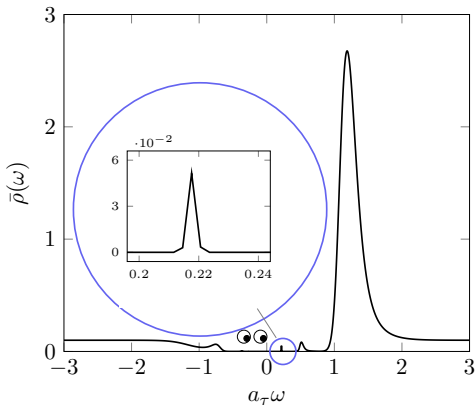
Nucleon (spin 1/2)

$$\bar{\rho}(\omega) \equiv \frac{1}{a_\tau} \frac{\rho(\omega)}{\langle C(\tau=0) \rangle_{\text{cfg}}}$$



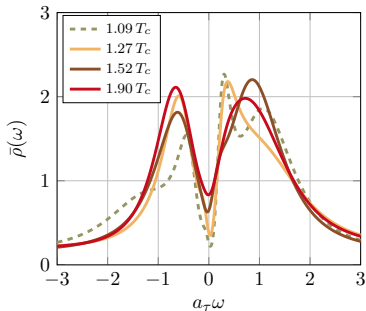
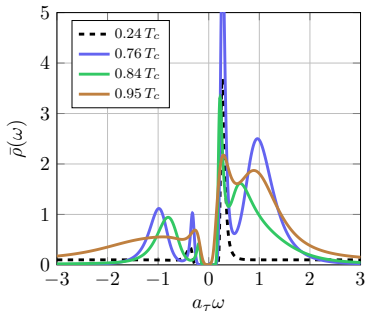
- $\omega > 0$ (+ parity): Very stable ground state below T_c
- $\omega < 0$ (- parity): Ground state moves inwards as $T \rightarrow T_c$
- Very symmetric spectral functions above T_c

Nucleon ground state without smearing ($T = 0.24T_c$)



The ground state is at the right place but we need smearing to enhance its signal

Δ -baryon (spin 3/2)

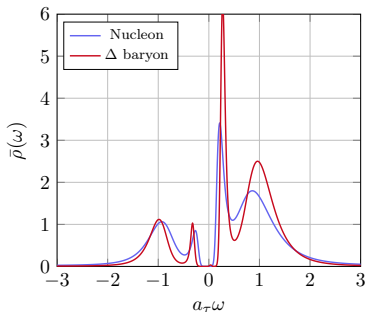
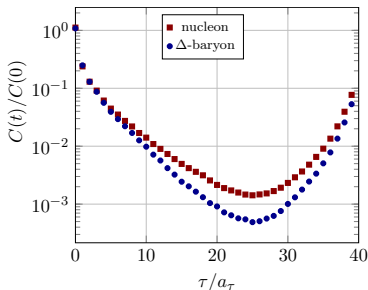


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N and Δ below T_c

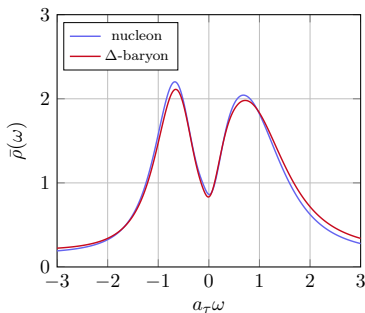
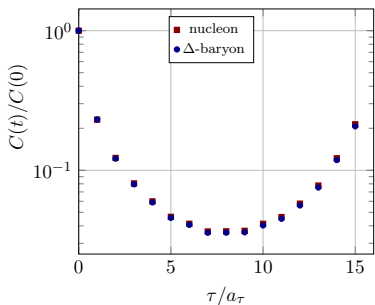
For ground-states ($I = 0$): $M_\Delta - M_N = \Delta M_{ss} = \frac{8}{3} \left(\frac{\hbar}{c}\right)^3 \frac{\pi \alpha_s}{m_{u,d}^2} |\psi(0)|^2$

$$\Delta M_{ss}^{+parity} = \begin{cases} (293 \pm 2) \text{ MeV} & \text{Nature } (m_\pi = 140 \text{ MeV}, T = 0) \\ (274 \pm 96) \text{ MeV} & \text{Lattice } (m_\pi = 384 \text{ MeV}, T = 44 \text{ MeV}) \end{cases}$$



Asymmetric correlators and spectral functions

N and Δ above T_c

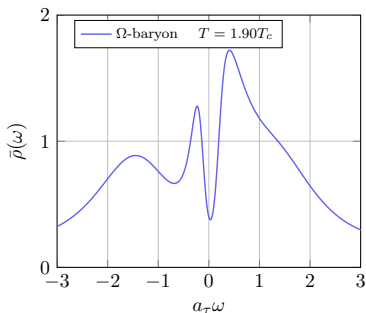
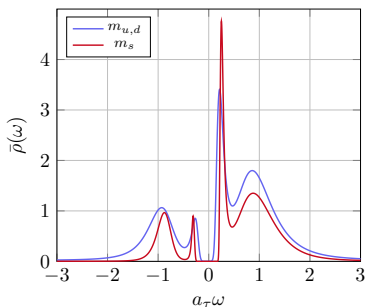


- Symmetric correlators and spectral functions (Parity restoration)
- Same correlators and spectral functions for N and Δ (GS melted)

N and Ω

below T_c

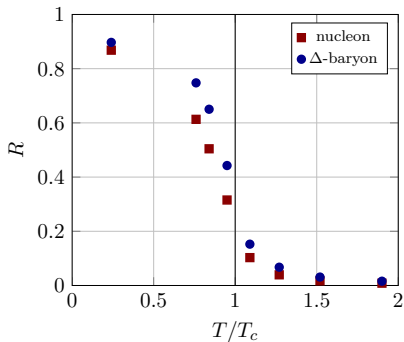
above $T_c \rightarrow$ Parity not yet restored



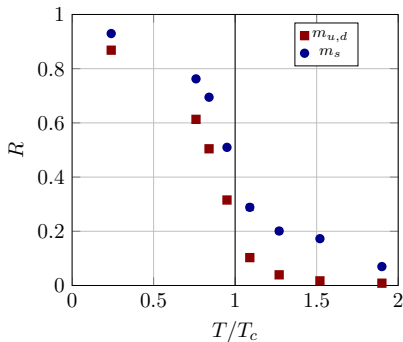
T/T_c	m^+ [MeV]	m^- [MeV]			
0	1672.4(0.3)	2250?	2380?	2470?	PDG
0.24	1703(159)	2232(380)			

[See also HadSpec Collaboration, 0810.3588v1]

Nucleon vs Δ



Nucleon vs Ω



Parity restoration

Signal of parity doubling

Both signals occur around T_c

Conclusions and perspectives

Summary

- Negative parity channel more affected by temperature than positive parity channel
- Parity restoration above T_c for N and Δ baryons
- Signal of parity doubling for Ω at T_c
- Chiral symmetry is strongly explicitly broken by m_s

Outlook

- To use chiral (overlap) fermions
- Finer lattice spacing