

QUARK MATTER 2015

October 2nd 2015



Experimental overview on quarkonium production

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INFN Torino



Outlook

2

Overview of recent/new results on charmonium and bottomonium production in pp, p-A and A-A collisions from RHIC to LHC energies

11 talks + 10 posters on quarkonium experimental results



Apologize for all the results which could not fit in this talk!!!

Quarkonium at RHIC and LHC³

Facility	Experiment	System	$\sqrt{s_{NN}}$ (GeV)	Data taking
RHIC	PHENIX/STAR	Au-Au, Cu-Cu, Cu-Au, U-U	200, 193, 62, 39	2000-2015
		p-A, d-Au	200	
		pp	200-500	
LHC	ALICE/ATLAS /CMS/LHCb	Pb-Pb	2760	2010-2012
		p-Pb	5020	2013
		pp	2760, 7000, 8000, 13000	2010-2015

Quarkonium at RHIC and LHC⁴

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				2010-2015

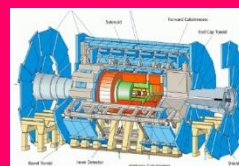
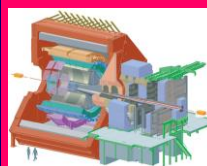
Quarkonium production
investigated via collisions:

- with different beam species
- at various energies

Quarkonium at RHIC and LHC ⁵

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RHIC	PHENIX/STAR	Au-Au, Cu-Cu, Cu-Au, U-U	200, 193, 62, 39	2000-2015
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			200-500	
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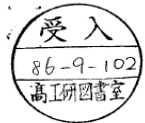
All LHC experiments investigate quarkonium production



- high luminosity and energies
- complementary results due to different kinematic coverages of the experiments

AA: from suppression...

6



PHYS. LETT. B, in press

BROOKHAVEN NATIONAL LABORATORY

June 1986

BNL-38344

J/ψ SUPPRESSION BY QUARK-GLUON PLASMA FORMATION

T. Matsui

Center for Theoretical Physics
Laboratory for Nuclear Science
Massachusetts Institute of Technology
Cambridge, MA 02139, USA

and

H. Satz

Fakultät für Physik
Universität Bielefeld, D-48 Bielefeld, F.R. Germany
and
Physics Department
Brookhaven National Laboratory, Upton, NY 11973, USA

ABSTRACT

If high energy heavy ion collisions lead to the formation of a hot quark-gluon plasma, then colour screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. To study this effect, we compare the temperature dependence of the screening radius, as obtained from lattice QCD, with the J/ψ radius calculated in charmonium models. The feasibility to detect this effect clearly in the dilepton mass spectrum is examined. We conclude that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark-gluon plasma formation.

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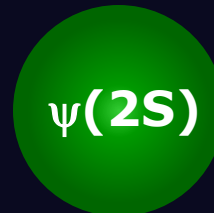
the original idea:

quarkonium production suppressed via color screening in the QGP



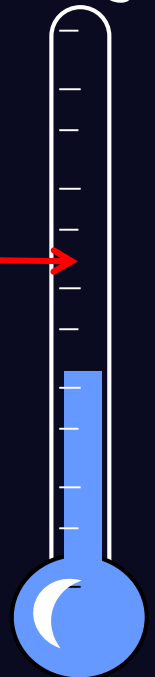
sequential melting

differences in the quarkonium binding energies lead to a sequential melting with increasing temperature



T_c →

$T < T_c$



Quarkonium is a QGP thermometer

AA: from suppression...

7



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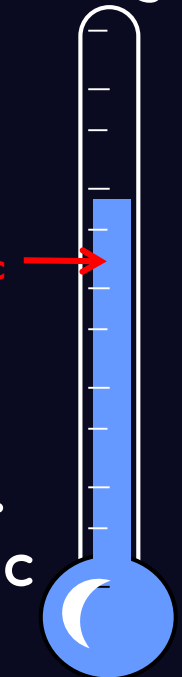
differences in the quarkonium binding energies lead to a sequential melting with increasing temperature

$\psi(2S)$

J/ψ

$\Upsilon(1S)$

T_c



$T \sim 1.5 T_c$

Quarkonium is a QGP thermometer

AA: from suppression...

8



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J/ψ SUPPRESSION BY QUARK-GLUON PLASMA FORMATION

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the original idea:

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sequential melting

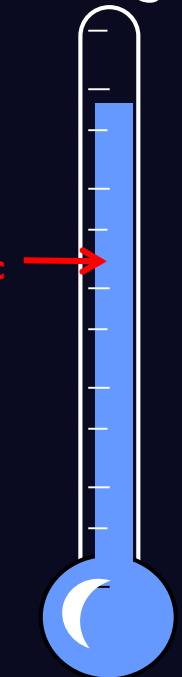
differences in the quarkonium binding energies lead to a sequential melting with increasing temperature

$\psi(2S)$

J/ψ



T_c



$T > T_c$

T.Matsui and H.Satz, Phys.Lett.B178 (1986) 416

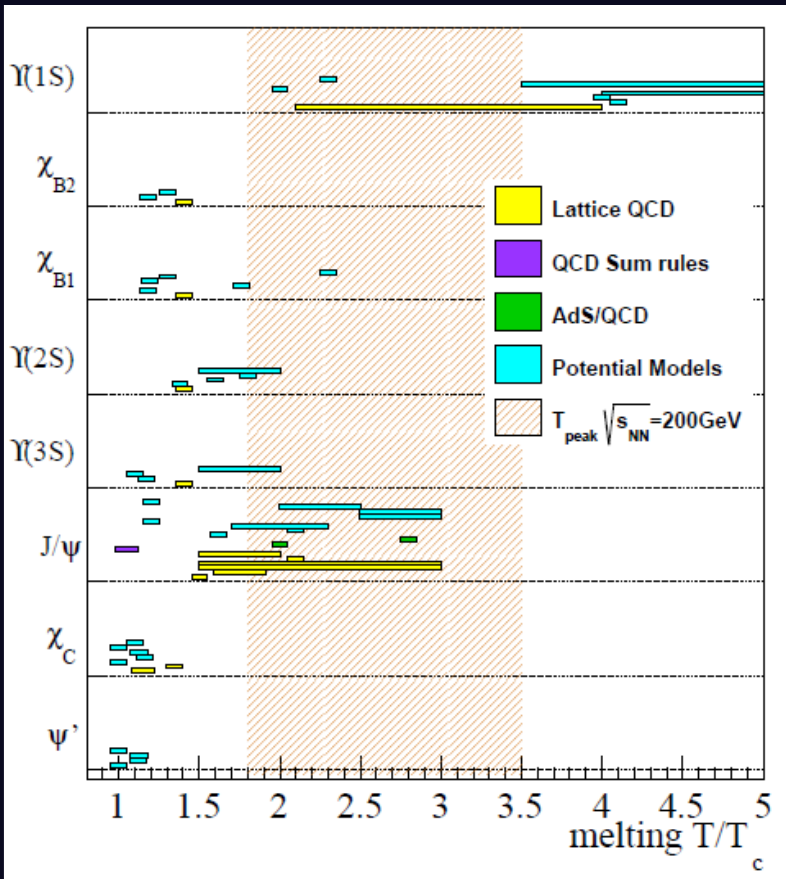
Quarkonium is a QGP thermometer

AA: from suppression...

9

→ **the original idea:**
quarkonium production suppressed
via color screening in the QGP

→ **sequential melting**
differences in the quarkonium binding
energies lead to a sequential melting
with increasing temperature



PHENIX, Phys.Rev C91, 024913

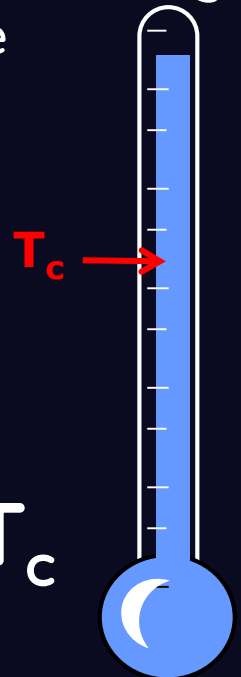
Quarkonium is a QGP thermometer

$\psi(2S)$

J/ψ

$\Upsilon(1S)$

$T \gg T_c$



...to recombination

10

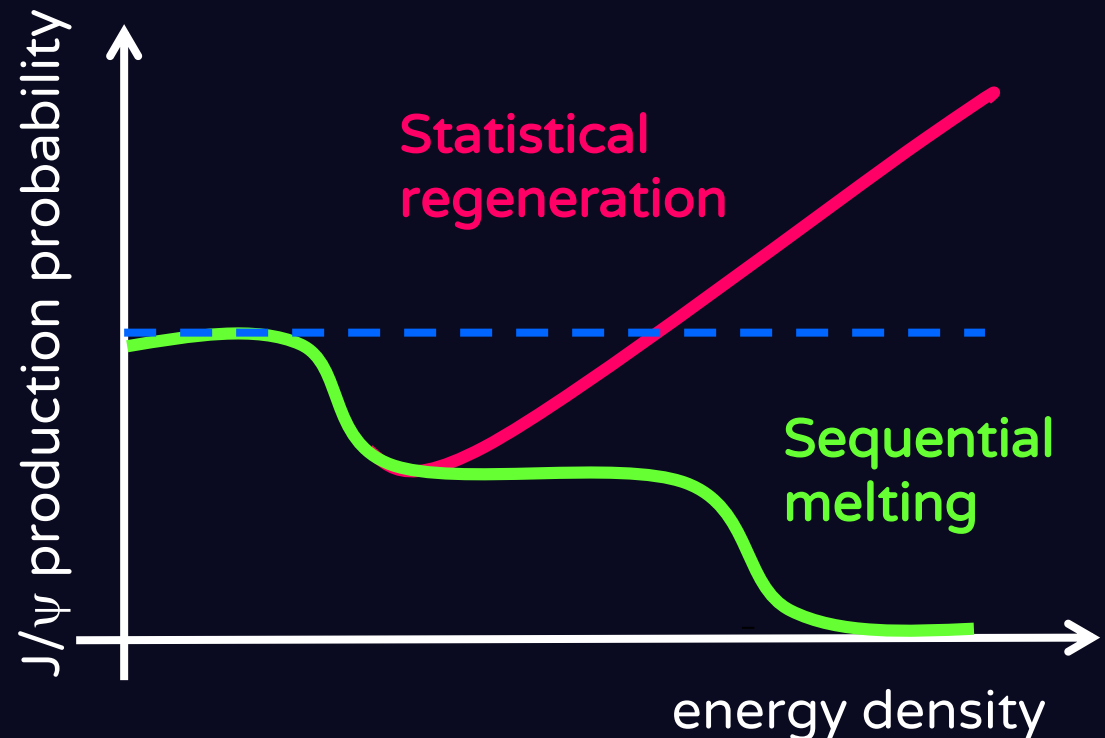
➡ (Re)combination

increasing the collision energy the cc pair multiplicity increases

Central AA collisions	SPS 20 GeV	RHIC 200 GeV	LHC 2.76TeV
N_{ccbar}/event	~ 0.2	~ 10	~ 85

enhanced quarkonia production via (re)combination at hadronization or during QGP phase

P. Braun-Muzinger, J. Stachel,
PLB 490(2000) 196
R. Thews et al,
Phys.Rev.C63:054905(2001)



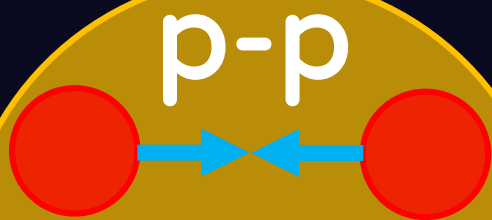
➔ on top of mechanisms related to hot matter, other cold matter effects might affect quarkonium production

- nuclear parton shadowing/color glass condensate
- energy loss
- $c\bar{c}$ in medium break-up

} investigated
in p-A (d-A)
collisions

➔ The assessment of the size of these effects is fundamental to interpret quarkonium A-A results

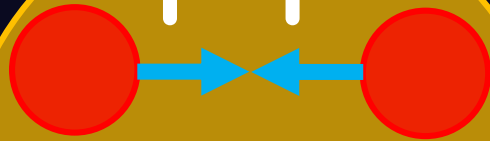
Summarizing quarkonium in HI¹²



“vacuum” reference
for A-A and p-A,
genuine pp physics
program

Summarizing quarkonium in HI¹³

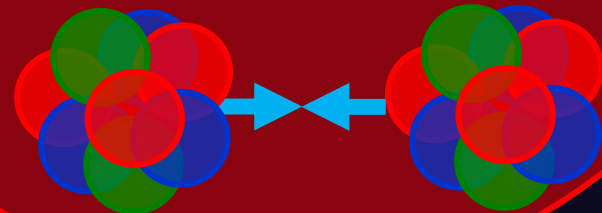
p-p



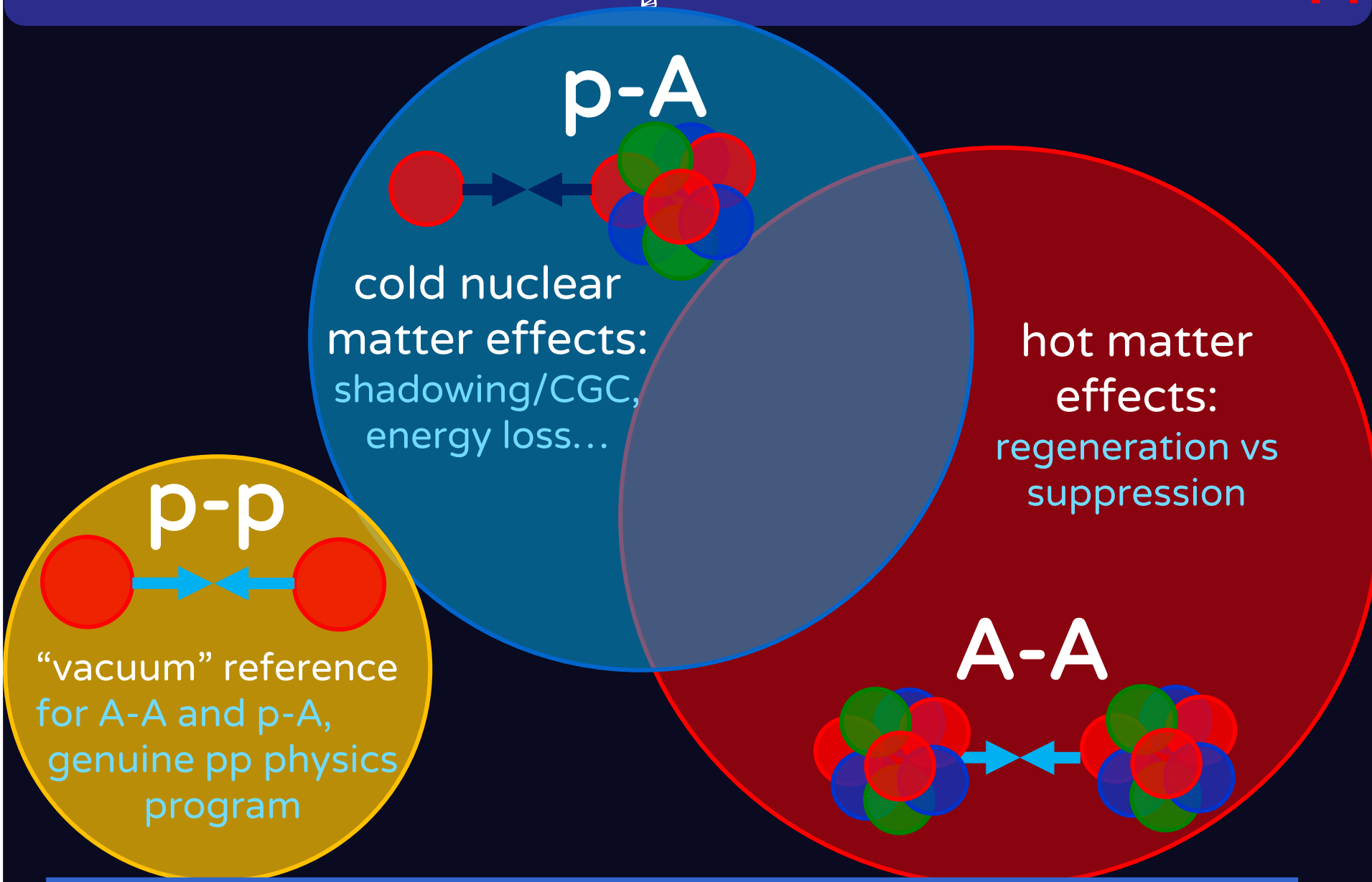
“vacuum” reference
for A-A and p-A,
genuine pp physics
program

hot matter
effects:
regeneration vs
suppression

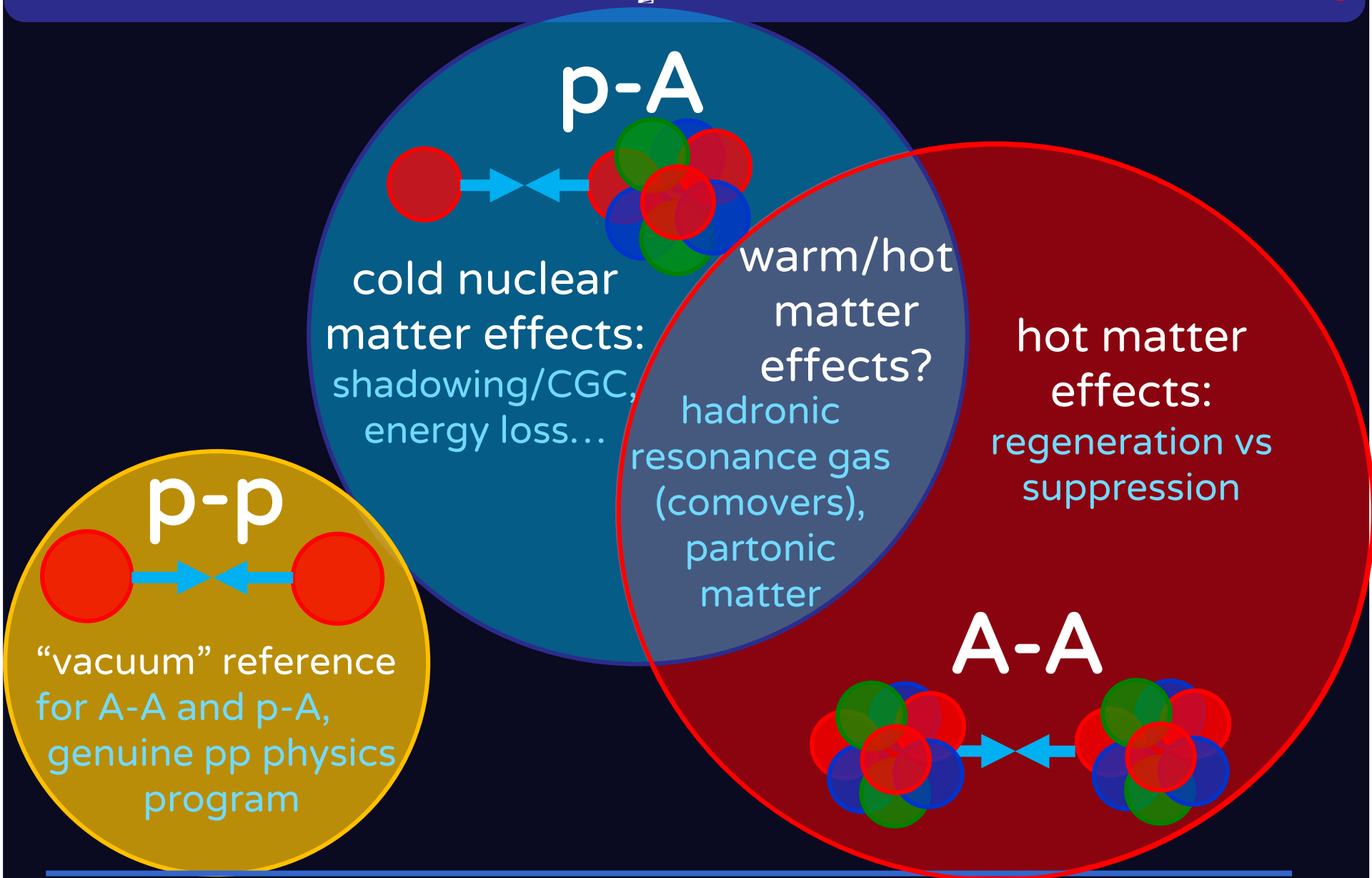
A-A



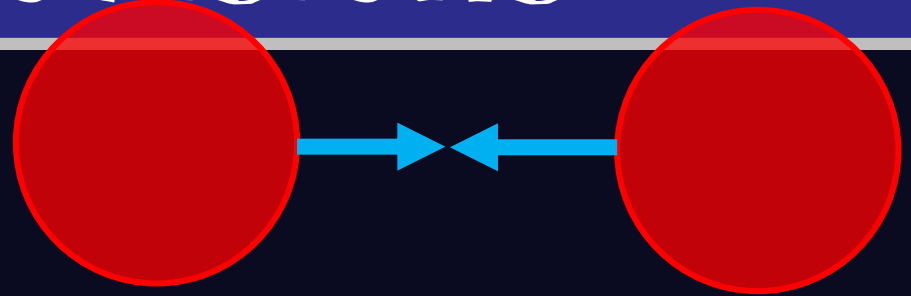
Summarizing quarkonium in HI₁₄



Summarizing quarkonium in HI¹⁵



Quarkonium in pp collisions



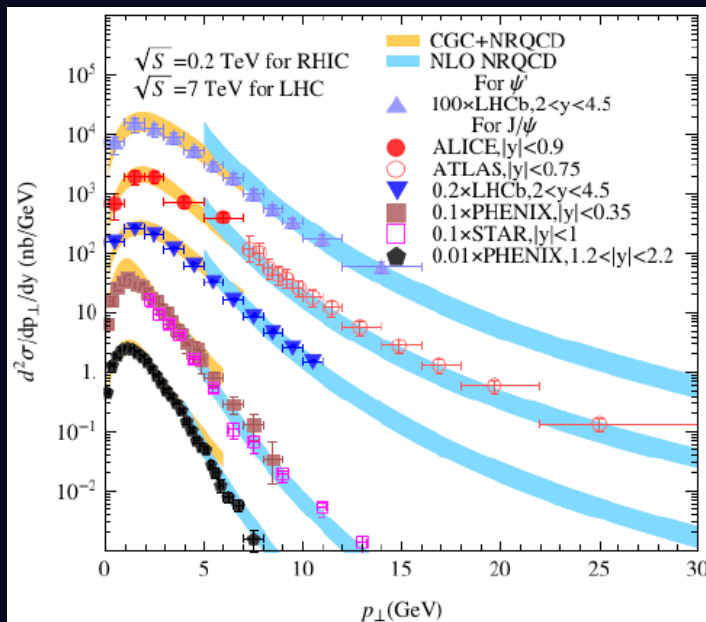
Recent news on J/ψ in pp

17

➔ pp as testing ground for quarkonium production mechanisms

- results now cover a large \sqrt{s} range
- very high p_T reached (ATLAS and CMS up to $\sim 100 \text{ GeV}/c$)

CMS. PRL 114.191802(2015)



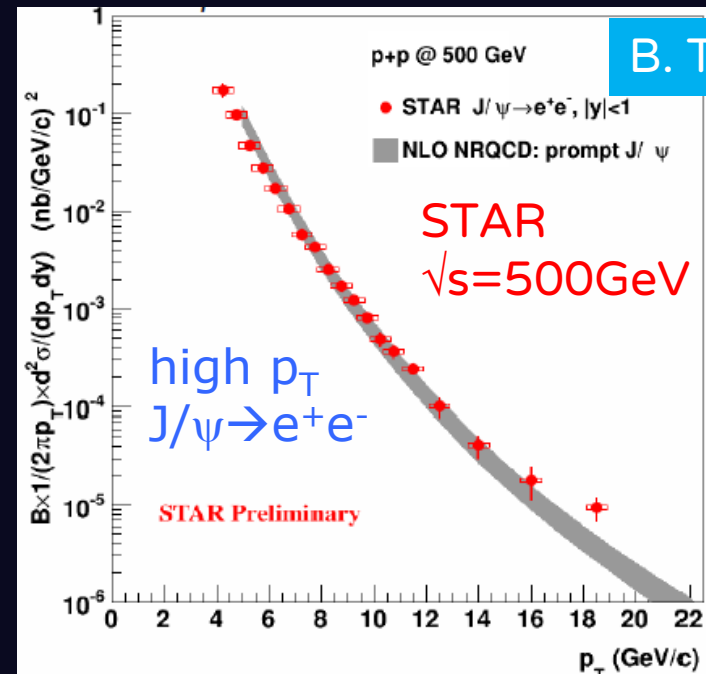
Ma et al. PRL 113 (2014)192301

NLO NRQCD

➔ describes production for $p_T > 4 \text{ GeV}/c$

CGC+NRQCD

➔ description of the low p_T range



B. Trzeciak

NLO NRQCD describes STAR results @ $\sqrt{s} = 500 \text{ GeV}$ ($p_T > 4 \text{ GeV}/c$)

Shao et al. JHEP 1505 (2015) 103

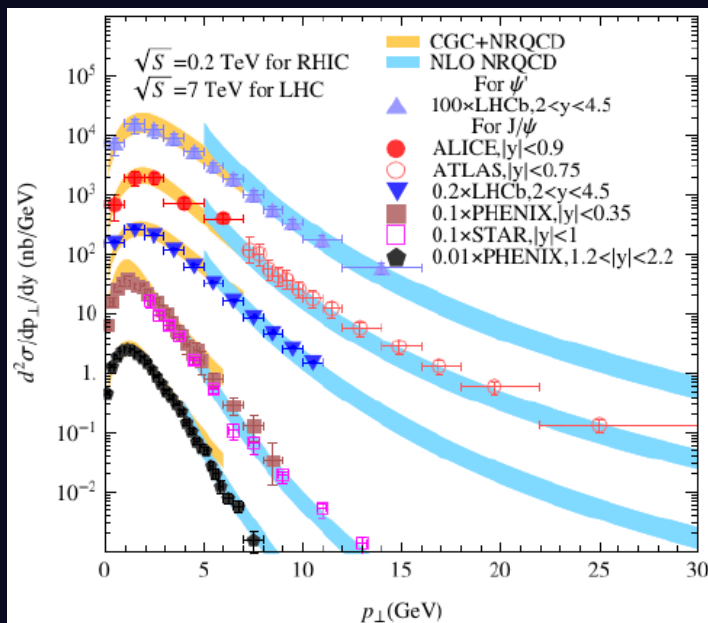
Recent news on J/ψ in pp

18

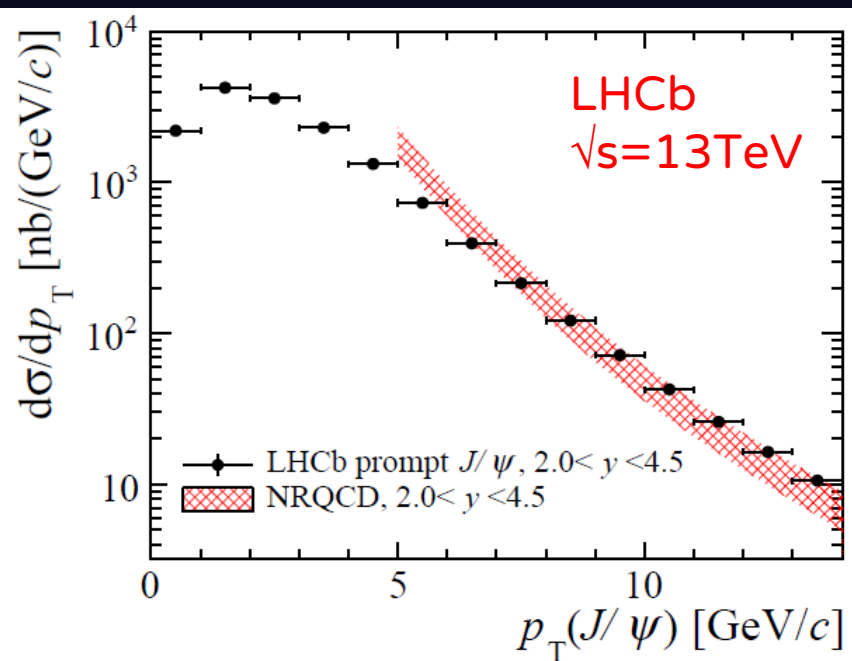
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CMS, PRL 114,191802(2015)



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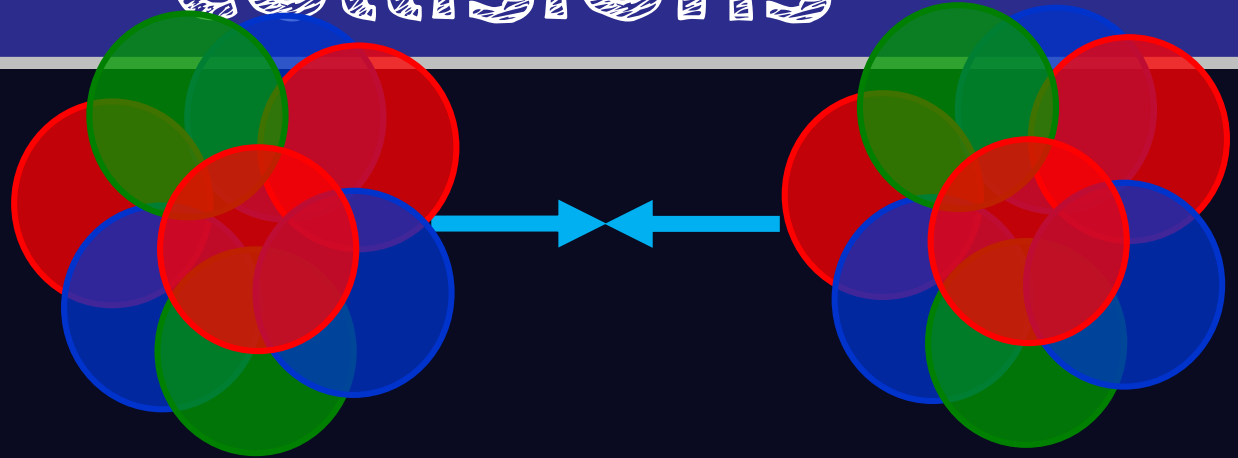
CGC+NRQCD

➔ description of the low p_T range

NLO NRQCD describes high p_T data up to the LHC top energy ($\sqrt{s} = 13 \text{ TeV}$)

LHCb, arXiv:1509.00771

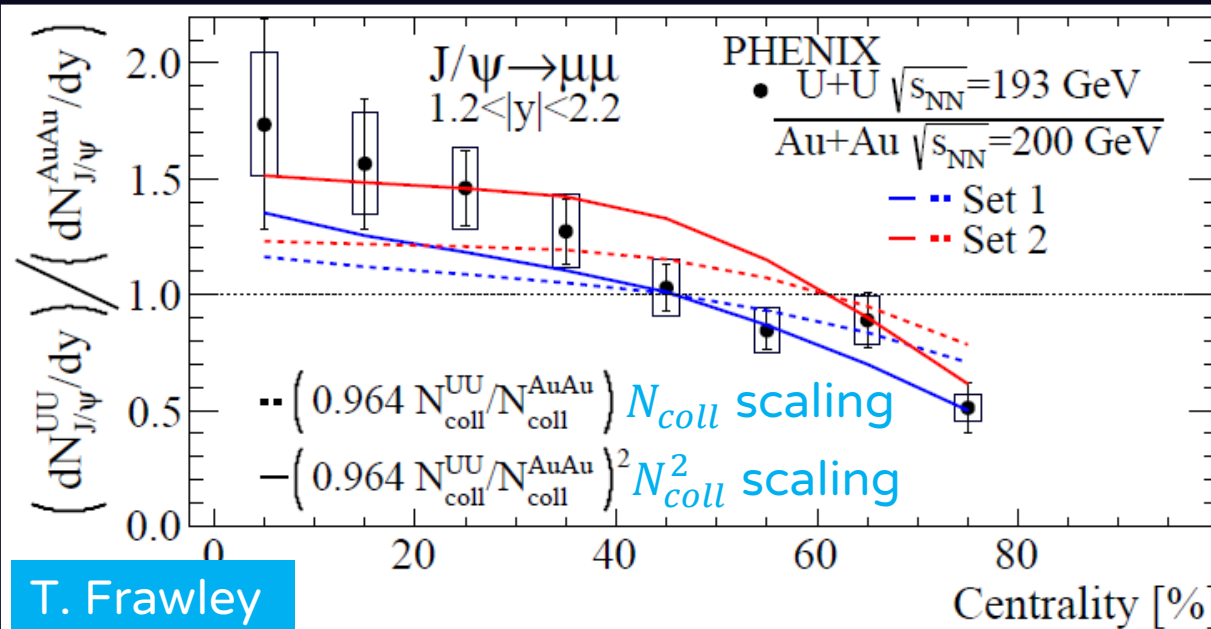
Quarkonium in AA collisions



J/ψ production at RHIC

20

➔ (re)combination/suppression role investigated comparing U-U and AuAu



in central U-U collisions:

1) stronger suppression
 due to color screening

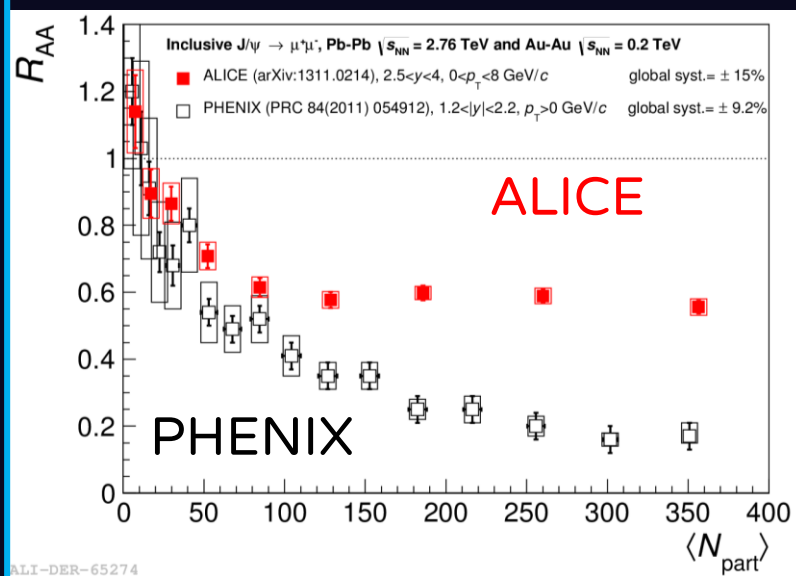
$$\varepsilon_{AuAu} \sim 80-85\% \varepsilon_{UU}$$

2) J/ψ recombination
 favoured by 25%
 larger N_{coll} in UU

$$N_{J/\psi}^{stat} \sim N_c^2 \sim N_{coll}^2$$

- ➔ results slightly favour N_{coll}^2 scaling \rightarrow dominant (re)combination over suppression when going from central U-U to Au-Au collisions
- ➔ quantitative comparison depends on the choice of the uranium Woods-Saxon parametrizations

evidence for recombination²21



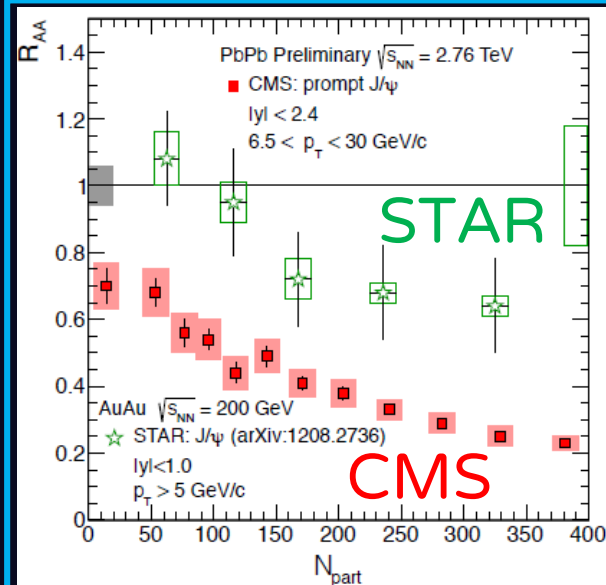
→ PHENIX vs. ALICE

low p_T J/ψ

in spite of LHC larger energy densities, stronger J/ψ suppression, vs centrality, in PHENIX, with respect to ALICE,

similar trend observed both at forward at mid rapidity

ALICE PLB 734 (2014) 314, arXiv:1505.08804



→ STAR vs. CMS

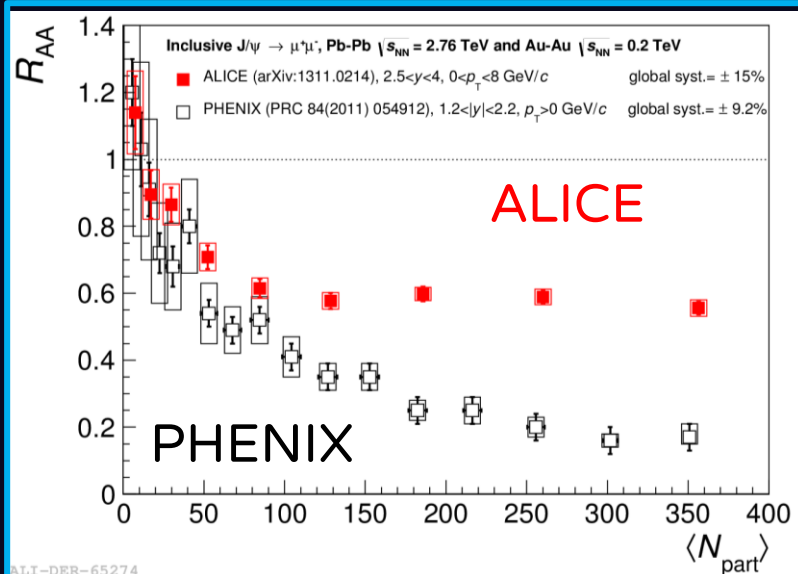
high p_T J/ψ

opposite J/ψ behavior compared to low- p_T results

→ suppression stronger at higher \sqrt{s} , as expected from QGP dissociation

CMS PAS-HIN-12-014

evidence for recombination²²



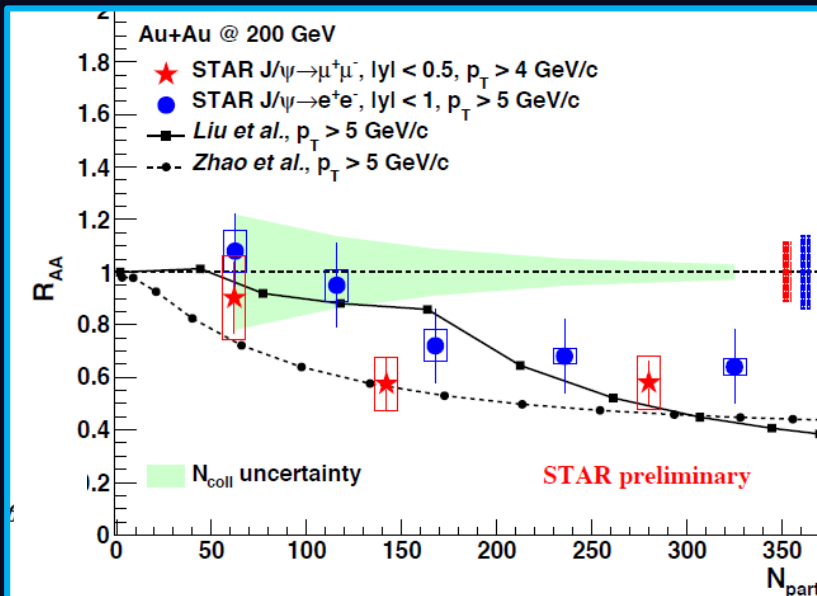
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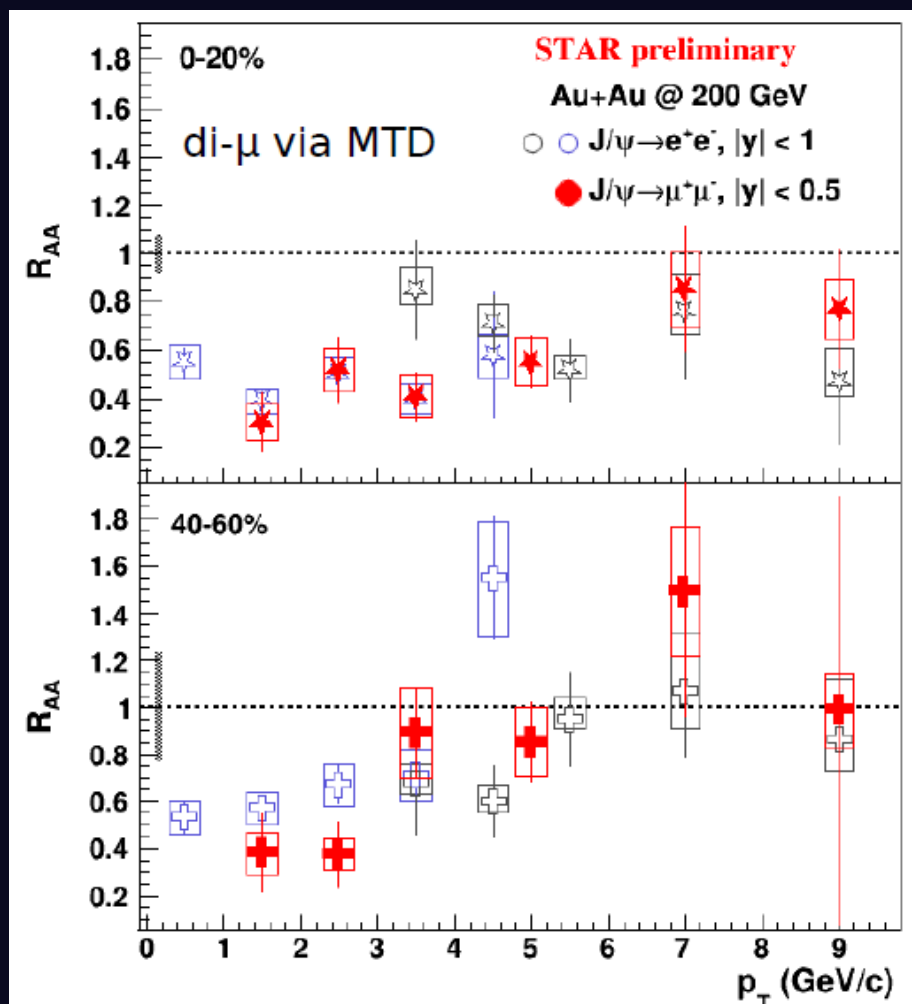
New STAR results based on MTD (J/ψ → μμ) confirm suppression in central collisions observed in the dielectron channel

R.Ma

J/ψ R_{AA} vs p_T at RHIC

23

R.Ma



→ J/ψ suppression stronger at low p_T

→ similar p_T dependence observed in the dimuon and dielectron decay channels

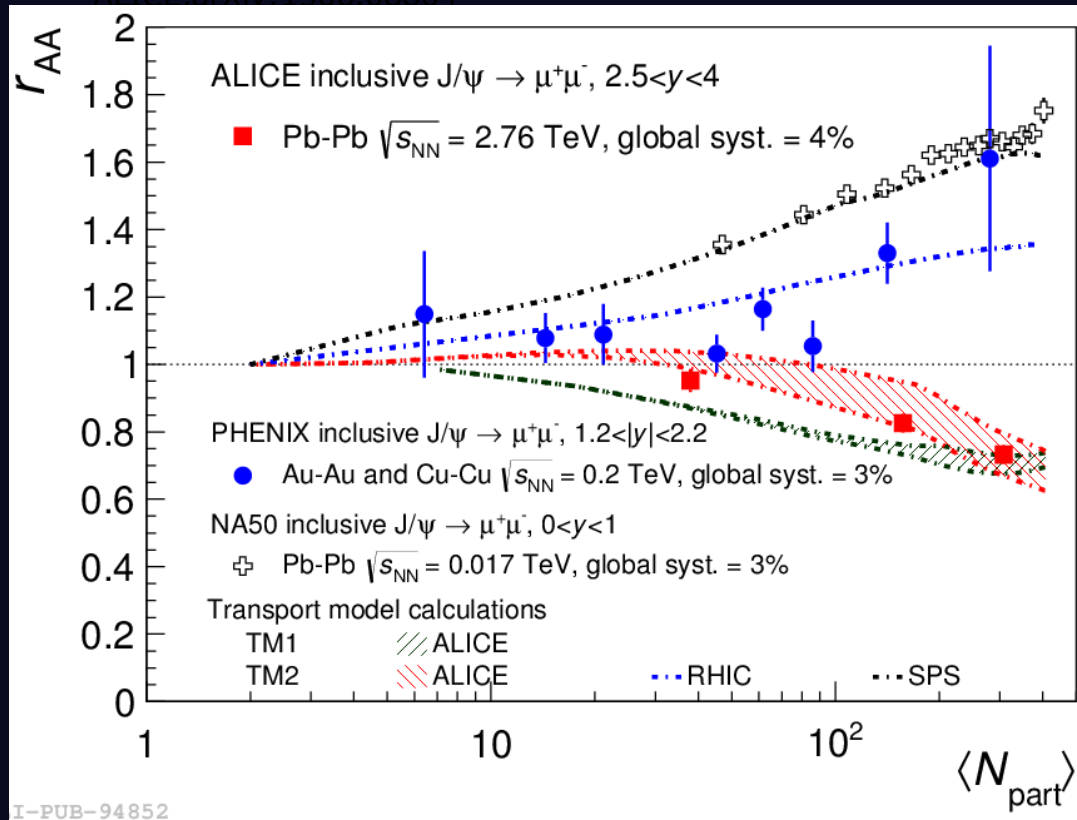
STAR PLB722(2013)55
STAR PR90(2014)024906
Liu et al. PLB678(2009)72
Zhao et al. PRC 82(2010) 064905

$\langle p_T \rangle$ increases from pp to AuAu

Evolution of $J/\psi \langle p_T^2 \rangle$

24

H. Pereira da Costa



$$r_{AA} = \frac{\langle p_T^2 \rangle_{AA}}{\langle p_T^2 \rangle_{pp}}$$

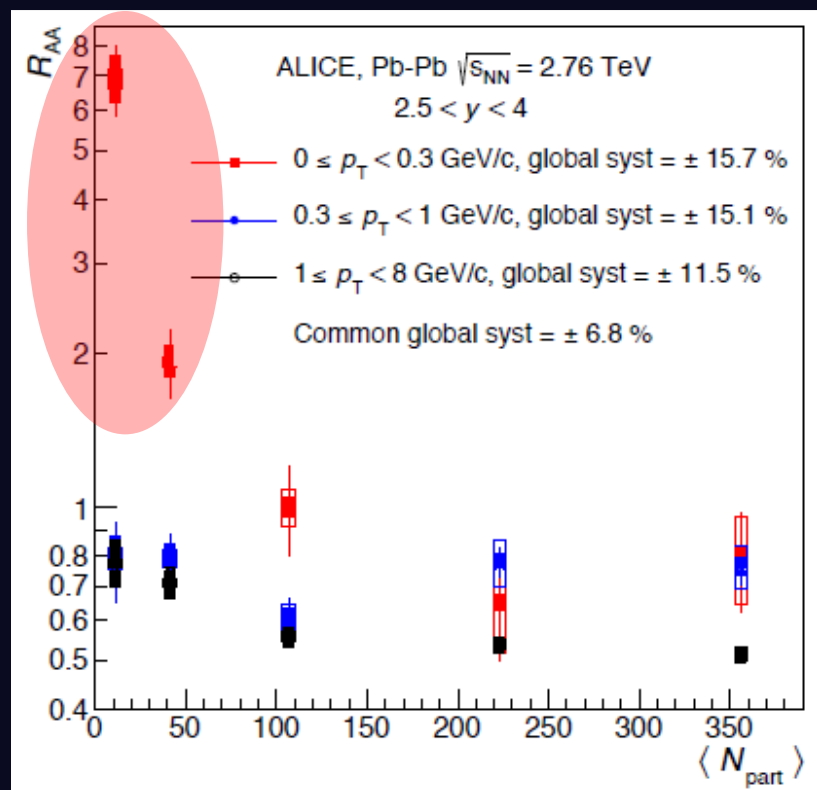
- r_{AA} centrality evolution strongly depends on \sqrt{s}
- decreasing r_{AA} trend, observed at LHC
→ due to (re)combination, which dominates J/ψ production at low p_T
- transport models, already describing $J/\psi R_{AA}$, also reproduce the r_{AA} evolution

J/ψ at very low p_T

25

Strong R_{AA} enhancement in peripheral collisions for $0 < p_T < 0.3$ GeV/c

G. Martinez

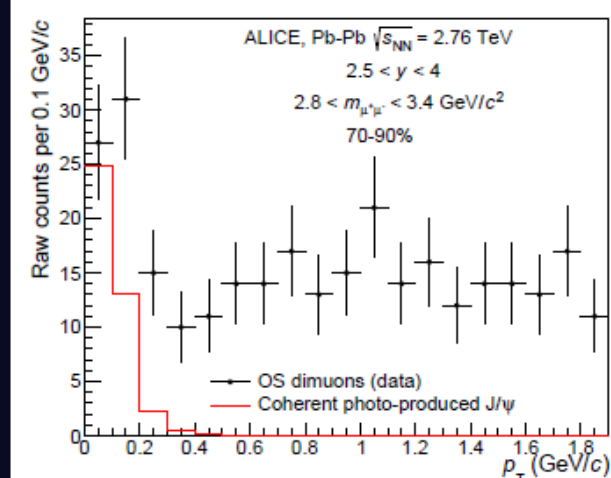


significance of the excess is 5.4 (3.4) σ in 70-90% (50-70%)

behaviour not predicted by transport models

excess might be due to coherent J/ψ photoproduction in PbPb (as measured also in UPC)

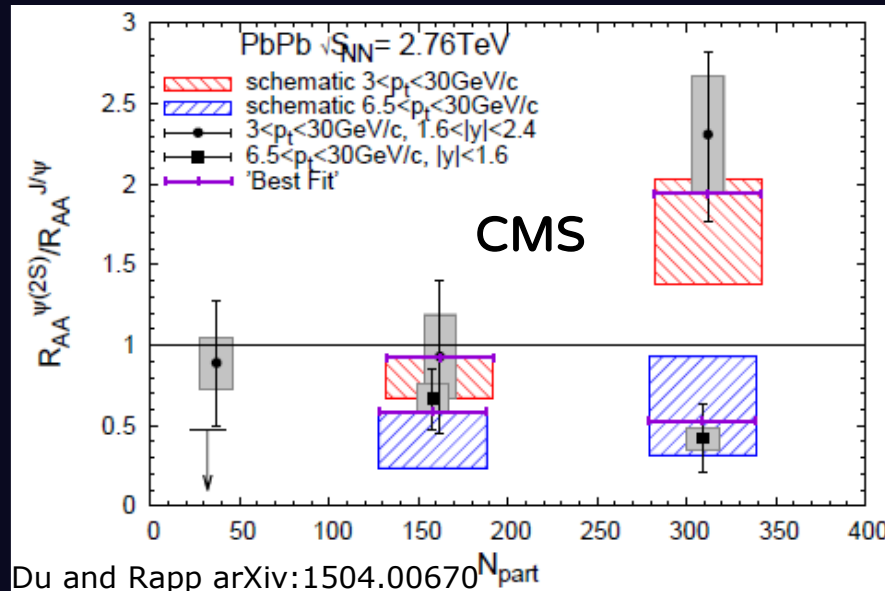
if excess is “removed” requiring $p_T^{J/\psi} > 0.3$ GeV/c \rightarrow ALICE R_{AA} lowers by 20% at maximum (in the most peripheral bin)



$\psi(2S)$ in AA collisions

26

➔ $\psi(2S)$ production modified in AA with a strong kinematic dependence



FW-y, $3 < p_T < 30 \text{ GeV/c} \rightarrow R_{AA}^{J/\psi} < R_{AA}^{\psi(2S)}$

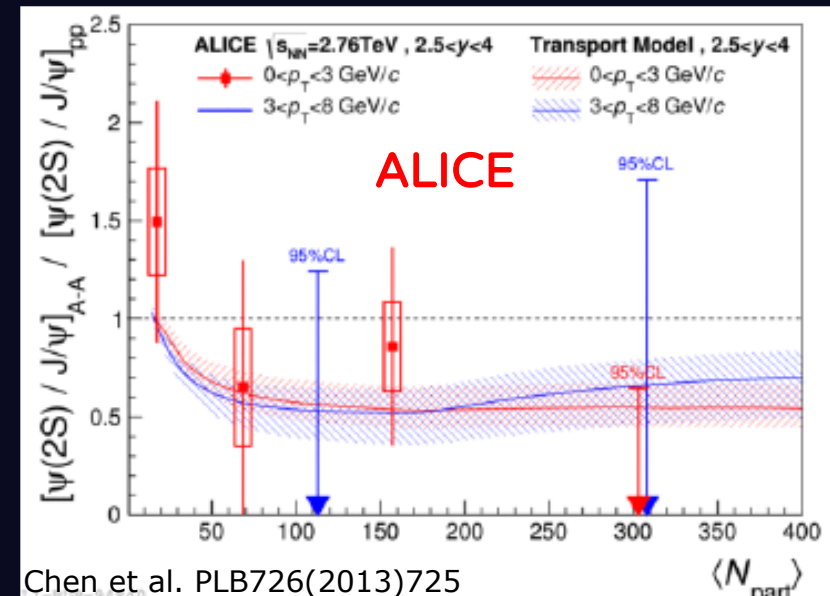
➔ later $\psi(2S)$ regeneration, when radial flow is stronger, might explain the rise

Mid-y $6.5 < p_T < 30 \text{ GeV/c} \rightarrow R_{AA}^{J/\psi} > R_{AA}^{\psi(2S)}$

➔ stronger suppression of weakly bound $\psi(2S)$ with respect to J/ψ

CMS, PRL 113(2014) 262301

Y. Kim



FW-y, $0 < p_T < 3 \text{ GeV/c} \rightarrow R_{AA}^{J/\psi} > R_{AA}^{\psi(2S)}$

➔ ALICE trend in agreement with transport models and with statistical hadronization approach

H. Pereira da Costa

ALICE arXiv:1506.08804

$\Upsilon(ns)$ production in AA

27

M. Jo

Reanalysis of 2011 CMS data:

- improved reconstruction
- high statistics pp reference (x20)

→ confirm the centrality dependent suppression for $\Upsilon(1S)$ and $\Upsilon(2S)$

→ Sequential suppression observed at LHC:

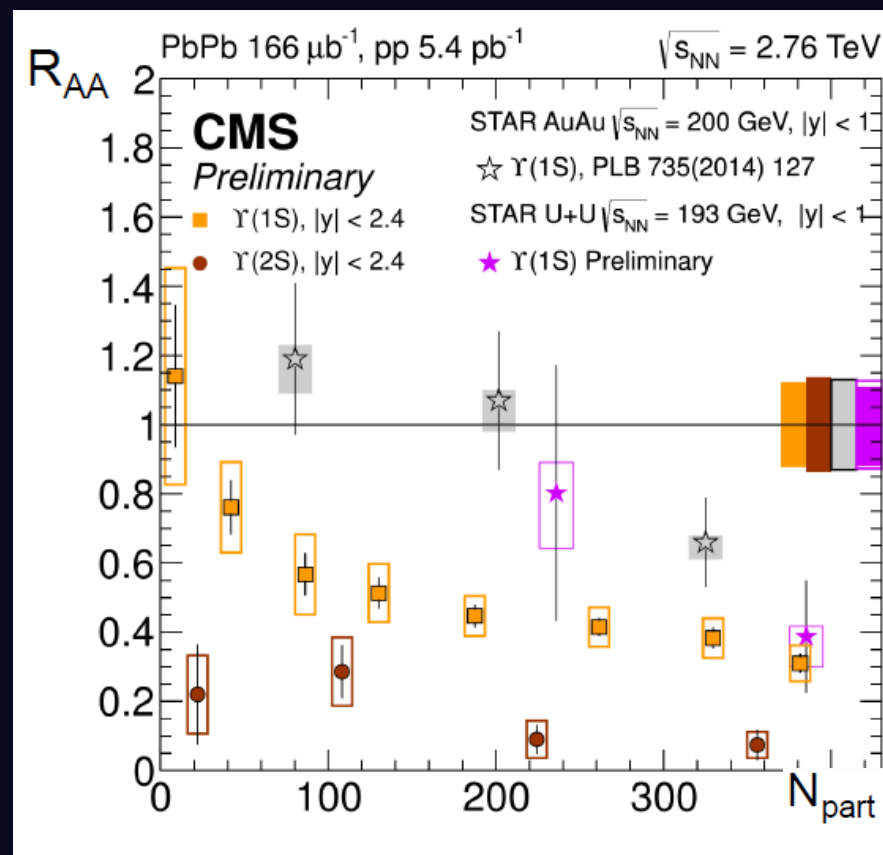
$$R_{AA}^{\Upsilon(3S)} < R_{AA}^{\Upsilon(2S)} < R_{AA}^{\Upsilon(1S)}$$

$$R_{AA}(\Upsilon(1S)) = 0.43 \pm 0.03 \pm 0.07$$

$$R_{AA}(\Upsilon(2S)) = 0.13 \pm 0.03 \pm 0.02$$

$$R_{AA}(\Upsilon(3S)) < 0.14 \text{ at } 95\% \text{ CL}$$

→ $\Upsilon(1S)$ suppressed also in central Au-Au and U-U collisions at RHIC



CMS, PRL109 (2012) 222301 and HIN-15-001
STAR, PLB735 (2014) 127 and preliminary U+U

Υ (ns) production in AA

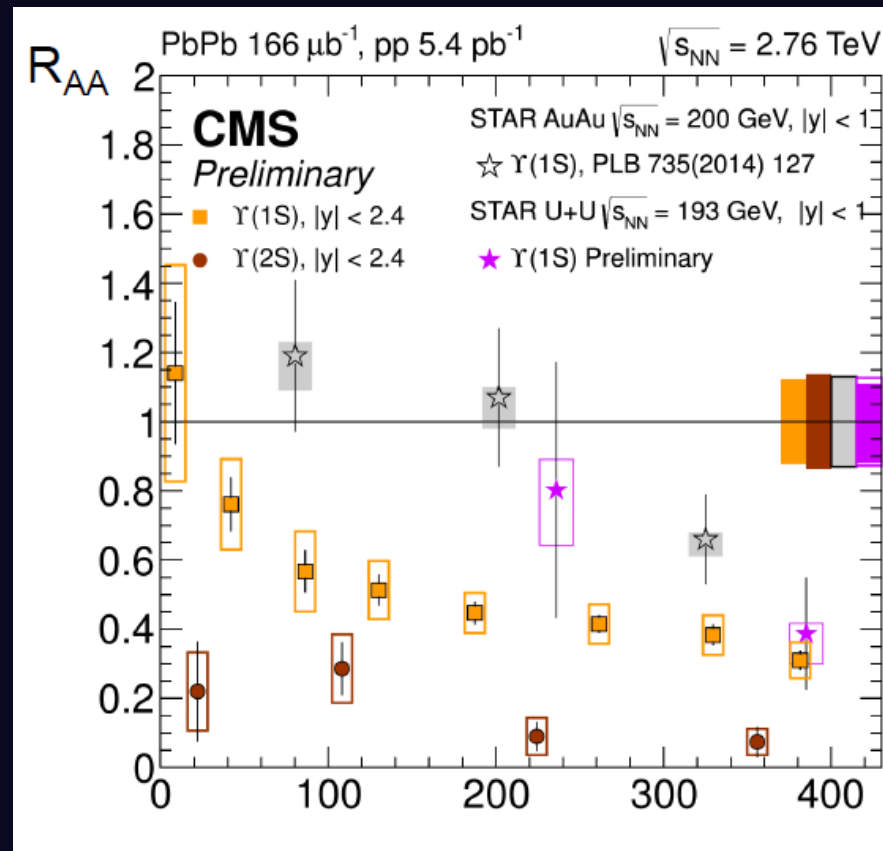
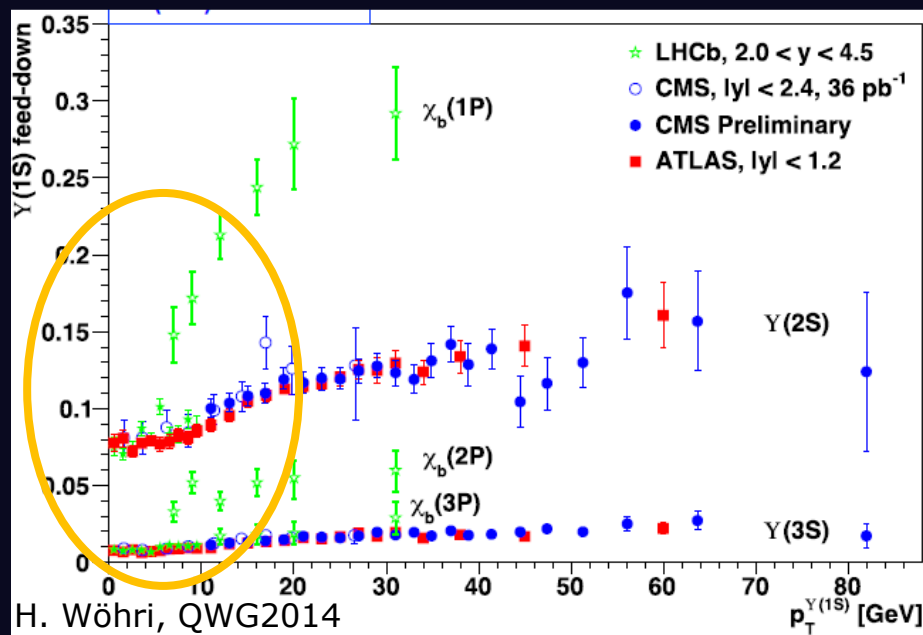
28

M. Jo

Reanalysis of 2011 CMS data:

- improved reconstruction
- high statistics pp reference (x20)

➔ Feed-down from excited states seems not enough to explain the observed $\Upsilon(1S)$ suppression

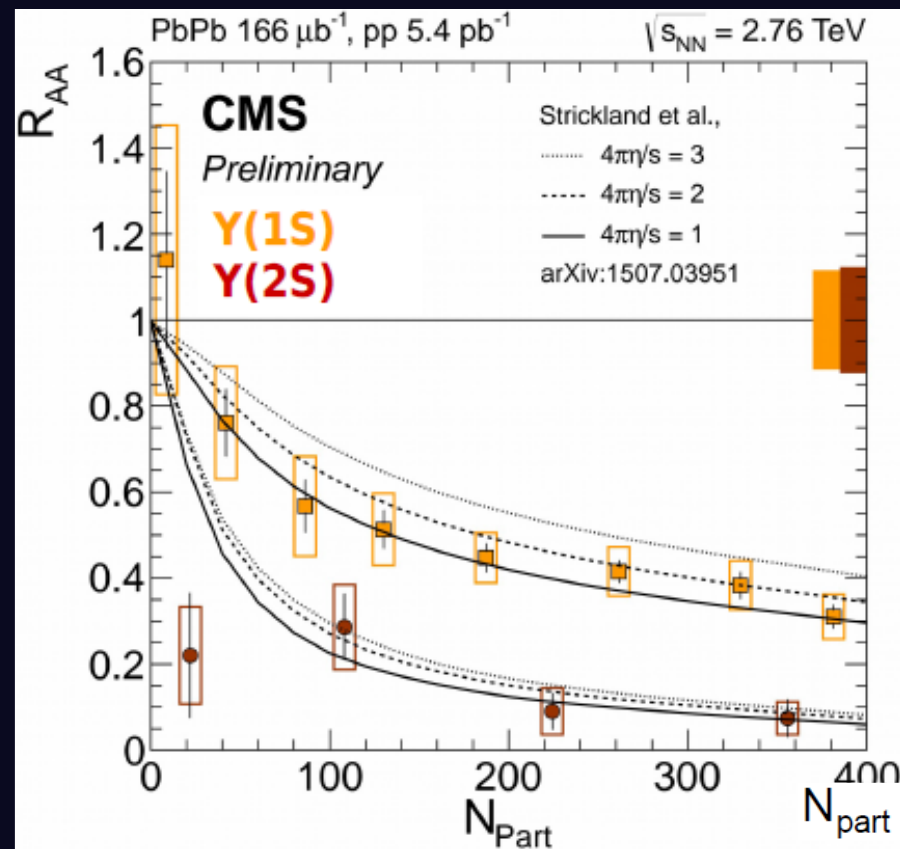
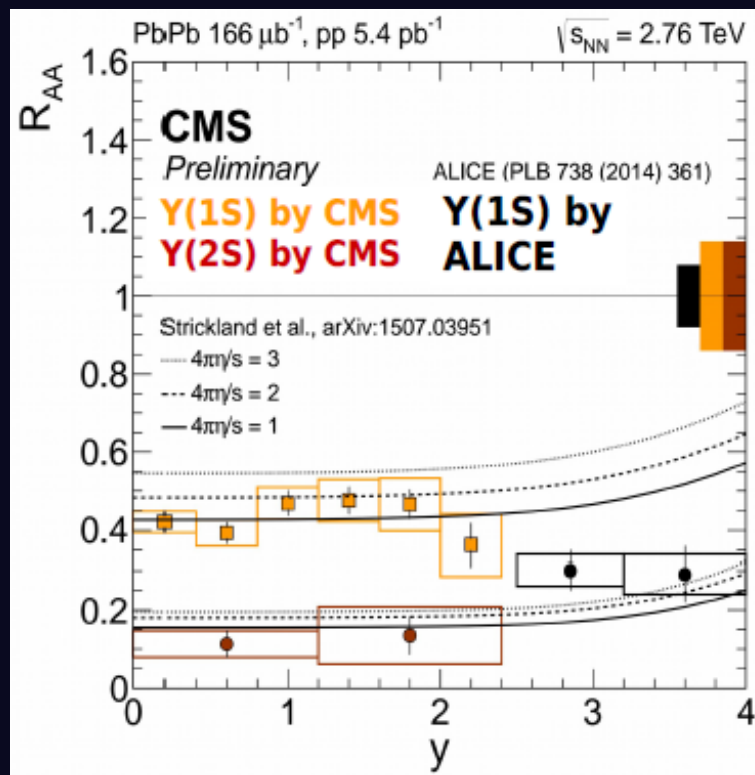


CMS, PRL109 (2012) 222301 and HIN-15-001
STAR, PLB735 (2014) 127 and preliminary U+U

Υ (ns) production in AA

29

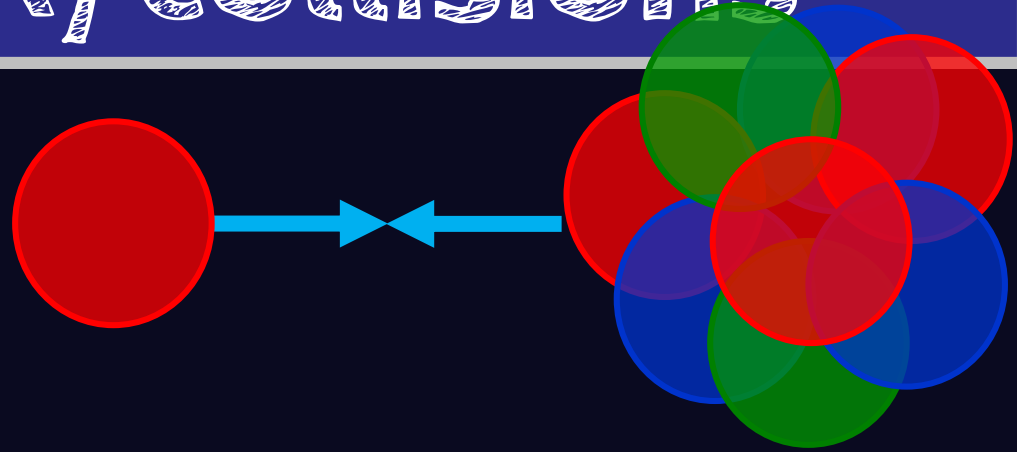
I. Das, M. Jo



- ➔ no p_T or y dependence of the $\Upsilon(1S)$ and $\Upsilon(2S)$ suppressions
- ➔ new CMS R_{AA} closer to ALICE forward- y results

- ➔ models reproduce the p_T and centrality dependence. Rapidity description still needs tuning

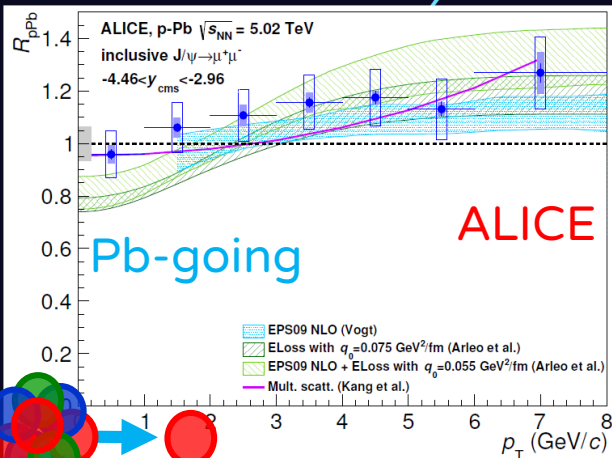
Quarkonium in p-A (d-A) collisions



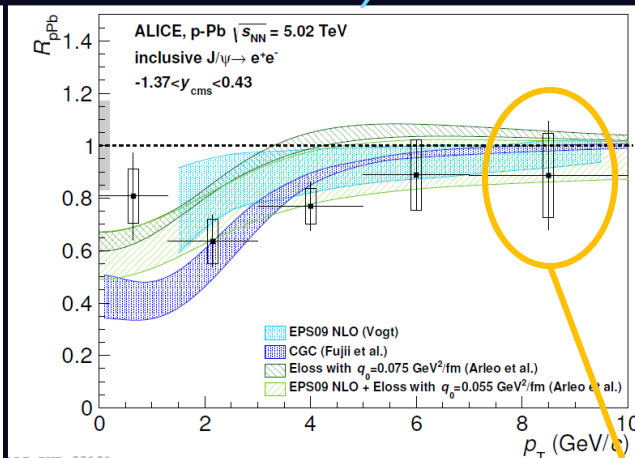
J/ψ R_{pA} vs p_T

31

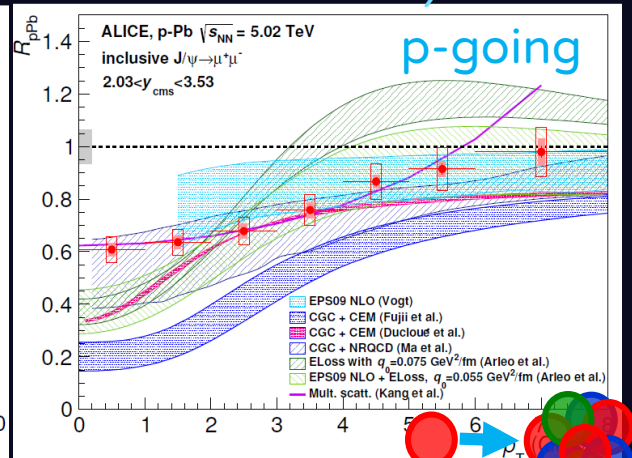
backward- γ



mid- γ



forward- γ



J/ψ production strongly depends on γ and p_T

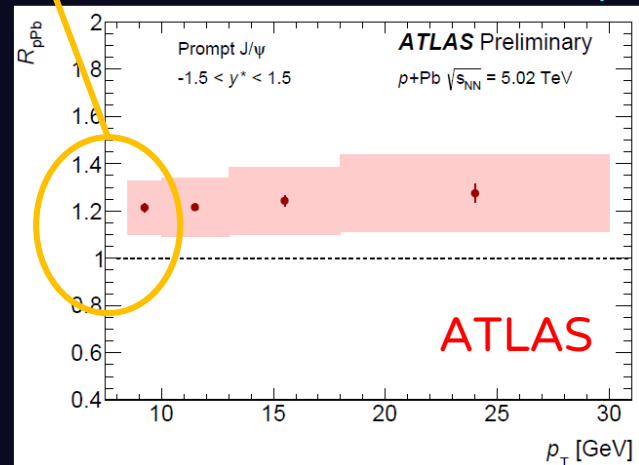
- backward- γ : $R_{pA} \sim 1$, negligible p_T dependence
- forward- γ : R_{pA} increases with p_T
- mid- γ : small p_T dependence, $R_{pA} \sim 1$ for $p_T > 3$ GeV/c (ALICE) and slightly larger for $p_T > 8$ GeV/c (ATLAS)

➔ shadowing and coherent parton energy loss models reasonably describe the data

➔ agreement with CGC depends on implementation

$\sim 0.7\sigma$ difference

mid- γ

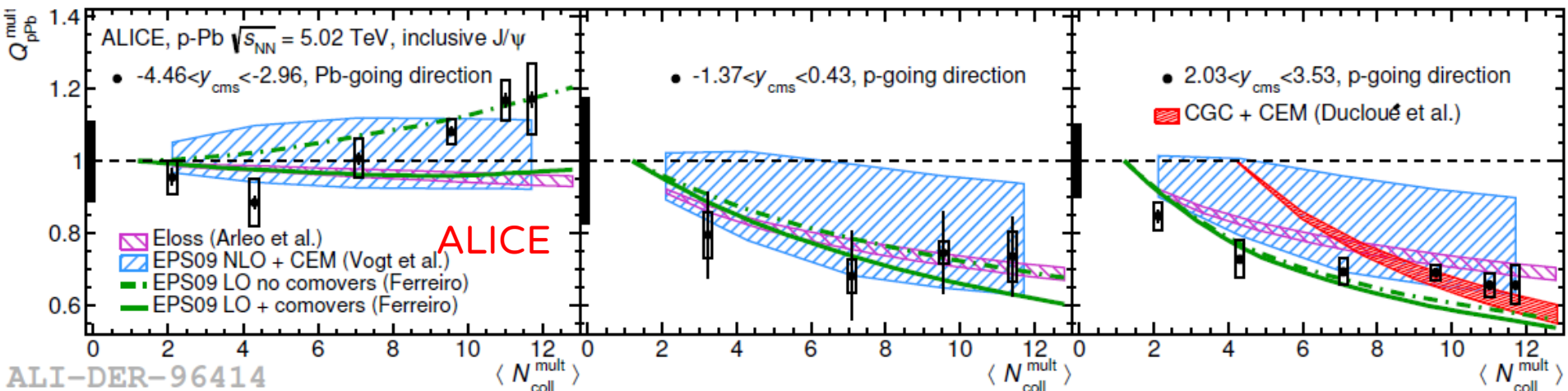


J/ψ centrality dependence 32

backward- y

mid- y

forward- y



➔ J/ψ production is investigated vs. centrality

ALICE:

mid and fw- y : suppression increases with centrality

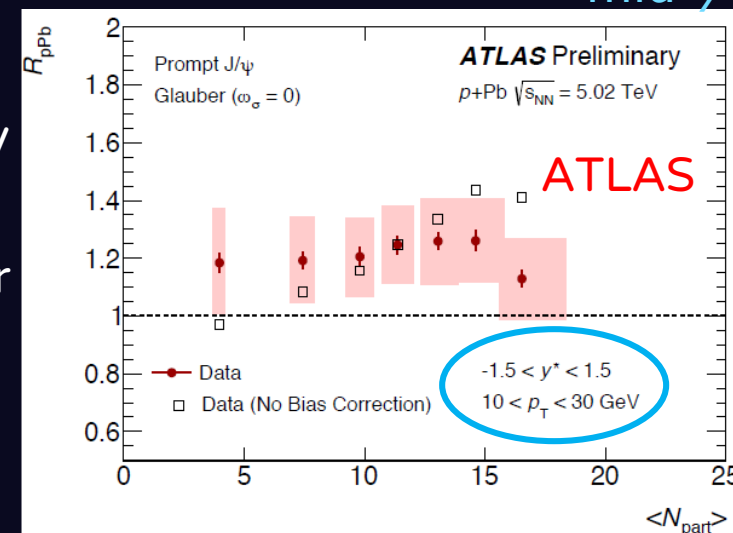
backward- y : hint for increasing Q_{pA} with centrality

Shadowing and coherent energy loss models in fair agreement with data

ATLAS

Flat centrality dependence in the high p_T range

mid- y



$\psi(2S)$ production in pA

33

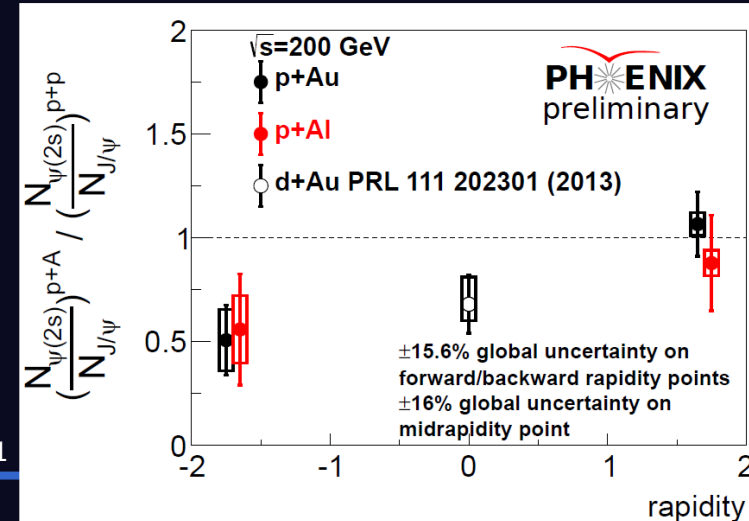
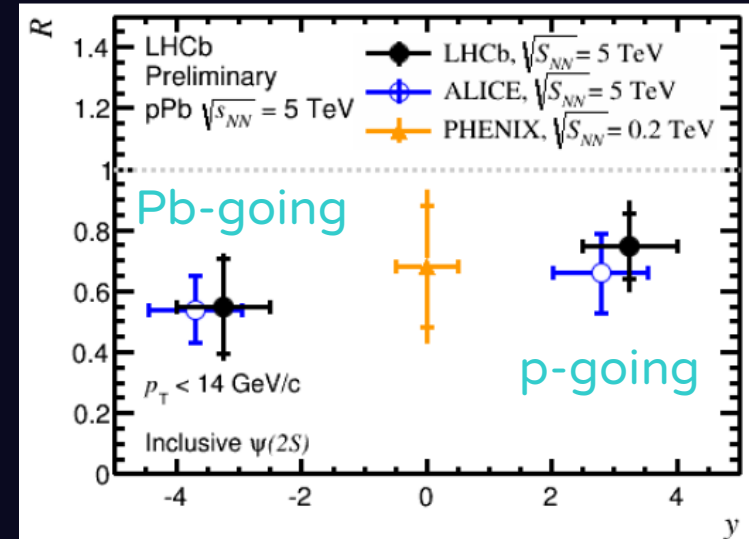
Being more weakly bound than the J/ψ , the $\psi(2S)$ is an interesting probe to have further insight on the charmonium behaviour in pA

→ $\psi(2S)$ suppression stronger than the J/ψ one at RHIC and LHC

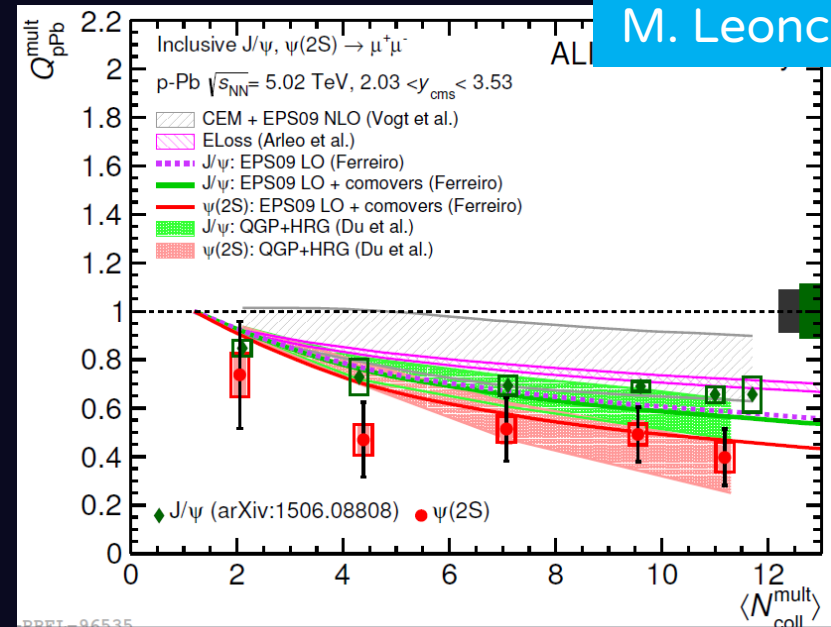
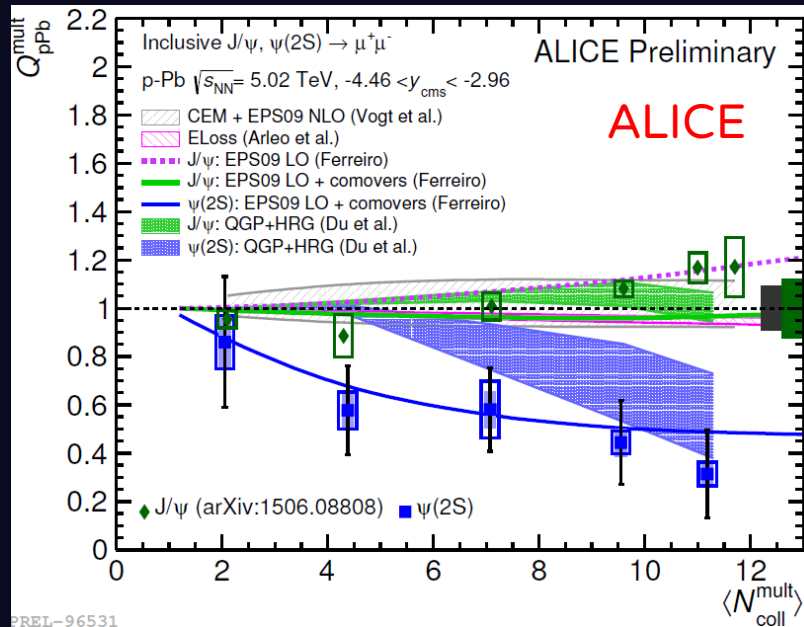
→ unexpected because time spent by the cc pair in the nucleus (τ_c) is shorter than charmonium formation time (τ_f)

→ shadowing and energy loss, almost identical for J/ψ and $\psi(2S)$, do not account for the different suppression

M. Brooks T. Frawley, M. Leoncino, Z. Yang



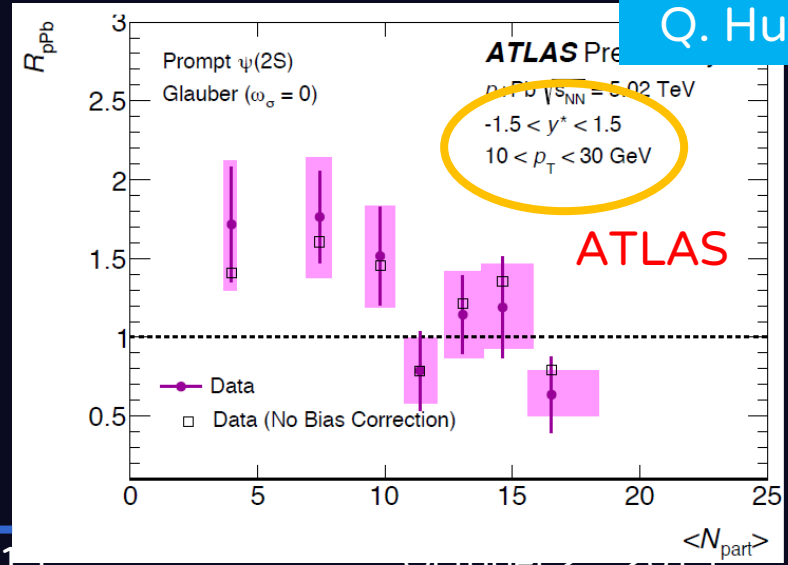
$\psi(2S)$ production vs centrality 34



M. Leoncino

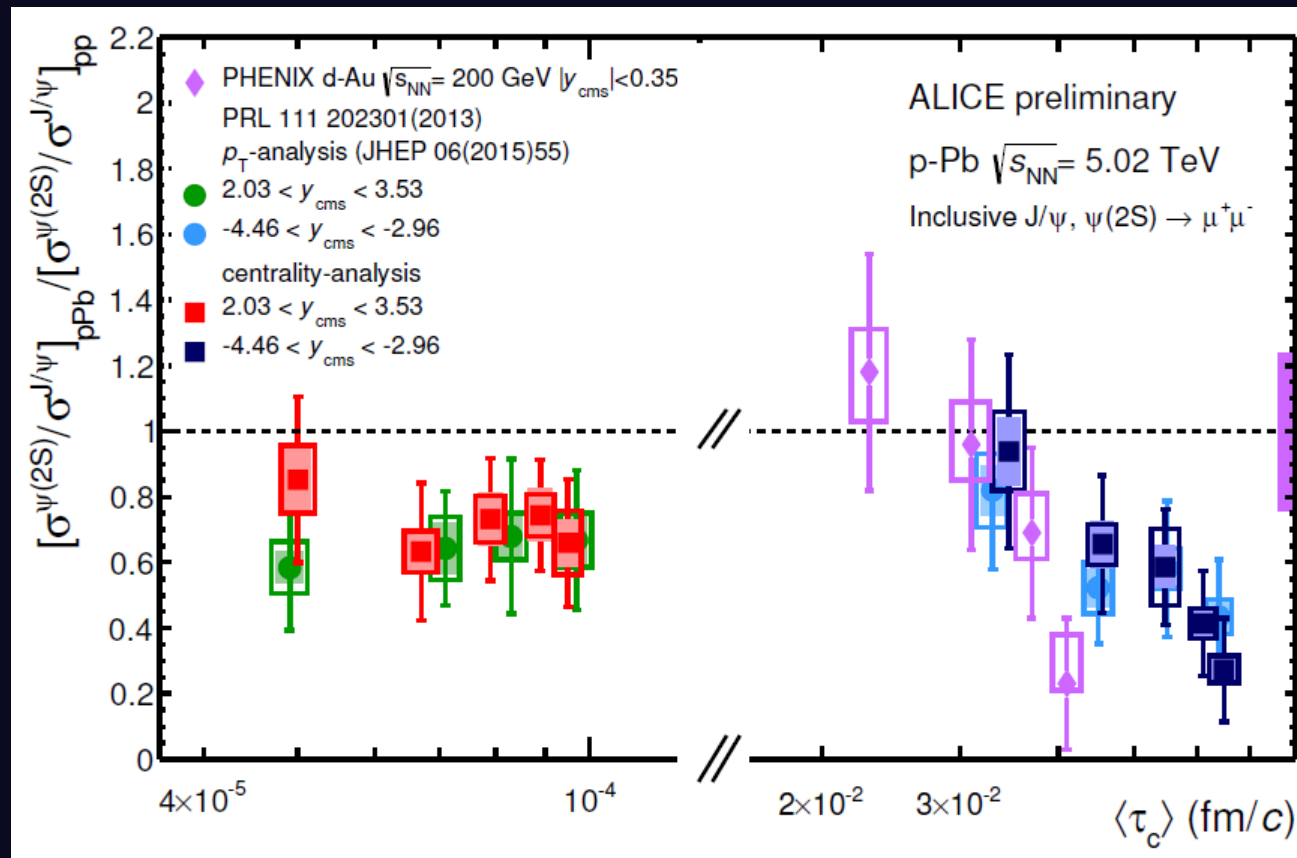
→ QGP+hadron resonance gas (Rapp) or comovers (Ferreiro) models describe the stronger $\psi(2S)$ suppression

→ Consistency between ALICE and ATLAS results?
It would be useful to have theory prediction for the higher p_T range



Q. Hu

$\psi(2S)$ versus crossing time 35

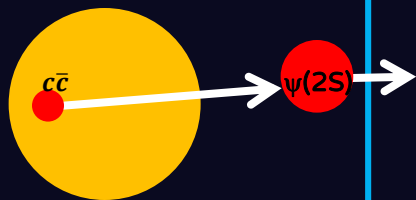


$$\tau_c = \frac{\langle L \rangle}{(\beta_z \gamma)}$$

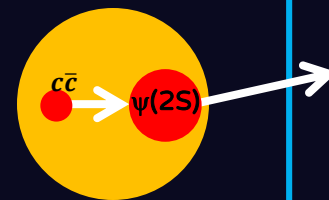
D. McGlinchey, A. Frawley and R. Vogt, PRC 87,054910 (2013)

M. Leoncino

Forward- y : $\tau_c \ll \tau_f$
interaction with
nuclear matter
cannot play a role



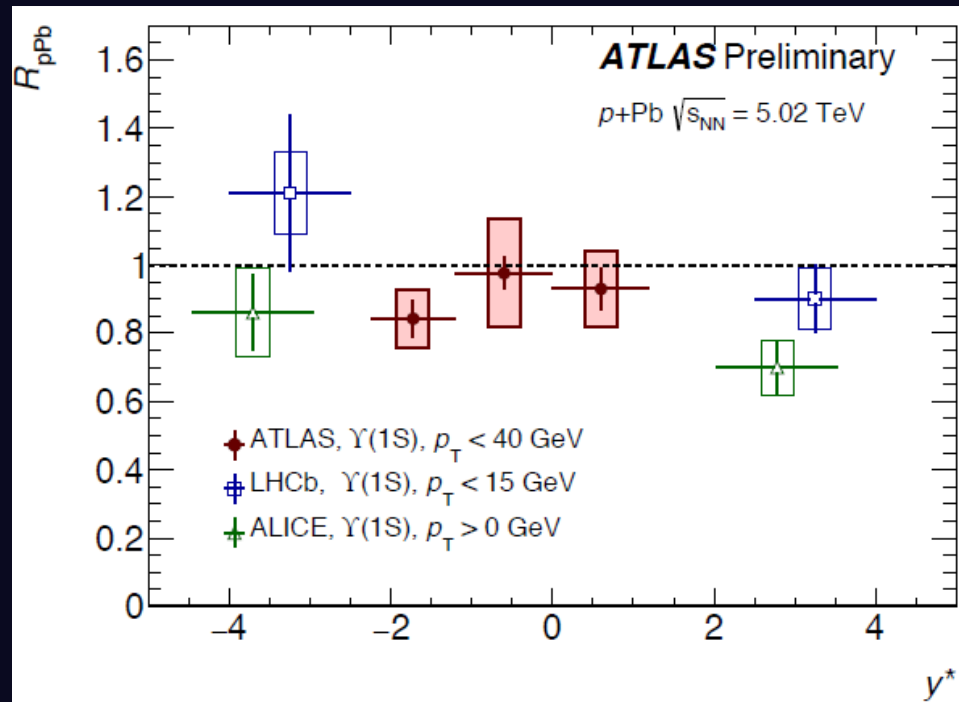
Backward- y : $\tau_c \lesssim \tau_f$
indication of effects
related to break-up
in the nucleus?



$\Upsilon(1S)$ in pA collisions

36

Q. Hu

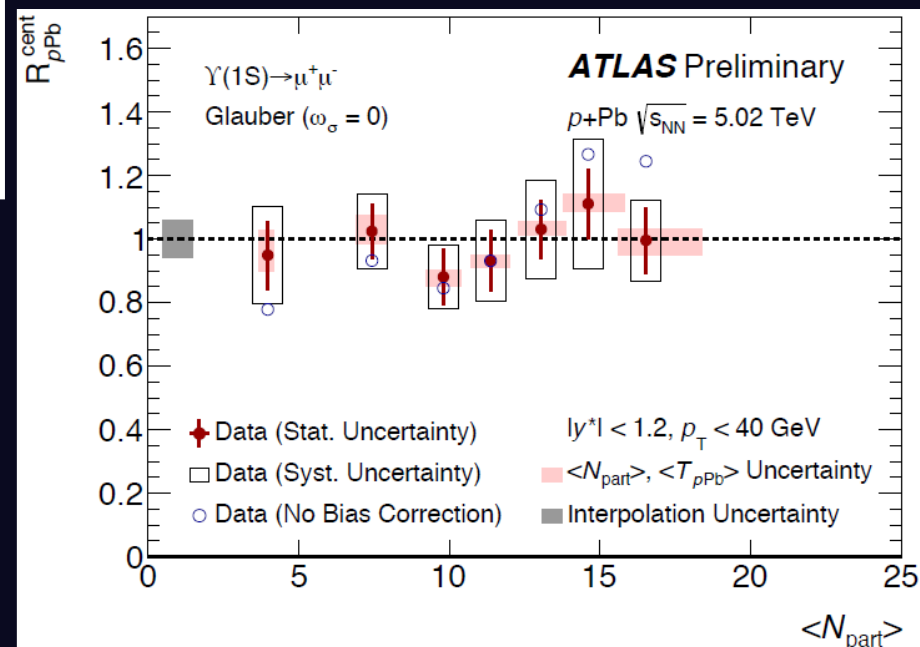


➔ No significant rapidity dependence of $\Upsilon(1S)$ R_{pA}

➔ Weak CNM effects on $\Upsilon(1S)$

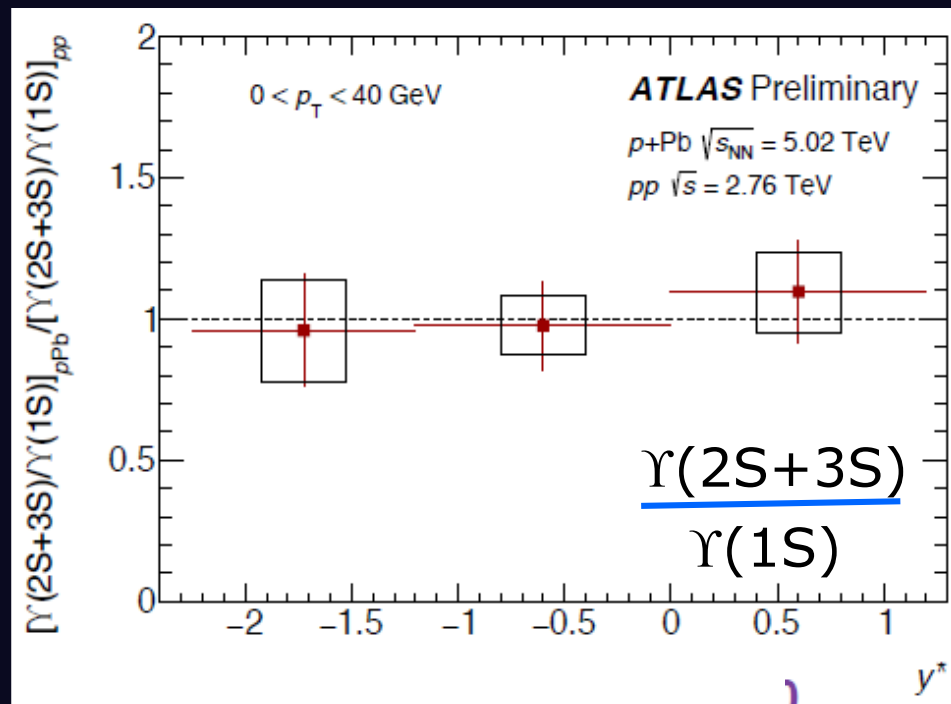
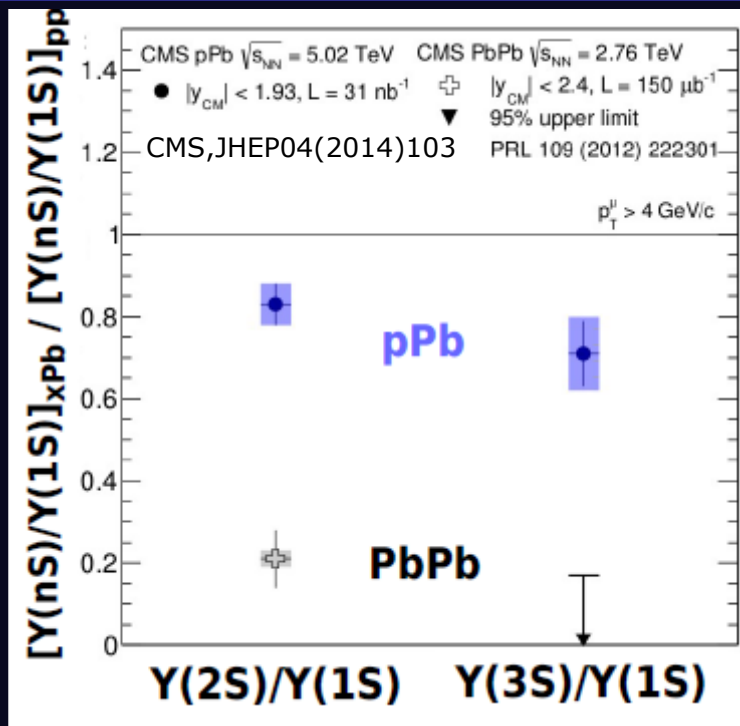
➔ $\Upsilon(1S)$ R_{pA} consistent with unity and rather flat versus centrality

ALICE, Phys. Lett. B 740 (2015) 105
 ATLAS-CONF-2015-050
 LHCb, JHEP 07(2014)094



$\Upsilon(ns)$ in pA collisions

37



→ CMS:

- excited states suppressed with respect to $\Upsilon(1S)$
- Initial state effects similar for the $\Upsilon(ns)$ states
→ Final states effects at play?

→ ATLAS:

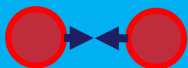
- no strong y (and p_T) dependence
- agreement with CMS within uncertainties

M. Jo, Q. Hu

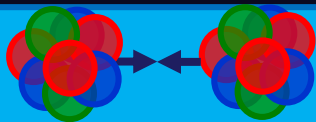
Conclusions

38

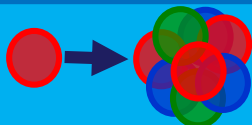
New results on quarkonium in pp, pA and AA shown in QM:

pp 

New differential J/ψ cross-section results at RHIC and LHC top energies will help constraining production models

AA 

New results confirm the role of suppression and recombination mechanisms at play on the various quarkonium states

pA 

Interplay of shadowing and energy loss describes J/ψ production. Comover-like effects seem to affect excited quarkonium states

Thanks!

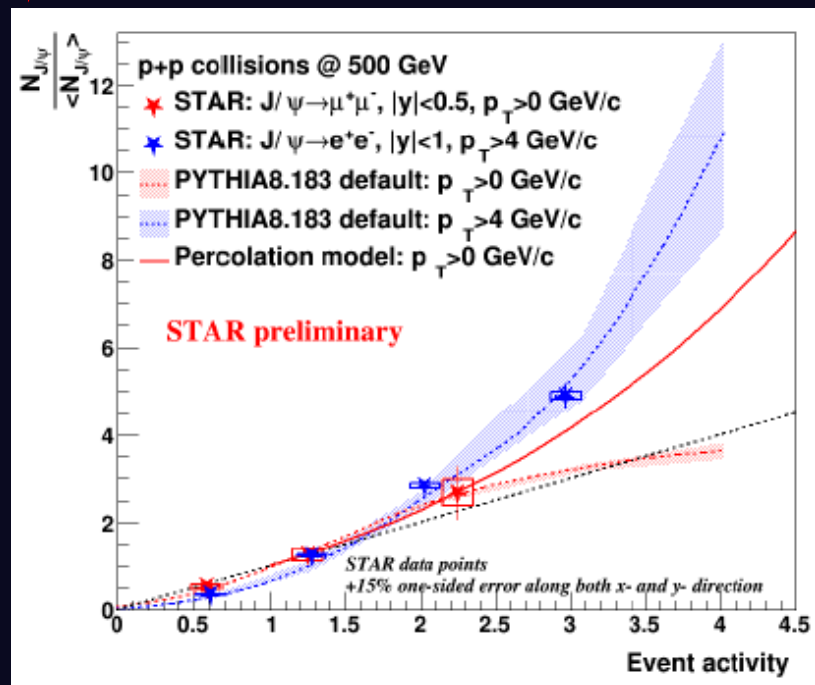


Backup slides pp

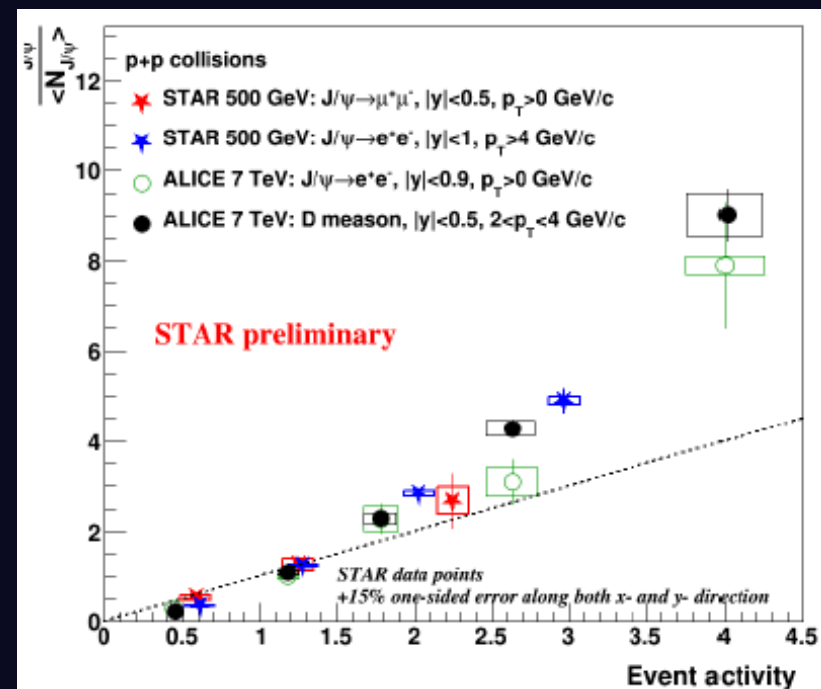
J/ψ production in pp

40

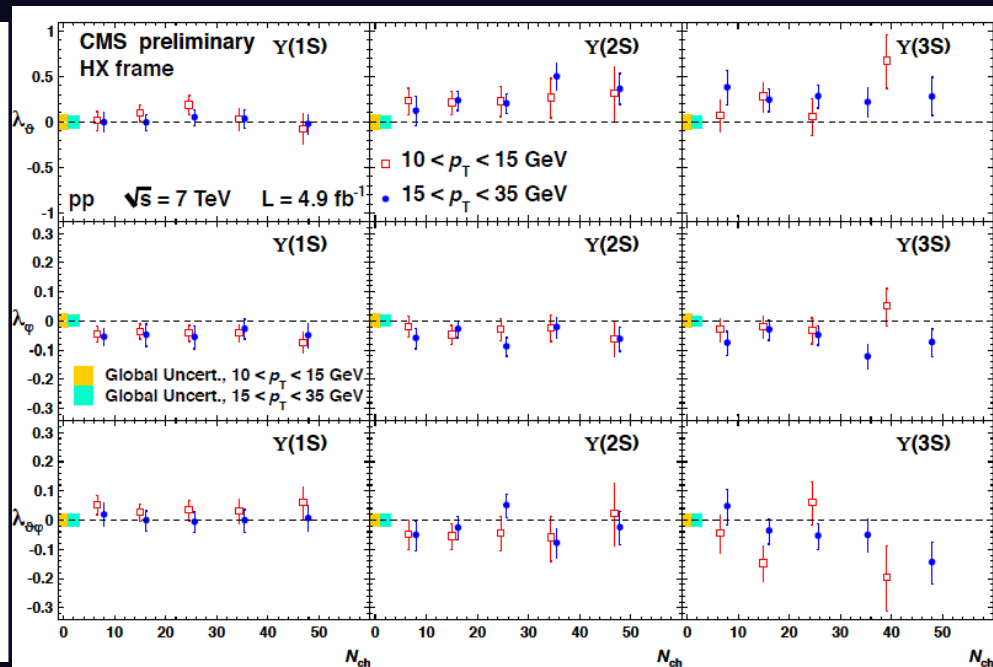
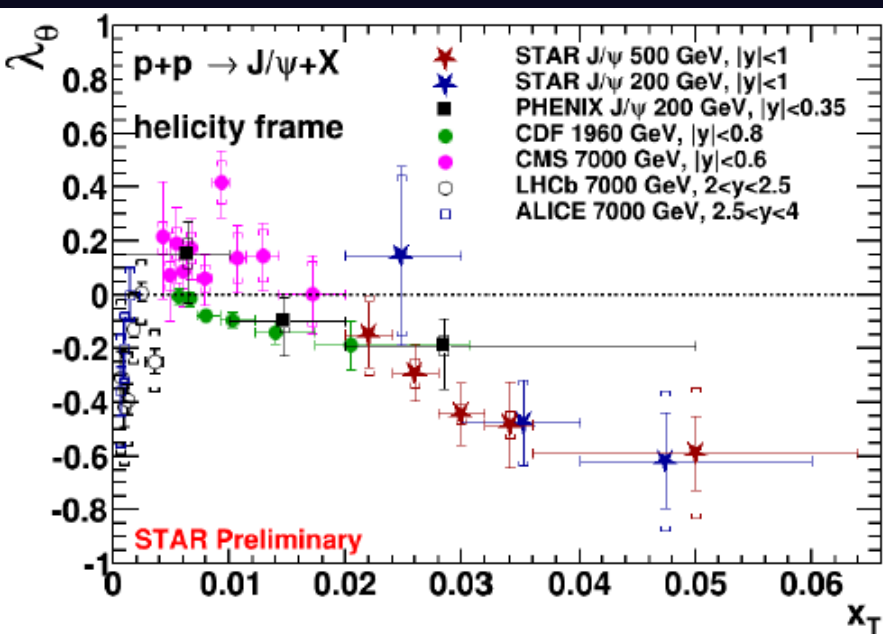
Correlation between J/ψ yields and multiplicity



- Stronger than linear increase for $p_T > 4$ GeV
- Possible interpretations: multi parton interactions (MPI) or percolation
- Similar trend observed in pp at $\sqrt{s} = 7$ TeV



J/ψ and Υ (ns) polarization 41



CMS PAS HIN-15-003

➡ Same trend in $\sqrt{s}=200$ and new STAR 500 GeV results

➡ Common x_T scaling towards negative λ_θ

$$x_T = 2p_T/\sqrt{s}$$

➡ Υ polarization in pp at $\sqrt{s}=7$ TeV: no trend observed as a function of the charged particle multiplicity

No changes in the production process between low and high multiplicity

Backup slides AA

Uranium parametrization

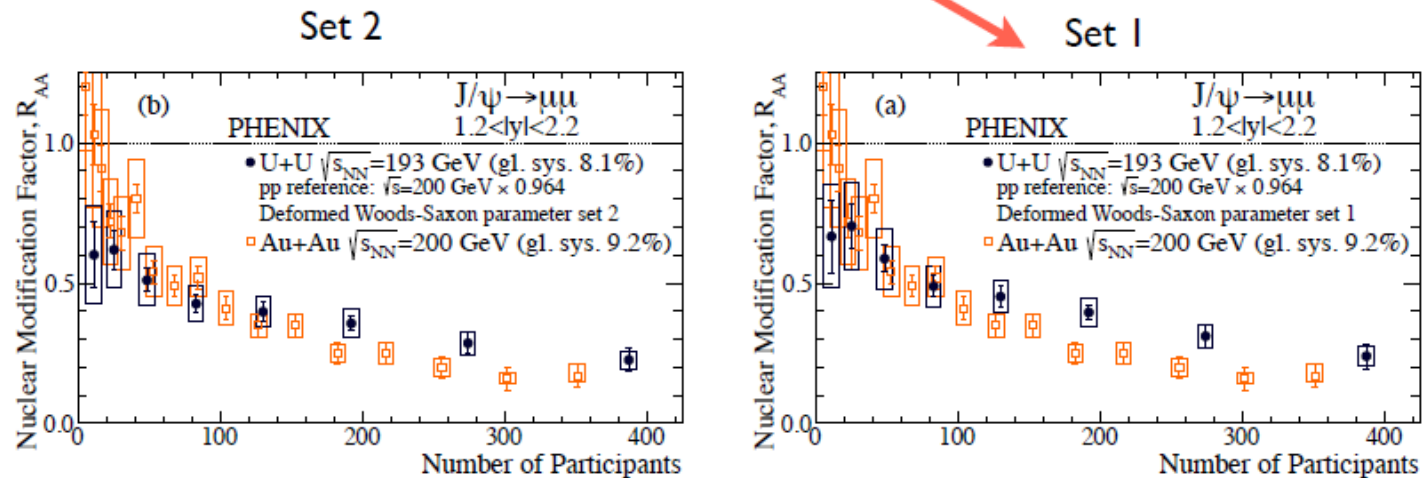
Effect of U deformation model

8

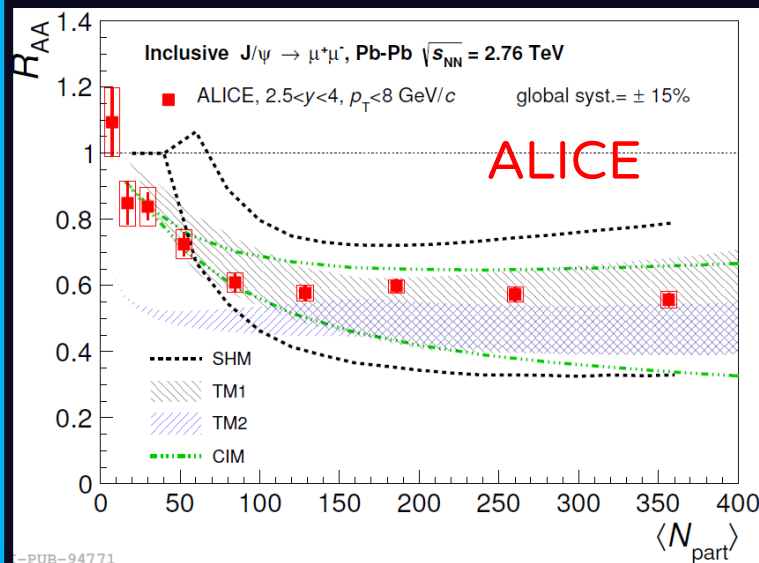
The parameters for set 1 are significantly different in their **surface diffuseness**:

Parameter	set 1	set 2
R (fm)	6.81	6.86
a (fm)	0.6	0.42
β_2	0.28	0.265
β_4	0.093	0

Larger surface diffuseness for set 1 results in a less compact nucleus, a larger reaction cross section by 12%, and **N_{coll} values that are smaller by 6 - 15%**



evidence for recombination² 44

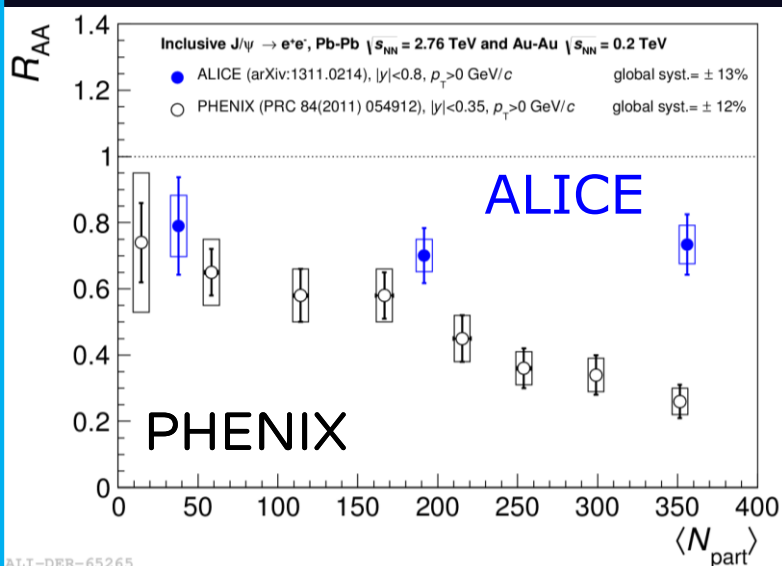


➔ **PHENIX vs. ALICE**

low p_T J/ψ

agreement with models including J/ψ (re)combination in QGP or in the hadronic phase

ALICE PLB 734 (2014) 314, arXiv:1505.08804



➔ **PHENIX vs. ALICE**

low p_T J/ψ

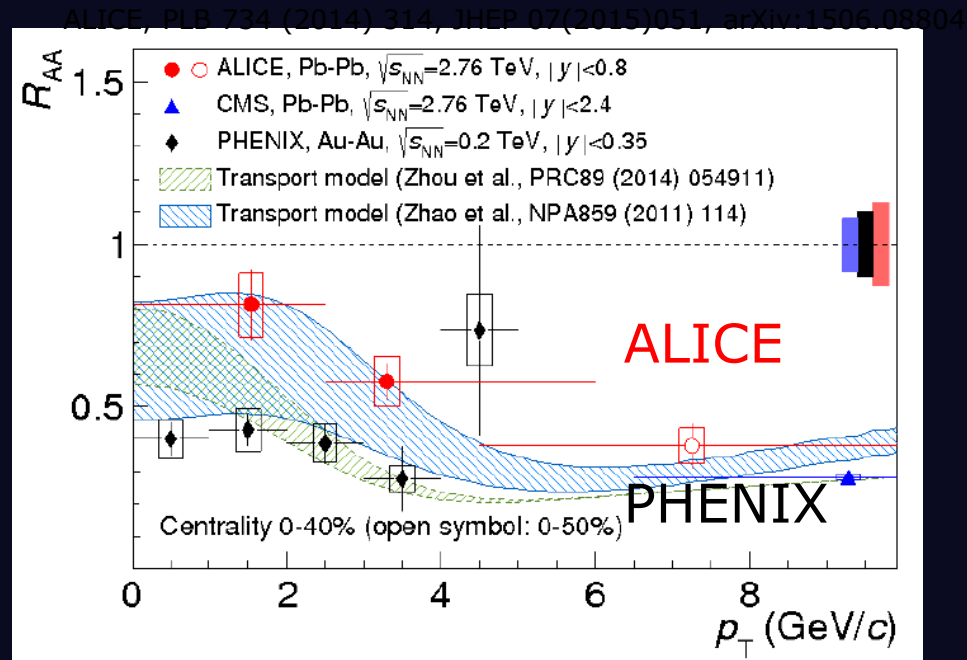
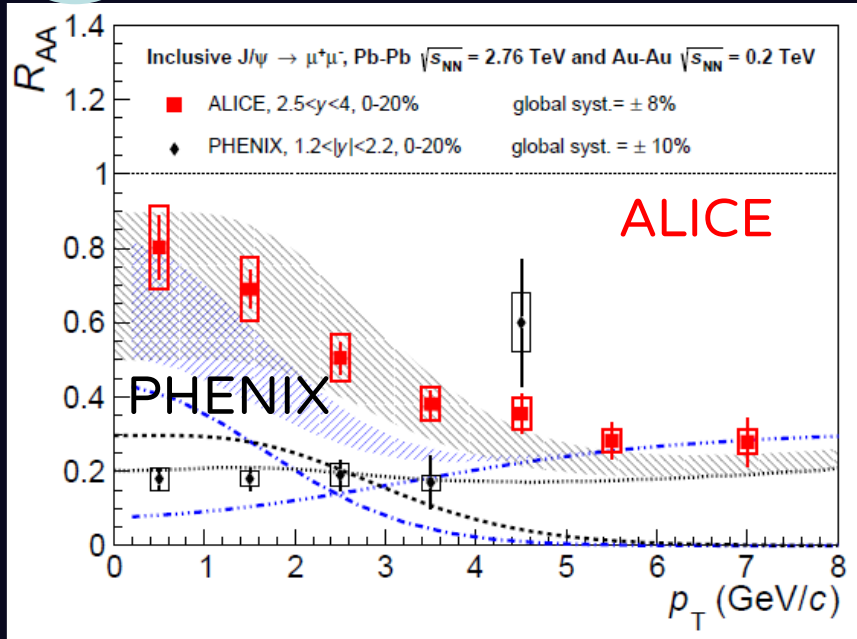
similar trend observed both at forward at mid rapidity

ALICE PLB 734 (2014) 314, arXiv:1505.08804

evidence for recombination²₄₅

2

Comparison LHC vs RHIC (R_{AA} p_T dependence)



- weaker J/ψ suppression at low p_T in agreement with models including J/ψ (re)combination in QGP or in the hadronic phase

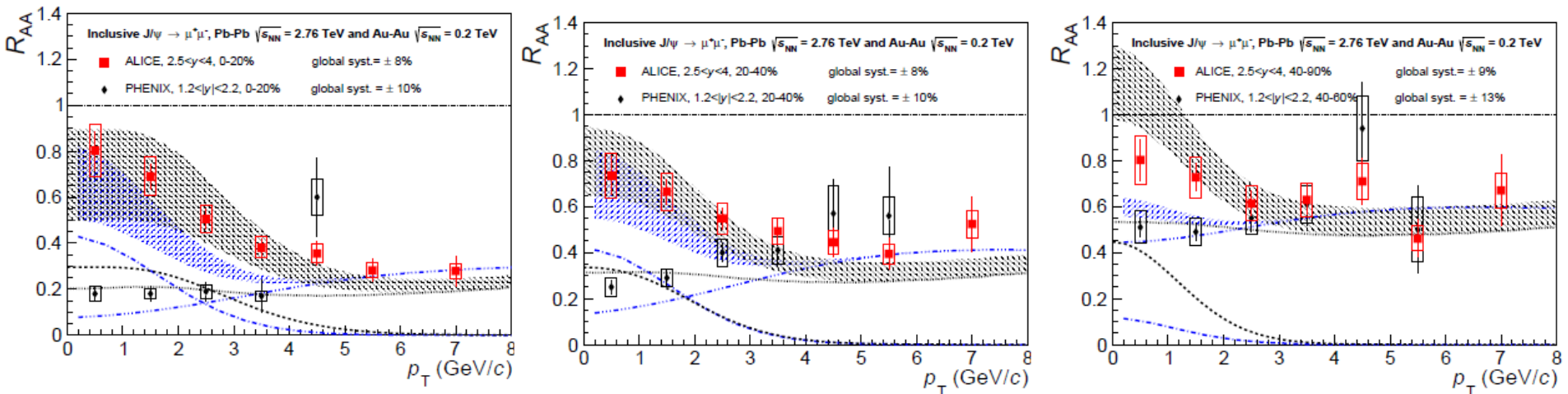
Multi-differential J/ψ studies 46

➔ p_T -centrality multi-differential studies allows detailed comparison with theory models

0-20%

20-40%

40-90%



TM1 Zhao et al., Nucl.Phys.A859 (2011) 114

TM2 Zhou et al. Phys.Rev.C89 (2014)054911

ALICE, arXiv:1506.08804

.... Primordial J/ψ (TM1)

--- Regenerated J/ψ (TM1)

... Primordial J/ψ (TM2)

... Regeneration J/ψ (TM2)

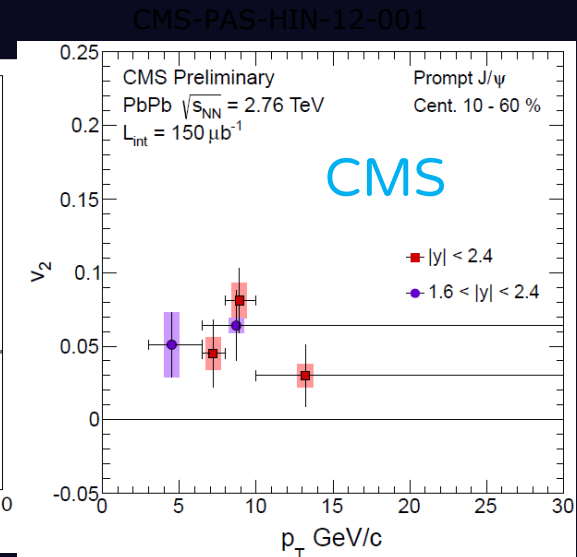
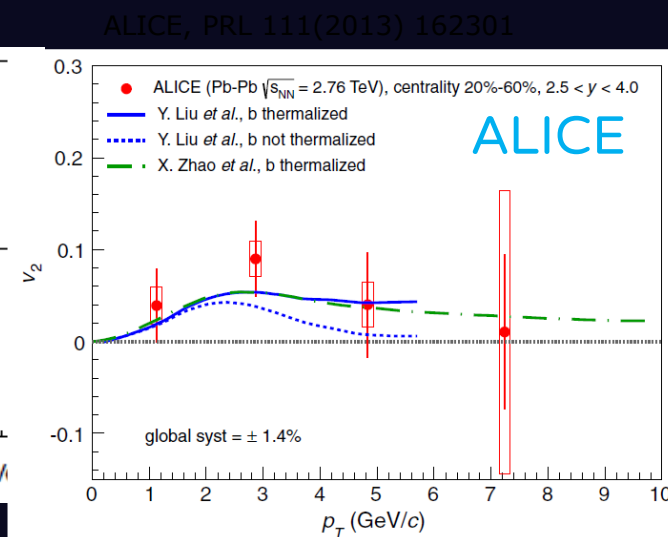
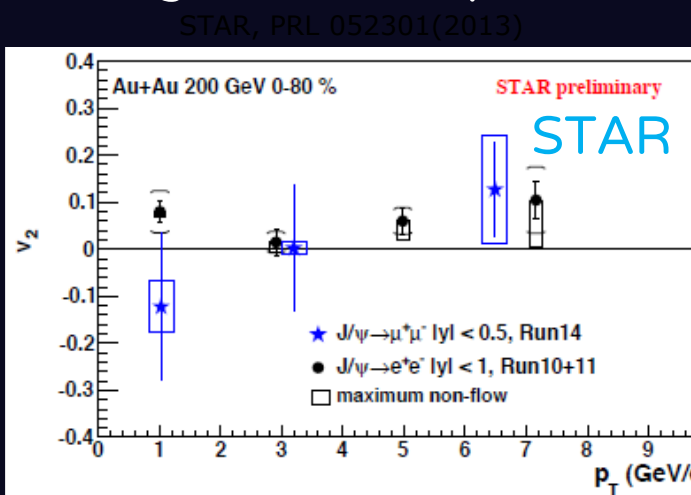
➔ Model provide a fair description of the data, even if with different balance of primordial/regeneration components

Still rather large theory uncertainties: models will benefit from precise measurement of σ_{cc} and CNM effects

4

J/ψ elliptic flow

➔ The contribution of J/ψ from (re)combination should lead to a significant elliptic flow

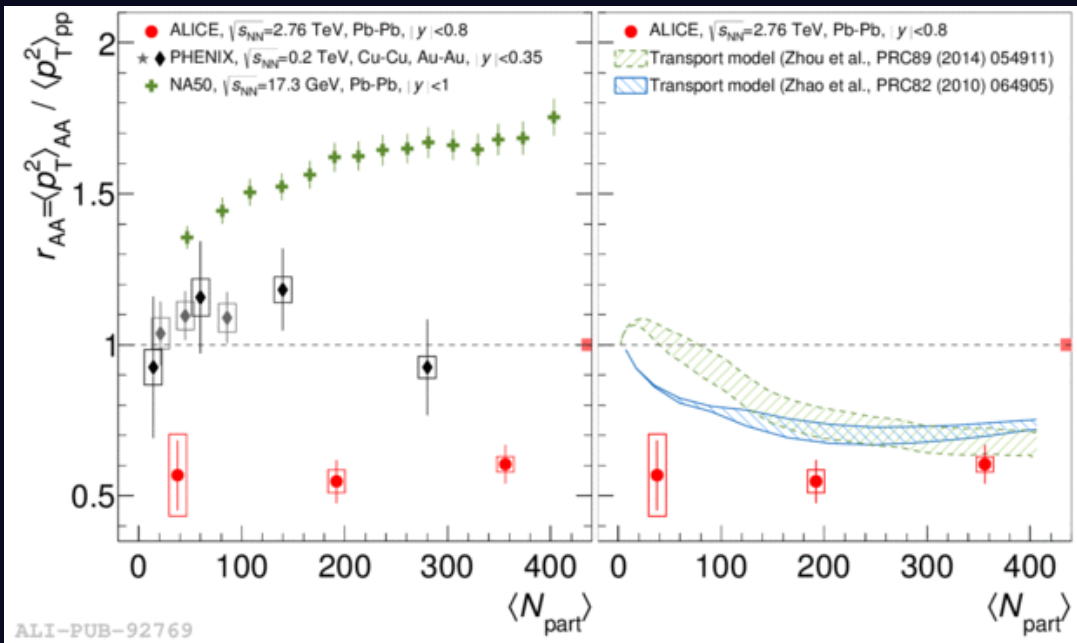


➔ Hint for J/ψ flow at LHC, contrary to $v_2 \sim 0$ observed at RHIC!

ALICE: qualitative agreement with transport models including regeneration
CMS: path-length dependence suppression effect?

$$r_{AA} = \frac{\langle p_T^2 \rangle_{AA}}{\langle p_T^2 \rangle_{pp}}$$

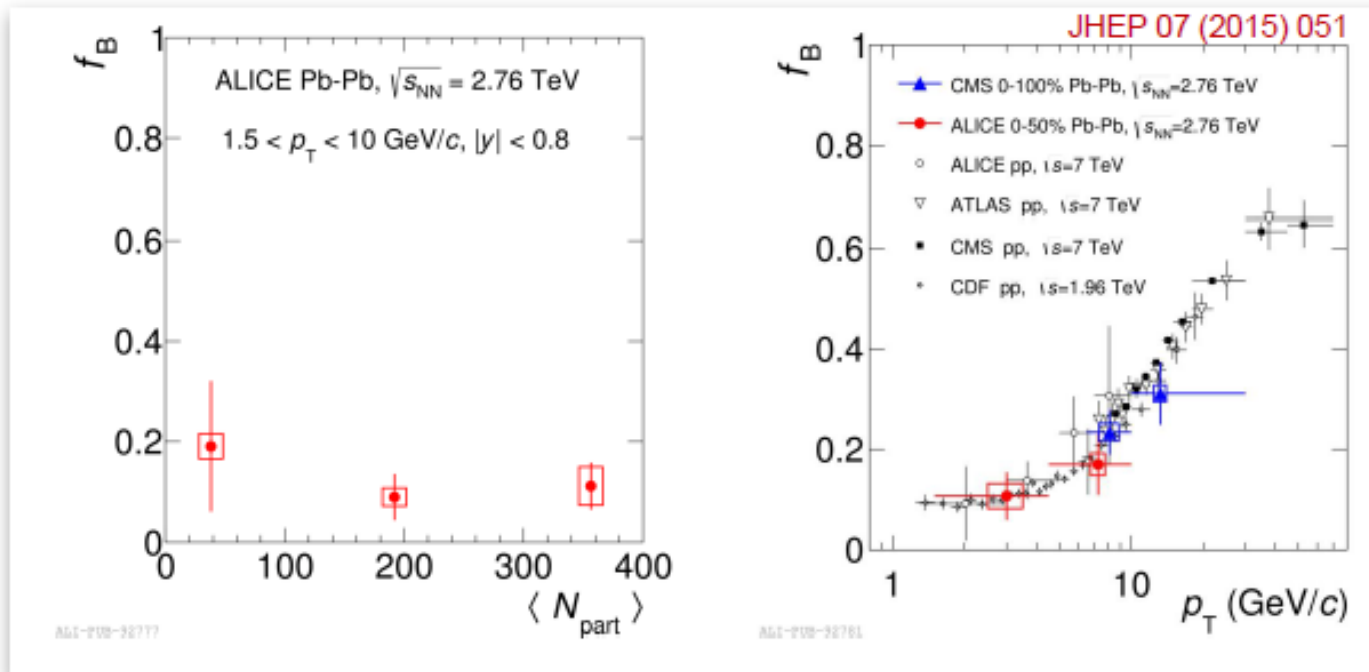
ALICE, JHEP07(2015)051



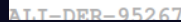
→ Strong \sqrt{s} -dependence as at forward- y but less pronounced centrality dependence observed by ALICE at mid-rapidity

→ Transport models, predicts a larger r_{AA} , with a stronger centrality dependence

$f_B = \text{non-prompt} / \text{inclusive}$



Y. Kim, H. Pereira da Costa

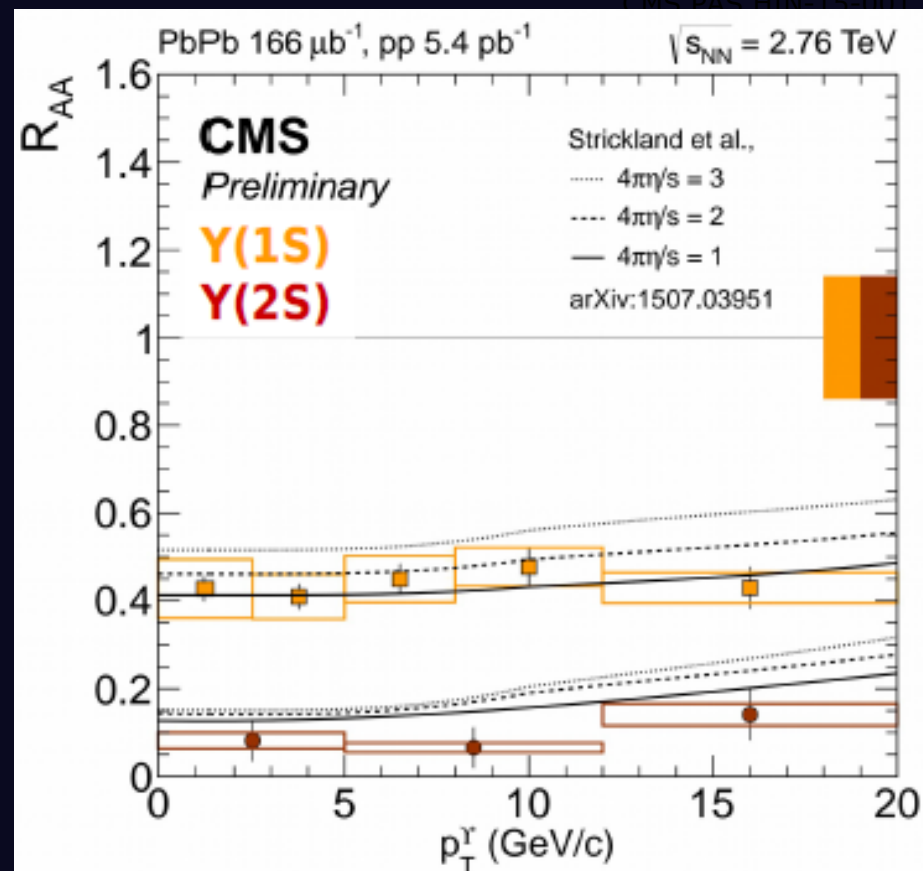
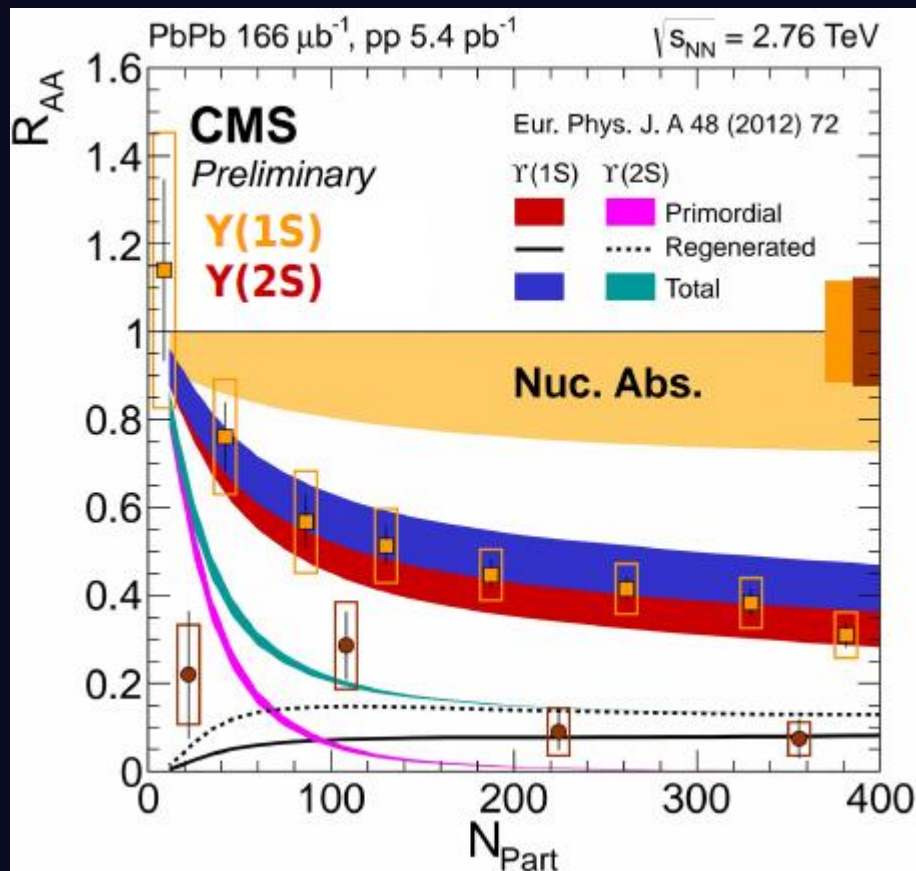


- CMS, PRL 113(2014) 262301, ALICE arXiv:1506.08804

$\Upsilon(ns)$ production in AA

51

CMS PAS HIN-15-001



J/ψ at very low p_T

52

Coherent photoproduction

$$N_{AA}^{\text{excess } J/\psi} / (1+f_D+f_V) \rightarrow d\sigma_{J/\psi}^{\text{coh}}/dy$$

Cent. (%)	$N_{AA}^{J/\psi}$	$N_{AA}^{\text{excess } J/\psi}$	$d\sigma_{J/\psi}^{\text{coh}}/dy (\mu\text{b})$
0-10	$339 \pm 85 \pm 78$	< 251	< 318
10-30	$373 \pm 87 \pm 75$	< 237	< 290
30-50	$187 \pm 37 \pm 15$	$62 \pm 37 \pm 21$	$73 \pm 44^{+26}_{-27} \pm 10$
50-70	$89 \pm 13 \pm 2$	$50 \pm 14 \pm 5$	$58 \pm 16^{+8}_{-10} \pm 8$
70-90	$59 \pm 9 \pm 3$	$51 \pm 9 \pm 3$	$59 \pm 11^{+7}_{-10} \pm 8$

Photon from the Pb EM field interacts with the Pb nucleus (coherent) or with a nucleon (incoherent).

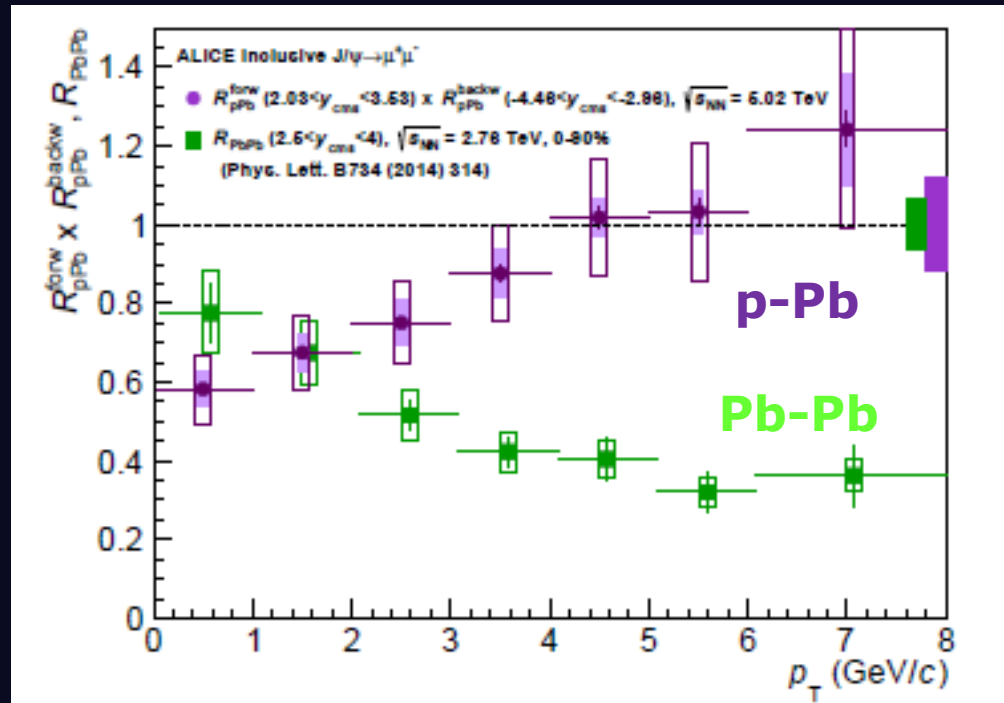
An extrapolation of measured J/ψ in UPC ($b_1=2R_{\text{Pb}}, b_2=\infty$) to 70-90% ($N_{\text{part}} \sim 11$) centrality class ($b_1^{70-90\%}, b_2^{70-90\%}$), provides a cross section of $\sim 40 \mu\text{b}$. Good qualitative agreement with the measured value $59 \pm 16 \mu\text{b}$.

From pA to AA

63

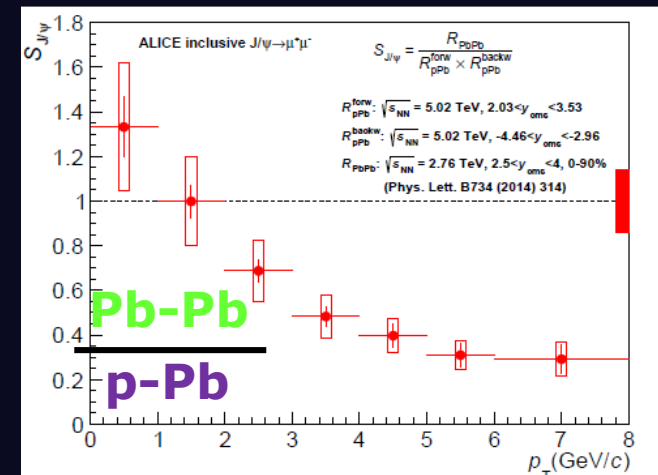
➔ Once CNM effects are measured in pA, what can we learn on J/ψ production in PbPb?

- Hypothesis:
- 2→1 kinematics for J/ψ production
 - CNM effects (dominated by shadowing) factorize in p-A
 - CNM obtained as $R_{pA} \times R_{Ap}$ (R_{pA}^2), similar x-coverage as PbPb



➔ Sizeable p_T dependent suppression still visible → CNM effects not enough to explain AA data at high p_T

➔ we get rid of CNM effects with
AA / pA



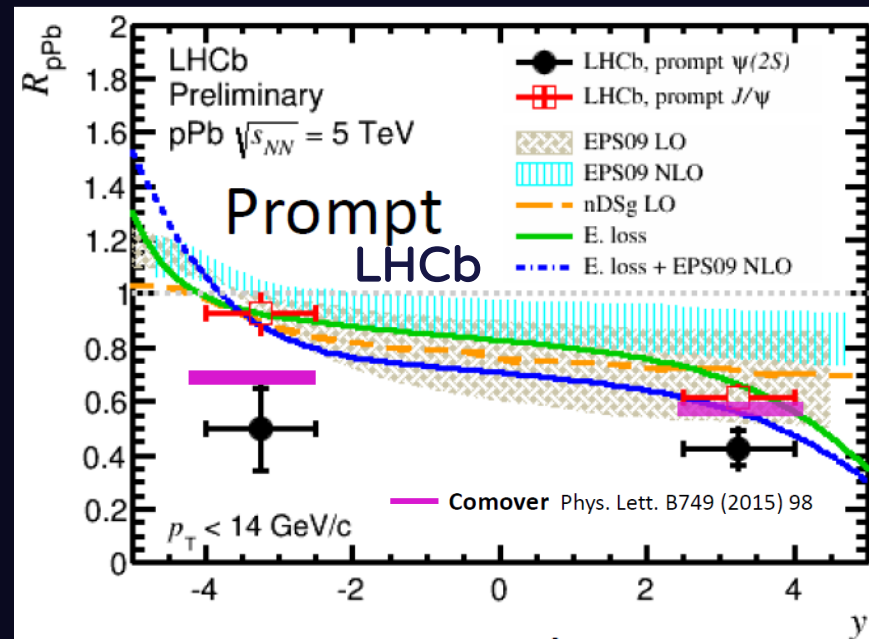
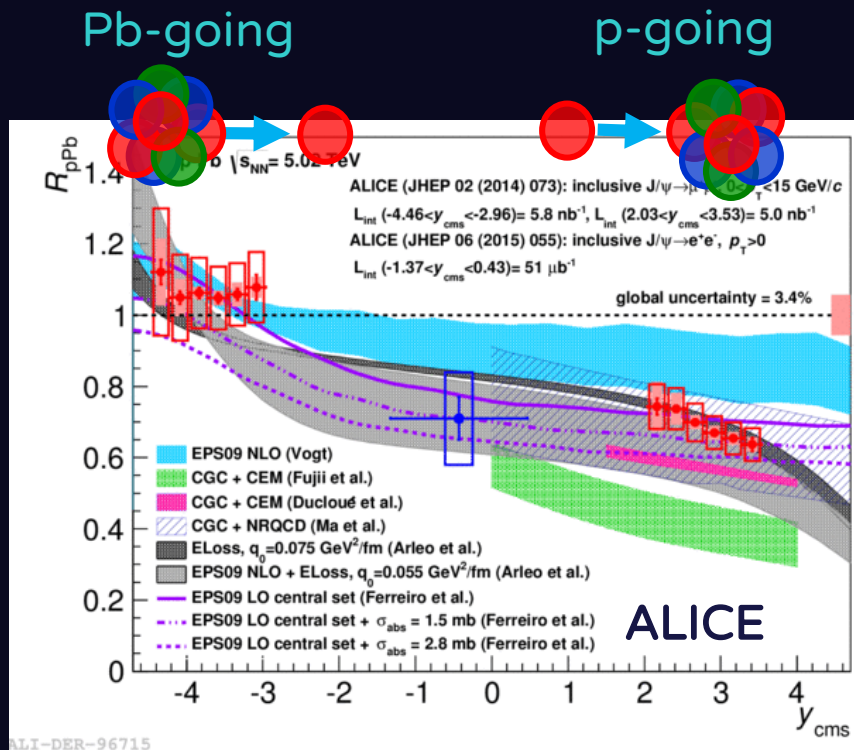
53

Backup slides pA

J/ψ production in pA

55

M. Leoncino, Z. Yang



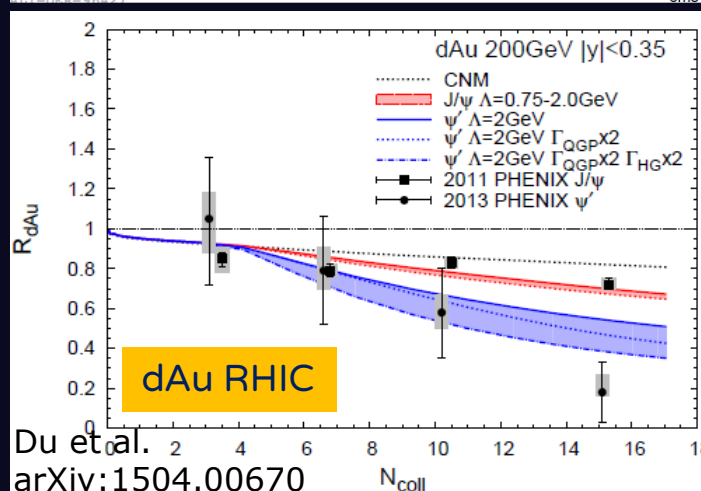
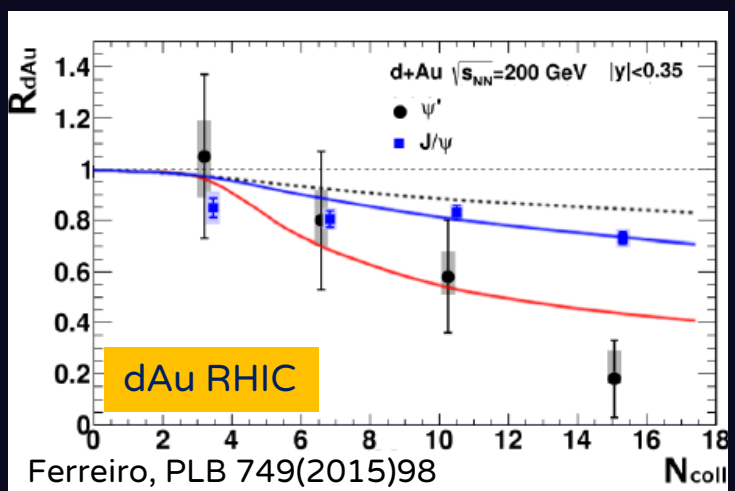
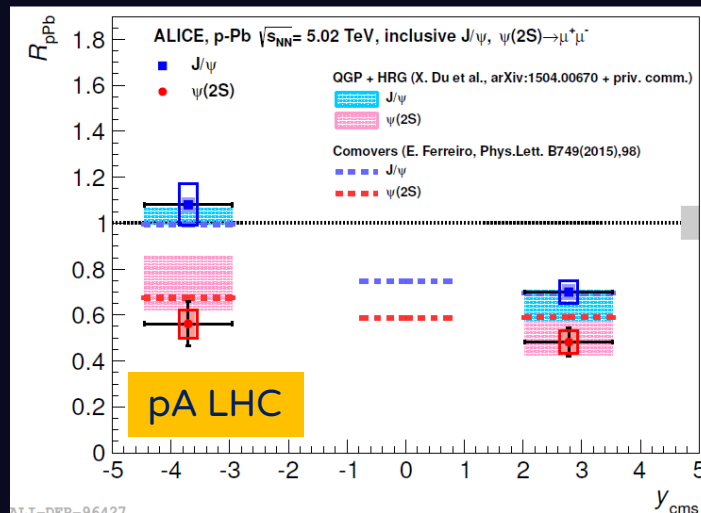
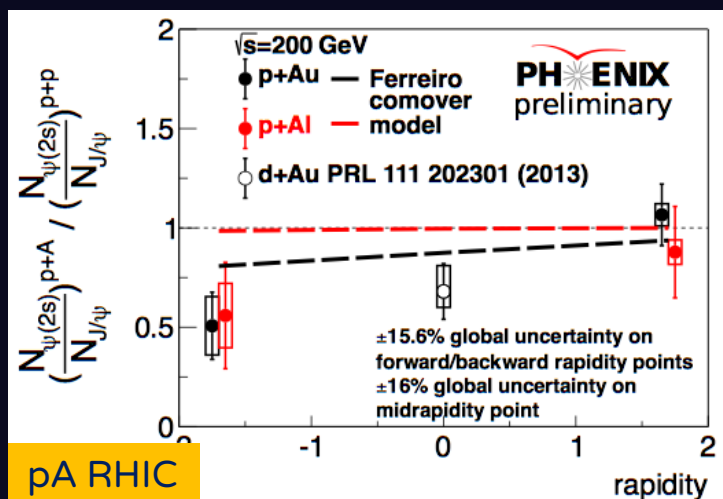
J/ψ production modified by CNM effects $\rightarrow R_{pA}$ decreases at forward y

Theoretical predictions:

- shadowing calculations and models including coherent parton energy loss reasonably describe the data
- agreement with CGC depends on the implementation

Comparison to theoretical models 56

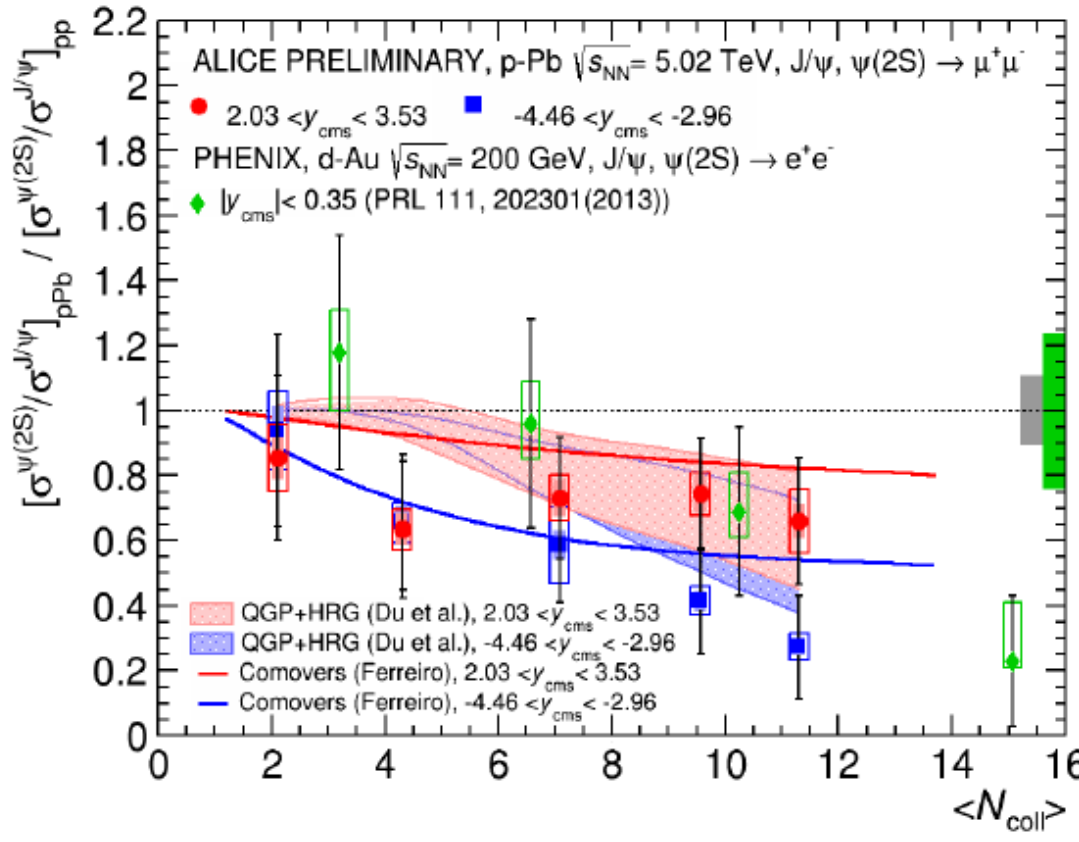
➔ QGP+hadron resonance gas (Rapp) or comovers models (Ferreiro) reasonably describe both J/ψ and $\psi(2S)$ suppression at RHIC and LHC



J/ψ
→ small suppression beyond CNM effects

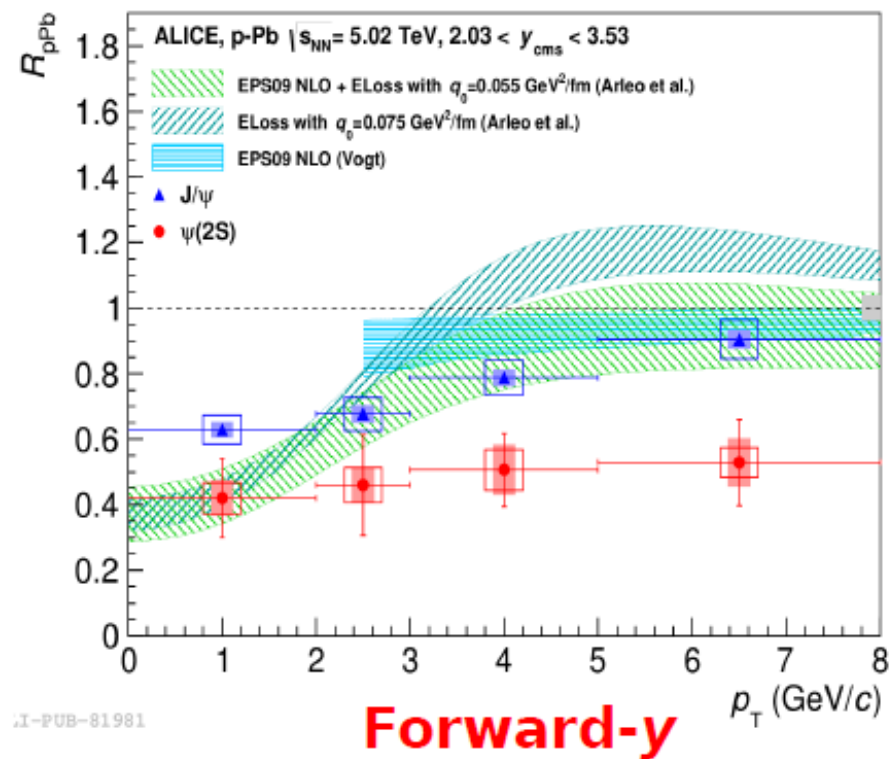
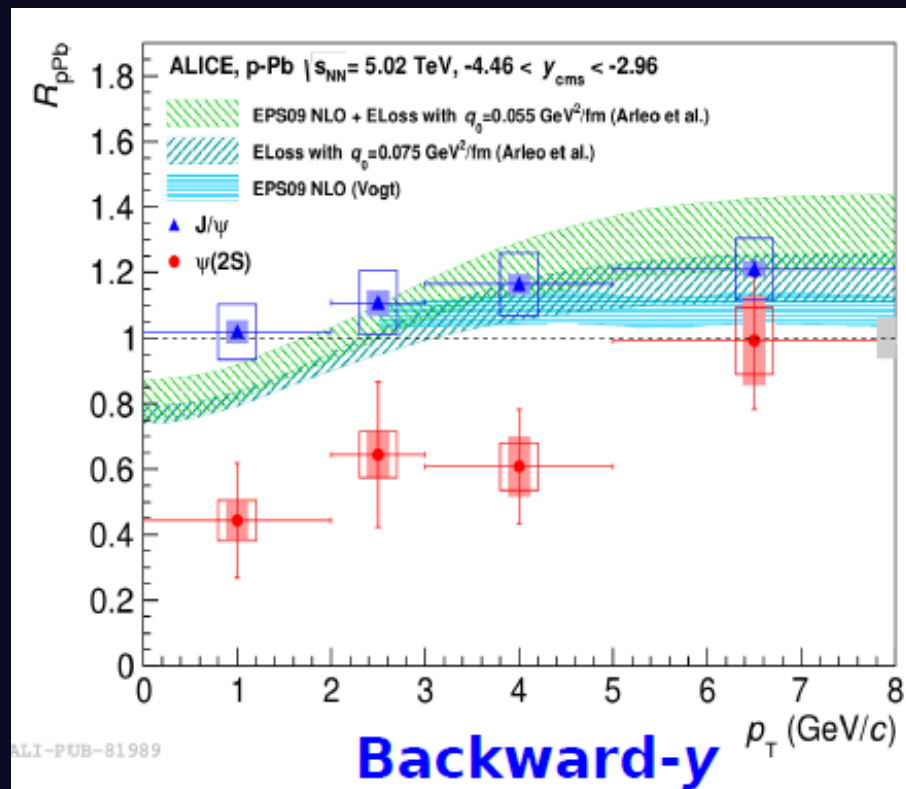
$\psi(2S)$
→ strongly affected by comovers due to its larger size
→ comovers more important in the A-going direction

$\psi(2S) / J/\psi$ double ratio 57



Similar suppression trend observed versus centrality, by both ALICE and PHENIX

→ QGP+hadron resonance gas (Rapp) or comovers models (Ferreiro) describe the observed suppression



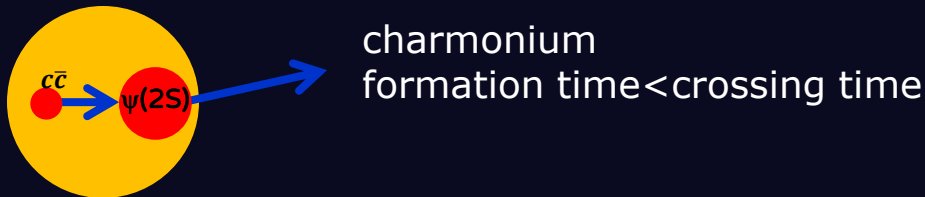
$\psi(2S)$ production in pA

59

- Being more weakly bound than the J/ψ , the $\psi(2S)$ is an interesting probe to have further insight on the charmonium behaviour in pA
- Low energy $\psi(2S)$ p-A results from NA50, E866 and HERA-B:

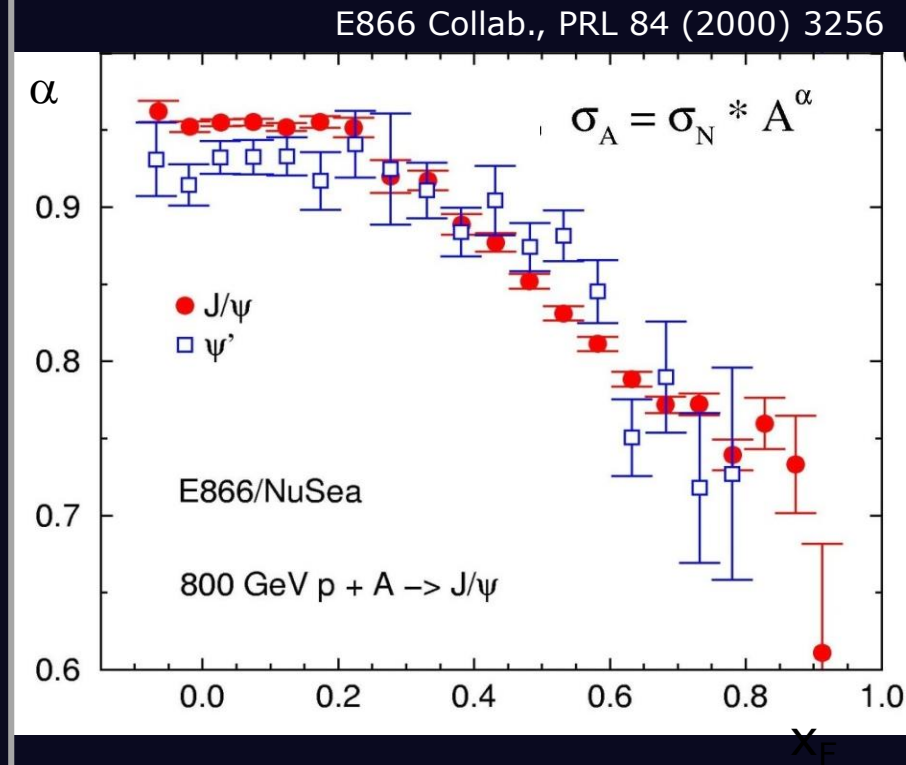
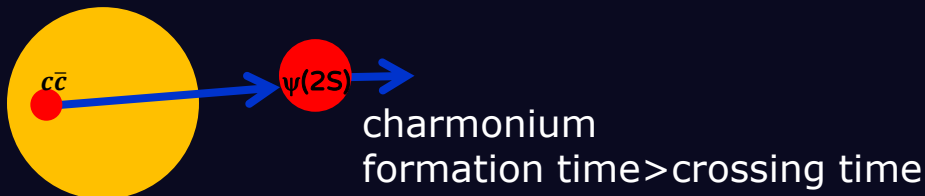
mid- y ($x_F \sim 0$):

$\psi(2S)$ suppression stronger than J/ψ one, interpreted via pair break-up
→ fully formed resonances traversing the nucleus



forward- y (high x_F):

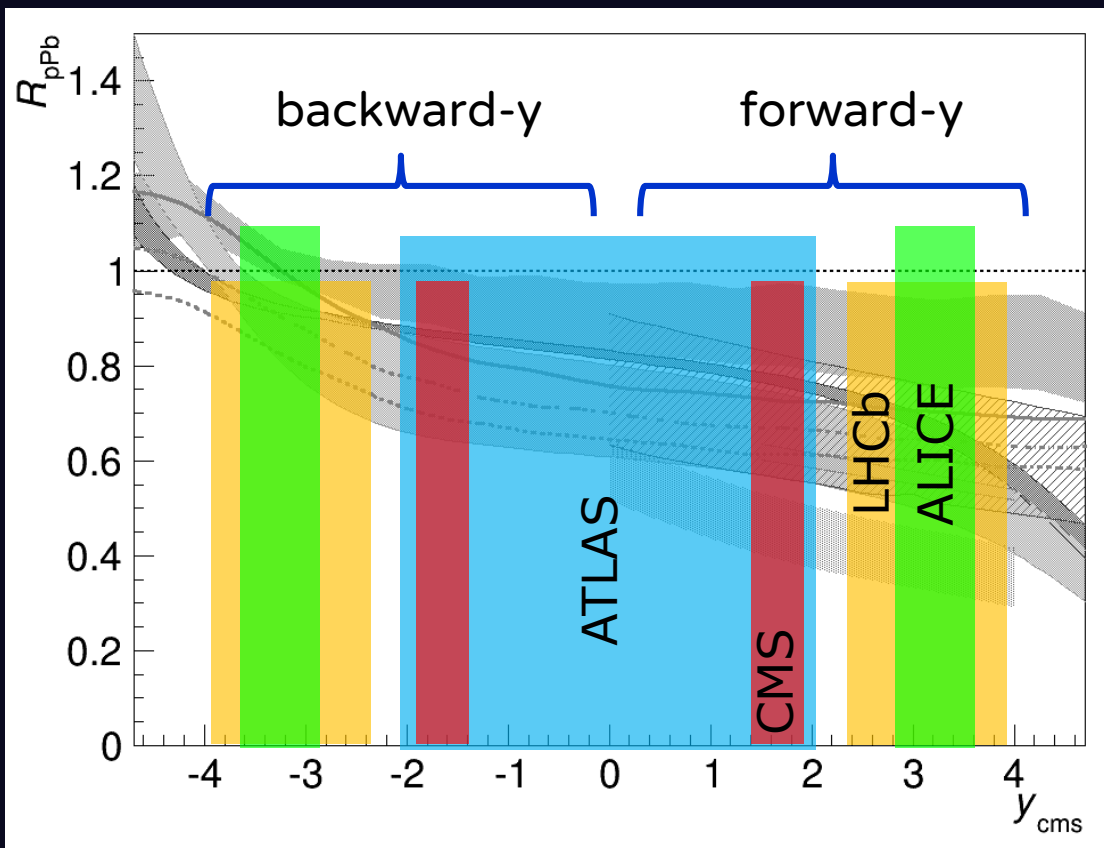
suppression becomes identical
→ dominated by energy loss



J/ψ forward-to-backward ratio 60

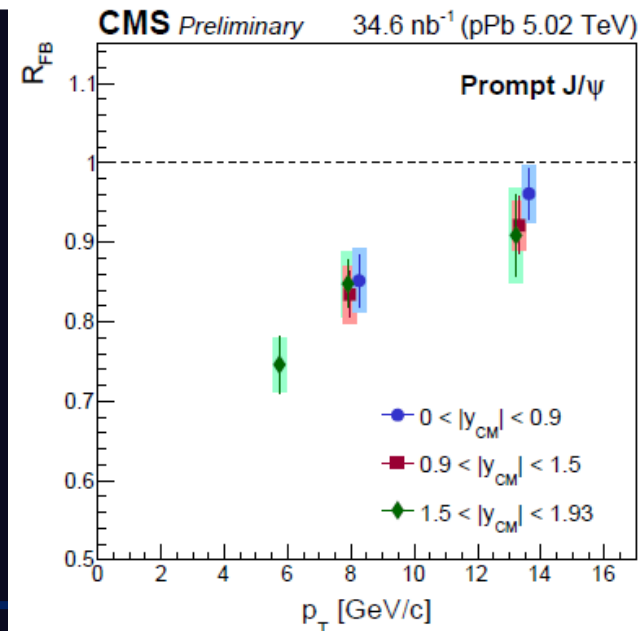
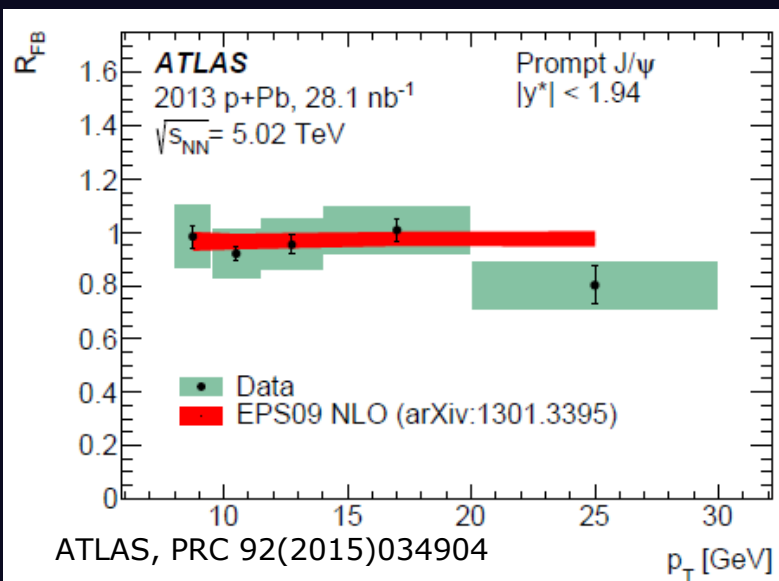
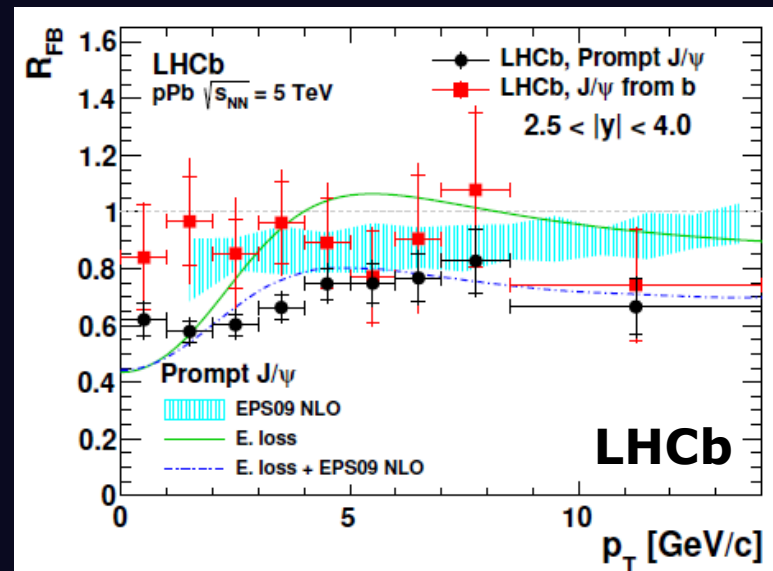
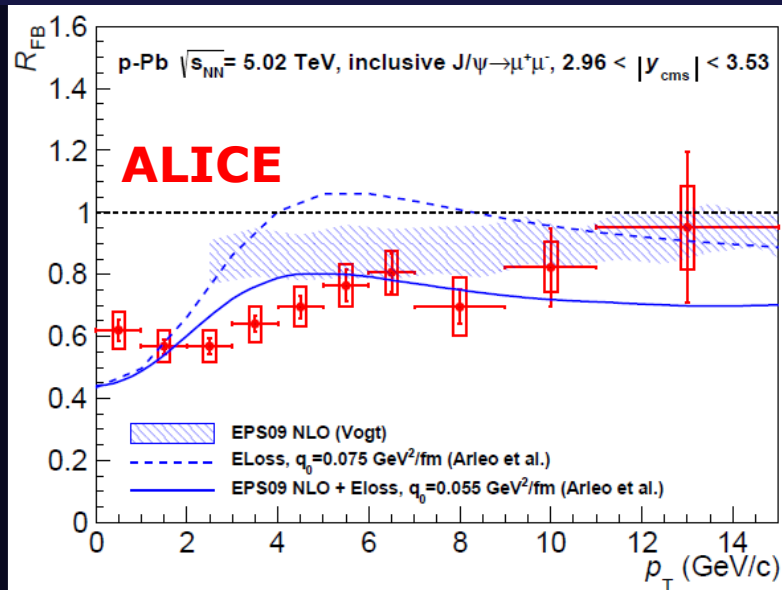
- Forward to backward ratio R_{FB} in a common y range
→ easier to evaluate wrt R_{pA} , since no pp reference is needed
→ but less straightforward to interpret

$$R_{FB} = \frac{Y_{J/\psi}^{forward}}{Y_{J/\psi}^{backward}}$$

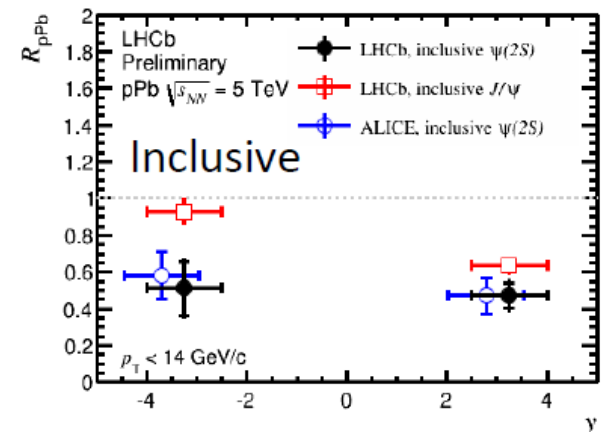
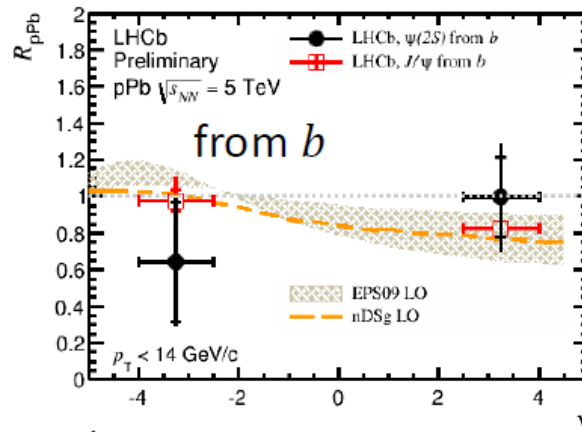
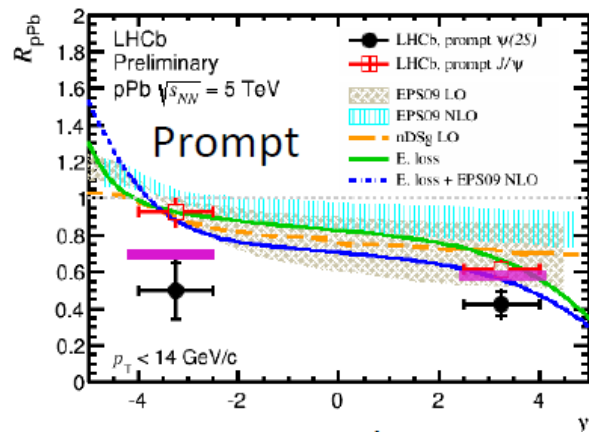


- Comparison of RFB results between experiments is not straightforward:
- 1) CNM effects have a strong rapidity dependence
 - 2) kinematic ranges explored by the experiments are rather different

J/ψ forward-to-backward ratio 61



J/ψ R_{pA} inclusive and prompt

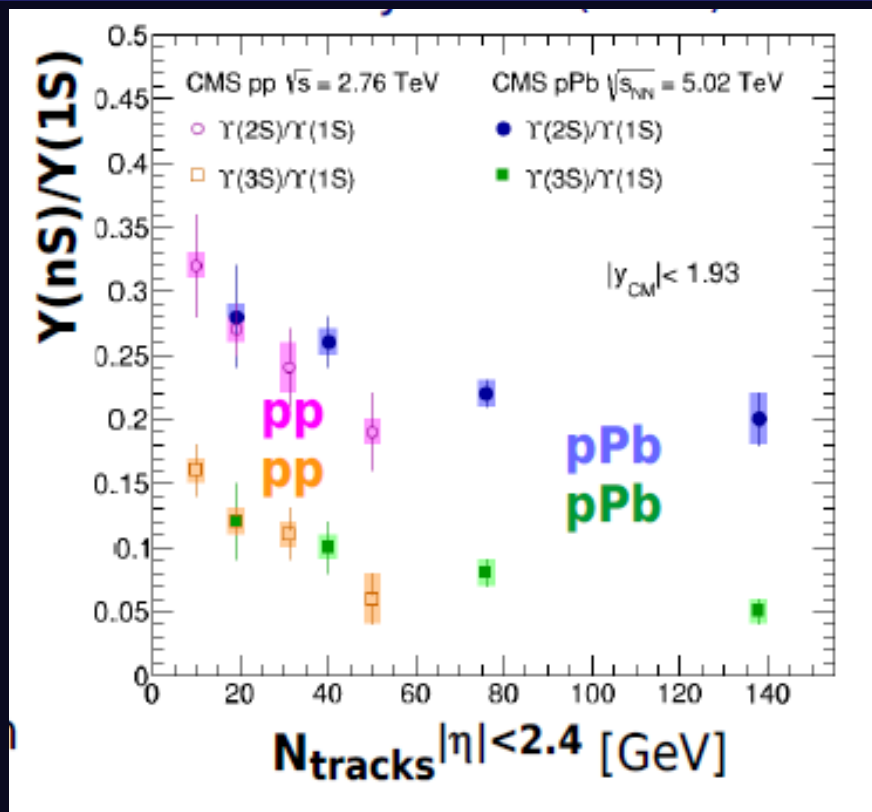


EPS09 LO Phys. Rev. C88 (2013) 047901
EPS09 NLO Int. J. Mod. Phys. E22 (2013) 1330007
nDSg LO Phys. Rev. C88 (2013) 047901
E. loss JHEP 03 (2013) 122
E. loss + EPS09 NLO JHEP 03 (2013) 122

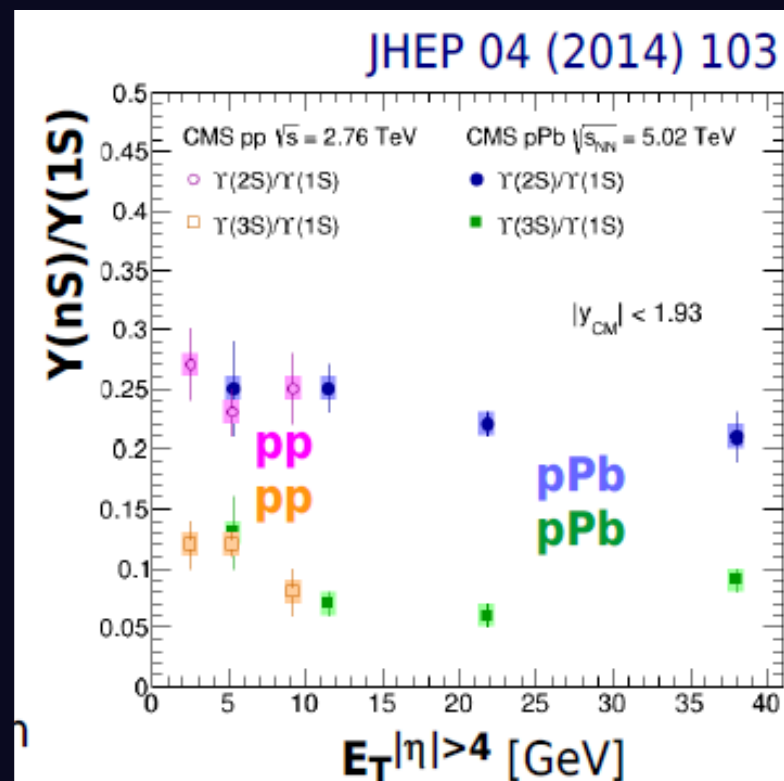
EPS09 LO Nucl.Phys.A926 (2014) 236
nDSg LO
Comover Phys. Lett. B749 (2015) 98

Y vs ev. activity

63



- $Y(nS)/Y(1S)$ ratios fall with event-activity CMS
 - Is the multiplicity affecting the $Y(nS)$?
 - Are the $Y(nS)$ produced differently with multiplicity?



$-\eta$			$+\eta$
HF	Y		HF
$[-5.2, -4]$	$[-1.93, 1.93]$		$[4, 5.2]$
N_{tracks}			
$[-2.4, 2.4]$			