

# An analytic hydrodynamical model of rotating 3D expansion

Márton NAGY

Eötvös University, Dept. of Atomic Physics

In collaboration with: Tamás Csörgő

MTA Wigner RCP, Budapest, Hungary

**2 December, 2014**

- Rotation in high-energy heavy-ion collision
- A new exact solution of fireball hydrodynamics
- Observables, signatures of rotation



Nemzeti  
Kiválóság  
Program

**TÁMOP 4.2.4.A**

# Introduction

## Motivation:

In non-central collisions:  
initial angular momentum  
leads to final state tilt of fireball

## Exact hydrodynamical solutions:

insight into the dynamics  
using simple analytic formulas

## Determination of rotation:

important question,  
theoretically, experimentally.

*Looking for rotating hydro solutions to follow time evolution, calculate observables, tell if there is rotation*

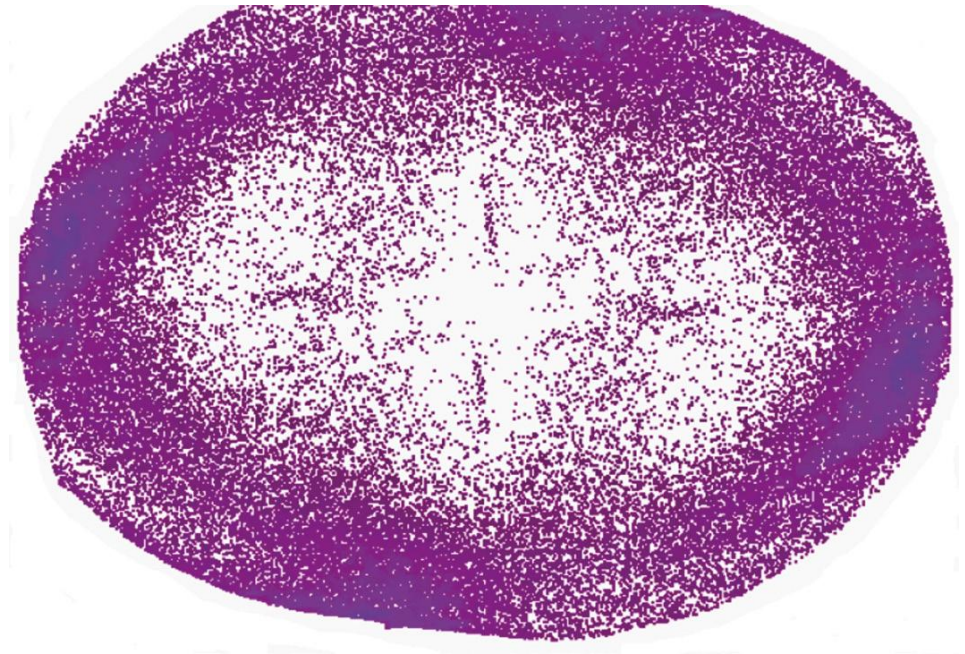
## This work follows the footsteps of:

*Csizmadia, Csörgő, Lukács, PLB443 (1998) 21*

*Csörgő, Acta Phys. Polon. B37 (2006) 483 & references therein*

*Csörgő, Akkelin, Hama, Lukács, Sinyukov, PRC67 (2003) 034094*

*Csörgő, Nagy, PRC89 (2014) 044901*



*Example from EPN 43/22 (2012) 91  
(L. Cifarelli, L.P. Csernai, H. Stöcker)*

# Equations of hydrodynamics

## Basic equations (non-relativistic, perfect fluid)

$$\partial_t n + \mathbf{v} \nabla n = -n \nabla \mathbf{v} \quad \text{mass conservation}$$

$$\partial_t \varepsilon + \mathbf{v} \nabla \varepsilon = -(\varepsilon + p) \nabla \mathbf{v} \quad \text{energy conservation}$$

$$\partial_t \mathbf{v} + (\mathbf{v} \nabla) \mathbf{v} = -\frac{1}{nm_0} \nabla p \quad \text{Euler equation}$$

Follow only from local thermal equilibrium & local conservation laws.

## Equation of State (EoS):

$$p = nT \quad \varepsilon = \kappa(T)p$$

**Frames:** inertial  $K$ , co-rotating:  $K'$

tilt angle (x-z plane):  $\mathcal{G}$

angular velocity: 
$$\boldsymbol{\omega} = \begin{pmatrix} 0 \\ \dot{\mathcal{G}} \\ 0 \end{pmatrix}$$

## Equations in rotating frame:

$$\left( \frac{\kappa}{T} + \frac{d\kappa}{dT} \right) (\partial_t + \mathbf{v}' \nabla') T + \nabla' \mathbf{v}' = 0$$

$$\partial_t n + \mathbf{v}' \nabla' n = -n \nabla' \mathbf{v}'$$

$$\partial_t \mathbf{v}' + (\mathbf{v}' \nabla') \mathbf{v}' = -\frac{1}{nm_0} \nabla' p + f'$$

$$f' = 2\mathbf{v}' \times \boldsymbol{\omega} + \mathbf{r}' \times \dot{\boldsymbol{\omega}} + \boldsymbol{\omega} \times (\mathbf{r}' \times \boldsymbol{\omega})$$

# A Hubble-like rotating solution

**Hubble-type velocity:**

$$\mathbf{v}' = \begin{pmatrix} \frac{\dot{X}}{X} r_x' + \dot{\psi} \frac{Z}{X} r_z' \\ \frac{\dot{Y}}{Y} r_y' \\ \frac{\dot{Z}}{Z} r_z' - \dot{\psi} \frac{X}{Z} r_x' \end{pmatrix}$$

**scaling variable:**

(rotating ellipsoids)

$$s = \frac{r_x'^2}{X^2} + \frac{r_y'^2}{Y^2} + \frac{r_z'^2}{Z^2}$$

„volume“:

$$V = XYZ$$

**coordinate transformation:**

$$r_x' = r_x \cos \mathcal{G} - r_z \sin \mathcal{G}$$

$$r_z' = r_x \sin \mathcal{G} + r_z \cos \mathcal{G}$$

$$r_y' = r_y$$

**Solution of the continuity equations:**

$$n(\mathbf{r}, t) = n_0 \frac{V_0}{V} e^{-s/2}$$

for constant  $\kappa$  :

$$T(\mathbf{r}, t) = T(t) = T_0 \left( \frac{V_0}{V(t)} \right)^{\frac{1}{\kappa}}$$

**for non-constant  $\kappa$ :**

$$T(\mathbf{r}, t) = T(t)$$

$$\frac{V_0}{V} = \exp \left( \int_{T_0}^T d\beta \left( \frac{d\kappa}{d\beta} + \frac{\kappa}{\beta} \right) \right)$$

**Similar to & generalization of:**

*Csörgő, Acta Phys. Polon. B37 (2006) 483 & references therein*

*Csörgő, Akkelin, Hama, Lukács, Sinyukov, PRC67 (2003) 034094*

*Csörgő, Nagy, PRC89 (2014) 044901*

# Arriving at a parametric solution

## Remaining functions:

Axes:  $X(t)$ ,  $Y(t)$ ,  $Z(t)$

Auxiliary function:  $\psi(t)$

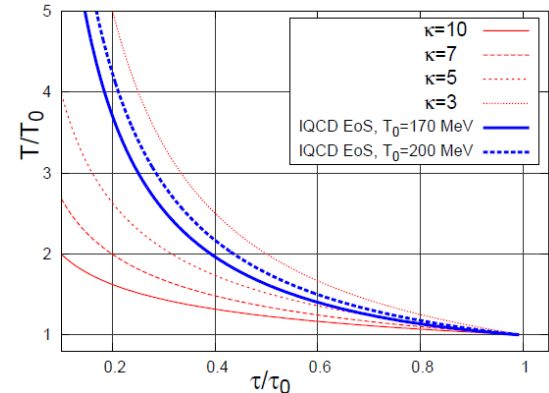
Angular velocity:  $\dot{\vartheta}(t)$

## Solution if

$$\psi(t) = \frac{\chi_0}{(X + Z)^2}$$

$$\dot{\vartheta}(t) = \frac{\chi_0}{(X + Z)^2}$$

*A slight (not quite physical) generalization is possible...*



## „Motion” of axes:

Governed by a Lagrangian!

In the constant  $\kappa$  case:

$$L = \frac{m}{2} (\dot{X}^2 + \dot{Y}^2 + \dot{Z}^2) - \frac{1}{\kappa} \left( \frac{V_0}{V} \right)^{1/\kappa} - \frac{m\chi_0^2}{(X + Z)^2}$$

Equations of motion:

Ordinary p.d.e.-s

Solution is easy numerically

**Conserved quantities:** total particle number:  $N_0 = n_0 V_0 (2\pi)^{3/2}$

total angular momentum (y component):  $M_y = m N_0 \chi_0$

total energy:  $E = \frac{m}{2} (\dot{X}^2 + \dot{Y}^2 + \dot{Z}^2) + \frac{1}{\kappa} \left( \frac{V_0}{V} \right)^{1/\kappa} + \frac{m\chi_0^2}{(X + Z)^2}$

# Observables

## Hydro evolution until freeze-out:

Dependent on EoS

## Our freeze-out criterion:

Same freeze-out temperature

Same time everywhere

Analytic results!

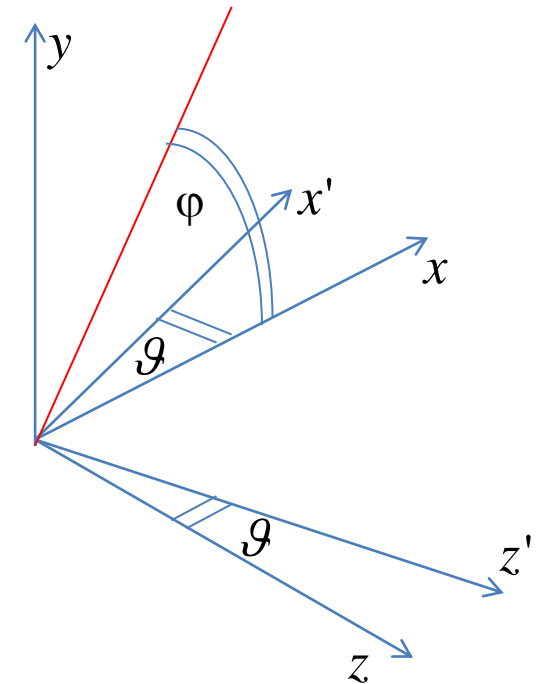
## Definitions & recipes:

Single-particle spectrum:

$$\frac{dn}{d^3\mathbf{p}} = \int d^3\mathbf{r} S(\mathbf{r}, \mathbf{p})$$

Bose-Einstein correlation:

$$C(\mathbf{K}, \mathbf{q}) = 1 + \lambda \frac{|\tilde{S}(\mathbf{K}, \mathbf{q})|^2}{|\tilde{S}(\mathbf{K}, 0)|^2}$$



## Source function:

$$S(\mathbf{r}, \mathbf{p}) = \frac{n(t_f, \mathbf{r})}{(2\pi m T_f)^{3/2}} \exp \left\{ -\frac{(\mathbf{p} - m\mathbf{v}(t_f, \mathbf{r}))^2}{2mT_f} \right\}$$

Important auxiliary quantity („speed of rotation“):

$$U = \frac{\chi_0}{X_f + Z_f}$$

# Observables: single-particle spectrum

**Definition:**

$$\frac{dn}{d^3\mathbf{p}} = \int d^3\mathbf{r} S(\mathbf{r}, \mathbf{p})$$

**Result:**

$$\frac{dn}{d^3\mathbf{p}} \propto \exp\left(-\frac{p_x^2}{2mT_x} - \frac{p_y^2}{2mT_y} - \frac{p_z^2}{2mT_z} - \frac{\beta_{xz}}{m} p_x p_z\right)$$

Auxiliary quantities:

$$T_y^* = T_f + m\dot{Y}^2$$

$$T_x^* = T_f + m(\dot{X}^2 + U^2), \quad T_z^* = T_f + m(\dot{Z}^2 + U^2), \quad \beta_* = mU(\dot{X} - \dot{Z})$$

**Coefficients:**

$$T'_x = T_x^* + \frac{\beta_*^2}{T_z^*}$$

$$T'_z = T_z^* + \frac{\beta_*^2}{T_x^*}$$

$$\beta'_{xz} = \frac{\beta_*}{T_x^* T_z^* - \beta_*^2}$$

**In lab frame:**

$$\frac{1}{T_x} = \frac{\cos^2 \vartheta_f}{T'_x} + \frac{\sin^2 \vartheta_f}{T'_z} + \beta'_{xz} \sin(2\vartheta_f)$$

$$\frac{1}{T_z} = \frac{\sin^2 \vartheta_f}{T'_x} + \frac{\cos^2 \vartheta_f}{T'_z} - \beta'_{xz} \sin(2\vartheta_f)$$

$$\beta_{xz} = \frac{\sin(2\vartheta_f)}{2} \left( \frac{1}{T'_z} - \frac{1}{T'_x} \right) + \beta'_{xz} \cos(2\vartheta_f)$$

# Observables: average $p_T$ -spectrum

## Definition:

$$E \frac{dn}{d^3\mathbf{p}} = \frac{1}{2\pi p_T} \frac{dn}{dp_T dy} \left\{ 1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\varphi - \Psi_n)] \right\}$$

## In this model:

***All event planes coincide!***

## Result

(same as in Csörgő et al., PRC67 (2003) 034094)

Azimuthal average, (up to 2nd order in  $v$ ):

$$\frac{1}{2\pi p_T} \frac{dn}{dp_T dy} \propto \exp\left(-\frac{p_z^2}{2mT_z} - \frac{p_T^2}{2mT_{\text{eff}}}\right) \times \left[ I_0(w) + \frac{v^2}{4} (I_0(w) + I_1(w)) \right]$$

## Auxiliary quantities:

$$w = \frac{p_T^2}{4m} \left( \frac{1}{T_y} - \frac{1}{T_x} \right)$$

$$T_{\text{eff}} = \frac{1}{2} \left( \frac{1}{T_y} + \frac{1}{T_x} \right)$$

$$v = -\frac{\beta_{xz}}{m} p_z p_T$$



# Observables: flow coefficients

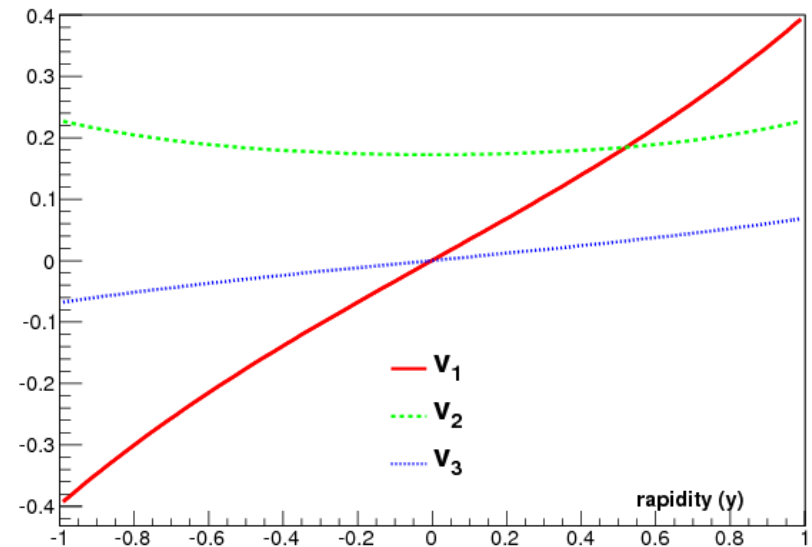
**Results:** (same as in Csörgő *et al.*, *PRC67* (2003) 034094):

Flow coefficients (up to 2nd order in  $v$ ):

$$v_1 = \frac{v}{2} \left[ 1 + \frac{I_1(w)}{I_0(w)} \right]$$
$$v_2 = \frac{I_1(w)}{I_0(w)} + \frac{v^2}{8} \left[ 1 + \frac{I_2(w)}{I_0(w)} - 2 \frac{I_1^2(w)}{I_0^2(w)} \right]$$
$$v_3 = \frac{v}{2} \left[ \frac{I_1(w)}{I_0(w)} + \frac{I_2(w)}{I_0(w)} \right]$$

Modified Bessel function:

$$I_\nu(x) = \int_0^\pi d\varphi \frac{\cos(\nu\varphi)}{\pi} e^{x \cos\varphi}$$



**Universal scaling of  $v_2$  is preserved**  
(Csanád *et al.*, *Nucl.Phys. A742* (2004) 80)

**Directed flow: opposite in two directions!**  
*Useful for rotation orientation determination*

# Observables: two-particle correlations

**Bose-Einstein correlations (raw):**

**Definition:**

$$C(\mathbf{K}, \mathbf{q}) = 1 + \lambda \frac{|\tilde{S}(\mathbf{K}, \mathbf{q})|^2}{|\tilde{S}(\mathbf{K}, 0)|^2}$$

$$\tilde{S}(\mathbf{K}, \mathbf{q}) = \int d^3\mathbf{r} e^{i\mathbf{q}\cdot\mathbf{r}} S(\mathbf{r}, \mathbf{K})$$

**Result: Gaussian**

$$C(\mathbf{K}, \mathbf{q}) = 1 + \lambda \exp\left(-\sum_{ij} q_i q_j R_{ij}^2\right)$$

**Radii in tilted frame:**

$$R_x'^2 = X^2 \frac{T_f}{T_x^*}$$

$$R_y'^2 = Y^2 \frac{T_f}{T_y^*}$$

$$R_z'^2 = Z^2 \frac{T_f}{T_z^*}$$

$$R_{xz}'^2 = XZT_f \beta'_{xz}$$

**In the lab frame:**

$$R_x^2 = R_x'^2 \cos^2 \vartheta_f + R_z'^2 \sin^2 \vartheta_f + R_{xz}'^2 \sin(2\vartheta_f)$$

$$R_y^2 = R_y'^2$$

$$R_x^2 = R_x'^2 \sin^2 \vartheta_f + R_z'^2 \cos^2 \vartheta_f - R_{xz}'^2 \sin(2\vartheta_f)$$

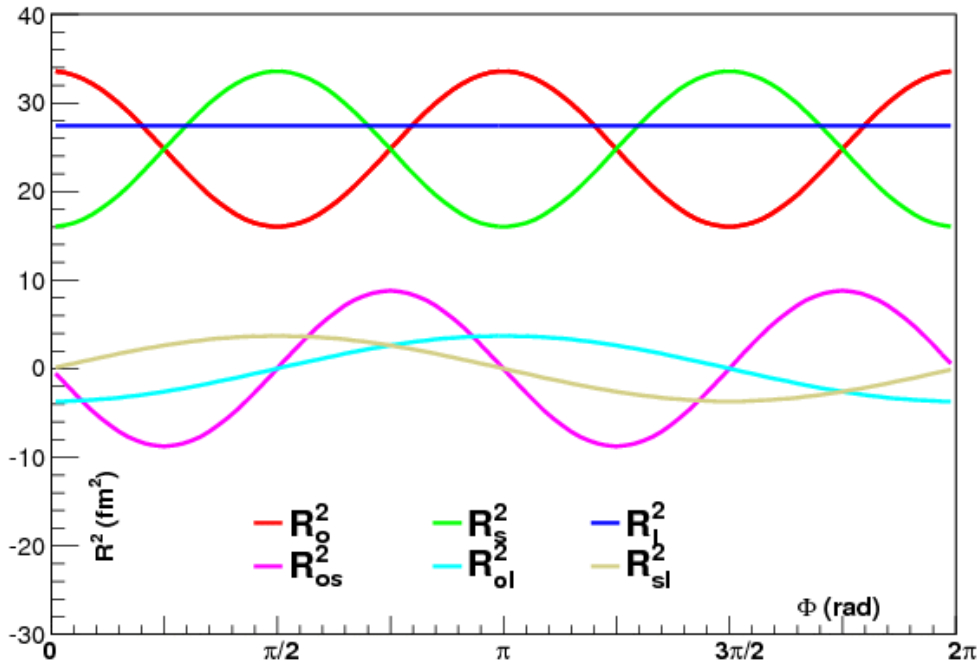
$$R_{xz}^2 = \frac{\sin(2\vartheta_f)}{2} (R_x'^2 - R_z'^2) + \cos(2\vartheta_f) R_{xz}'^2$$

*Coulomb correction treatment important!*

# Observables: two-particle correlations

**Bertsch-Pratt parametrization:**

$$C(\mathbf{K}, \mathbf{q}) = 1 + \lambda \exp\left(-\sum_{i,j=o,s,l} q_i q_j R_{ij}^2\right)$$



**HBT Bertsch-Pratt radii:**

$$R_o^2 = R_x^2 \cos^2 \varphi + R_y^2 \sin^2 \varphi$$

$$R_s^2 = R_x^2 \sin^2 \varphi + R_y^2 \cos^2 \varphi$$

$$R_l^2 = R_z^2$$

$$R_{os}^2 = (R_y^2 - R_x^2) \sin \varphi \cos \varphi$$

$$R_{ol}^2 = R_{xz}^2 \cos \varphi$$

$$R_{sl}^2 = -R_{xz}^2 \sin \varphi$$

**Oscillations: s, o, os, ol, sl !**

ol, sl values vanish for no tilt:  $\mathcal{G} = 0$

No 2nd order oscillations of ol, sl!

# Signatures of rotation

---

## Observables that appear because of tilt ( $\mathcal{G} = 0$ ):

Flow coefficients vs. rapidity:  $v_3$  and  $v_1$  have characteristic rapidity dependence!

## **HBt-radii**: azimuthal oscillation vs. reaction plane:

out-side: known, measured multiple times (characteristic to ellipsoid-like sources)

- eg. STAR beam energy scan HBT paper: extensive oscillation measurements

(*arXiv:1403.4972*)

out-long, side-long: not measured yet above AGS energies (as far as I know)

Interesting ***ol, sl oscillations*** depend on  $\cos \varphi$  and  $\sin \varphi$ :

1th order event plane needed! (fluctuations may interplay...)

*Rotation detection with HBT measurements is possible!*

Flow coefficients vs. rapidity: may help to establish 1th order event plane

\*\*\*

## Effect of physical rotation (not jus final tilt):

***Eigen-frames*** of the rotated ***coordinate-space*** ellipsoid, of the ***single particle***

***spectrum*** and that of the ***HBT correlation*** functions are ***three different*** frames!

Presented hydro solution allows a dynamical investigation of the rotation vs time!

# Summary and outlook

---

## **Hydrodynamical models:**

dynamically connect initial state with final state

## **Non-relativistic hydrodamical equations:**

Extended the class of known parametric solutions

Found rotating, 3-axis ellipsoidal solutions, for arbitrary EoS

Calculated observables

## **Simplified treatment**

Needs at least some relativistic generalization (eg. HBT radii now do not depend on transverse momentum) Relativistic generalization of solutions: hard task...

## **Signals of rotation:**

Directed & third flow vs. Rapidity

1th order oscillations of out-long, side-long HBT radii parameters

## **Measurement of tilt angle:**

A final state variable; initial angle known to be zero?

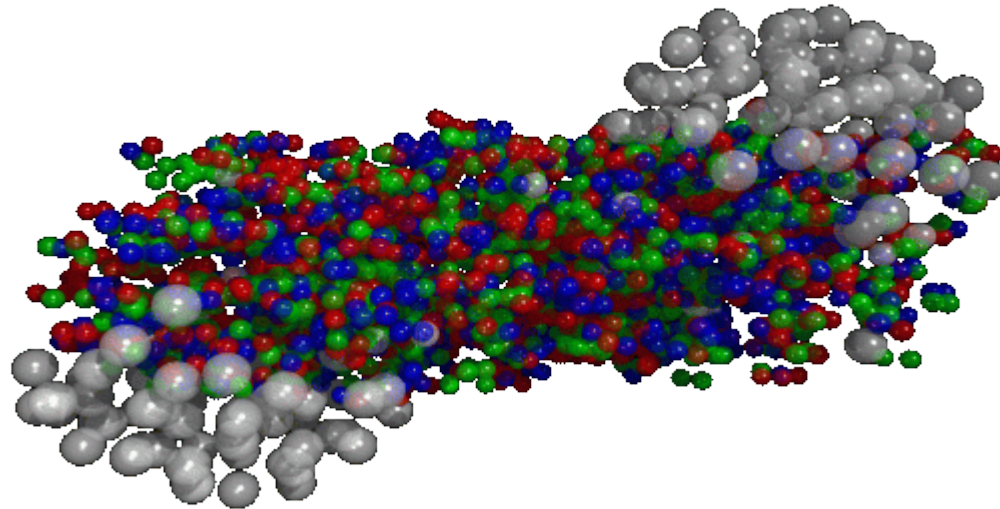
Possible to investigate EoS!

## **Further work needed:**

Extension of the Buda-Lund model to this rotating case: a semi-analytic relativistic generalization, *coming soon...*

---

Thank you for your attention!



This research was partly supported by the European Union and the State of Hungary, co-financed by the European Social Fund in the framework of TÁMOP 4.2.4. A/1-11-1-2012-0001 'National Excellence Program'.

---

Spare slides coming...  
*(soon)*