

Yang-Mills correlation functions at non-zero temperature



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XIIth Quark Confinement and the Hadron Spectrum - Thessaloniki, Greece

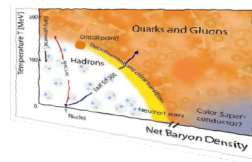
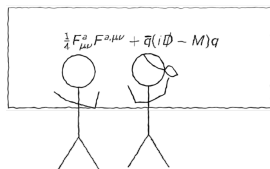
Sept. 2, 2016



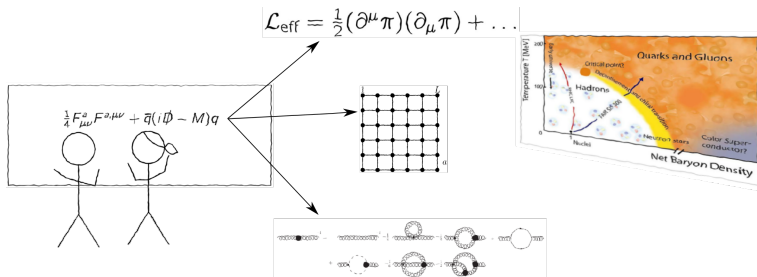
QCD: A simple theory?

$$\frac{1}{4} F_{\mu\nu}^a F^{a,\mu\nu} + \bar{q}(i\cancel{D} - M)q$$

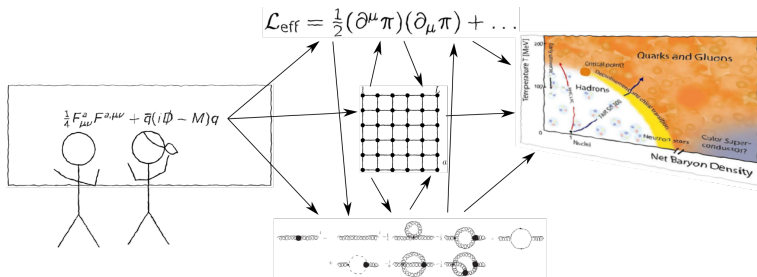
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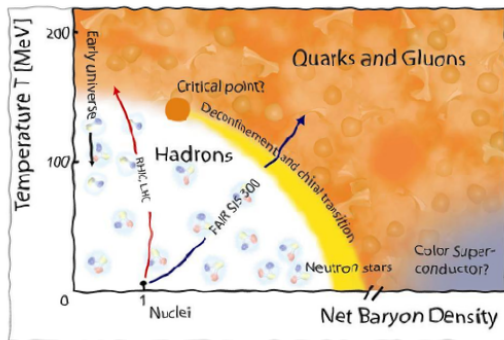
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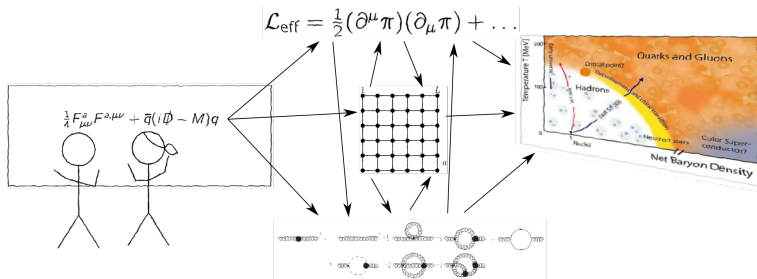


QCD: A simple theory?



- **Phases:** hadronic phase, quark-gluon plasma, color superconductor, quarkyonic?
- **Transitions:** first order line, crossover at $\mu = 0$
- **Critical point:** existence? position?

QCD: A simple theory?



- Challenges for all methods at $\mu > T$, e.g.
 - Lattice QCD: complex action problem
 - Models: parameters
 - Functional methods: reliability of truncations

Functional methods

Functional equations: Exact equations derived from QCD action.

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Dyson-Schwinger eqs.:
eqs. of motion of corr.
functions



$$i\frac{\partial}{\partial m} \text{[quark propagator]} = \frac{1}{2} \text{[gluon loop]} + \text{[ghost loop]} - 2 \text{[ghost-gluon loop]}$$

funct. renorm. group eqs.



eqs. of motion from 3PI eff. action

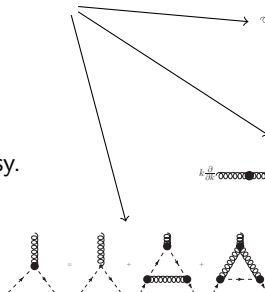
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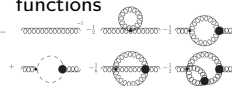
- Chiral limit accessible.
- No sign problem.
- Large scale separations easy.
- Real-time quantities accessible

→ [talk by Pawłowski]

$$\frac{1}{4} F_{\mu\nu}^a F^{a,\mu\nu} + \bar{q}(i\not{D} - M)q$$



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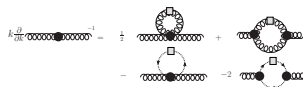
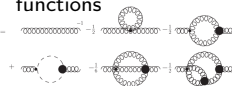
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Dyson-Schwinger eqs.:
eqs. of motion of corr.
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funct. renorm. group eqs.

Difficulty

Infinitely large systems of equations without obvious ordering scheme.

Outline

- 1 Introduction
- 2 Dyson-Schwinger equations and truncations
- 3 Testing truncations in $d = 3$
- 4 Non-vanishing temperature results

Landau gauge QCD

$$\mathcal{L} = \bar{\mathbf{q}}(-\not{D} + m)\mathbf{q} + \frac{1}{2}F^2 + \mathcal{L}_{gf} + \mathcal{L}_{gh}$$

$$F_{\mu\nu} = \partial_\mu \mathbf{A}_\nu - \partial_\nu \mathbf{A}_\mu + ig[\mathbf{A}_\mu, \mathbf{A}_\nu]$$



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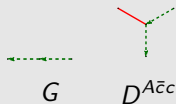


Landau gauge

- simplest one for functional equations

- $\partial_\mu \mathbf{A}_\mu = 0$: $\mathcal{L}_{gf} = \frac{1}{2\xi}(\partial_\mu \mathbf{A}_\mu)^2$, $\xi \rightarrow 0$

- requires ghost fields: $\mathcal{L}_{gh} = \bar{c}(-\square + g \mathbf{A} \times) c$



The tower of DSEs

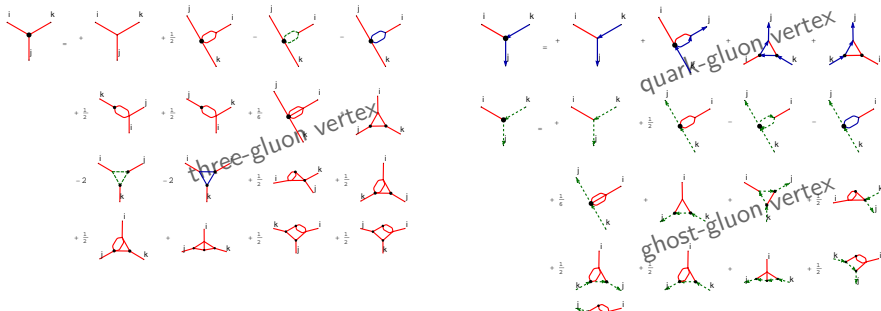
$$\begin{aligned}
 & \text{red line with dot} \stackrel{-1}{=} + \text{red line} \stackrel{-1}{=} - \frac{1}{2} \text{red line with red loop} - \frac{1}{2} \text{red line with red loop and dot} + \text{red line with green loop} \\
 & \quad + \text{red line with blue loop} \stackrel{-1}{=} - \frac{1}{6} \text{red line with red loop and dot} - \frac{1}{2} \text{red line with red loop and dot} \quad \text{gluon propagator} \\
 & \text{dashed green line with dot} \stackrel{-1}{=} + \text{dashed green line} \stackrel{-1}{=} - \text{dashed green line with red loop and dot} \quad \text{ghost propagator} \\
 & \text{blue line with dot} \stackrel{-1}{=} + \text{blue line} \stackrel{-1}{=} - \text{blue line with red loop and dot} \quad \text{quark propagator}
 \end{aligned}$$

The tower of DSEs

$$\begin{aligned}
 \text{gluon propagator} &= \text{tree} + \text{self-energy} - \frac{1}{2} \text{ghost loop} - \frac{1}{2} \text{quark loop} + \text{ghost-gluon vertex} \\
 &+ \text{ghost-gluon vertex} - \frac{1}{6} \text{ghost loop} - \frac{1}{2} \text{quark loop}
 \end{aligned}$$

$$\text{ghost propagator} = \text{tree} - \text{ghost loop}$$

$$\text{quark propagator} = \text{tree} - \text{quark loop}$$

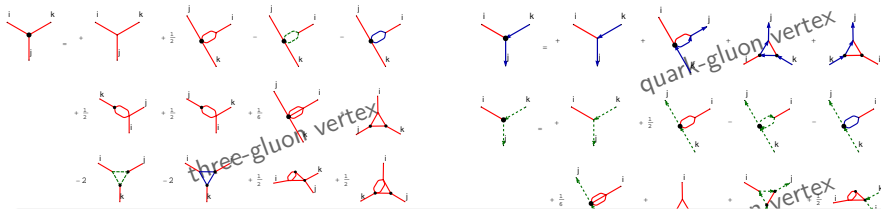


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 &+ \text{ghost-gluon-gluon loop} - \frac{1}{6} \text{ghost-gluon-gluon-gluon loop} - \frac{1}{2} \text{ghost-gluon-gluon-gluon-gluon loop}
 \end{aligned}$$

$$\text{ghost propagator} = \text{tree} - \text{ghost loop}$$

$$\text{quark propagator} = \text{tree} - \text{quark-gluon loop}$$



Infinitely many equations. In QCD, every n -point function depends on $(n + 1)$ - and possibly $(n + 2)$ -point functions.



Truncating the equations

Guides

- Perturbation theory
- Symmetries
- Lattice
- Analytic results

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Truncation

- Drop quantities (unimportant?)
- Use fits
- Model quantities (good models available? 'true' or 'effective'?)

Ideally: Find a truncation that has (I) **no parameters** and yields (II) **quantitative results**.

But...

...how do we know that the results are trustworthy?

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Lattice results

Available for

- Vacuum

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- Four-point functions?

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Available for

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- Propagators
- $T > 0$
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- $\mu > 0$?
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→ Comparison with lattice is helpful, but finally self-consistent checks are required.

Two words of caution:

- One cannot assume naturally that the hierarchy is the same for all T and μ .
- In simple truncations, the effect of a single correlation function is difficult to estimate.

Top-down for Yang-Mills theory

Top-down vs. bottom up



Green functions from QCD action vs. effective models

Top-down for Yang-Mills theory

Neglect all non-primitively divergent Green functions. \rightarrow Self-contained.

Full propagator equations (two-loop diagrams!):

$$\begin{aligned}
 & \text{Red line with dot} \stackrel{-1}{=} + \text{Red line} \stackrel{-1}{=} - \frac{1}{2} \text{Red line with self-energy loop} - \frac{1}{2} \text{Red line with ghost loop} + \text{Red line with ghost loop} \\
 & \quad - \frac{1}{6} \text{Red line with ghost loop} - \frac{1}{2} \text{Red triangle} \\
 & \text{Green dashed line with dot} \stackrel{-1}{=} + \text{Green dashed line} \stackrel{-1}{=} - \text{Green dashed line with self-energy loop}
 \end{aligned}$$

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 &\quad - \frac{1}{6} \text{Red line with ghost loop} - \frac{1}{2} \text{Red line with ghost loop}
 \end{aligned}$$

$$\text{Dashed line with dot} = \text{Dashed line} - \text{Dashed line with ghost loop}$$

Truncated three-point functions:

$$\begin{aligned}
 \text{Truncated 3-point} &= \text{Truncated 3-point} + \text{Truncated 3-point} + \text{Truncated 3-point} + \text{Truncated 3-point} \\
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 &\quad + \frac{1}{2} \text{Truncated 3-point} + \text{Truncated 3-point} - 2 \text{Truncated 3-point}
 \end{aligned}$$

Truncated four-gluon vertex:

$$\begin{aligned}
 \text{Truncated 4-gluon vertex} &= \text{Truncated 4-gluon vertex} + \frac{1}{2} \text{Truncated 4-gluon vertex} + 3 \text{Truncated 4-gluon vertex} \\
 &\quad + 3 \text{Truncated 4-gluon vertex} + 3 \text{Truncated 4-gluon vertex} - 6 \text{Truncated 4-gluon vertex}
 \end{aligned}$$

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 &- \frac{1}{6} \text{---} \text{---} - \frac{1}{2} \text{---} \text{---}
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 &+ \frac{1}{2} \text{---} \text{---} + \text{---} \text{---} - 2 \text{---} \text{---}
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Truncated four-gluon vertex:

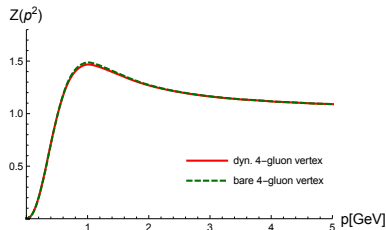
$$\begin{aligned}
 \text{---} \text{---} &= + \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} + 3 \text{---} \text{---} \\
 &+ 3 \text{---} \text{---} + 3 \text{---} \text{---} - 6 \text{---} \text{---}
 \end{aligned}$$

Testing top-down for Yang-Mills in $d = 3$

[MQH '16]

Quantitative study of truncation effects possible

- Varying the four-gluon vertex: bare vs. dynamic

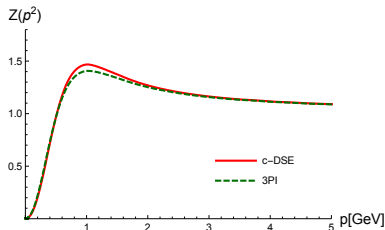


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Conclusions from $d = 3$

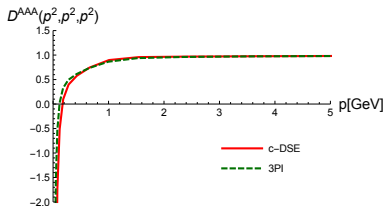
- Importance of two-loop diagrams in propagator, less in vertices

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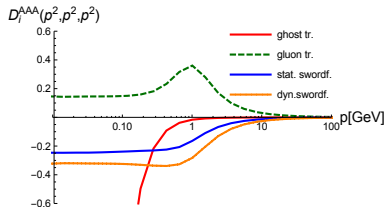
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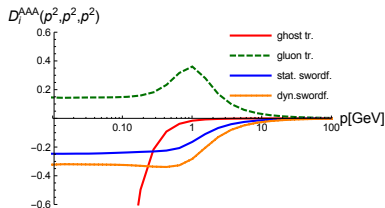
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Similar truncation in FRG ($d = 4$):

[Cyrol, Mitter, Strodthoff, Pawłowski '16; talks by Cyrol, Mitter]



Conclusions from $d = 3$

- Importance of two-loop diagrams in propagator, less in vertices
- Small deviations of vertices from tree-level
- Cancellations in gluonic vertices: large+large=small

What do we need to go beyond modern QCD phase diagram calculations?

Beyond effective interaction approximation: ✓ [Fischer, Lücker, Welzbacher '14]

Input for DSEs:

- model for quark-gluon vertex
- fits for gluon propagators at $\mu = 0$ from the lattice

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Possible improvements:

- fully dynamical propagators
- fully dynamical quark-gluon vertex

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- fully dynamical quark-gluon vertex \rightarrow requires propagators & other vertices

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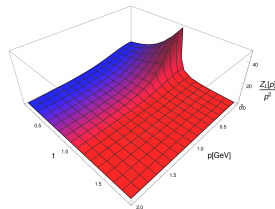
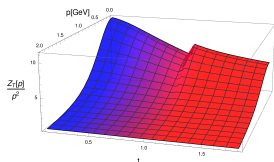
Ultimately, **full control** over Yang-Mills part required!

Non-vanishing temperature

Elementary Green functions \rightarrow (dual) quark condensate, Polyakov loop potential, ...

Propagators

Lattice results \rightarrow input to calculate other quantities [Fischer, Maas, Müller '10].



Fits based on [Maas, Pawłowski, von Smekal, Spielmann '12].

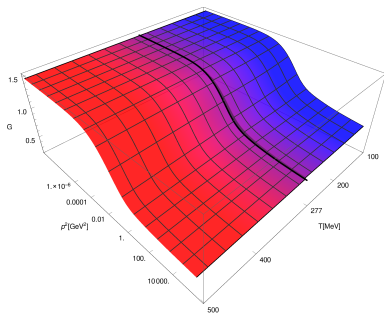
Note

Gluon propagator from lattice has no truncation artifacts.

Example: Ghost DSE

Ghost dressing $G(p^2)$ from DSE [MQH, von Smekal '13]:

$$\text{---}\bullet\text{---} \stackrel{-1}{=} \text{---}\text{---} \stackrel{-1}{-} \left(\text{---}\text{---}\text{---} \right) \text{---} - \left(\text{---}\text{---}\text{---} \right) \text{---}$$

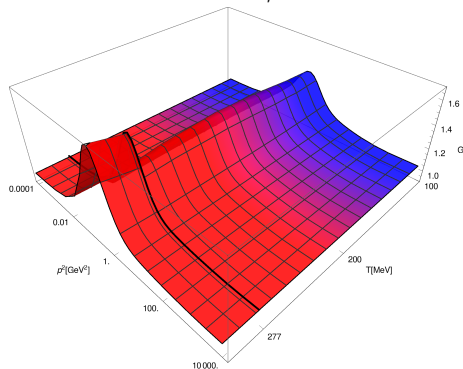


- Ghost insensitive to phase transition.

Ghost-gluon vertex

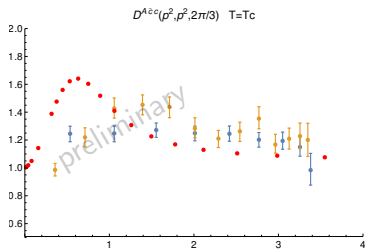
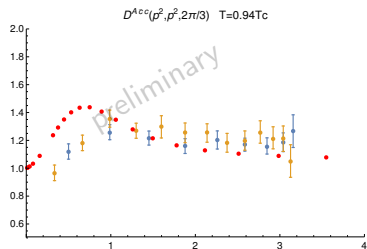
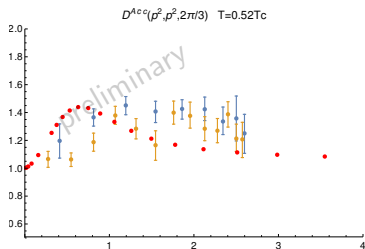
- Vertices on lattice more difficult than propagators.
- Full momentum dependence from functional equations.

Self-consistent solution, zeroth Matsubara only



Vertex from FRG: [Fister, Pawłowski '11; talk by Cyrol]

Ghost-gluon vertex: Continuum and lattice

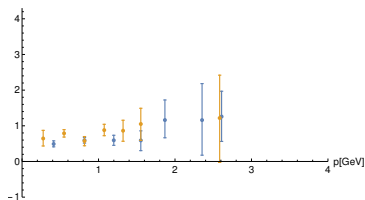
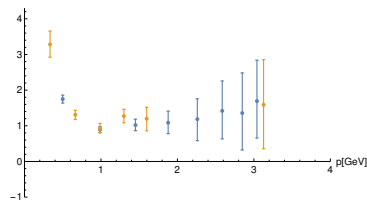
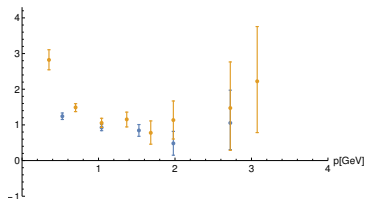
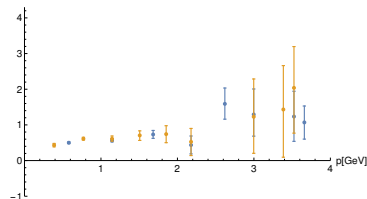


Lattice: [Fister, Maas '14]

Three-gluon vertex: Continuum and lattice

Note: IR suppression observed at $T = 0$. Zero crossing? [talk by Sternbeck]

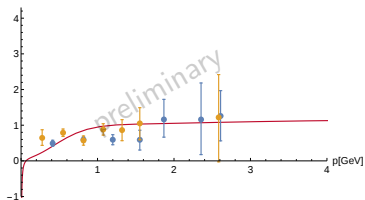
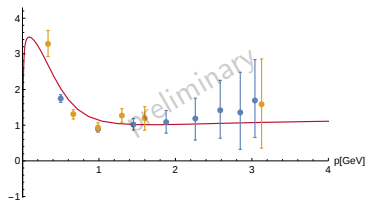
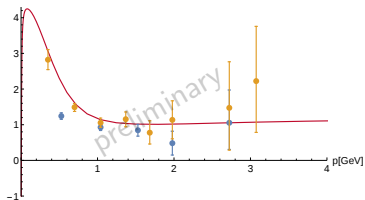
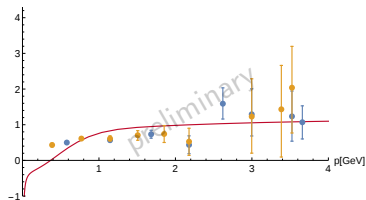
Three-gluon vertex: Continuum and lattice

 $D^{AAA}(p^2, p^2, 2\pi/3)$ $T=0.52T_c$

 $D^{AAA}(p^2, p^2, 2\pi/3)$ $T=0.94T_c$

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 $D^{AAA}(p^2, p^2, 2\pi/3)$ $T=1.08T_c$


Lattice: [Fister, Maas '14]

→ Enhancement around T_c , but still zero crossing.

Three-gluon vertex: Continuum and lattice

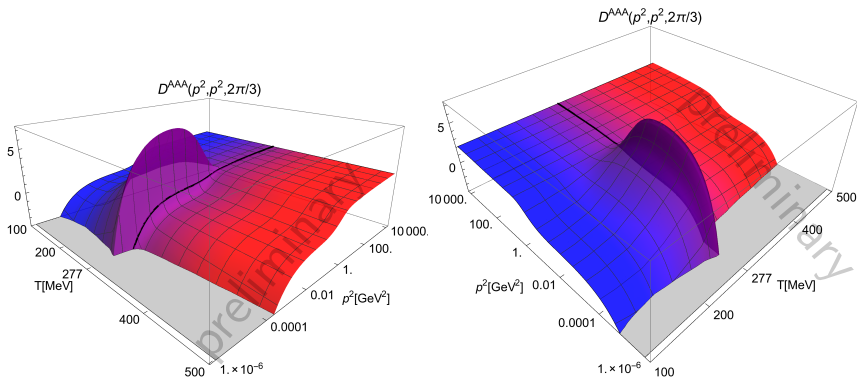
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Three-gluon vertex

DSE calculation: semi-perturbative approximation (first iteration only)



Summary and conclusions

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- Calculations of propagators, vertices and partially mixed systems show a **coherent picture** at $T = 0$.
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Thank you for your attention.