

Crystalline chiral condensate in an external magnetic field

Kazuya Nishiyama

Collaborator :Ryo Yoshiike, Shintaro Karasawa, Toshitaka Tatsumi (Kyoto University)

Motivation

Elucidate QCD phase structure by taking into account both of inhomogeneity and magnetic field

Preceding Study and problems

Inhomogeneous Chiral Condensate

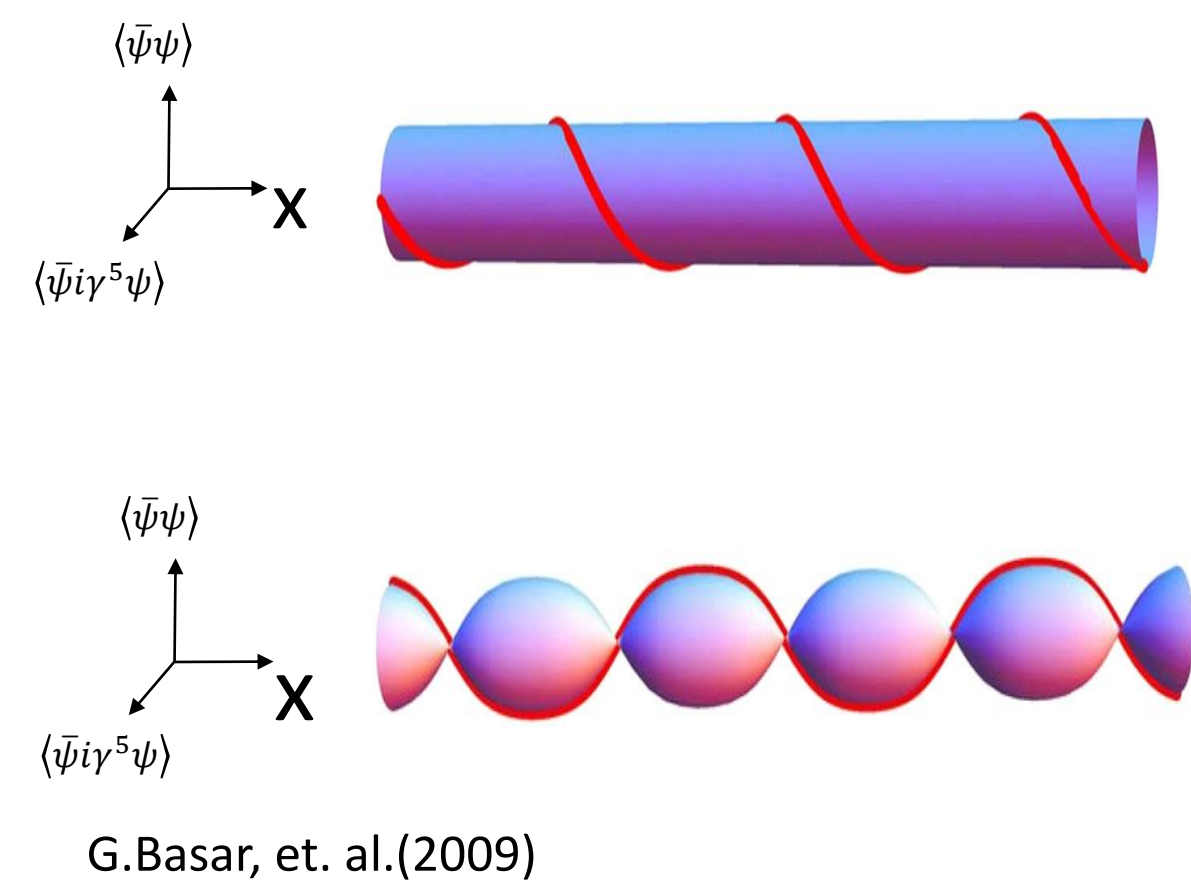
$$\Delta(x) = -2G[\langle\bar{\psi}\psi\rangle + i\langle\bar{\psi}i\gamma^5\tau^3\psi\rangle]$$

- Dual Chiral Density Wave(DCDW)
Plane wave type condensate

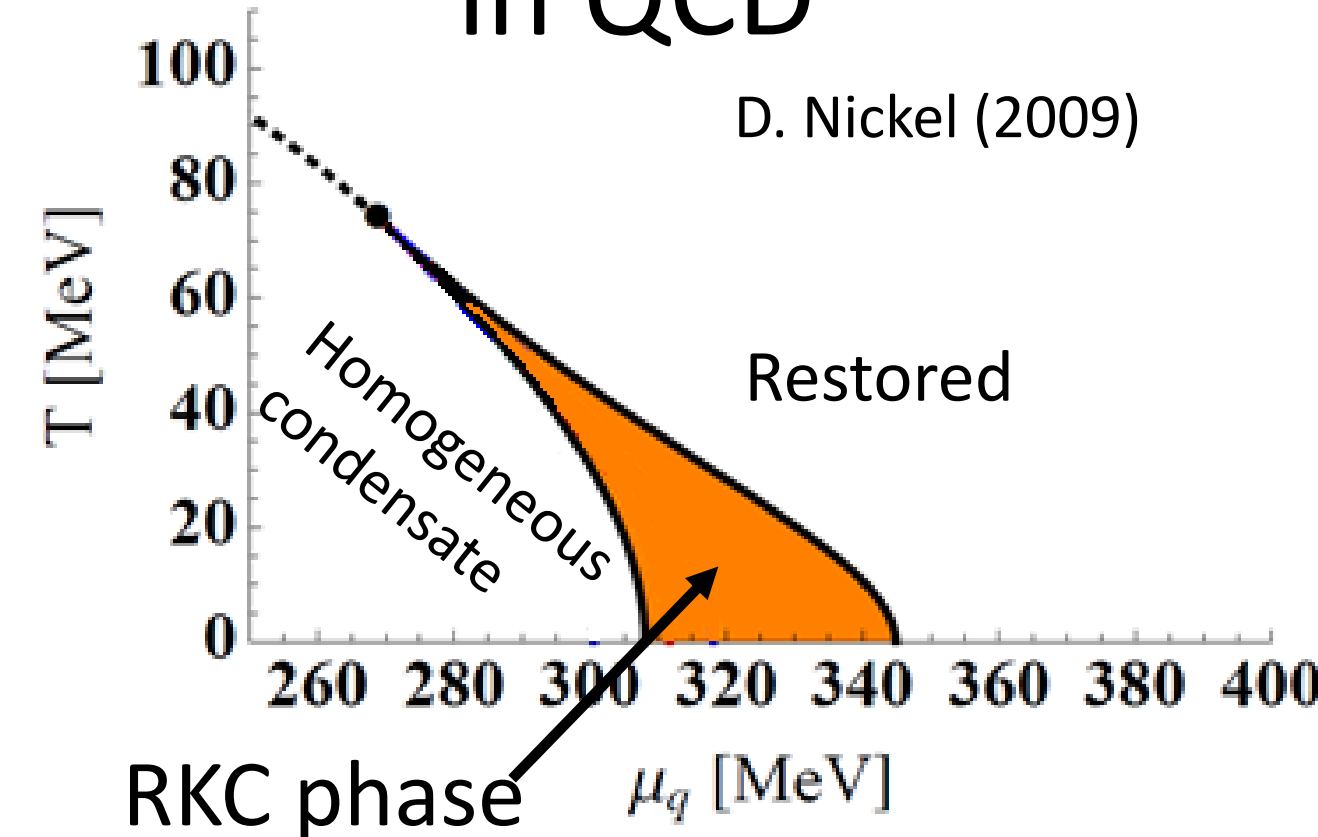
$$\Delta_{\text{DCDW}}(x) = \lambda e^{iqx}$$

- Real kink Crystal(RKC)
Multi soliton type condensate

$$\Delta_{\text{RKC}}(x) = \lambda \left(\frac{2\sqrt{\nu}}{1+\sqrt{\nu}} \right) \text{sn} \left(\frac{2\lambda x}{1+\sqrt{\nu}}; \nu \right)$$



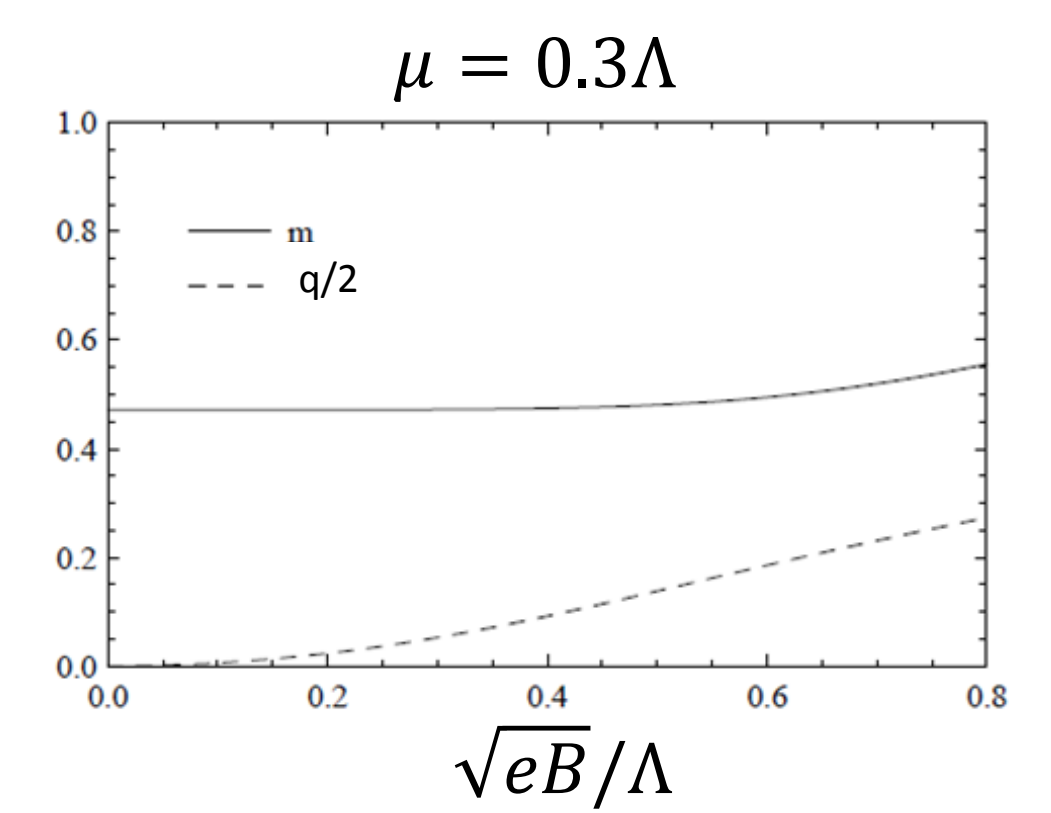
Inhomogeneous chiral Phase
in QCD



Without magnetic field, RKC is favored

DCDW in Magnetic Field

I.E. Frolov, et. al. (2010)



Inhomogeneity grows with the increase of B

- Which phase is favored in external magnetic field?
- Why does magnetic field drive DCDW?

Hybrid condensate

$$\Delta(x) := M(x)e^{iqx} = \frac{2m\sqrt{\nu}}{1+\sqrt{\nu}} \text{sn} \left(\frac{2mx}{1+\sqrt{\nu}}; \nu \right) \times e^{iqx}$$

$$\Delta_{\text{DCDW}}(x) \xleftarrow{\nu \rightarrow 1} \Delta(x) \xrightarrow{q \rightarrow 0} \Delta_{\text{RKC}}(x)$$

Free energy calculation

• 1 Particle Energy Spectrum

$$E_{n,\zeta,\alpha} = \left(F_\alpha + \zeta \frac{q}{2} \right) \sqrt{1 + \frac{2eBn}{\left(F_\alpha + \zeta \frac{q}{2} \right)^2}} \quad n=1,2,\dots \quad \zeta = \pm 1$$

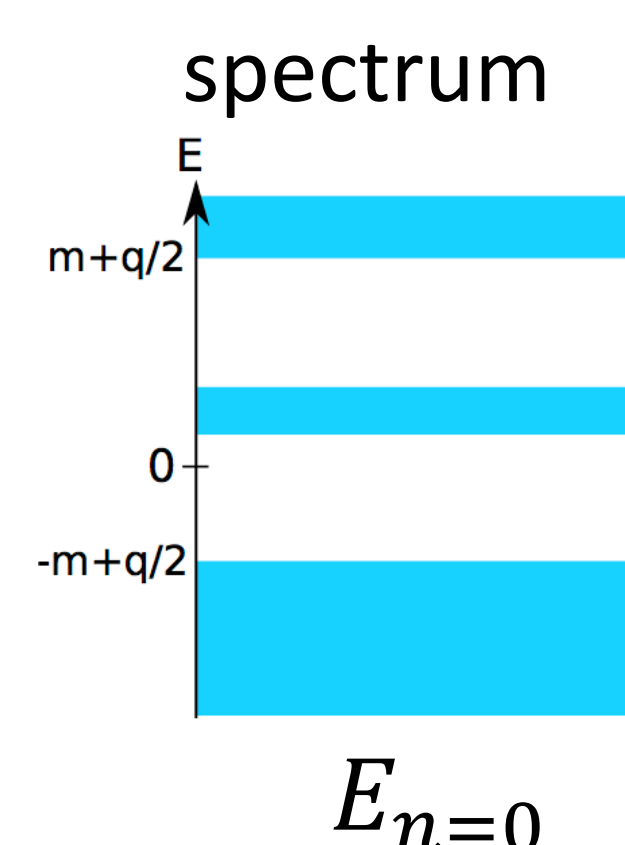
n : Landau level
 F_α : 1+1dim RKC spectrum (symmetric with respect to F=0)

$$E_{n=0,\alpha} = F_\alpha + \frac{q}{2} \quad n=0$$

$E_{n \geq 1, \zeta, \alpha}$: symmetric

$E_{n=0,\alpha}$: asymmetric (for $q \neq 0$)

$B \neq 0$ and $q \neq 0$, Energy spectrum is asymmetric.



• Anomalous Quark Number Density

$$\left. \frac{N}{V} \right|_{T=0} = - \left. \frac{\partial \Omega}{\partial \mu} \right|_{T=0} = N_c \sum_f \frac{|q_f B|}{2\pi} \sum_{\zeta} \sum_{n=0}^{\infty} \int dF \rho(F) \theta(E_{n\zeta F}) \theta(\mu - E_{n\zeta F})$$

$$+ N_c \sum_f \frac{1}{2} \frac{|q_f B|}{2\pi} \sum_F \text{sgn}(E_{n=0,F})$$

anomalous quark number density : ρ_{anom}

For DCDW

$$\rho_{anom} = N_c \sum_f \frac{1}{2} \frac{|q_f B|}{2\pi} \frac{q}{\pi}$$

$$\Omega_{anom} = -\mu \rho_{anom} \propto \mu B q$$

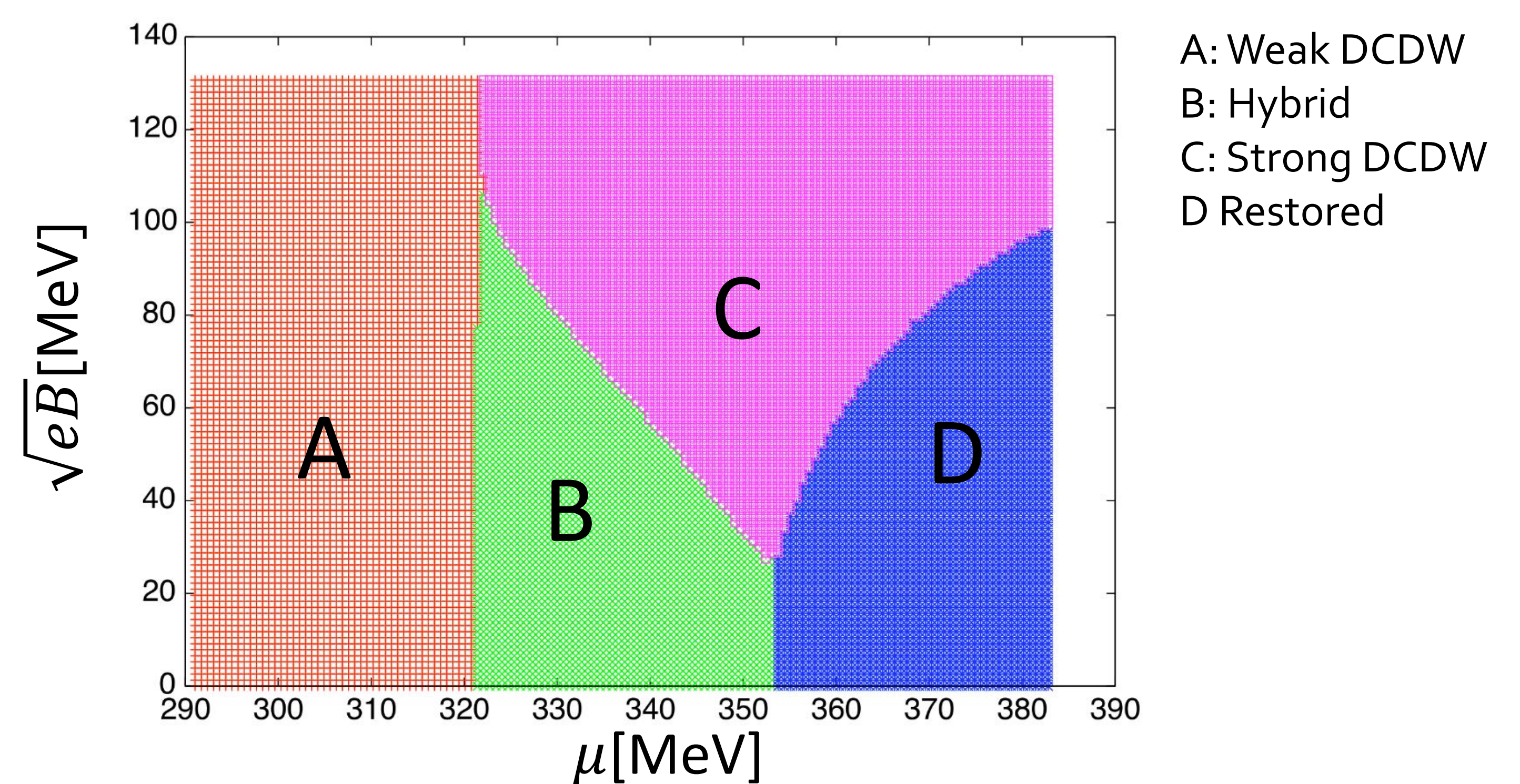
$B \neq 0$, homogeneous phase is unstable

• Free Energy

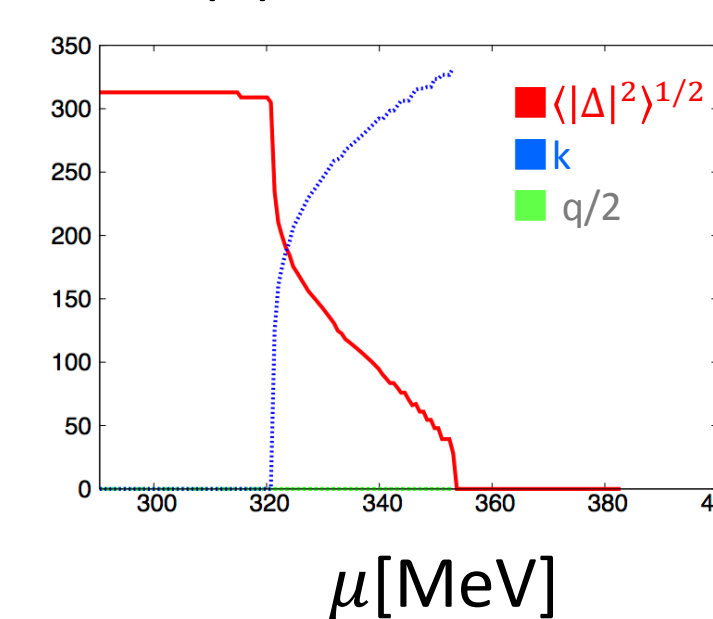
$$\Omega = \frac{\langle |\Delta(z)|^2 \rangle}{4G} - \frac{1}{\beta} N_c \sum_f \frac{|q_f B|}{4\pi^2} \sum_{\zeta} \sum_{n=0}^{\infty} \int dF \left[\frac{2F^2 - 2 \left(m \frac{1-\sqrt{\nu}}{1+\sqrt{\nu}} \right)^2 - m^2 + \langle |\Delta(z)|^2 \rangle}{\sqrt{(F^2 - m^2) \left(F^2 - \left(m \frac{1-\sqrt{\nu}}{1+\sqrt{\nu}} \right)^2 \right)}} \right] \ln \left(1 + e^{-\beta(E_{n\zeta}(\alpha) - \mu)} \right)$$

Phase structure is determined by $\frac{\partial \Omega}{\partial m} = \frac{\partial \Omega}{\partial \nu} = \frac{\partial \Omega}{\partial q} = 0$

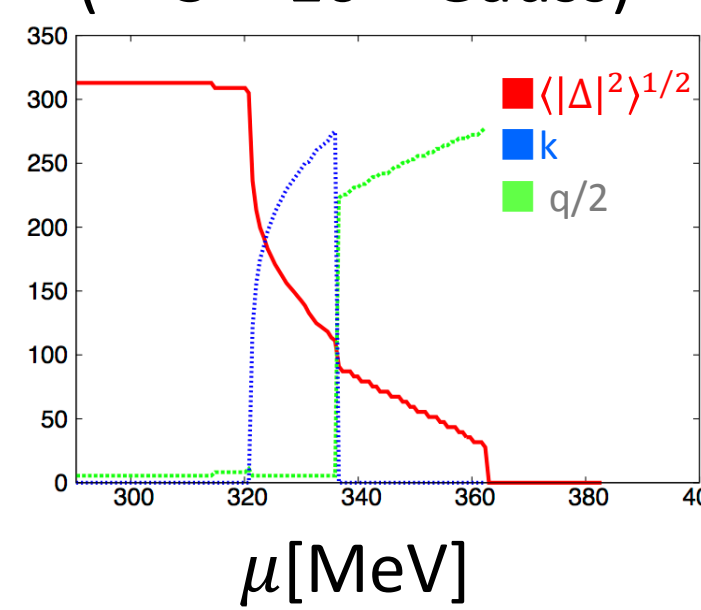
• Phase Diagram (μ -B plane, $T=0$)



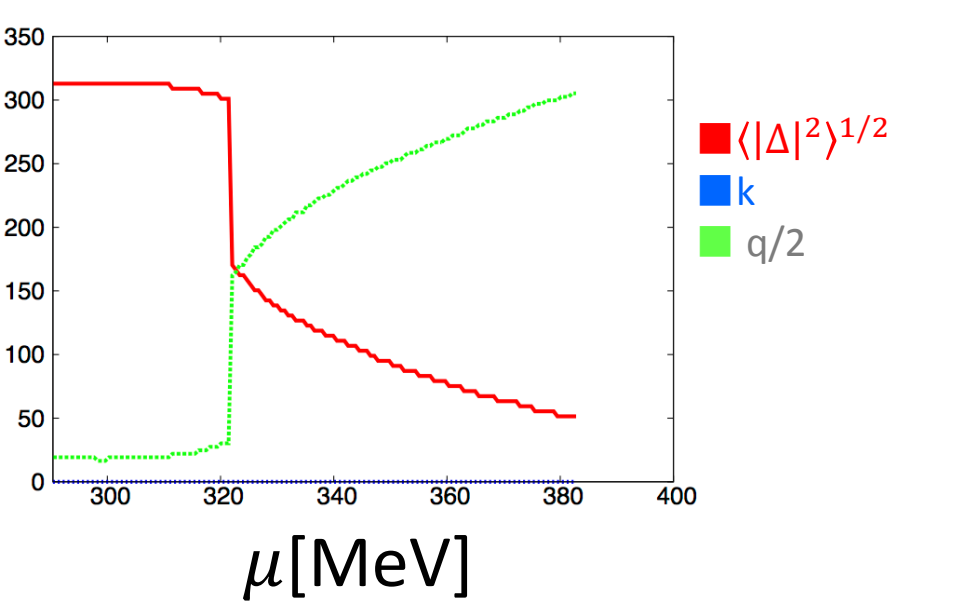
(a) $B = 0$



(b) $\sqrt{eB} = 70 \text{ MeV}$
($\sim 5 \times 10^{16}$ Gauss)



(c) $\sqrt{eB} = 120 \text{ MeV}$
($\sim 1.4 \times 10^{17}$ Gauss)



k is wavenumber of amplitude modulation
 $k = m / (2(1 + \sqrt{\nu})K(\nu))$

Summary and Outlook

- DCDW is favored in strong B due to spectral asymmetry
- Twisted kink crystal appears in weak B