

# Introduction: collective modes and exotic nuclei

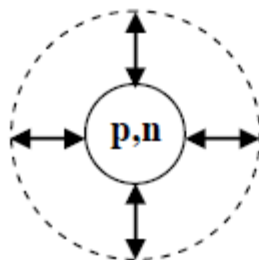
---

**Muhsin N. Harakeh**

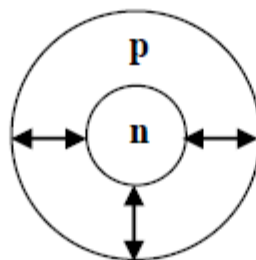
**KVI-CART, Groningen**

**Workshop on “Collective Mode Study  
through Beta-Decay Measurements”  
SUBATECH, Nantes; 19-20 January 2015**

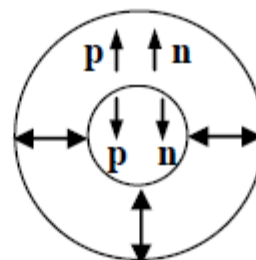
$\Delta L = 0$



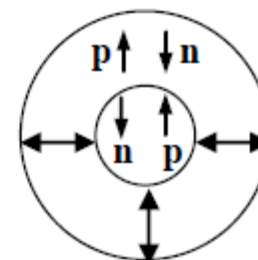
**ISGMR**



**IVGMR**



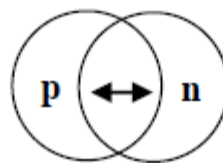
**ISSGMR**



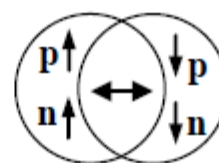
**IVSGMR**

$\Delta L = 1$

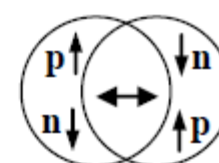
**ISGDR**  
**??**



**IVGDR**

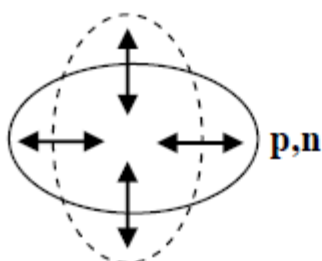


**ISSGDR**



**IVSGDR**

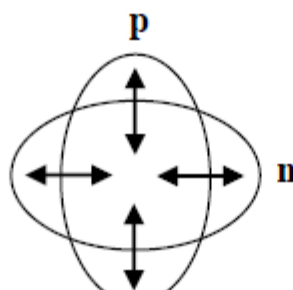
$\Delta L = 2$



**ISGQR**

$\Delta T = 0$

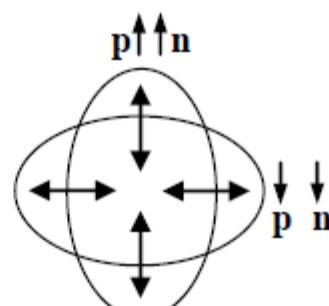
$\Delta S = 0$



**IVGQR**

$\Delta T = 1$

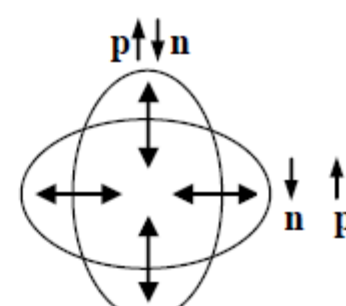
$\Delta S = 0$



**ISSGQR**

$\Delta T = 0$

$\Delta S = 1$



**IVSGQR**

$\Delta T = 1$

$\Delta S = 1$

**Microscopic picture: GRs are coherent (1p-1h) excitations induced by single-particle operators.**

- **Excitation energy depends on**
  - i) multipole  $L$  ( $L\hbar\omega$ , since radial operator  $\propto r^L$ ; except for ISGMR and ISGDR,  $2\hbar\omega$  &  $3\hbar\omega$ , respectively),*
  - ii) strength of effective interaction and*
  - iii) collectivity.*
- **Exhaust appreciable % of EWSR**
- **In addition to Landau damping, they acquire a width due to coupling to continuum and to underlying 2p-2h configurations.**

Nucleus  $\longrightarrow$  Many-body system with a finite size

Vibrations  $\longrightarrow$  Multipole expansion with  $r$ ,  $Y_{lm}$ ,  $\tau$ ,  $\sigma$

$\Delta S=0, \Delta T=0$     $\Delta S=0, \Delta T=1$     $\Delta S=0, \Delta T=1$     $\Delta S=1, \Delta T=1$     $\Delta S=1, \Delta T=1$

L=0: Monopole	<b>ISGMR</b> $r^2 Y_0$	<b>IAS</b> $\tau Y_0$	<b>IVGMR</b> $\tau r^2 Y_0$	<b>GTR</b> $\tau \sigma Y_0$	<b>IVSGMR</b> $\tau \sigma r^2 Y_0$
L=1: Dipole	<b>ISGDR</b> $(r^3 - 5/3 \langle r^2 \rangle r) Y_1$		<b>IVGDR</b> $\tau r Y_1$		<b>IVSGDR</b> $\tau \sigma r Y_1$
L=2: Quadrupole	<b>ISGQR</b> $r^2 Y_2$		<b>IVGQR</b> $\tau r^2 Y_2$		<b>IVSGQR</b> $\tau \sigma r^2 Y_2$
L=3: Octupole	<b>LEOR, HEOR</b> $r^3 Y_3$				



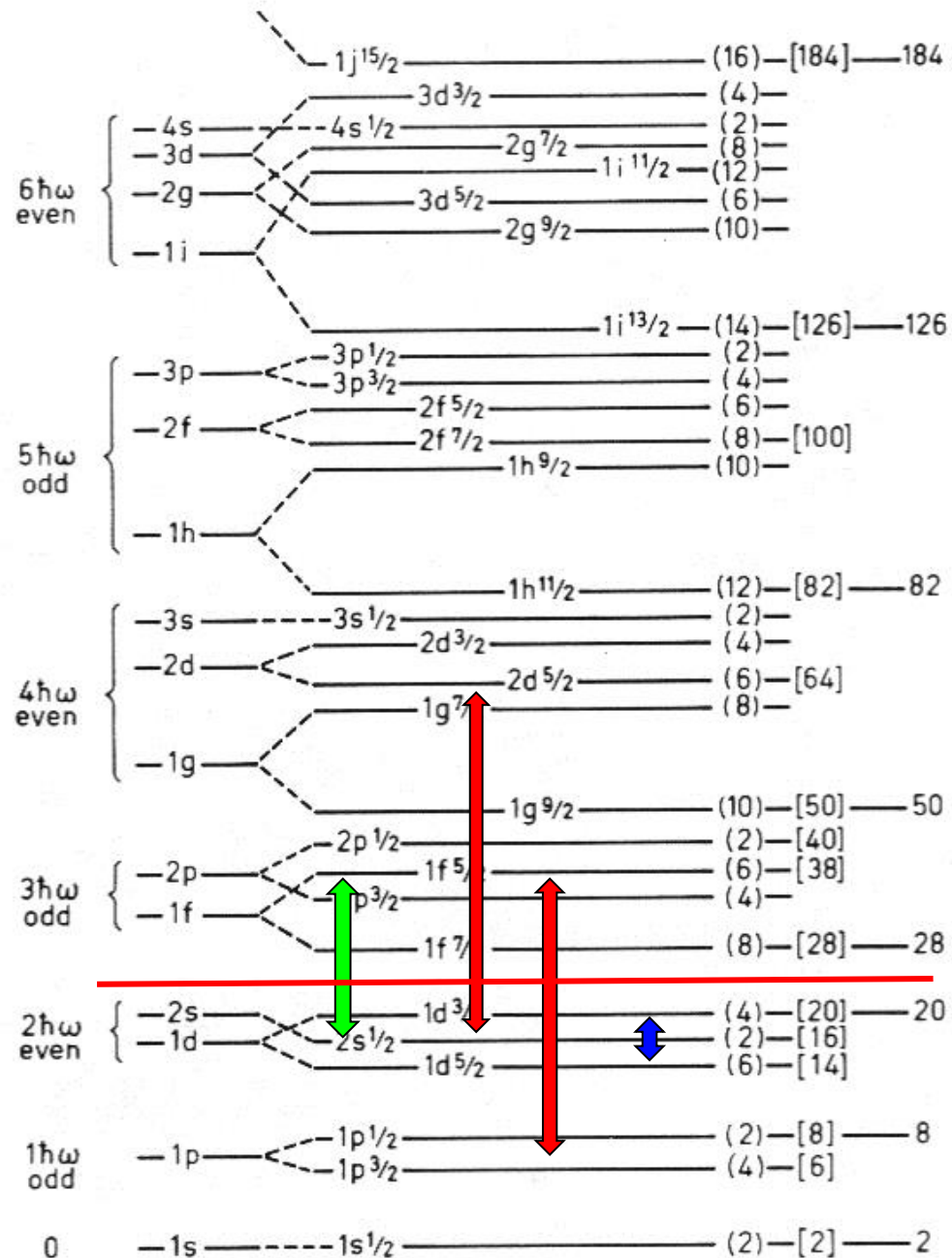
$\Delta N = 1$  E1 (IVGDR)



$\Delta N = 2$  E2 (ISGQR)  
&



$\Delta N = 0$  E0 (ISGMR)



# The collective response of the nucleus

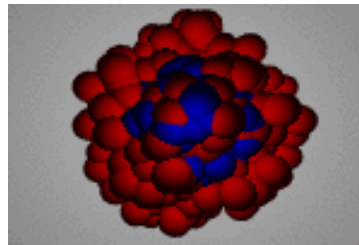
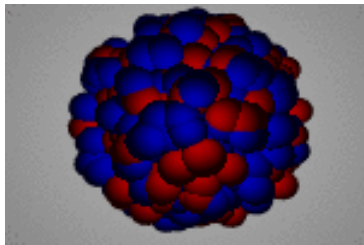
## Giant Resonances

### Electric giant resonances

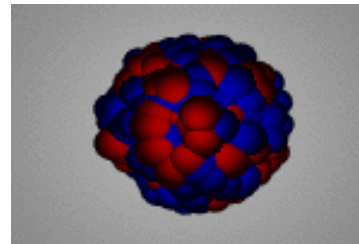
Isoscalar

Isovector

Monopole  
(GMR)



Dipole  
(GDR)



Quadrupole  
(GQR)

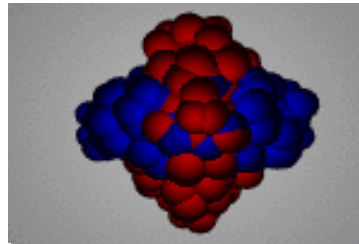
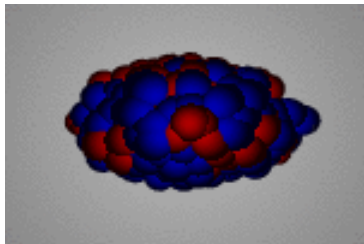
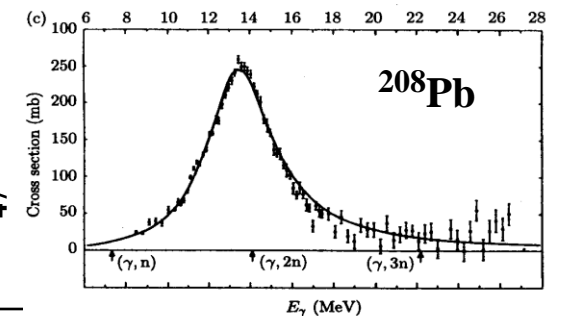
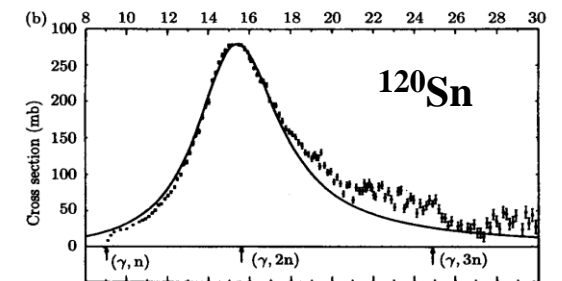
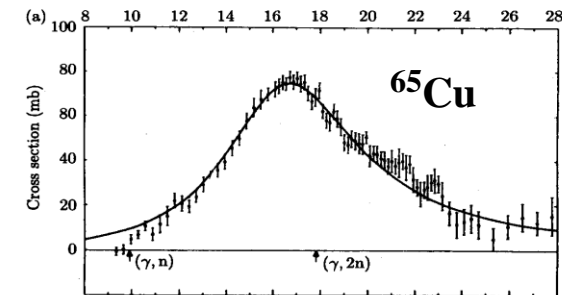


Photo-neutron  
cross sections



Berman and Fultz, Rev. Mod. Phys. 47 (1975)

47

Nantes; 19-20 January 2015

# Spin-isospin excitations

**Neutral ( $\nu, \nu'$ ) and charged ( $\nu_e, e^-$ ), ( $\nu_e, e^+$ ) currents**

**NC  $\Rightarrow$  Inelastic electron and proton scattering**

**$\Rightarrow$  M0, M1, M2**

**CC  $\Rightarrow$  Charge-exchange reactions**

**Isovector charge-exchange modes**

**$\Rightarrow$  IAS, GTR, IVSGMR, IVSGDR, etc.**

**Importance for nuclear astrophysics,**

**$\nu$ -physics,  $2\beta$ -decay, n-skin thickness, etc.**

**( $p, n$ ), ( ${}^3\text{He}, t$ ) {GT $^-$ }; ( $n, p$ ), ( $d, {}^2\text{He}$ ) & ( $t, {}^3\text{He}$ ) {GT $^+$ }**

# Charge-exchange probes

**$(p,n)$ -type ( $\Delta T_z = -1$ )**

- $\beta^-$ -decay
- $(p,n)$
- $(^3\text{He},t)$
- heavy ion

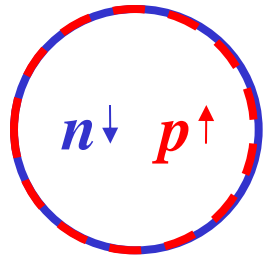
**$(n,p)$ -type ( $\Delta T_z = +1$ )**

- $\beta^+$ -decay
- $(n,p)$
- $(d,^2\text{He})$
- $(t,^3\text{He})$
- heavy ion; ( $^7\text{Li}, ^7\text{Be}$ )

- Energy per nucleon ( $>100$  MeV/u)
- Spin-flip versus non-spin-flip
- Complexity of reaction mechanism
- Experimental considerations



# Spin-isospin excitations



$\Delta L=0$   $\Delta S=1$   $\Delta T=1$   
GTR

- Gamow-Teller transitions;  
Isospin ( $\Delta T=1$ ); Spin ( $\Delta S=1$ )

## Advantages

- Cross section peaks at  
 $\theta = 0^\circ$  ( $\Delta L=0$ )
- Strong excitation of GT states  
at  $E/A=100\text{-}500$  MeV/u

⇒ At RCNP, Osaka

$E(^3\text{He}) \approx 150$  MeV/u

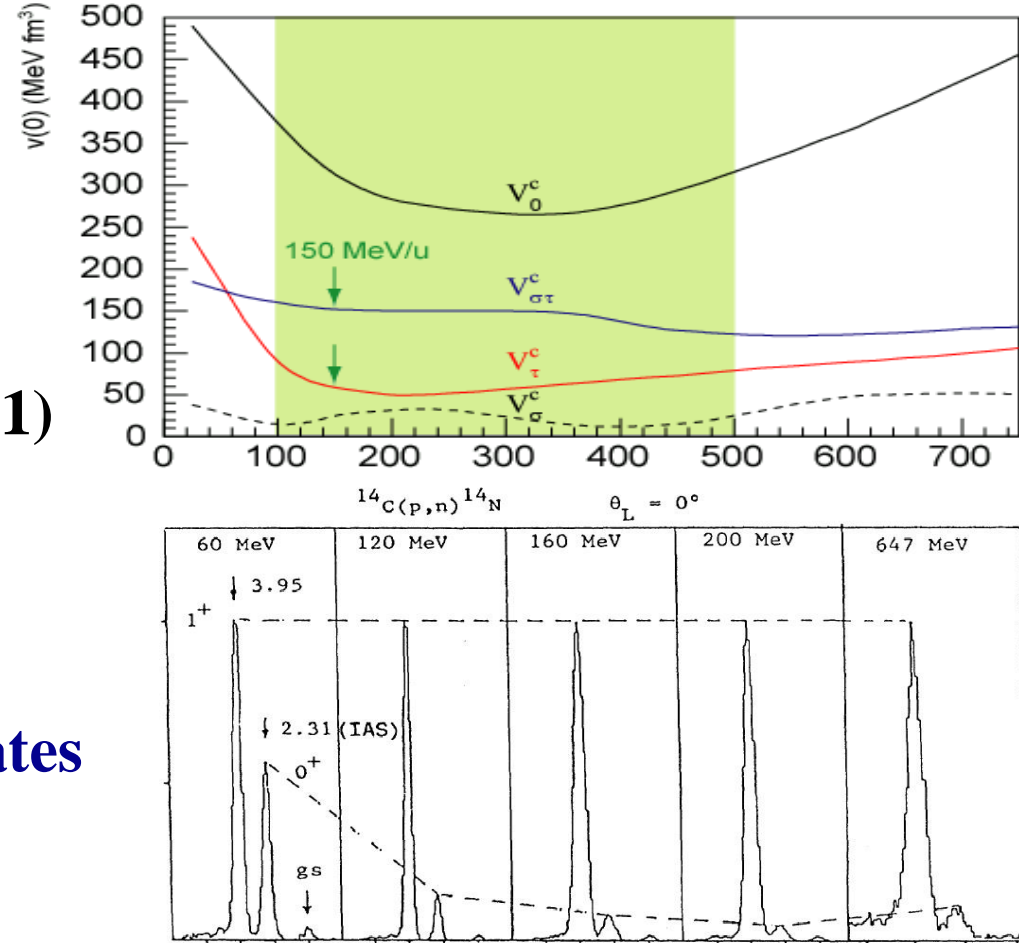
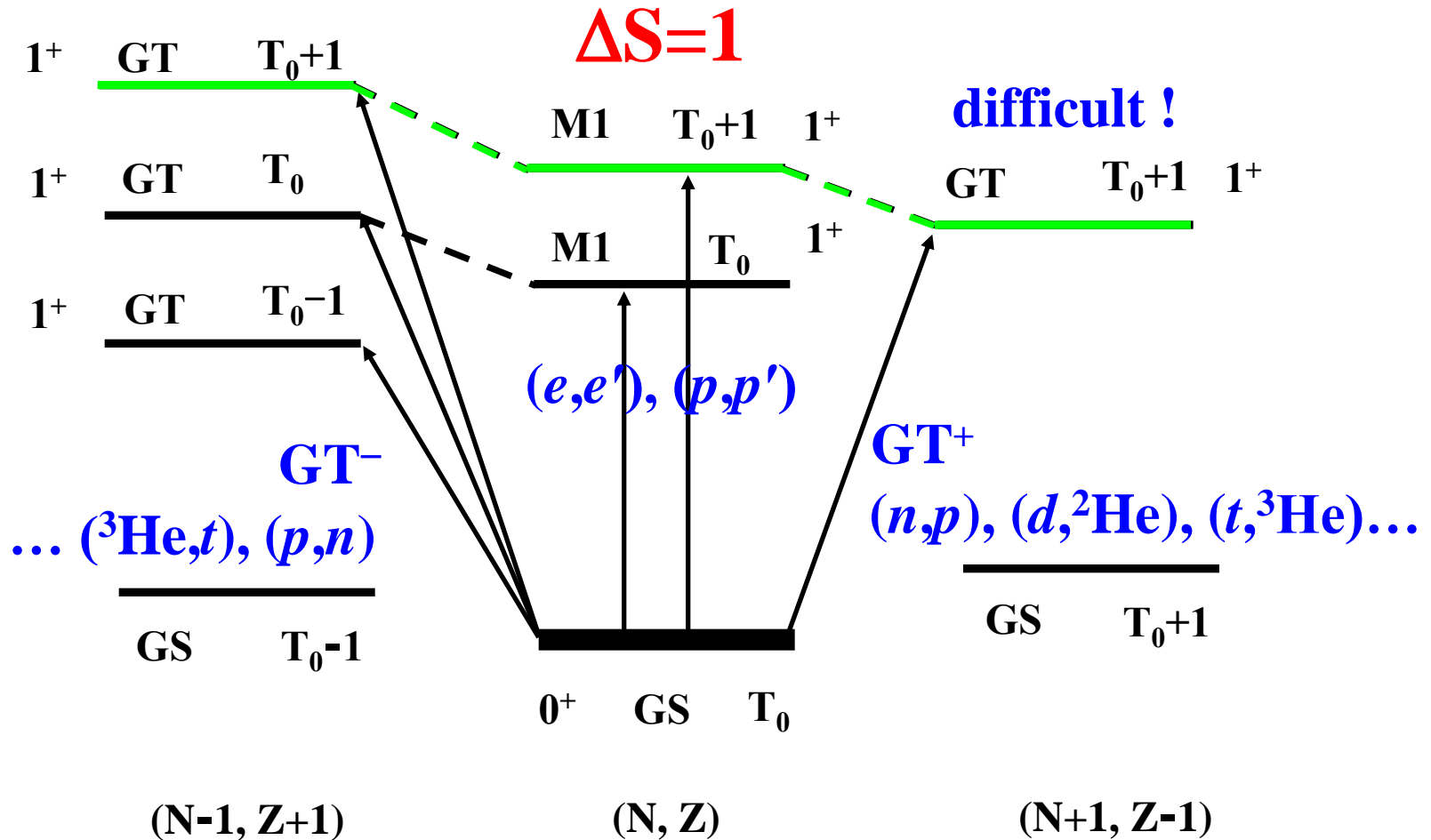


FIG. 4. Zero-degree cross-section spectra for the  $^{14}\text{C}(p,n)^{14}\text{N}$  reactions at the indicated bombarding energies. The spectra have been arbitrarily normalized. From Gaarde (1985) and Rapaport (1989).

# Spin-flip & GT transitions



# The ( $^3\text{He}, t$ ) reaction at 0 degree

- Cross sections at  $E(^3\text{He})=450$  MeV,  $q=0$  for ( $^3\text{He}, t$ ) reactions

$$\frac{d\sigma}{d\Omega} = \frac{\mu_i \mu_f}{(\pi \hbar^2)^2} \left( \frac{k_f}{k_i} \right) (N_\tau^D | J_\tau |^2 B(F) + N_{\sigma\tau}^D | J_{\sigma\tau} |^2 B(GT))$$

T. N. Taddeucci *et al.*, Nucl. Phys. A469, 125 (1987)

I. Bergqvist *et al.*, Nucl. Phys. A469, 648 (1987)

- Neutrino absorption cross sections

$$\sigma = \frac{1}{\pi \hbar^4 c^3} \left[ G_V^2 B(F) + G_A^2 B(GT) \right] \times F(Z, E_e) p_e E_e$$

$F(Z, E_e)$  is the relativistic Coulomb barrier factor

**Importance of charge-exchange reactions at intermediate energies**

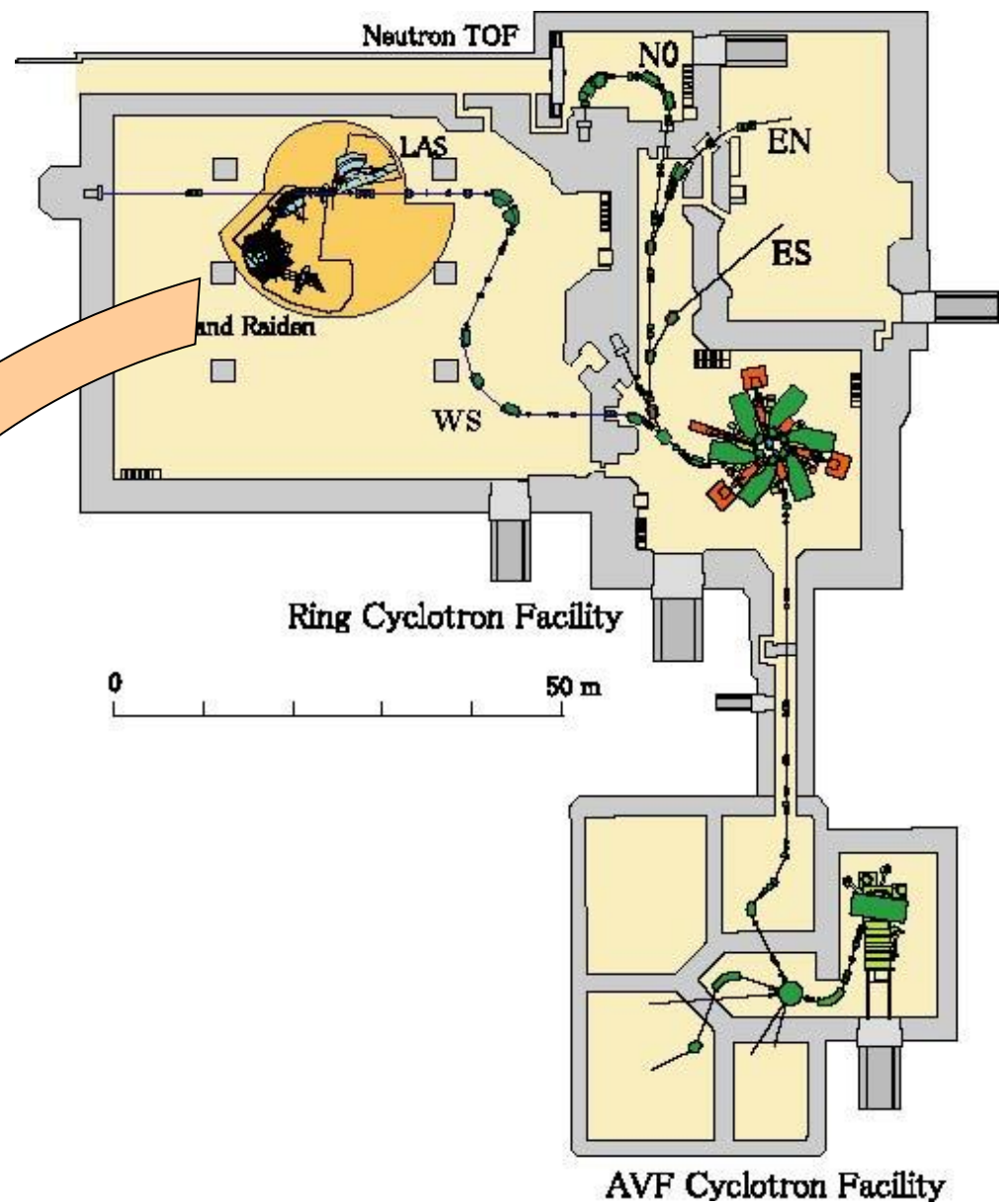
# Experiments

- RCNP facility

- K=400 MeV ring cyclotron

- Grand Raiden spectrometer

- Beam:  $^3\text{He}^{++}$ , 450 MeV

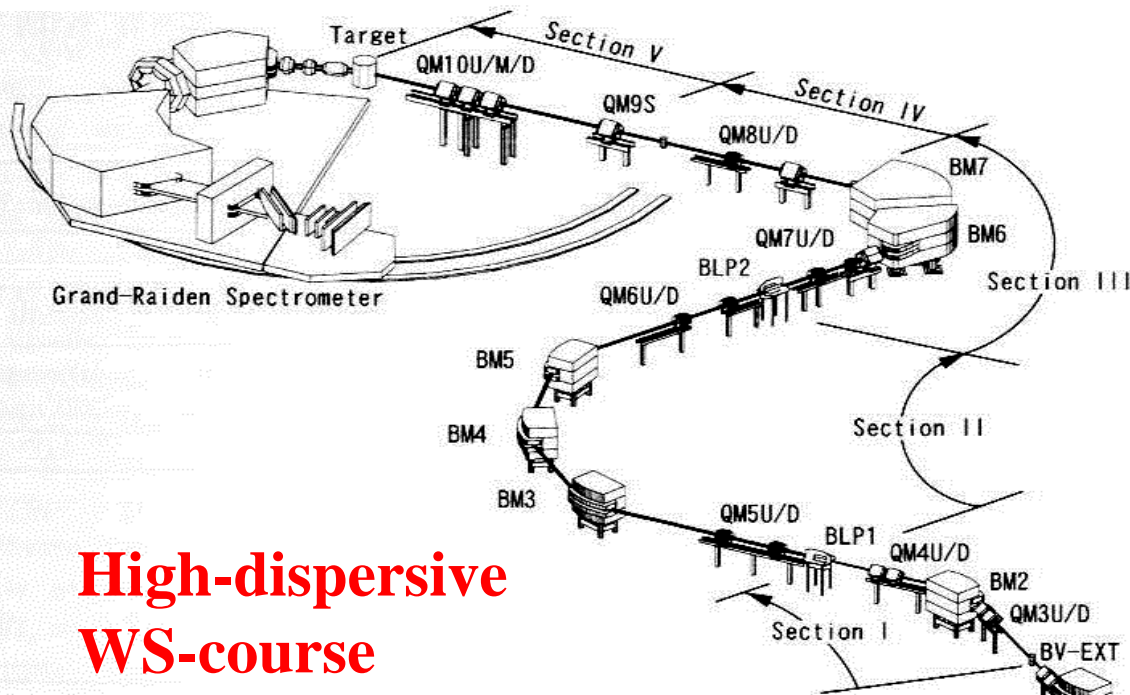


M. Fujiwara *et al.*, NIM A422 (1999) 484

# Beam line WS-course

## Grand-Raiden Spectrometer

M. Fujiwara *et al.*, NIM A422 (1999) 484



**High-dispersive  
WS-course**

T. Wakasa *et al.*, NIM A482 (2002) 79



# Evolution of Resolution in Charge-Exchange Reactions at Intermediate Energies

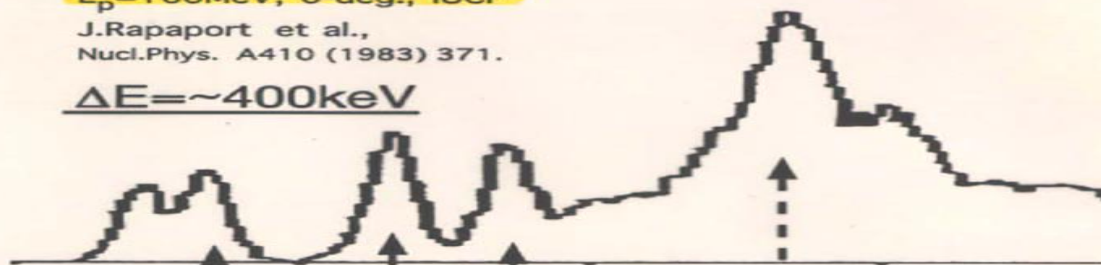
IUCF

$^{58}\text{Ni}(p,n)$

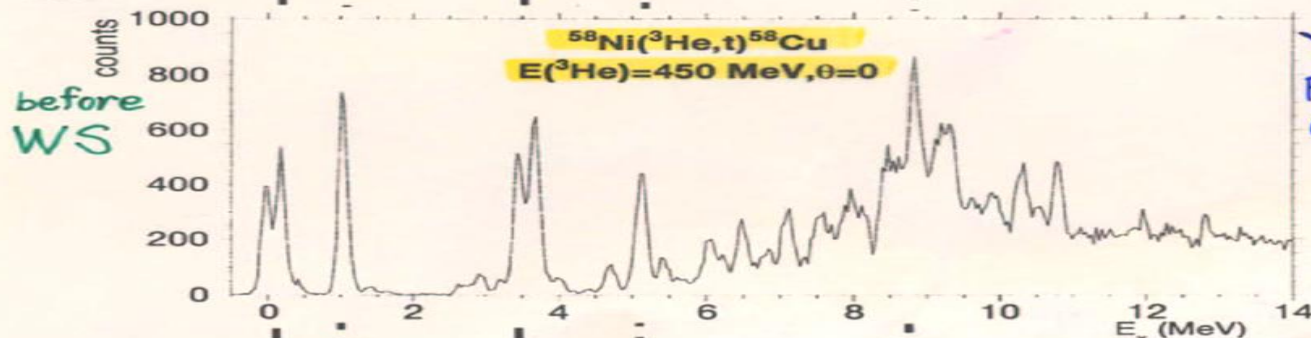
$E_p = 160\text{ MeV}$ , 0-deg., IUCF

J. Rapapaort et al.,  
Nucl. Phys. A410 (1983) 371.

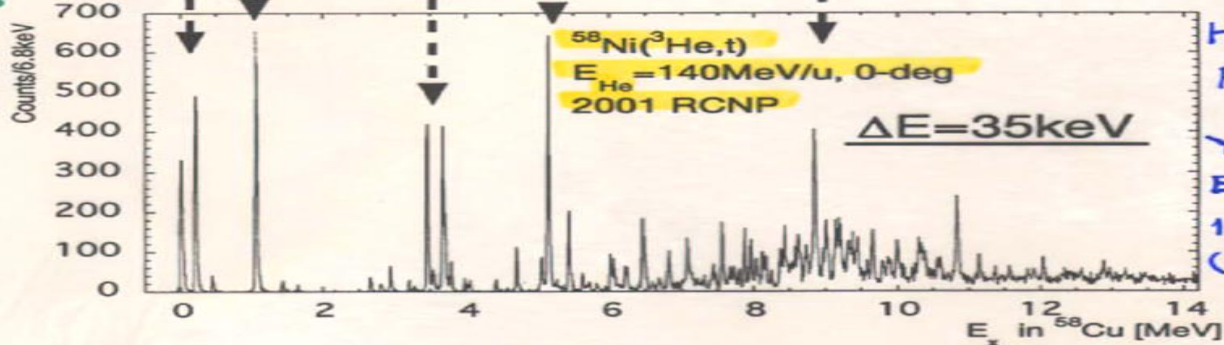
$\Delta E \sim 400\text{ keV}$



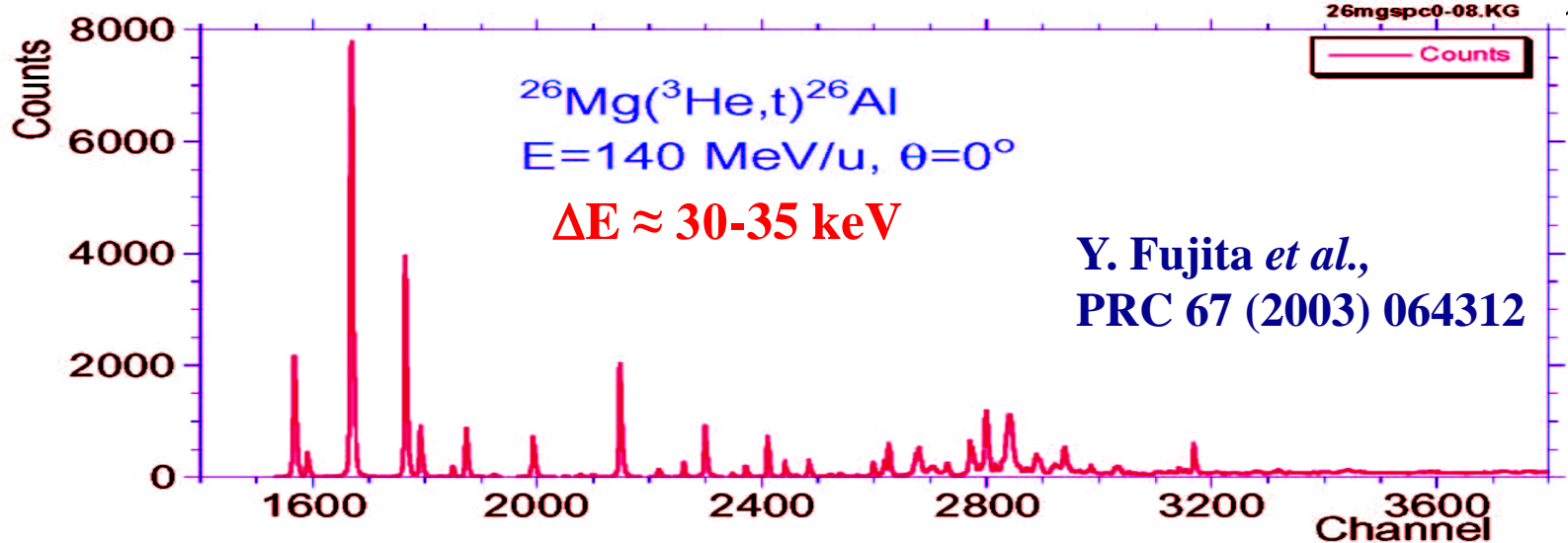
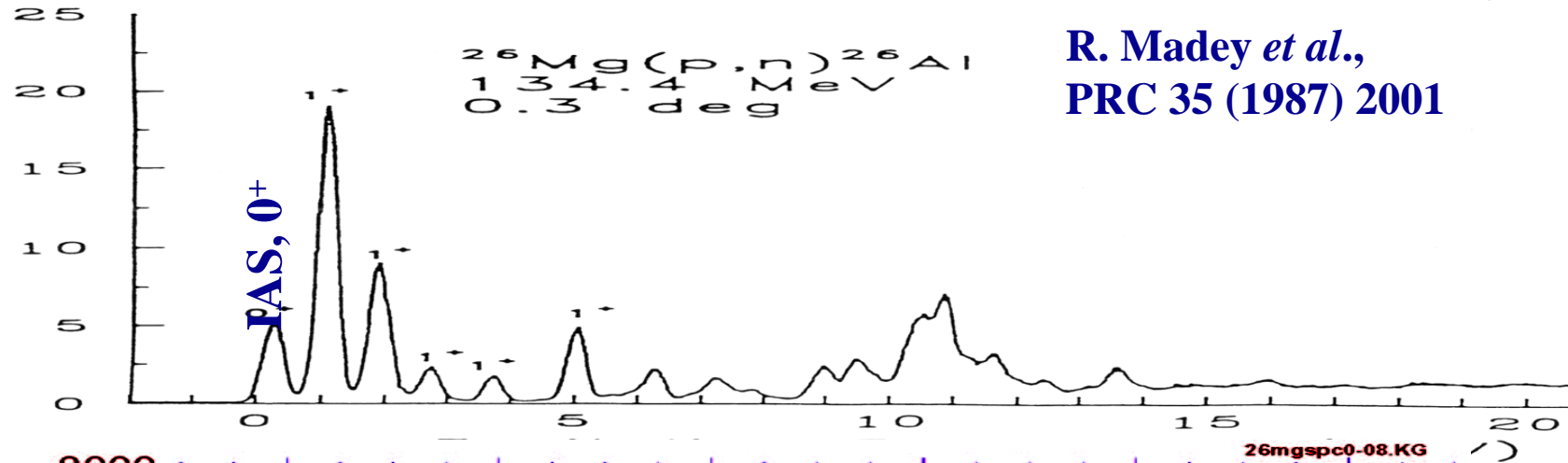
RCNP



WS

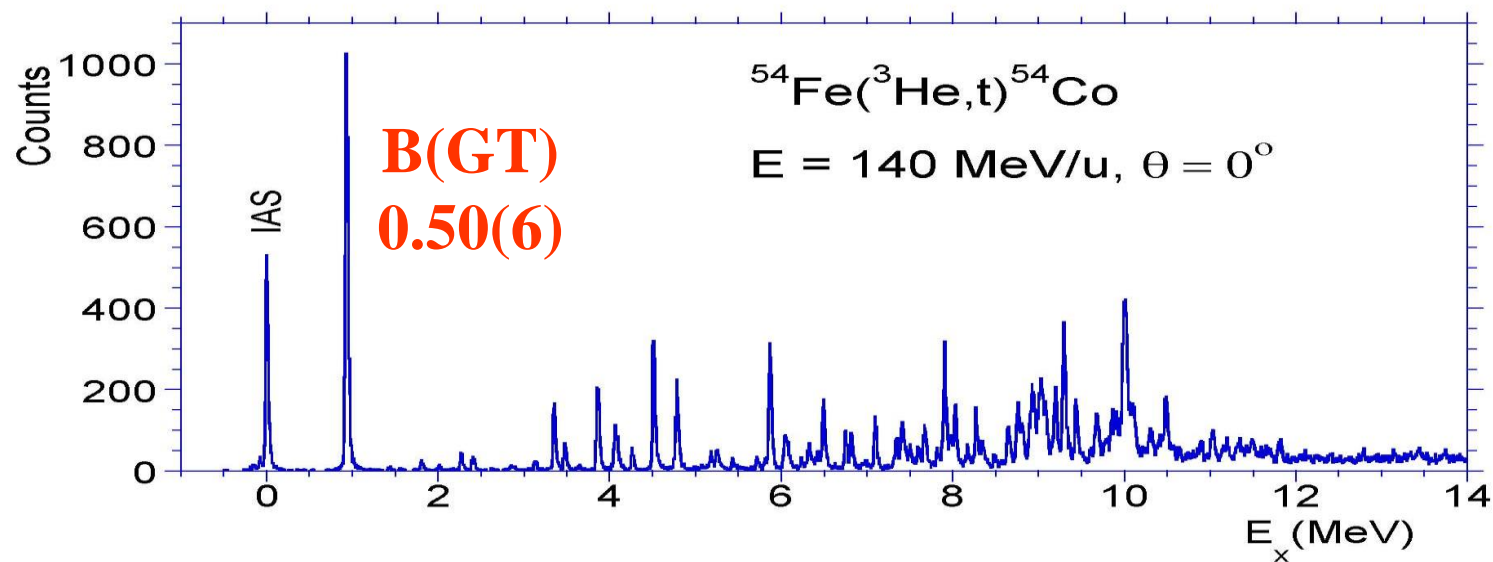
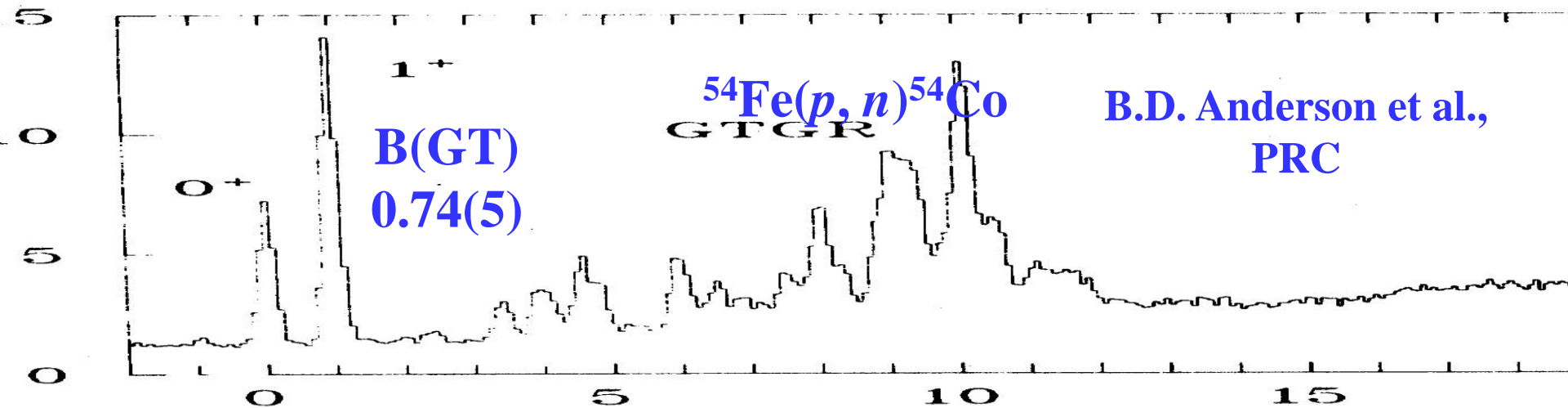


# $^{26}\text{Mg}(p,n)^{26}\text{Al}$ & $^{26}\text{Mg}(^3\text{He},t)^{26}\text{Al}$ spectra



# Prominent states are GT states and the IAS !

# $^{54}\text{Fe}(p,n)$ & $^{54}\text{Fe}(^3\text{He},t)$





$^{136}\text{Xe}(^3\text{He},t)^{136}\text{Cs}$

$E(^3\text{He}) = 420 \text{ MeV}$

$\Delta E = 42 \text{ keV}$

$B_{\text{exp}}(\text{GT}+) =$

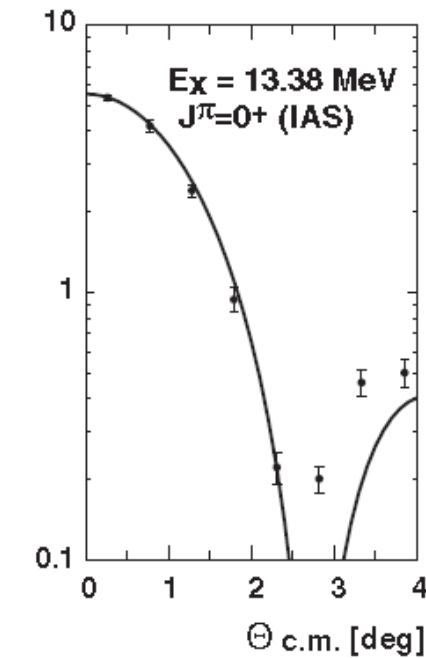
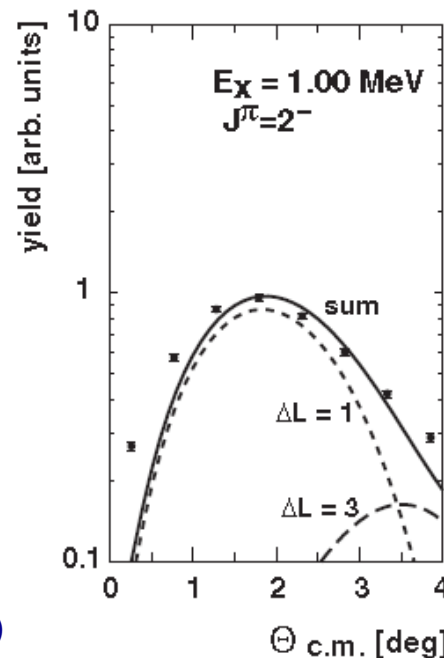
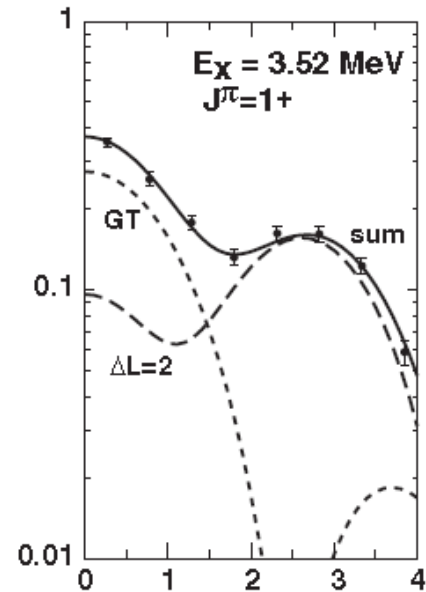
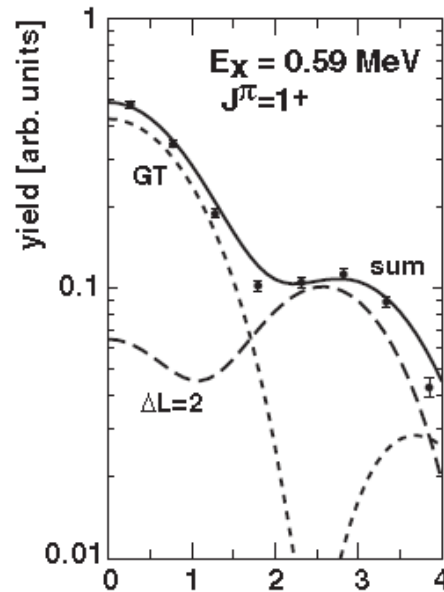
$$\frac{d\sigma(q=0)}{d\Omega} \cdot \left[ \frac{d\hat{\sigma}(\text{GT})}{d\Omega} \right]^{-1}$$

extrapolated  
(DWBA)

unit cross section

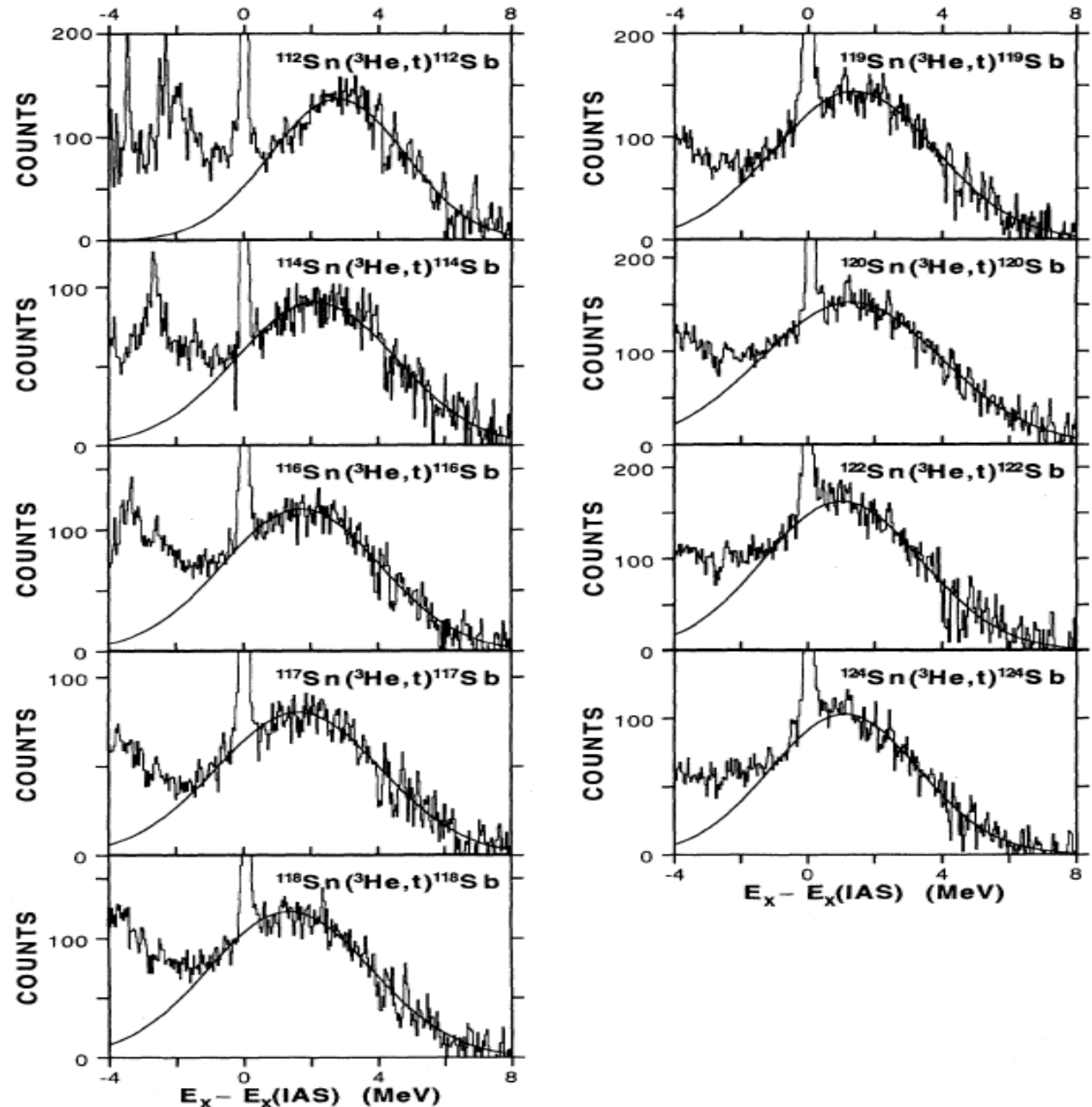
$\Delta L = 2$  &  $\Delta L = 0$  incoherent

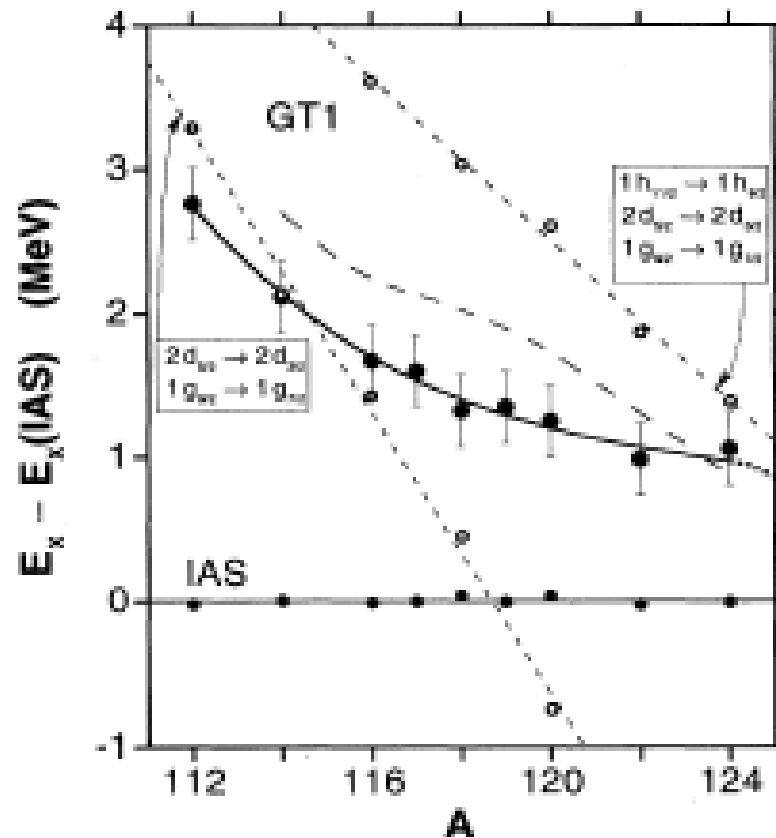
P. Puppe *et al.*, PRC 84 (2011) 051305(R)



$(^3\text{He}, t)$  charge-exchange reaction on all stable Sn nuclei at IUCF, Bloomington  
 $E(^3\text{He}) = 200 \text{ MeV}$   
 Excitation energy spectra are plotted relative to IAS.

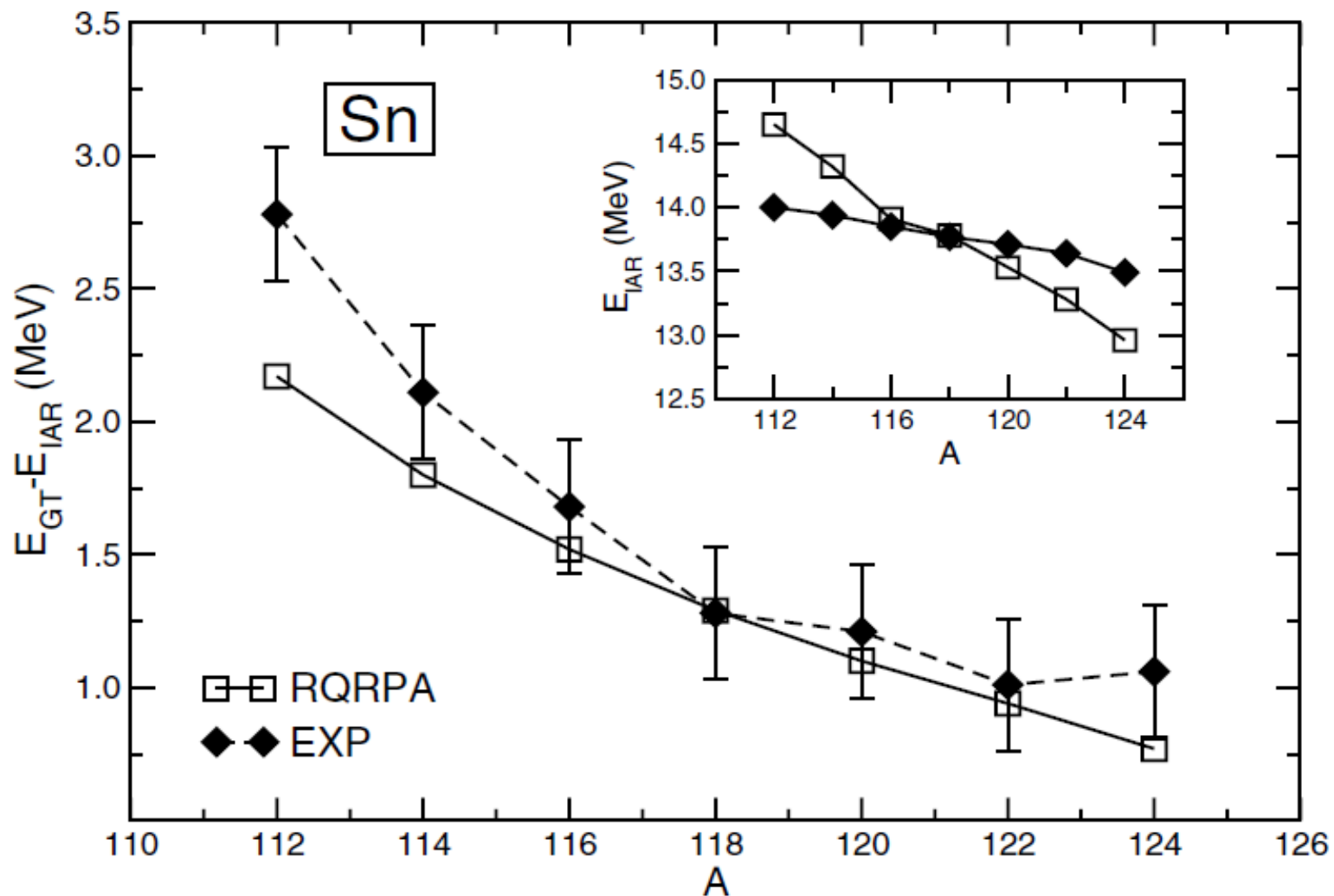
K. Pham *et al.*,  
 PRC **51** (1995) 526





Excitation energy of main component  
of GTGR relative to IAS.

K. Pham *et al.*, PRC **51** (1995) 526

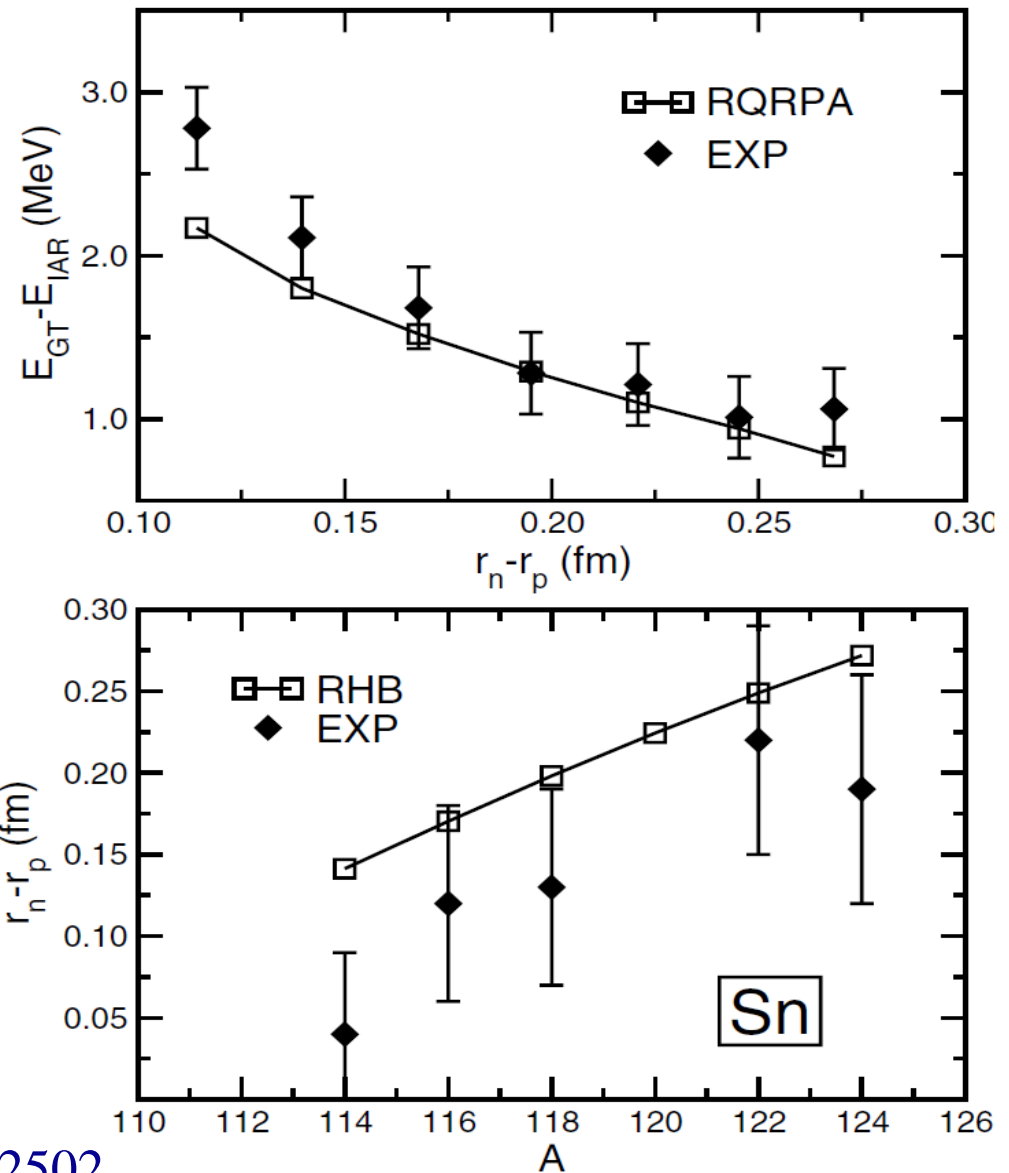


Comparison of theoretical calculations to experimental results for excitation energy of main component of GTGR relative to IAS.

Inset shows IAS energies

D. Vretenar *et al.*, PRL **91** (2003) 262502

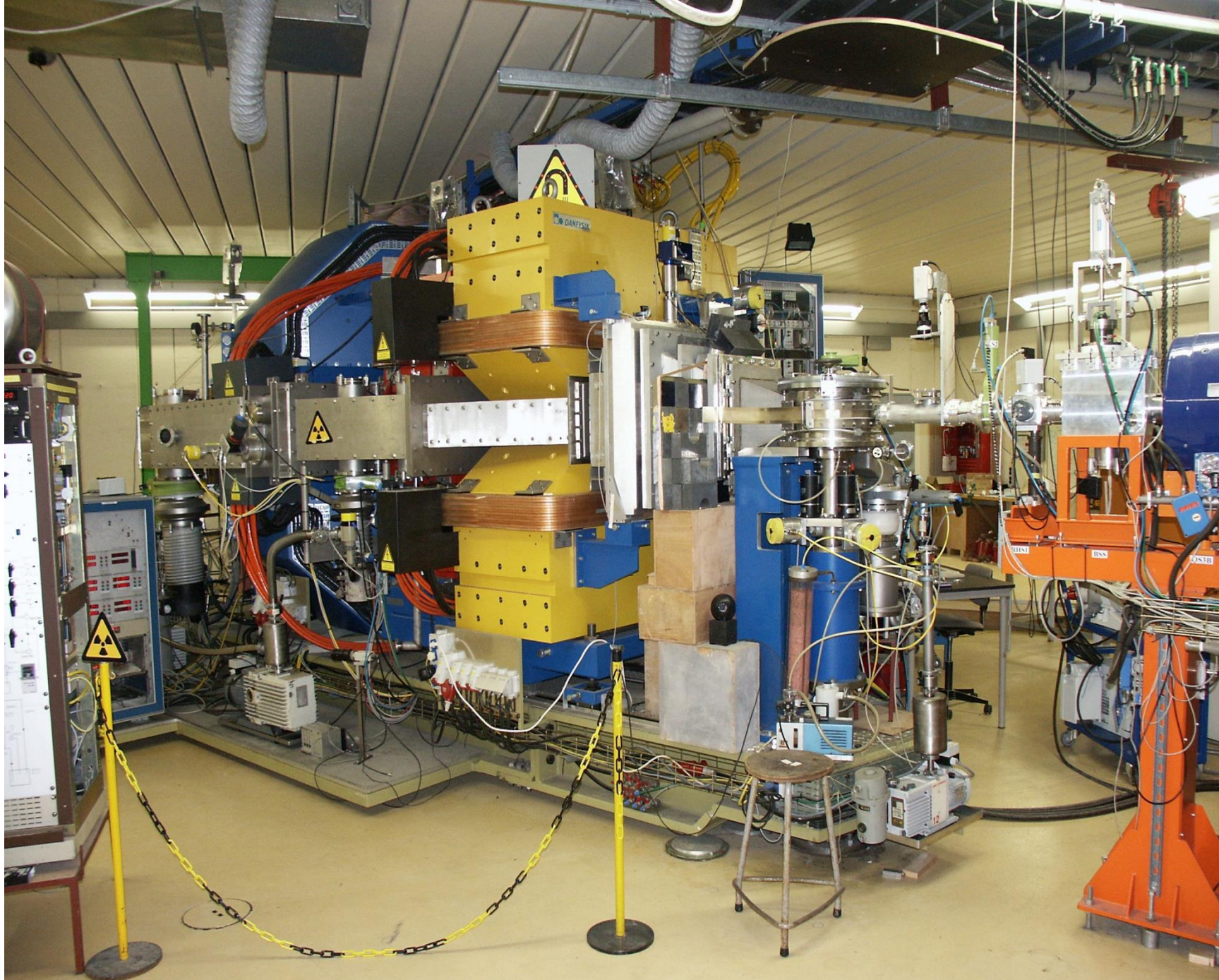
Theoretical pn-RQRPA and experimental differences of GTGR and IAS excitation energies as function of neutron-skin thickness (data from K. Pham). Lower panel shows comparison between theoretical neutron-skin thickness and experimental data (data from A. Krasznahorkay).



D. Vretenar *et al.*, PRL **91** (2003) 262502

A. Krasznahorkay *et al.*, PRL **83** (1999) 3216;  $r_n - r_p$





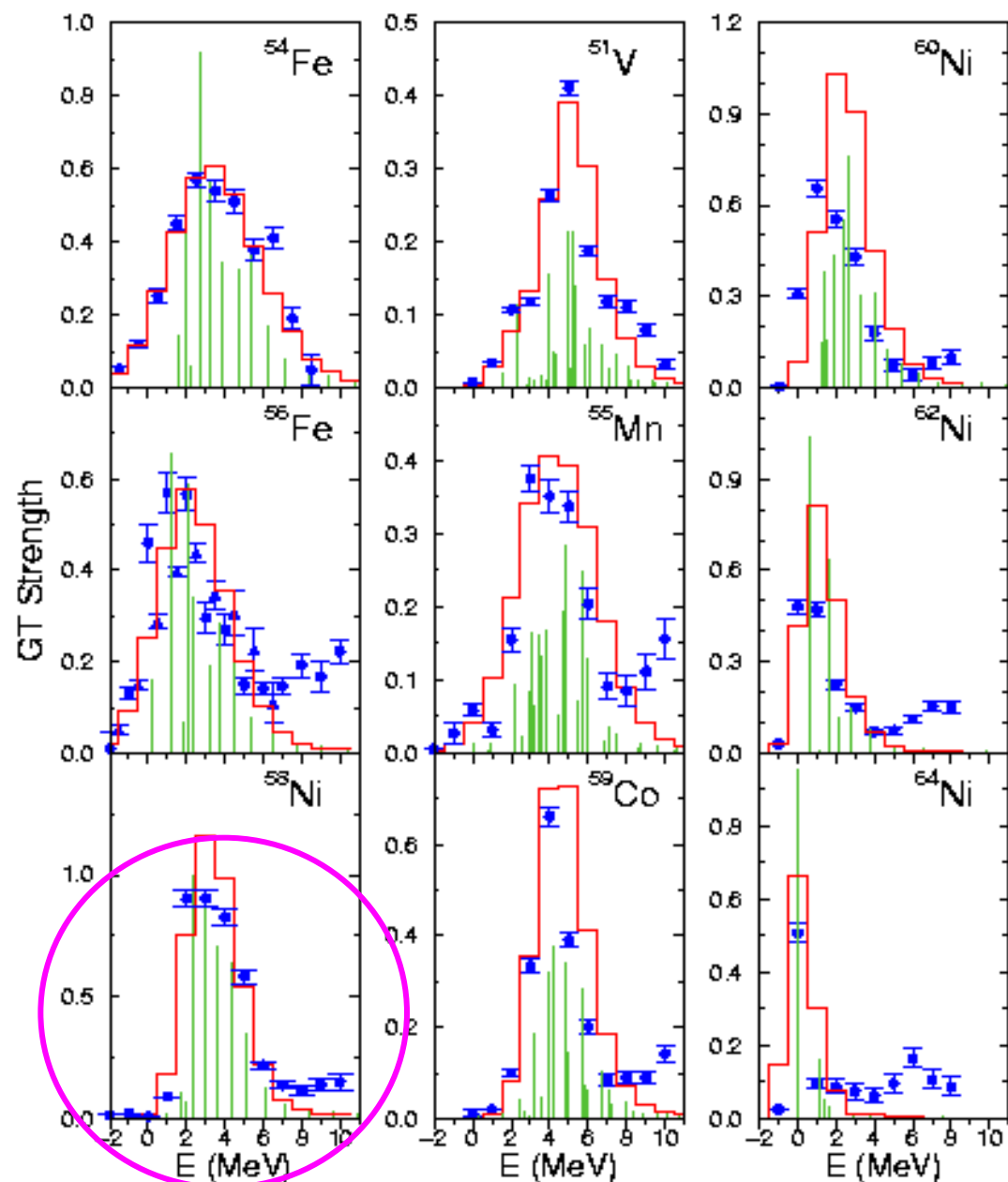
# *fp*-shell nuclei: large scale shell model calculations

E. Caurier *et al.*

NPA 653 (1999) 439

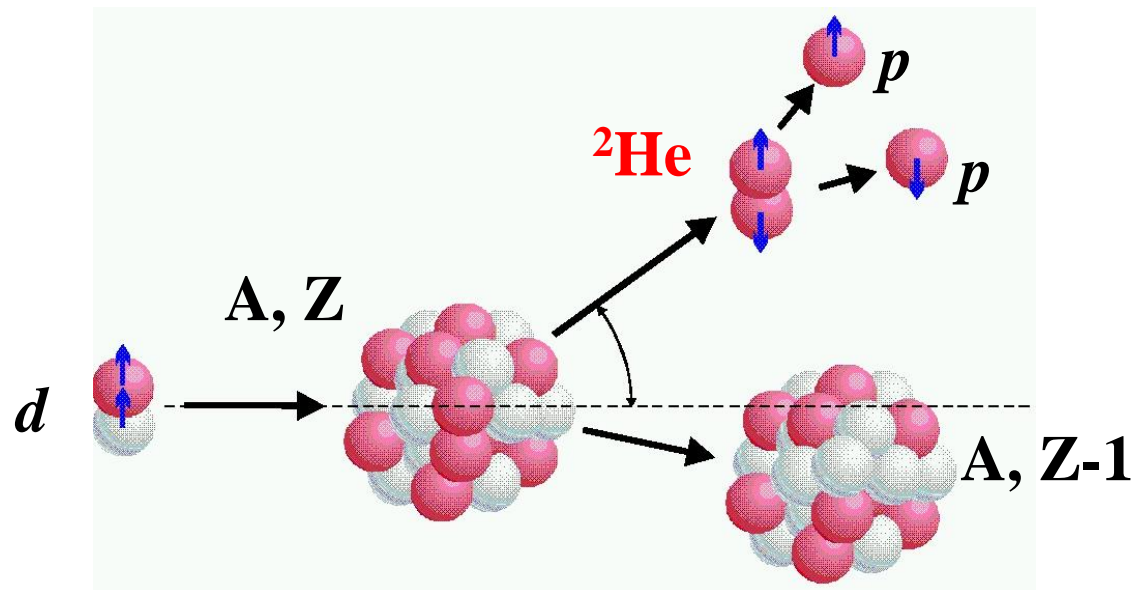
- Stellar weak reaction rates with improved reliability
- Large scale shell model (SM) calculations
- **Tuned to reproduce  $GT^+$  strength measured in  $(n,p)$**
- $(n,p)$  data from TRIUMF
- $GT^+$  strength from SM
- **Folded with energy resolution**

Case study:  $^{58}\text{Ni}$





# Exclusive excitations $\Delta S=\Delta T=1$ : ( $d, {}^2\text{He}$ )



${}^3\text{S}_1$  deuteron  $\Rightarrow$   ${}^1\text{S}_0$  di-proton ( ${}^2\text{He}$ )

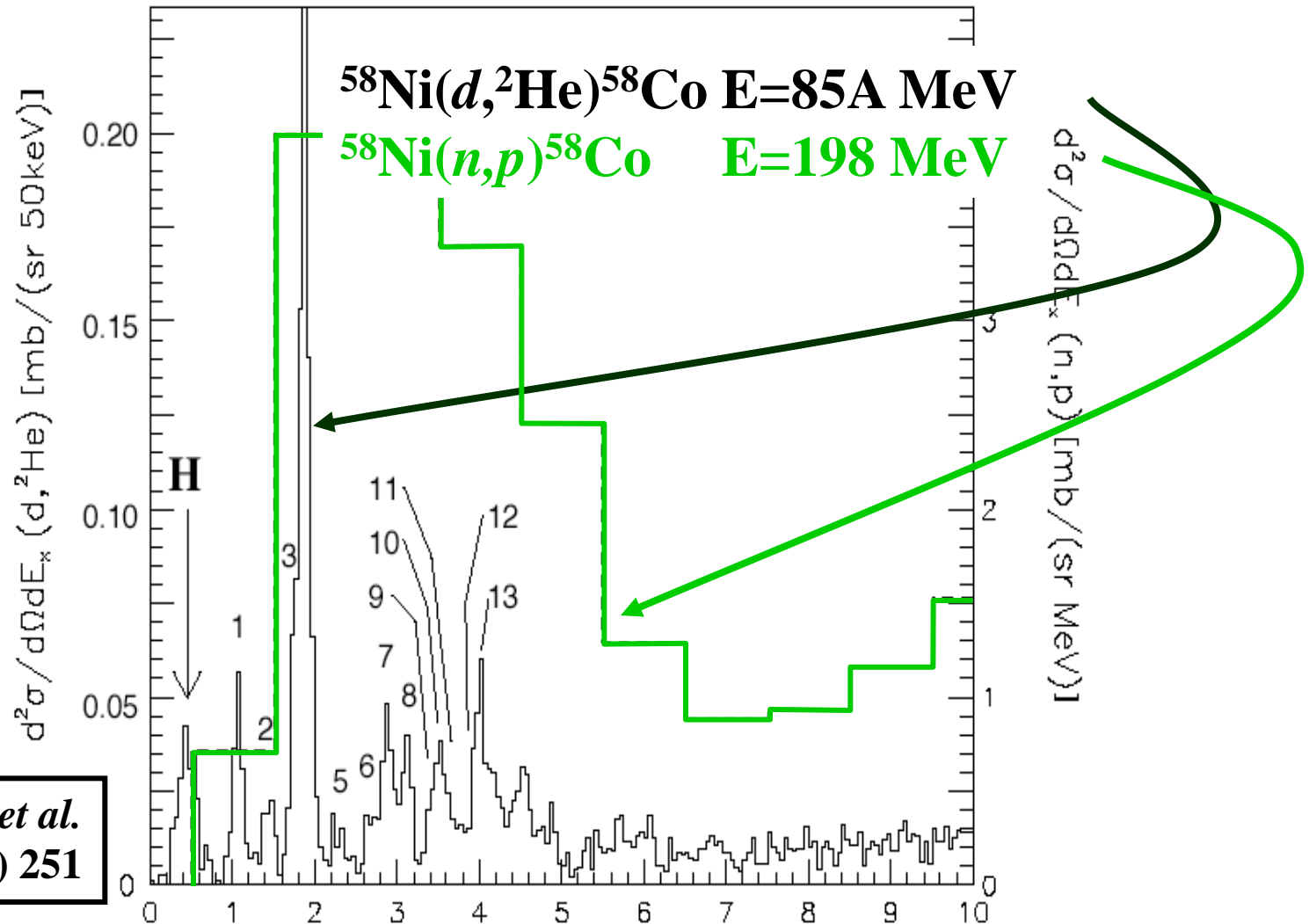
${}^1\text{S}_0$  dominates if (relative) 2-proton kinetic energy  $\varepsilon < 1$  MeV

$(n,p)$ -type probe with exclusive  $\Delta S=1$  character ( $\text{GT}^+$  transitions)

But near  $0^\circ$ : tremendous background from  $d$ -breakup

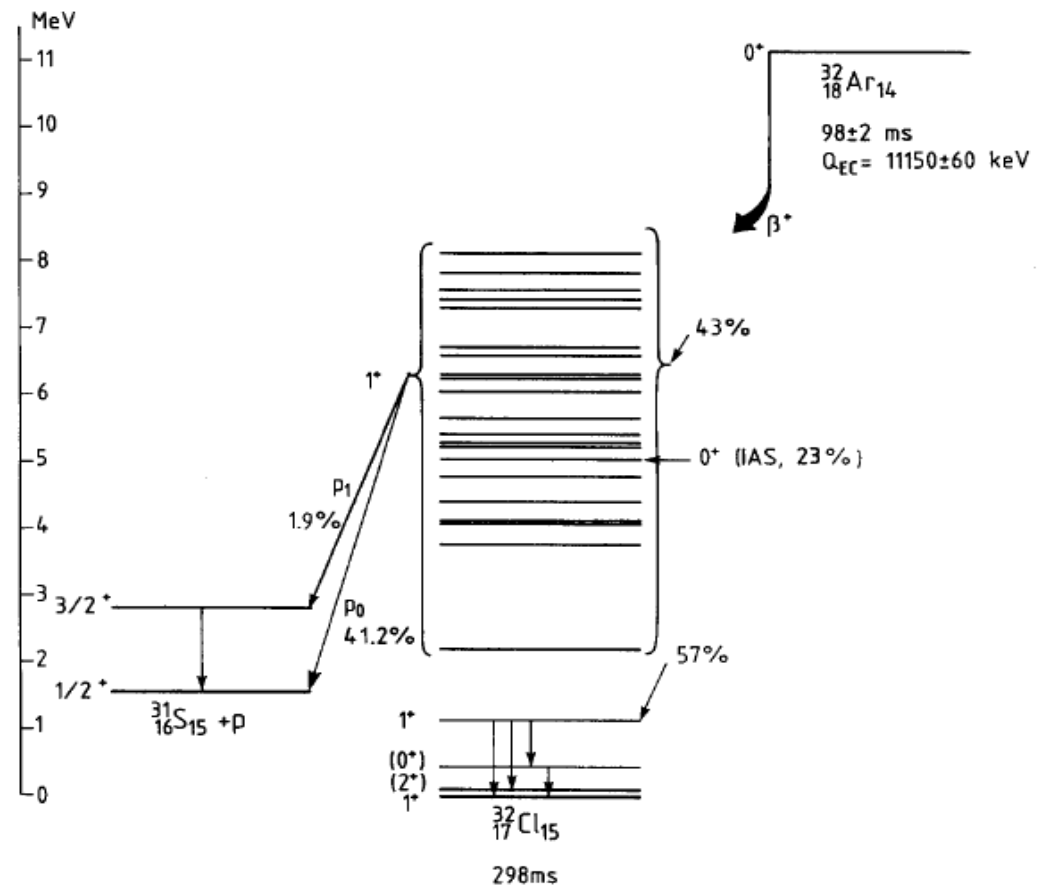


# $(d, {}^2\text{He})$ as $\text{GT}^+$ probe in $fp$ -shell nuclei



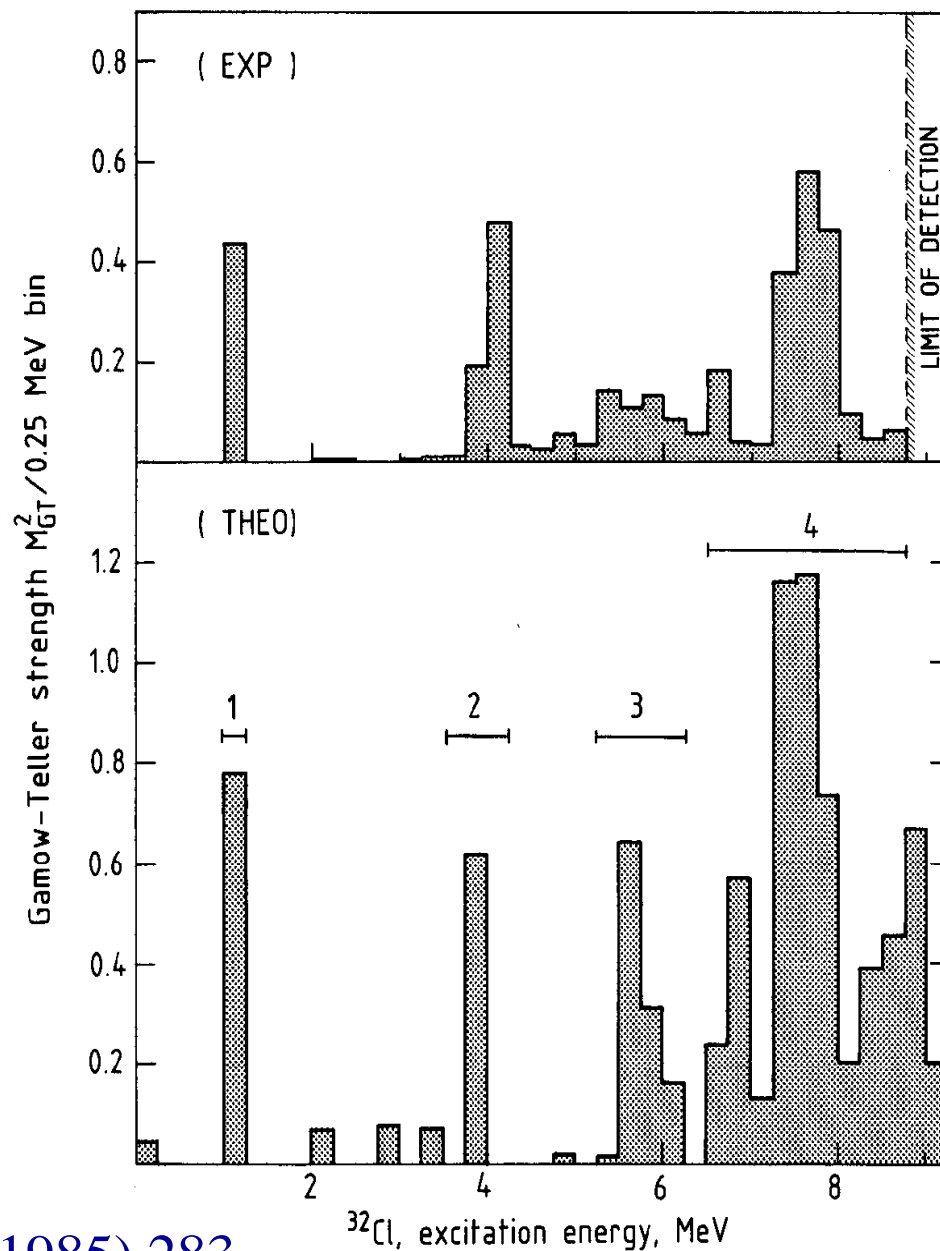
M. Hagemann *et al.*  
PLB 579 (2004) 251

Decay scheme of  $^{32}\text{Ar}$ .  
The limit of detection  
corresponded to 8.75 MeV  
excitation in  $^{32}\text{Cl}$ .



T. Björnstad *et al.*, NPA **443** (1985) 283

The  $^{32}\text{Ar}$  Gamow-Teller strength function beta decay and beta-delayed proton decay. The GTR is observed at about 7.5 MeV, in agreement with shell-model predictions by Müller *et al.* [PRC **3** (1971) 700] shown in lower part of the figure. Quenching in the four regions indicated in figure are 0.6, 1.0, 0.5, and 0.4, respectively. The overall renormalization of the axial-vector strength is  $(g'_A/g_A)^2=0.49$ .



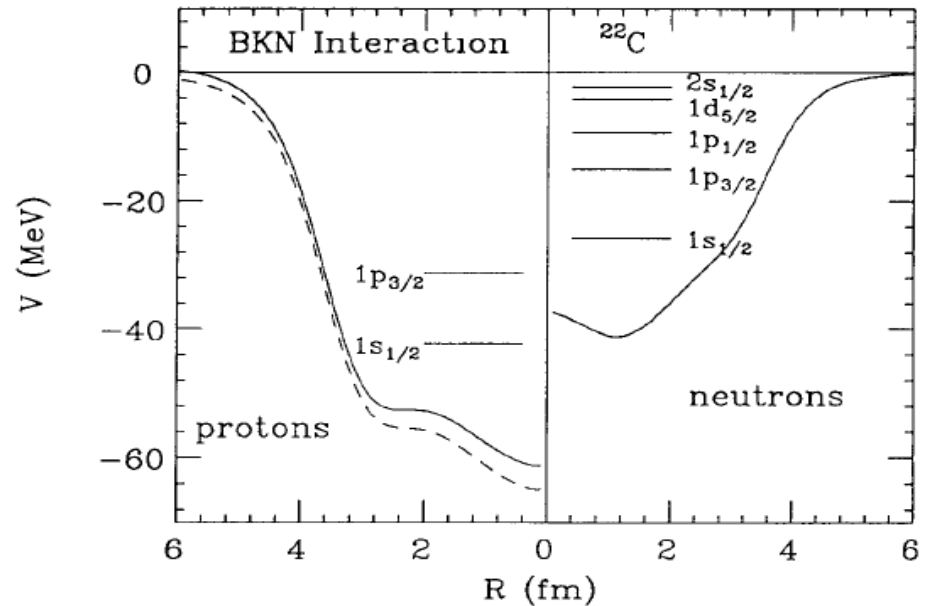
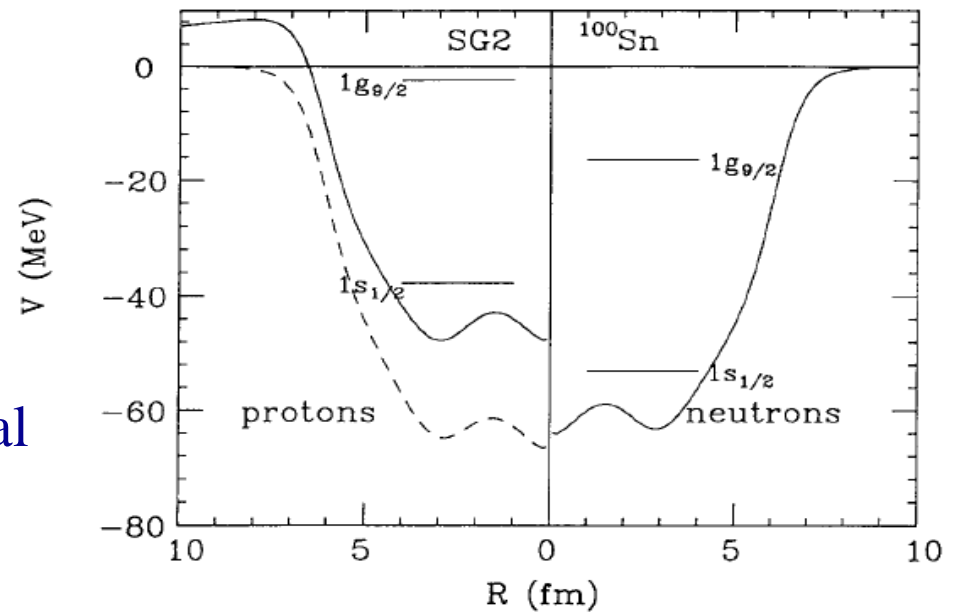
T. Björnstad *et al.*, NPA **443** (1985) 283

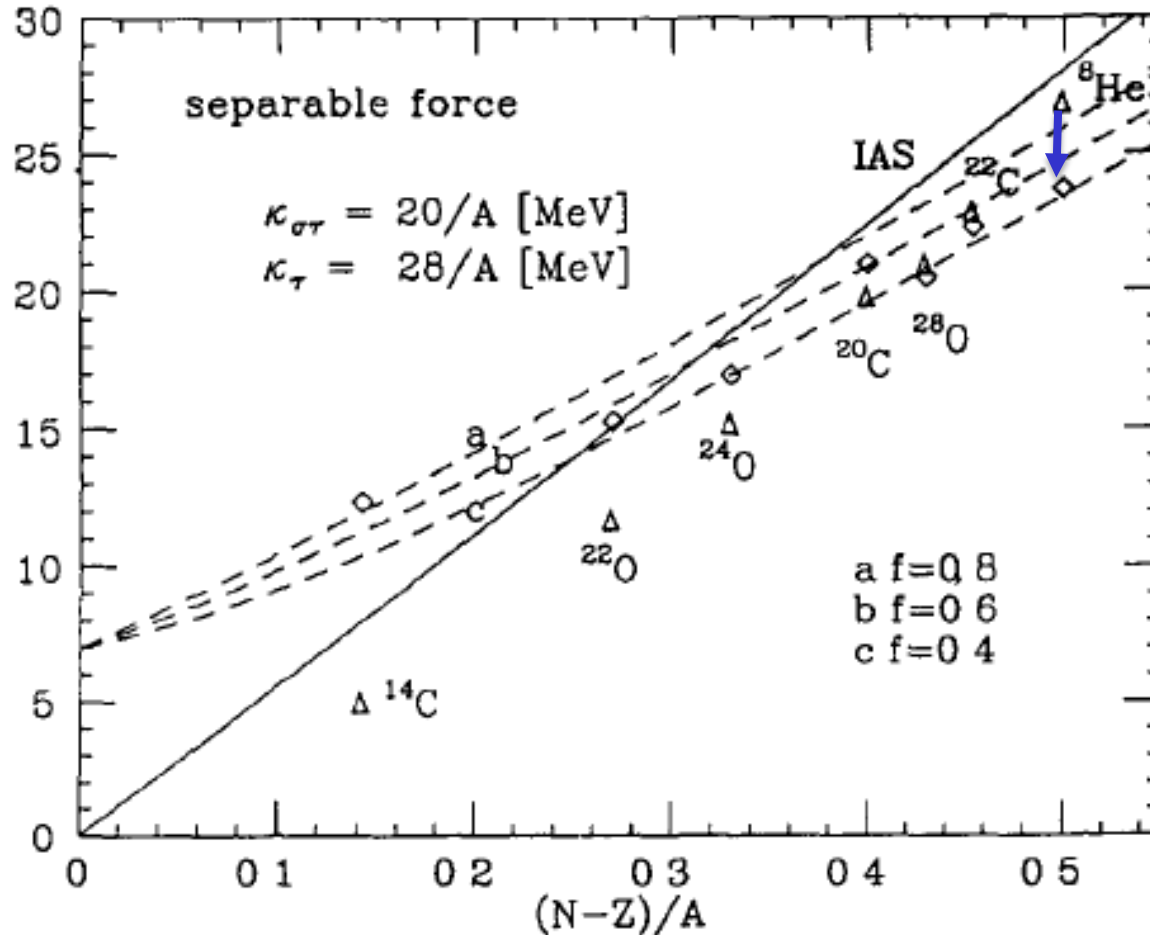
## Hartree-Fock potentials

Dashed: Without Coulomb potential

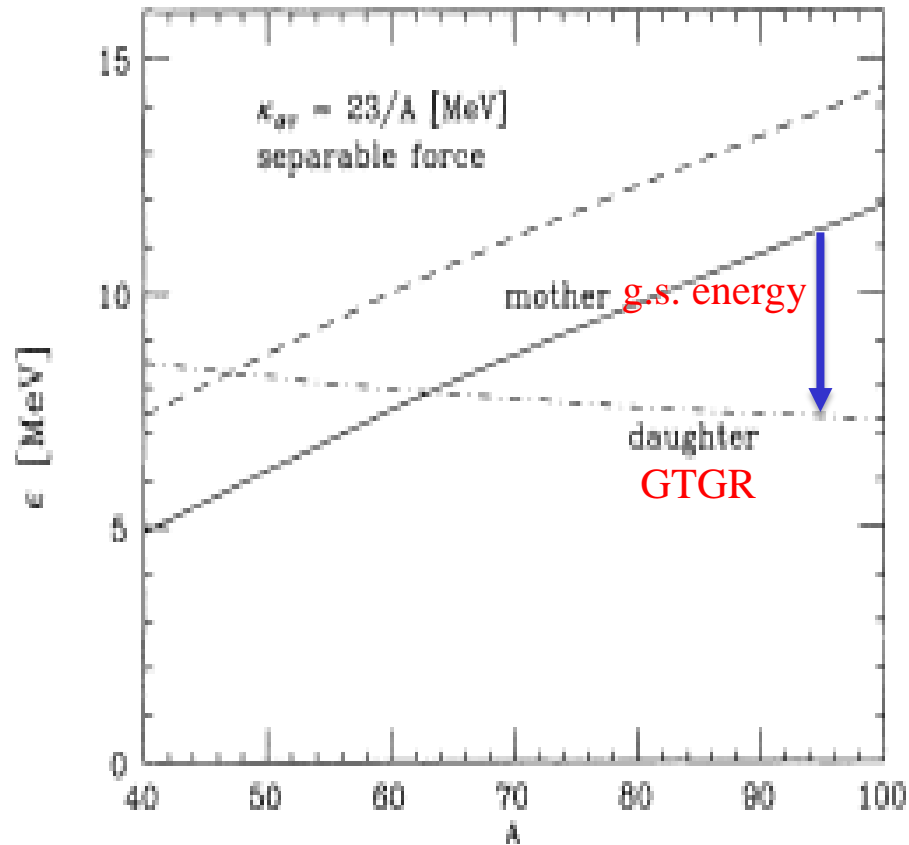
Solid: With Coulomb potential

I. Hamamoto & H. Sagawa,  
 PRC **48** (1993) R960;  
 I. Hamamoto,  
 NPA A577 (1994) 19c



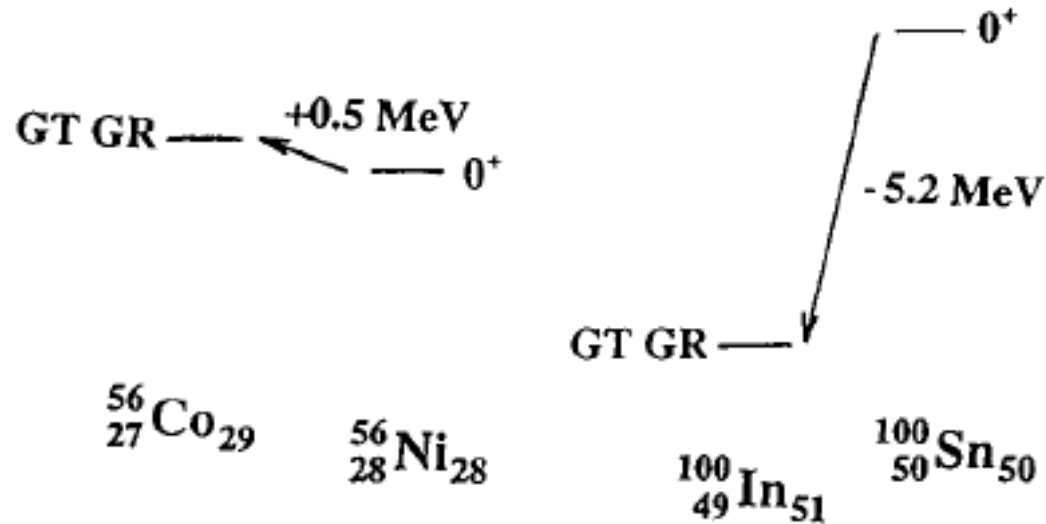


Energies of GTGR (dashed lines) and IAS (solid line) obtained in a schematic model with  $f$  the quenching factor of GT strength. **Diamonds are GTGR based on mother nuclei and triangles are g.s. of mother nuclei**



Energies of GTGR (dash-dotted line) of  $N=Z$  nuclei calculated in a schematic model.

Dashed (solid) line is energy of mother nucleus without (with) correction term  $\alpha = 2.5$  MeV.

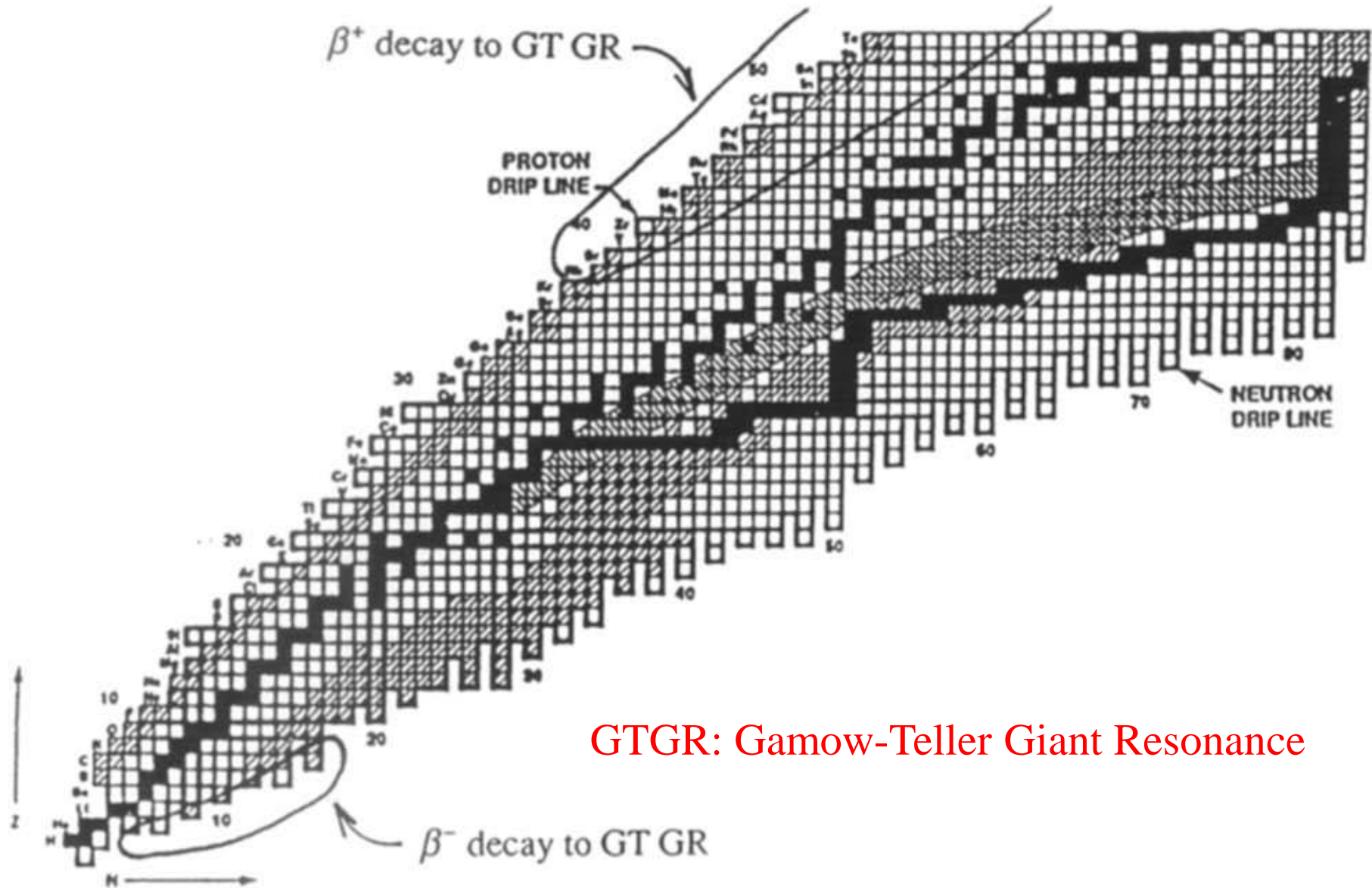


## HF-TDA

Energy of GTGR relative to g.s. of  $N=Z$  nuclei

$^{56}\text{Ni}$  &  $^{100}\text{Sn}$

I. Hamamoto, NPA A577 (1994) 19c



GTGR: Gamow-Teller Giant Resonance

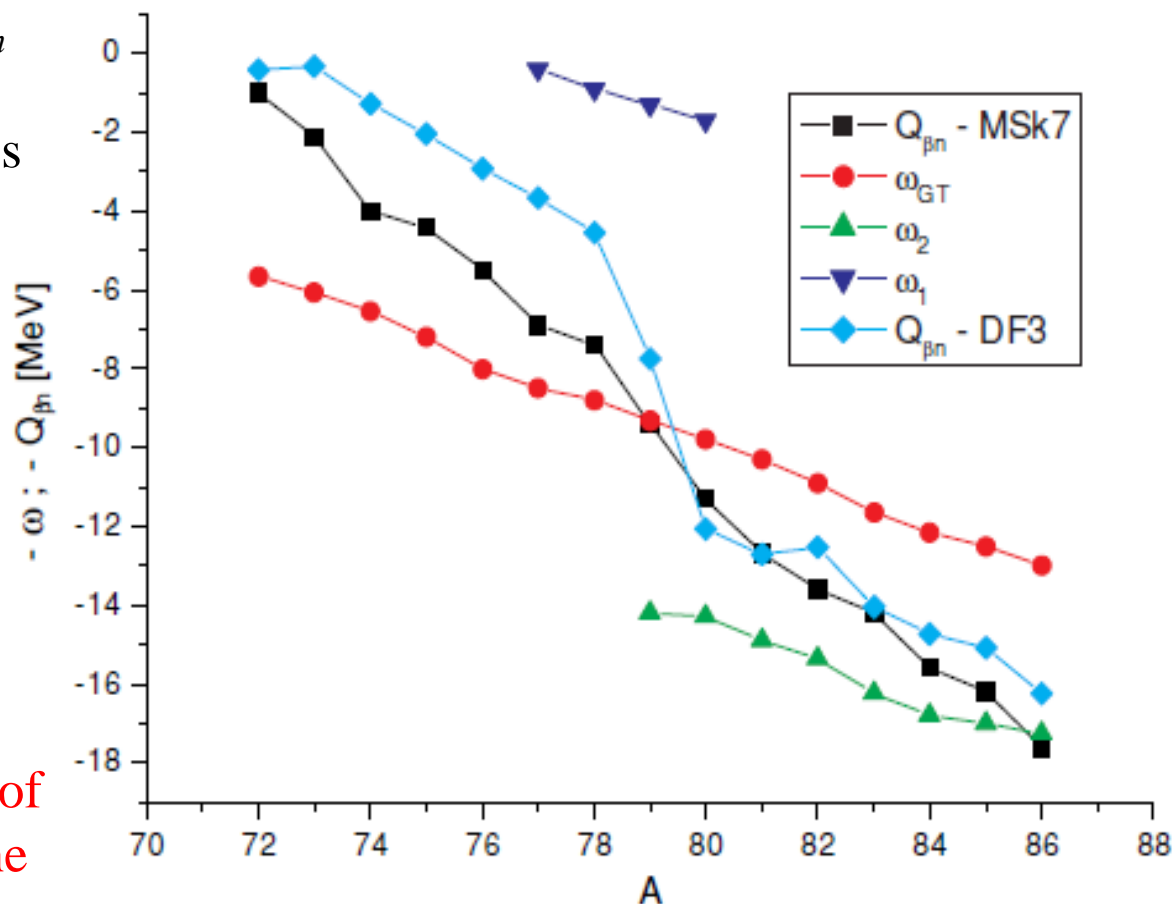
I. Hamamoto, NPA A577 (1994) 19c



The calculated position of  $Q_{\beta n}$  window for delayed neutron emission compared to energies of the GT pygmy-resonance ( $\omega_{GT}$ ) and main  $0^-$  transitions ( $\omega_{1,2}$ ) in **Ni isotopes**.

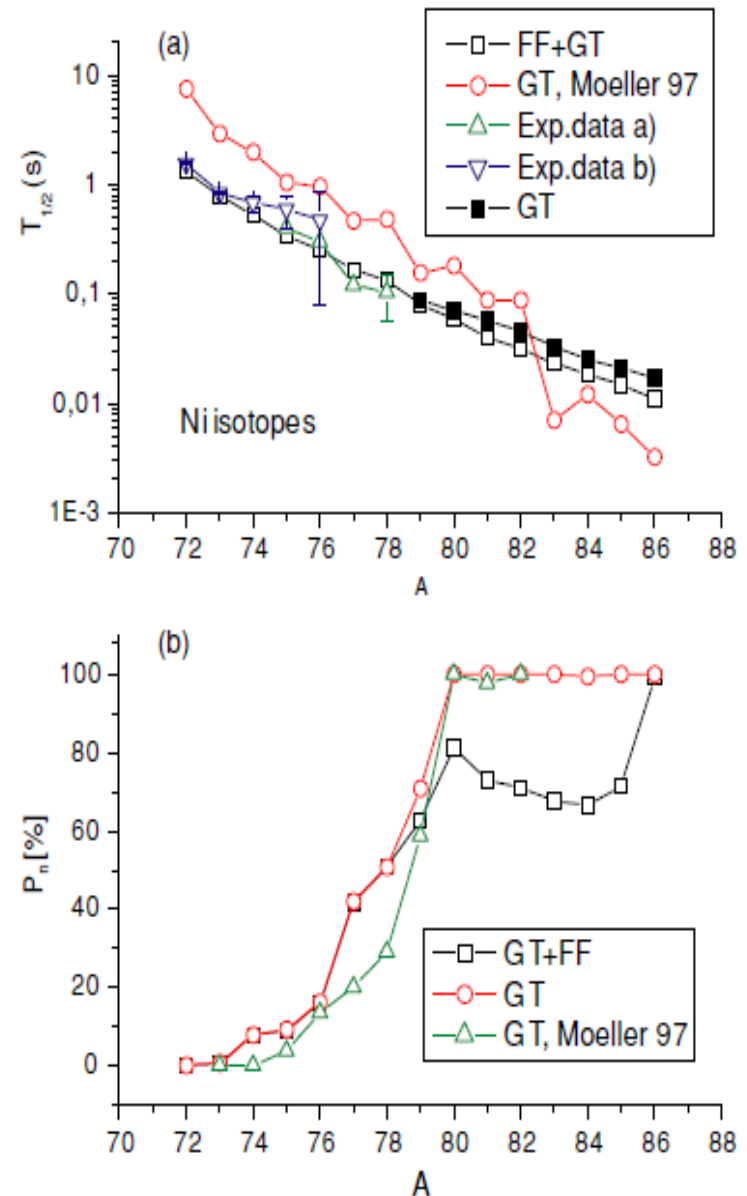
Transition energies and neutron-emission windows are plotted with the opposite sign.

Zero energy corresponds to mothers ground states.  
For the lowest energy branch of the  $0^-$  transitions ( $\omega_1$ ), only the points for  $A = 77-81$  are shown.



I.N. Borzov, PRC **71** (2005) 065801

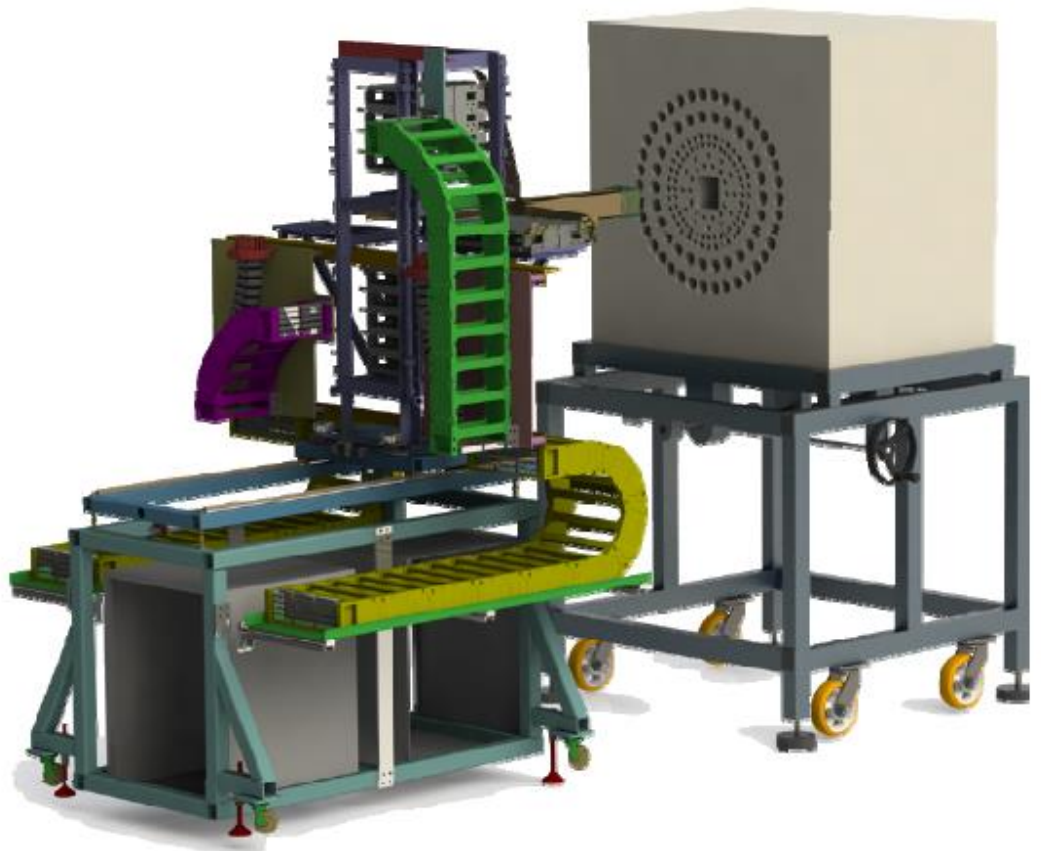
(a) Total  $\beta$ -decay half-lives and  
 (b) delayed neutron emission  
 probabilities for Ni isotopes predicted  
 from density functional of Fayans (DF3)  
 + continuum QRPA including the allowed  
 (GT) plus first-forbidden  
 (FF) transitions, and the allowed  
 transitions, in comparison with the  
 Finite-range droplet model (FRDM) +  
 RPA calculation for allowed transitions  
 (Möller) and experimental data a) and b).



I.N. Borzov, PRC **71** (2005) 065801

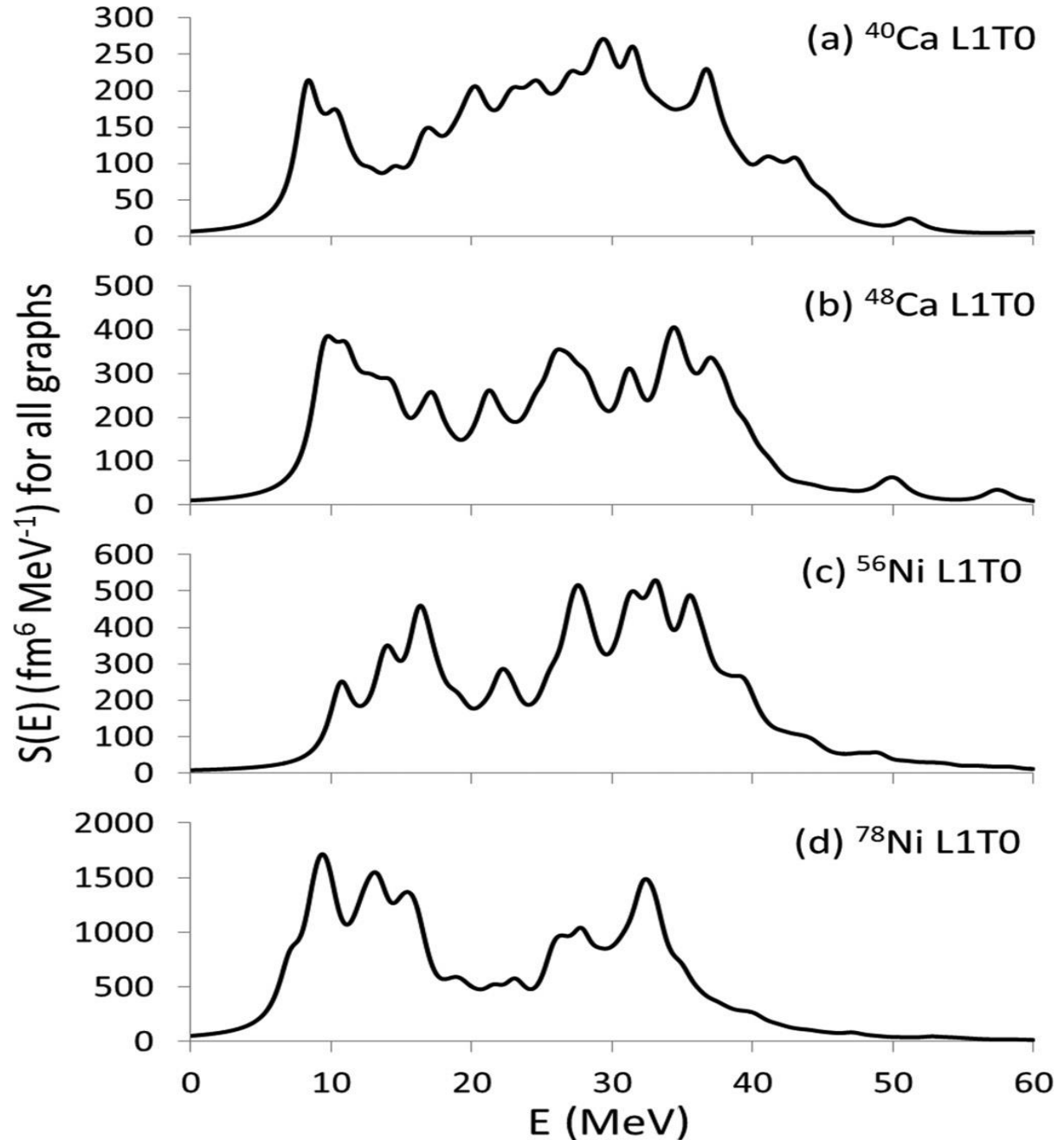
Neutron detector BELEN consisting of 48 counters will merge with other neutron detection setups from Japan, Russia, and USA and the AIDA particle detector from the UK into the “BRIKEN” detector to perform experiments at RIKEN in 2015/2016.

**BRIKEN** collaboration is intending to measure  $\beta$ -delayed neutron emission in the “ $^{78}\text{Ni}$  and beyond” part of nuclear chart.



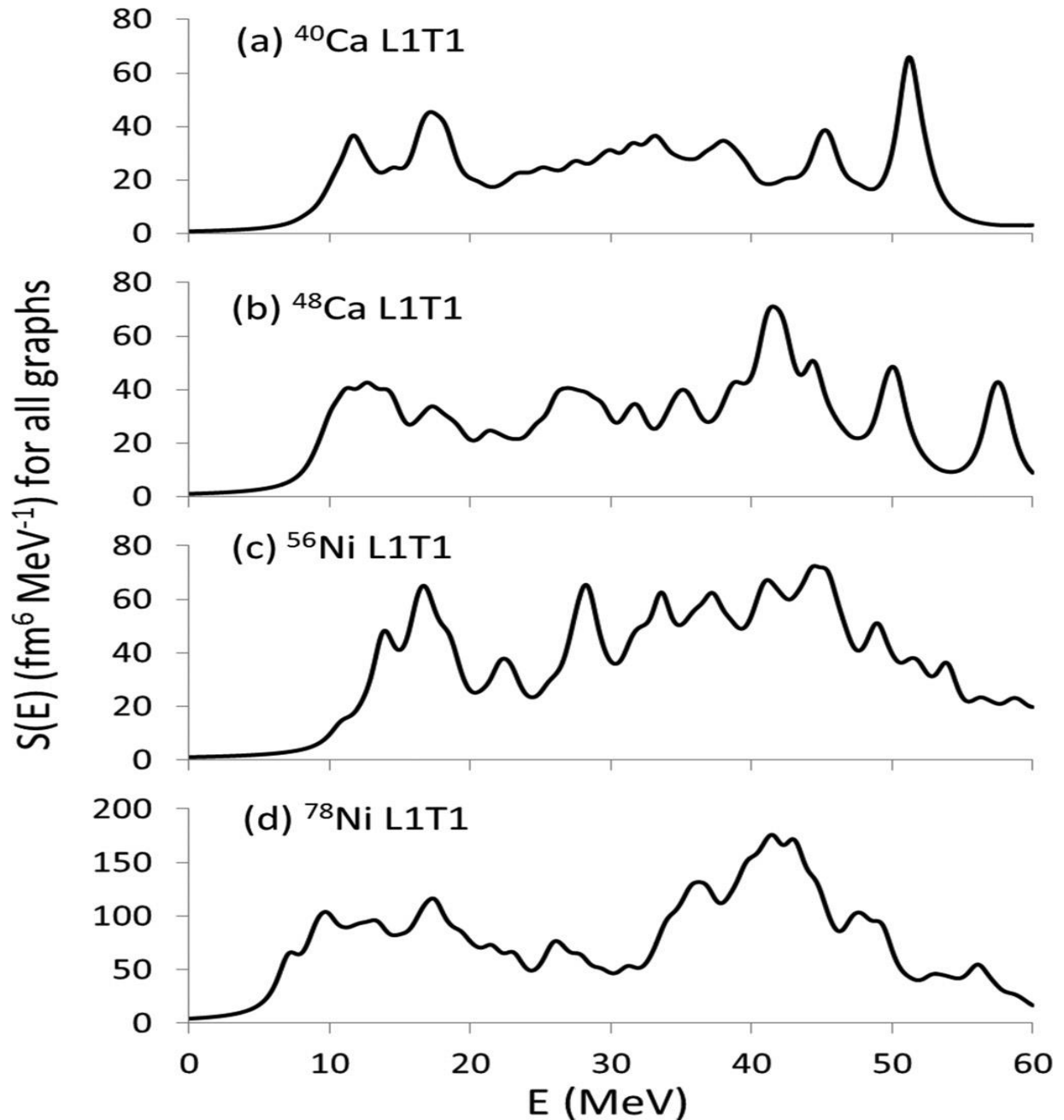
Self-consistent  
HF-based RPA  
results for the  
strength function  
 $S(E)$  distribution  
of the isoscalar  
dipole response

N. Auerbach *et al.*,  
PRC **89** (2014) 014335

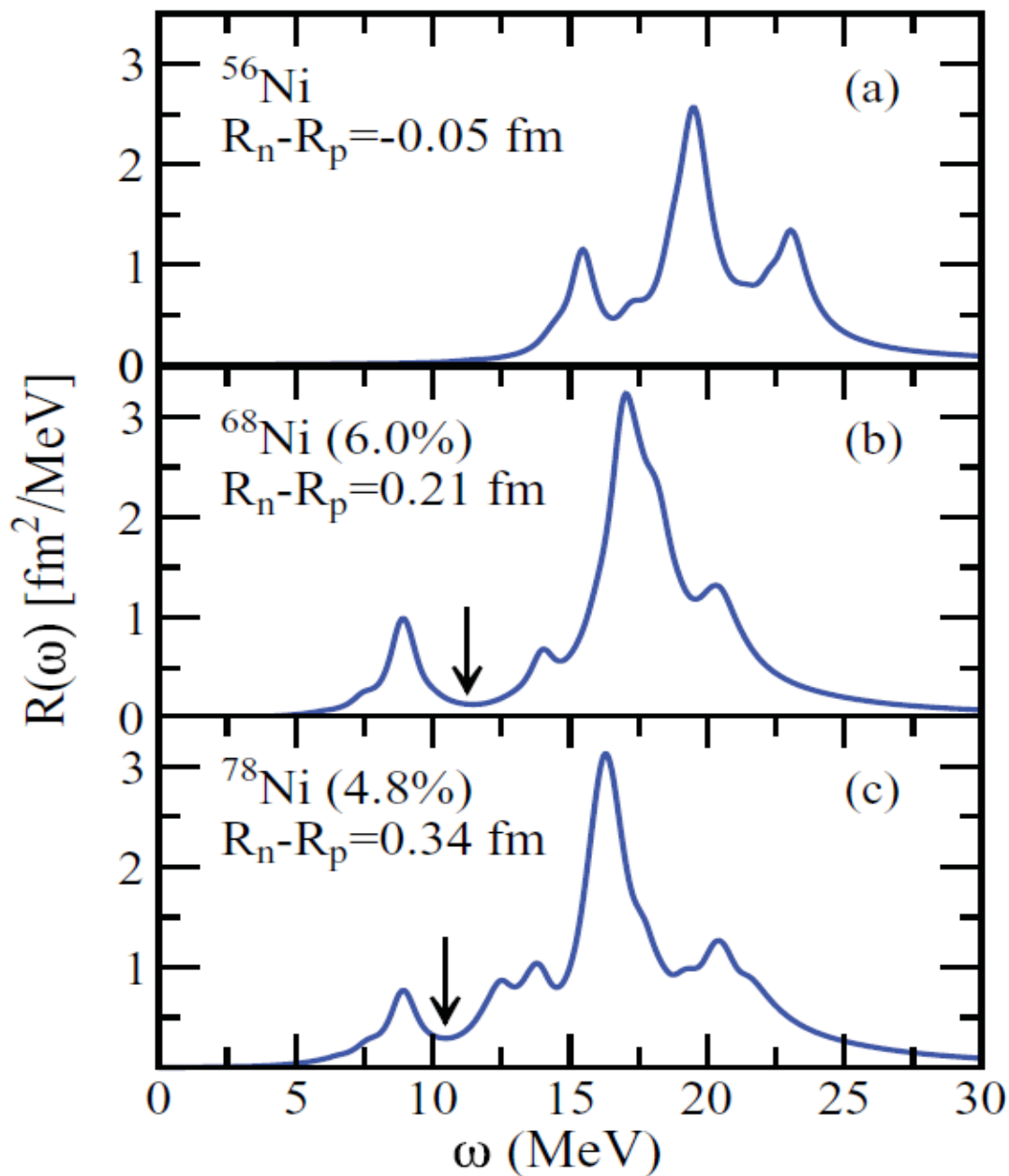


Self-consistent  
HF-based RPA  
results for the  
strength function  
 $S(E)$  distribution  
of the isovector  
dipole response

N. Auerbach *et al.*,  
PRC **89** (2014) 014335



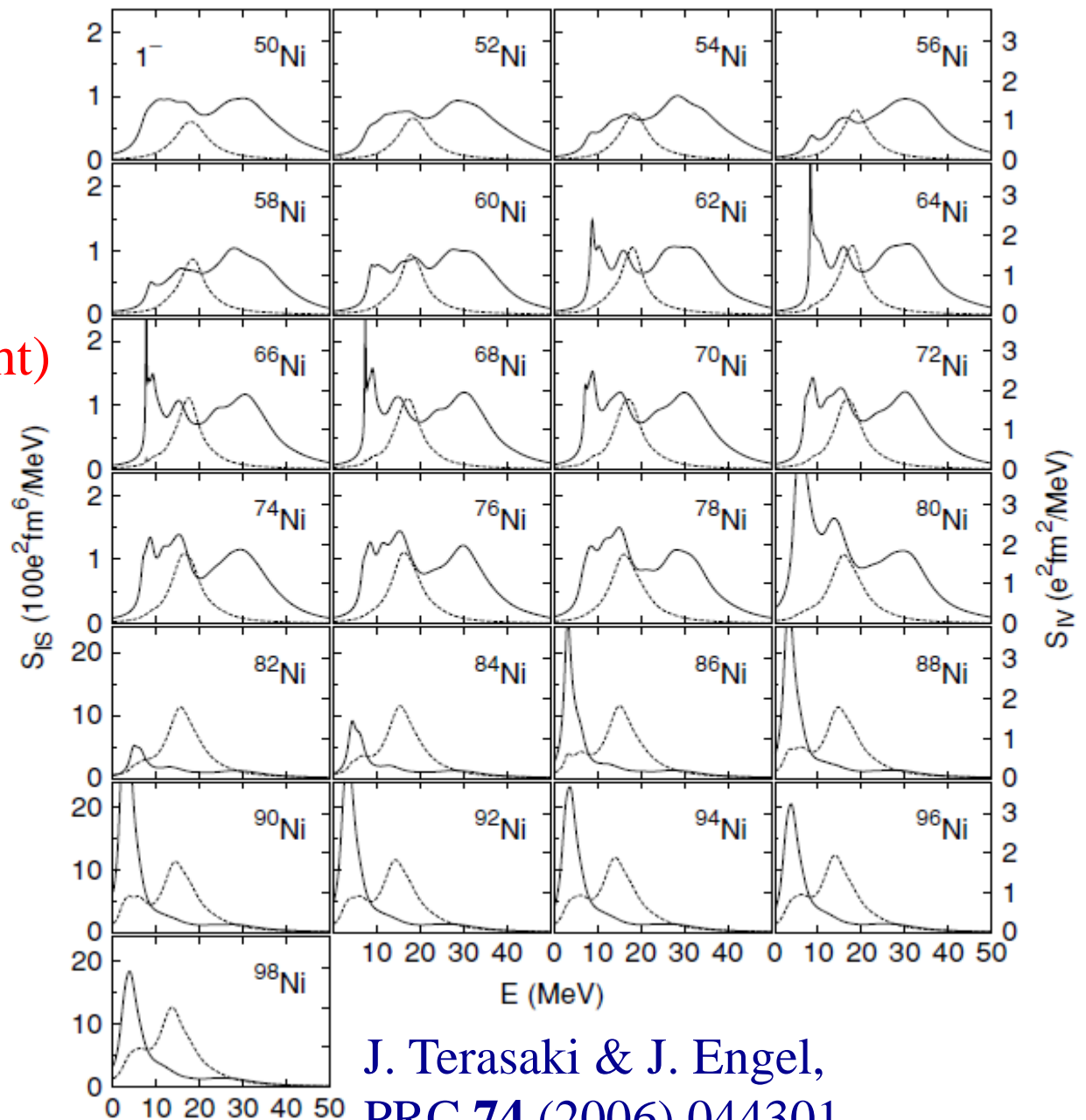
Distribution of isovector dipole strength for the three closed-(sub)shell nickel isotopes  $^{56}\text{Ni}$ ,  $^{68}\text{Ni}$ , and  $^{78}\text{Ni}$  calculated in HF-plus-RPA using the FSUGold interaction parameter set.



J. Piekarewicz,  
PRC **83** (2011) 034319

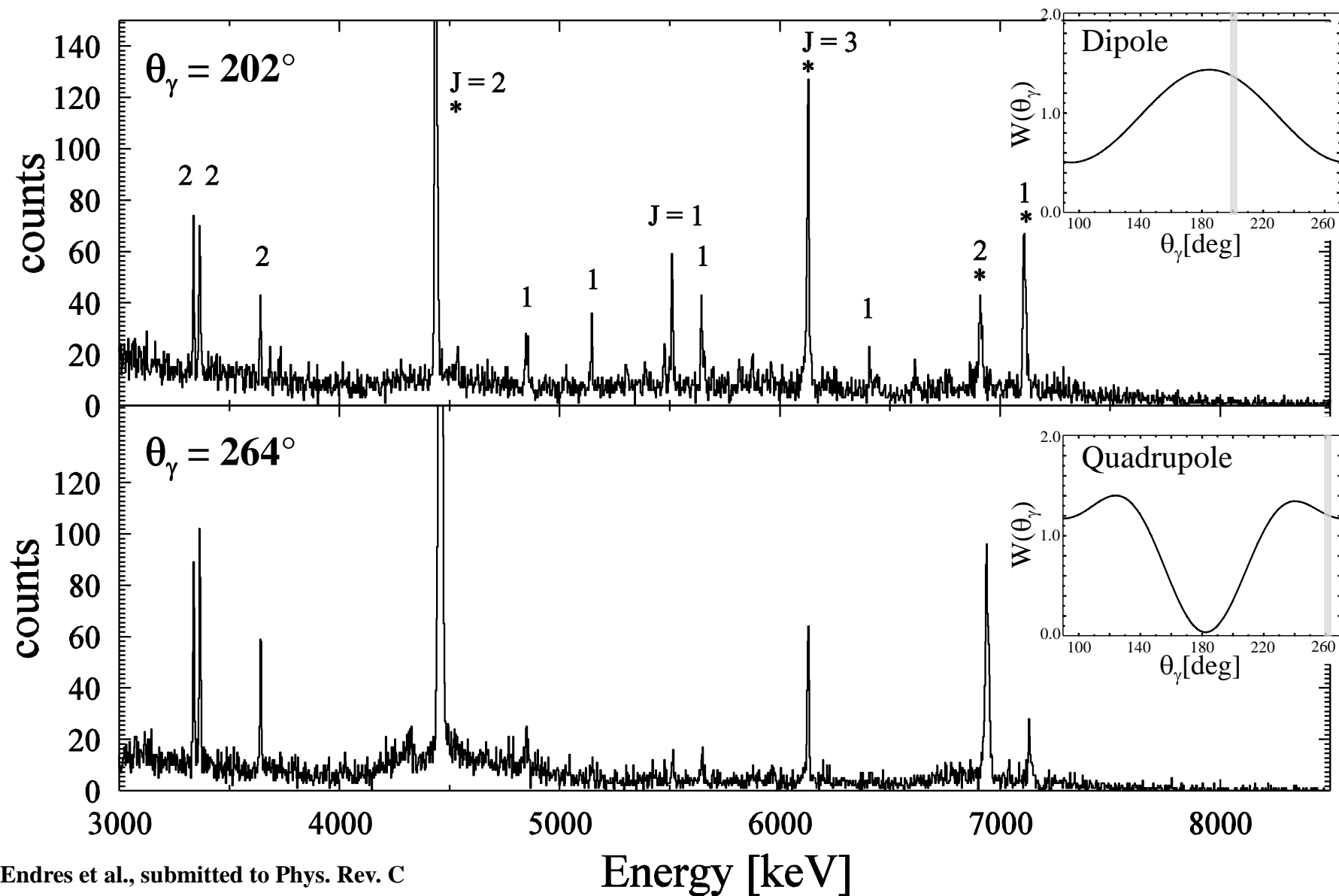


Isoscalar  
(solid, scale on left)  
and Iovector  
(dashed, scale on right)  
 $1^-$  strength functions  
for even Ni isotopes  
(SkM\* interaction).



J. Terasaki & J. Engel,  
PRC 74 (2006) 044301

# Multipole assignment with $\alpha$ - $\gamma$ angular correlation

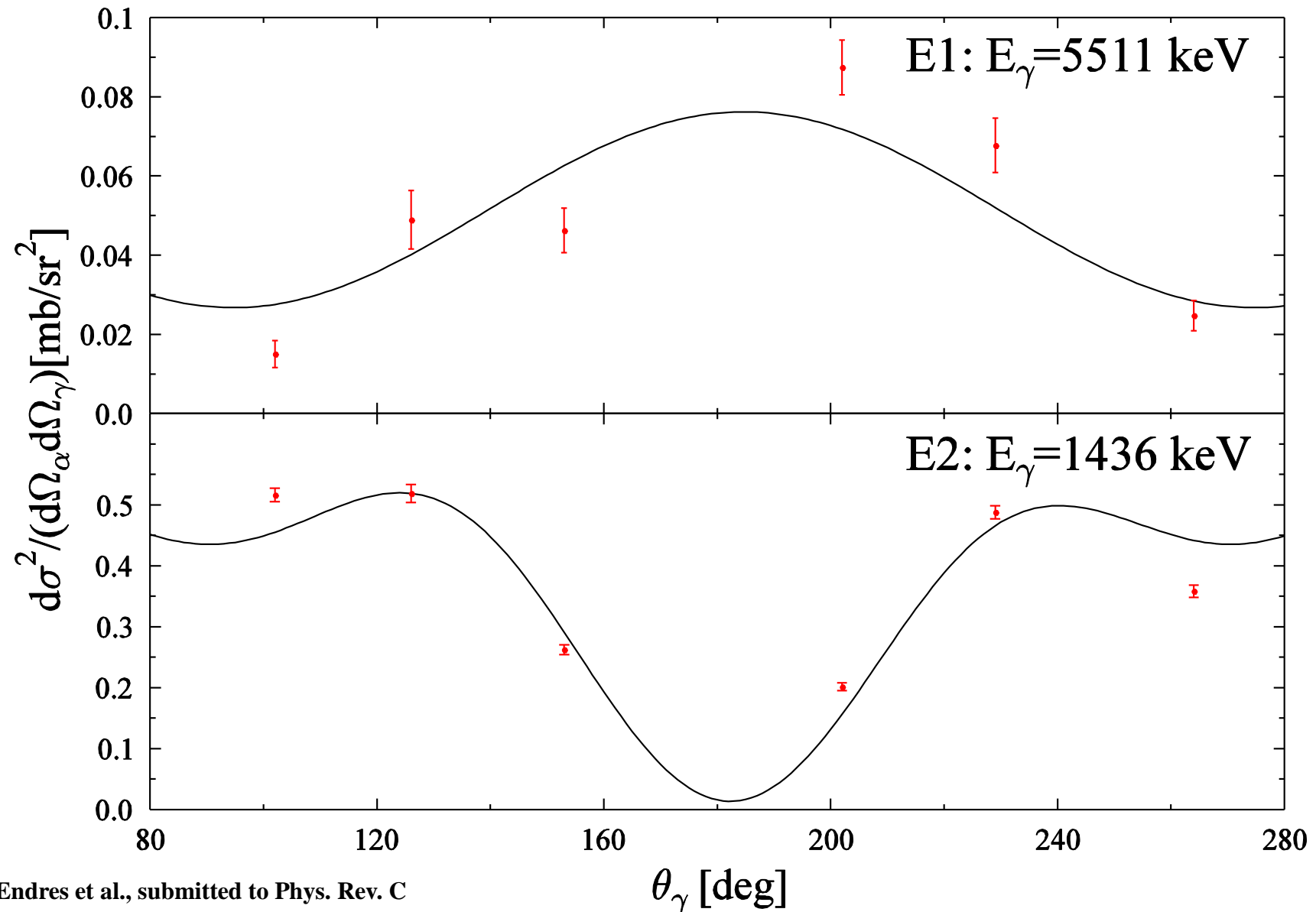


J. Endres et al., submitted to Phys. Rev. C

Nantes; 19-20 January 2015



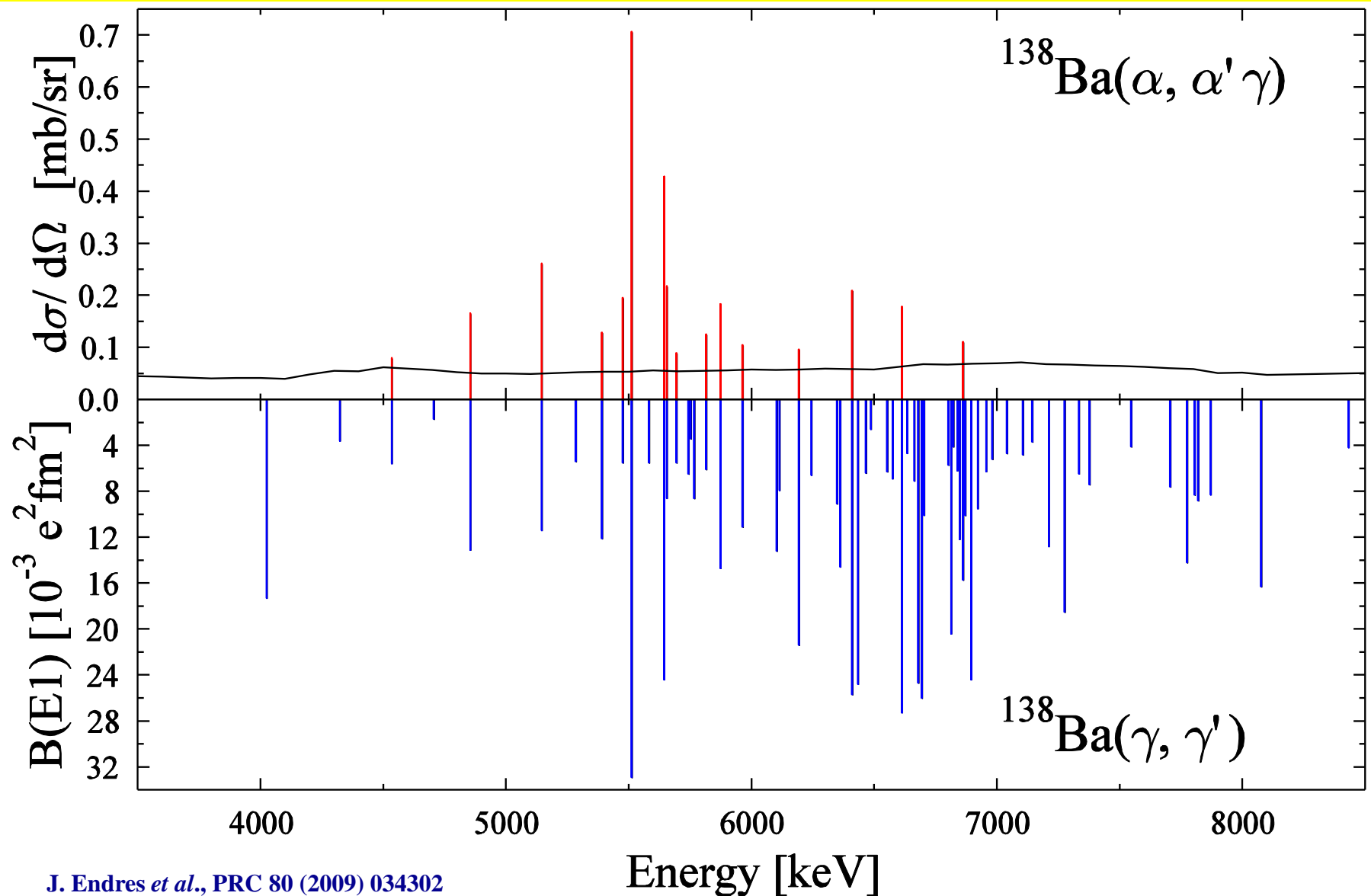
# Multipole assignment with $\alpha$ - $\gamma$ angular correlation



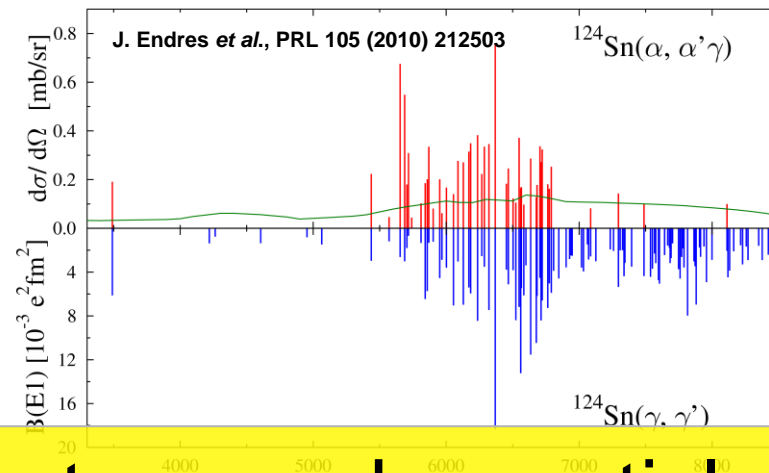
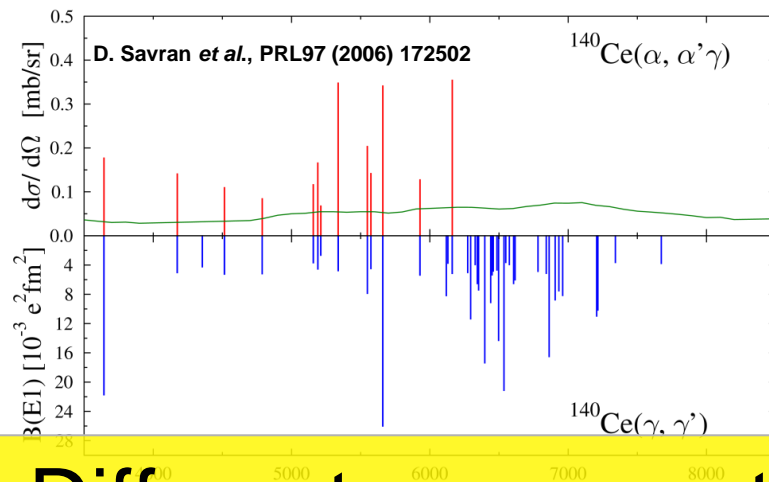
J. Endres et al., submitted to Phys. Rev. C

Nantes; 19-20 January 2015

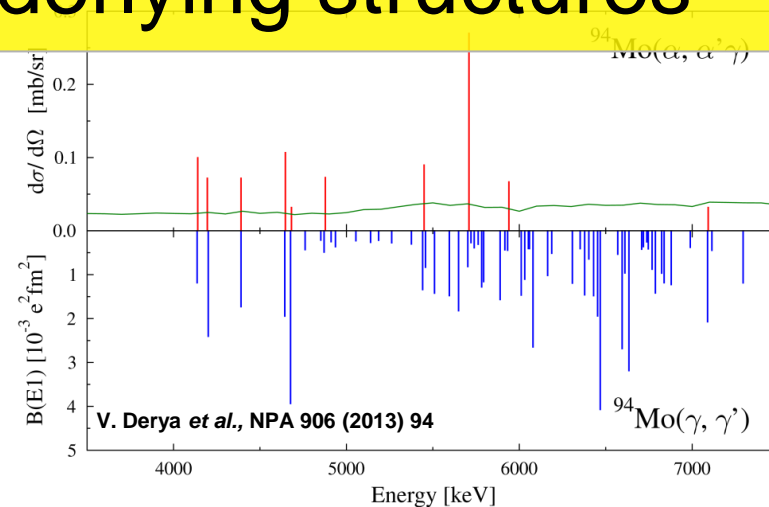
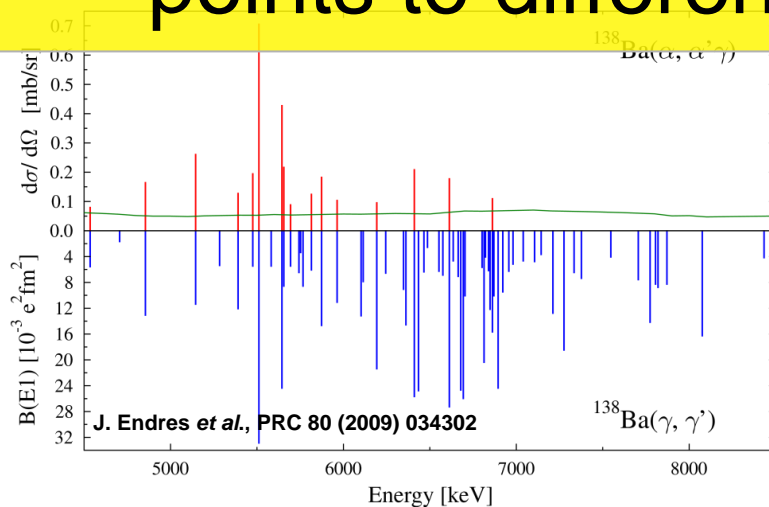
# Comparison of $(\alpha, \alpha'\gamma)$ with $(\gamma, \gamma')$ on $^{138}\text{Ba}$



# Systematics



Different response to photons and  $\alpha$ -particles points to different underlying structures



*Thank you for your attention*