

Baryon number cumulant ratios at finite density in the strong-coupling lattice QCD

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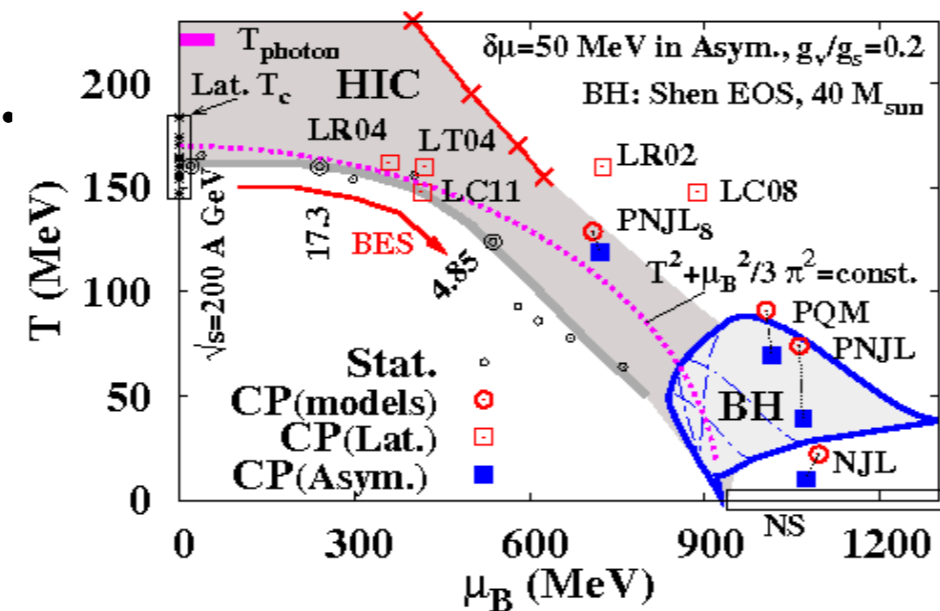
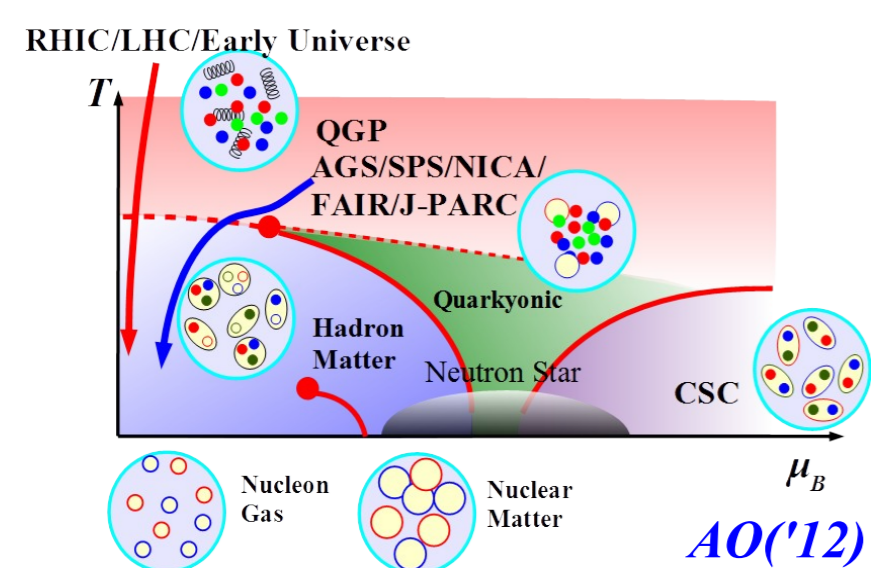
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Introduction

- Critical point in QCD phase diagram connects
Crossover trans. at low μ and first order trans. at high μ

Asakawa, Yazaki ('89)

→ CP could be accessible in BES.

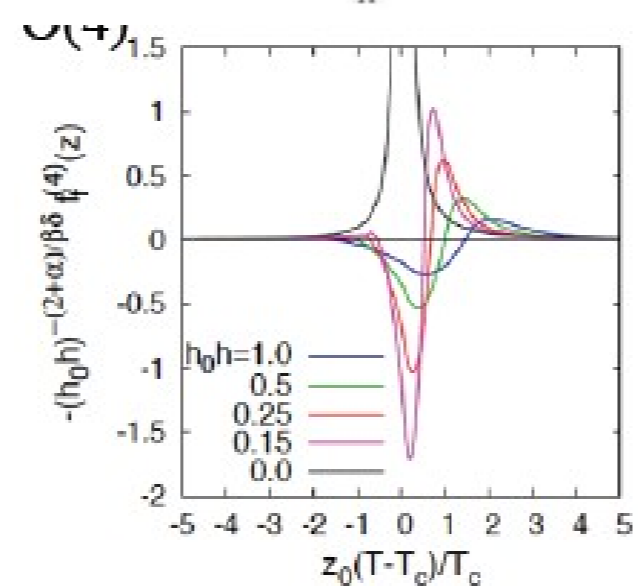
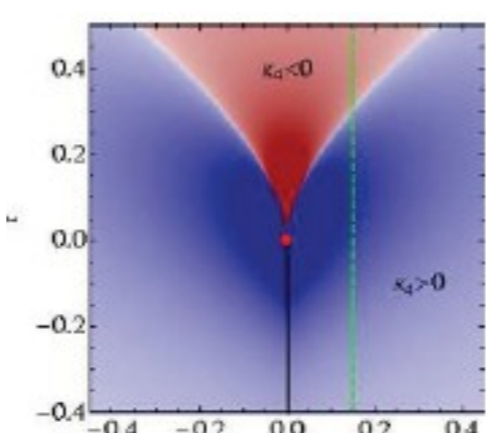
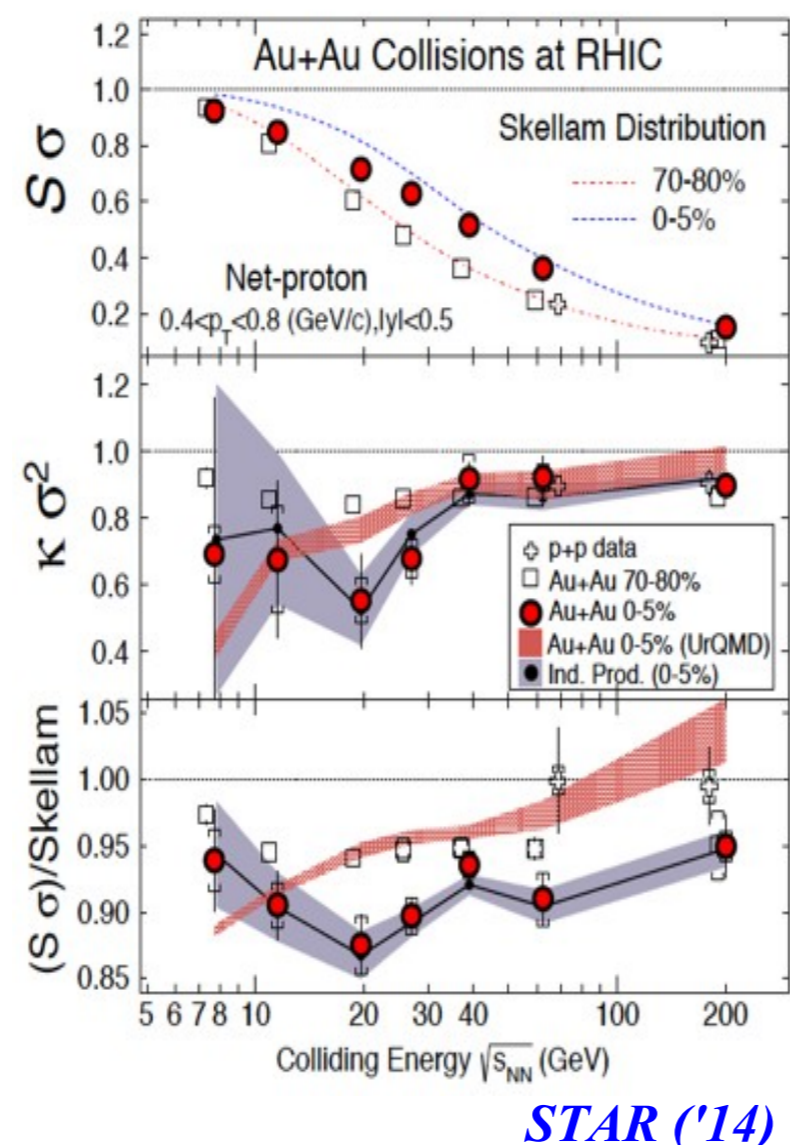


- Fluctuations of conserved charge = Promising signal of CP

$$S\sigma = \chi_\mu^{(3)}/\chi_\mu^{(2)} \quad \kappa\sigma^2 = \chi_\mu^{(4)}/\chi_\mu^{(2)}$$

$$\chi_\mu^{(n)} = \frac{1}{VT^3} \frac{\partial^n \log Z}{\partial (N_c \mu/T)^n}$$

- Contain info. on earlier stage.
- Remnant of O(4) criticality may cause non-monotonic behavior of the singular part.
- Larger fluctuation of order par. around CP → Negative κ
- Stephanov ('11) [Z(2)]
- Non-monotonic behavior of cumulant ratio is observed @ BES → CP signal?



We need calculation including regular part at finite μ !
This work:
Net baryon number cumulants at finite T and μ in the strong-coupling and chiral limit of lattice QCD

Strong-Coupling Lattice QCD

- Lattice QCD action (unrooted staggered fermion)

$$L = \frac{1}{2} \sum_x [V_x^+ - V_x^-] + \frac{1}{2\gamma} \sum_{x,j} \eta_j(x) [\bar{\chi}_x U_j(x) \chi_{x+j} - \chi_{x+j}^* U_j^*(x) \bar{\chi}_x] + \frac{m_0}{\gamma} \sum_x \bar{\chi}_x \chi_x + \frac{2N_c}{g^2} S^{\text{plaq}}$$

$$V_x^+ = \bar{\chi}_x U_0(x) e^{\mu/\gamma^2} \chi_{x+\hat{0}}, \quad V_x^- = \chi_{x+\hat{0}}^* U_0^*(x) e^{-\mu/\gamma^2} \bar{\chi}_x$$

- $\eta_j(x) = (-1)^{x_0 + \dots + x_{j-1}}$
- U(1)_L x U(1)_R chiral sym.
- $\chi_x \rightarrow \exp[i\theta \varepsilon(x)] \chi_x$, $\varepsilon(x) = (-1)^{x_0 + x_1 + x_2 + x_3}$
- Anisotropy parameter γ ($T = \gamma^2/N_c$) Bilic et al. ('92)

- Strong coupling limit

Damgaard, Kawamoto, Shigemoto ('84),
Jolicœur, Kluberg-Stern, Lev, Morel, Petersson ('84).

Spatial link integral

→ Fermion action with four-Fermi int. (LO in 1/d expansion)

$$S_{\text{eff}}^{(\text{SCL})} = \frac{1}{2} \sum_x [V_x^+ - V_x^-] + \frac{m_0}{\gamma} \sum_x M_x - \frac{1}{4N_c \gamma^2} \sum_{x,j} M_x M_{x+j} \quad (M_x = \bar{\chi}_x \chi_x)$$

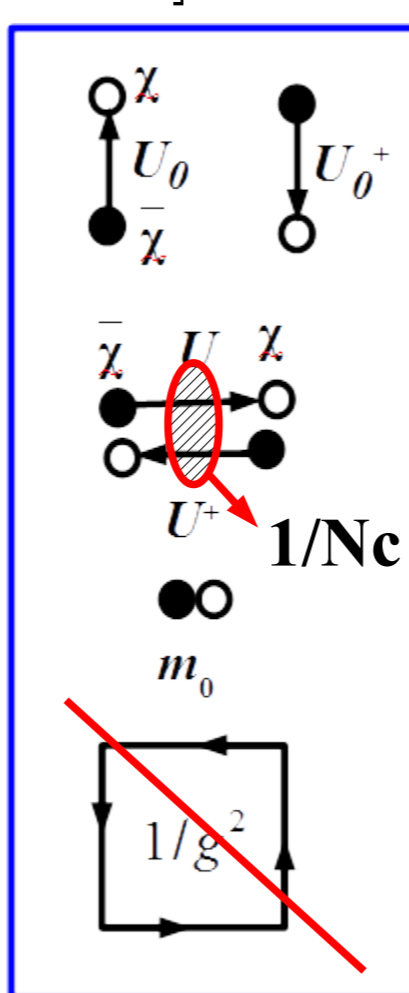
- Extended Hubbard Stratonovich transf.

→ Eff. Action of Auxiliary Field

$$\exp(\alpha AB) = \int d\varphi d\varphi^* \exp(-\alpha[\varphi^* \varphi + \varphi^* A + \varphi B])$$

$$S_{\text{eff}}^{(\text{EHS})} = \frac{1}{2} \sum_x [V_x^+ - V_x^-] + \frac{1}{\gamma} \sum_x m_x M_x + S_{\text{AF}}$$

$$S_{\text{AF}} = \frac{L^3}{4N_c} \sum_{k\tau} f(k) [\sigma_{k\tau}^* \sigma_{k\tau} + \pi_{k\tau}^* \pi_{k\tau}] \quad \text{Ichihara, Nakano, AO ('14)}$$



MC integral over AF
→ Fluctuation Obsv.

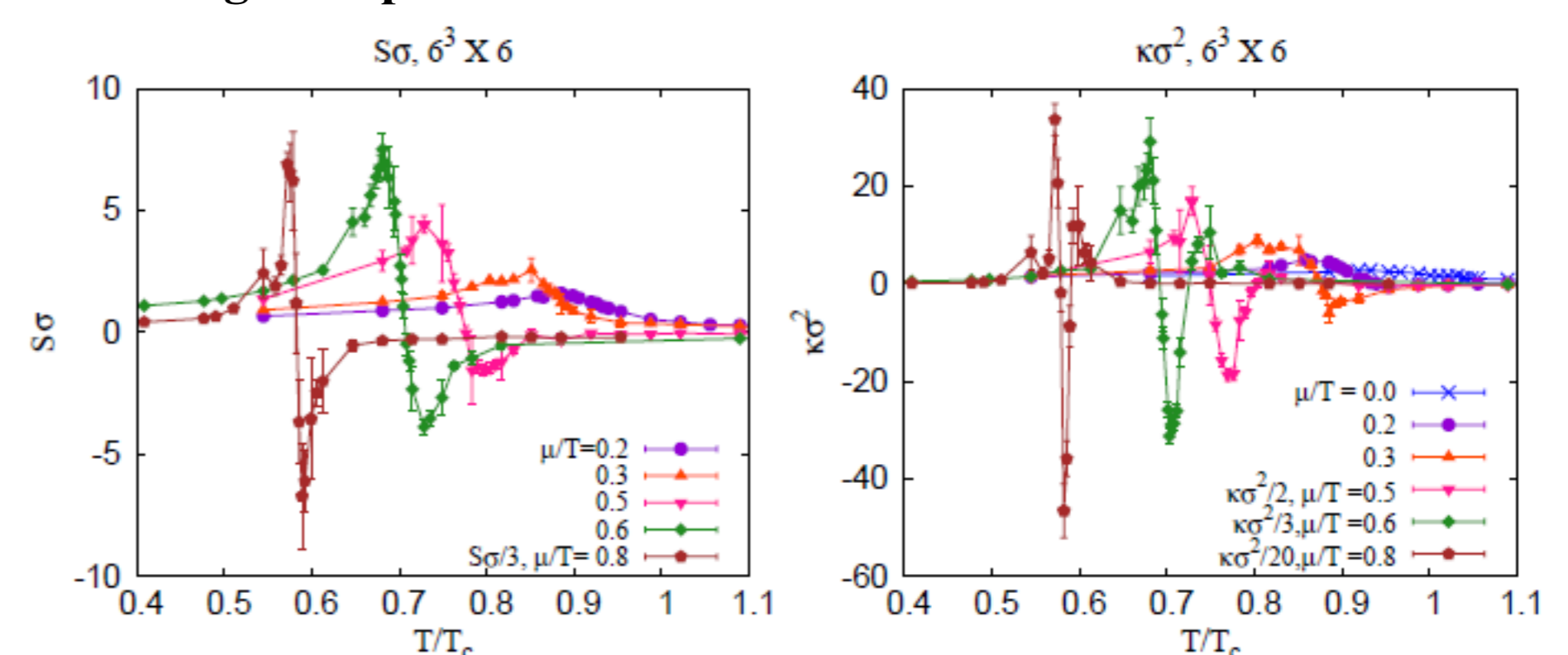
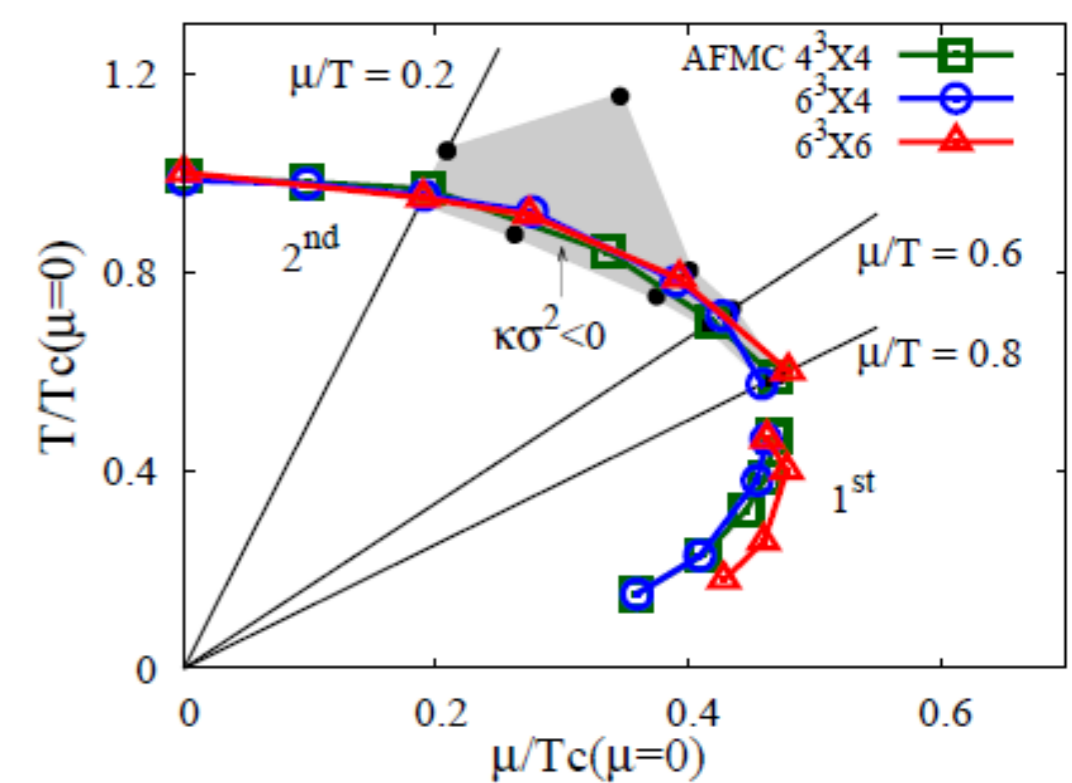
Results

- Lattice set up

- unrooted staggered fermion [O(2) symmetry]
- Strong-coupling and chiral limit
- Auxiliary Field MC method on 4⁴, 6³×4, 6⁴, 8⁴ lattices Ichihara, Nakano, AO ('14)
- $\mu/T = 0.2, 0.3, 0.5, 0.6, 0.8$

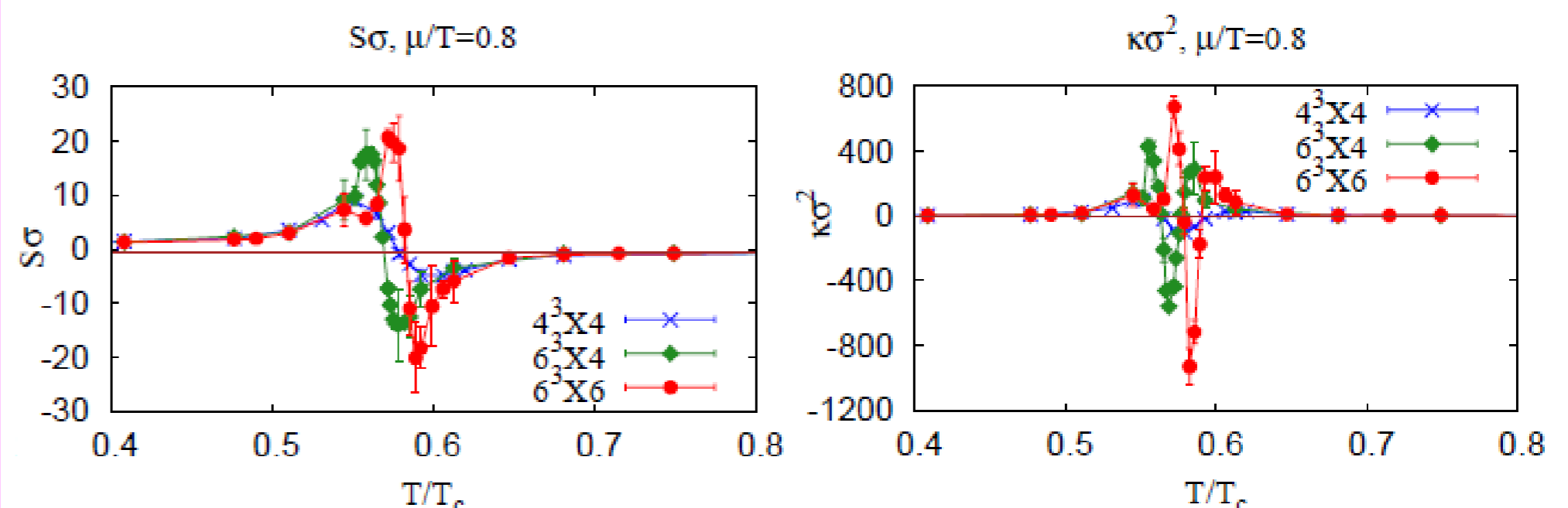
- (μ, T) dependence

- Small μ/T :
Peak around boundary.
- Large μ/T :
Oscillatory behavior ($S\sigma$: +-, $\kappa\sigma^2$: +-+)
Larger amplitude towards TCP.



- Size dependence

- Amplitude show divergent behavior and Negative valley in $\kappa\sigma^2$ narrows with increasing lattice size. [Singular part dominance]
- Consistent with the O(4) scaling analysis Friman, Karsch, Redlich, Skokov ('11) [Qualitative behavior is similar in O(4) and O(2)]



- Negative normalized kurtosis region in the QCD phase diagram

- Negative kurtosis area exists around the phase boundary.
- Expected to shrink in the chiral and thermodynamic limit.

Summary & Discussion

- We investigate normalized skewness ($S\sigma$) and kurtosis ($\kappa\sigma^2$) in the chiral and strong coupling limit ($\sim 6^4$ lattice).
- We find oscillatory behavior at large μ/T and negative kurtosis valley due to the finite size effect.
- Peak heights of skewness and kurtosis increases and negative valley of kurtosis shrinks on larger lattices, as suggested in O(4) scaling analysis in the chiral limit.
- Finite lattice size is found to smear the critical behavior of cumulants as the finite quark mass does.
- Important next steps
 - Cumulants on larger lattices [Finite size scaling],
 - Finite mass effects [Negative region should survive on large lattice].

Refs: M. Asakawa, K. Yazaki, NPA504 ('89) 668 / M. A. Stephanov, PRL107 ('11) 052301 / B. Friman, F. Karsch, K. Redlich, V. Skokov, EPJC71 ('11) 1694 / L. Adamczyk, et al. (STAR Collab.), PRL 112 ('14) 032302 / T. Ichihara, A. Ohnishi, T.Z. Nakano, PTEP 2014('14),123D02 / T. Ichihara, K. Morita, A. Ohnishi, PTEP 2015 ('15), in press / A. Ohnishi, PTPS 193 ('12), 1.