

New horizons for MCAS: heavier masses and α -particle scattering

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and collaborators

Outline

The MCAS formulation, brief reviews: K. Amos, *et al*,
Nuclear Physics, **A728** (2003) 65; **A912** (2013) 7

Mirror nuclei method

Application of MCAS method to ^{16}O ,
using the vibrational model

Final results for Nucleon + ^{16}O

Nucleon + mass-18 preliminary results

α -particle scattering from even-mass nuclei

MCAS:

Multichannel Algebraic Scattering Formalism

1. Discretization of the coupled-channel equations by separable expansion of the channel interactions.
2. Pauli principle inclusion by use of orthogonalizing pseudo-potentials.
3. Fast, effective search procedure for resonances and bound states.
4. Can use rotational or vibrational models for the structure of the target nucleus.

Why mirror nuclei

- Two nuclei are called “mirror nuclei” if one changes into the other by interchanging all protons and neutrons
- Example: ^{14}C , the isotope of carbon used in carbon dating (half-life \approx 5700 years) has as its mirror ^{14}O , a short-lived isotope of oxygen (half life = 70.6 sec)
- Nuclei with a proton excess tend to be less stable than those with a neutron excess
- Current MCAS role: analyze bound and resonant spectra to support and interpret experimental work

Nucleon – ^{16}O scattering

The mirror concept cannot be used to get information on ^{16}O , since it is its own mirror.

However, energy levels of ^{16}O , as well as those of ^{17}O , the compound system of $n+^{16}\text{O}$, are well known.

So, we carry out MCAS calculations on $n+^{16}\text{O}$ scattering, to get accurate fits to the spectrum, including resonant states, of ^{17}O . From these we extract neutron scattering “data”.

Using the same parameters, but adding a Coulomb force, we obtain a spectrum for ^{17}F , as well as proton scattering cross sections. ^{17}F is the mirror system to ^{17}O .

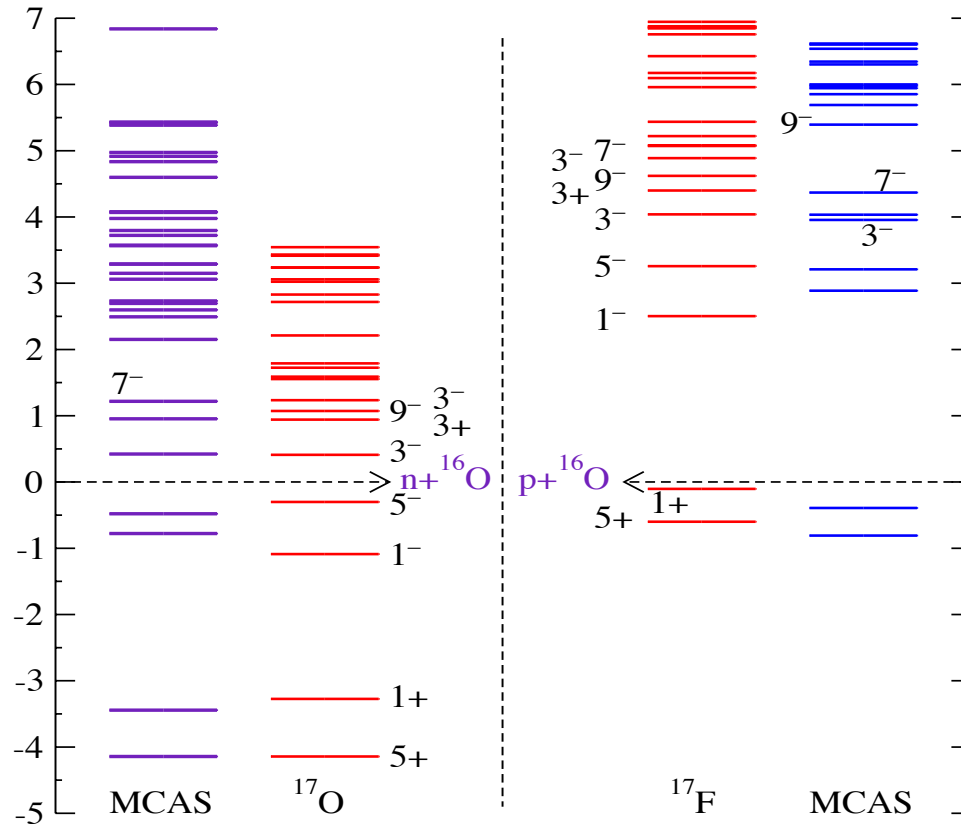
Difficulties with MCAS for ^{16}O

- ^{16}O is a doubly-magic nucleus: 8 protons in the $0s_{1/2}$, $0p_{3/2}$ and $0p_{1/2}$ states: filled s and p shells.
- That means ^{16}O is spherical in its ground state, which causes difficulties for us, since we need the assumption of a deformed target nucleus to which the incoming neutron or proton is coupled. The rotational model does not work well with a spherical ground state.
- So, instead, we use the vibrational model, for the first time with MCAS. Now the deformation is dynamic, and coupling to the projectile works better.
- Results shown here are obtained with the vibrational model. This is a more complicated model, and obtaining good results has required much work.

The parameters

MeV	parity -	parity +	geometry	value	Coulomb
V_0	- 47.15	- 50.6	R_0	3.15 fm	2.608 fm
V_{LL}	2.55	0.0	a	0.65 fm	0.513 fm
V_{Ls}	6.9	7.2	β_2	0.21	w= 0.051
V_{ss}	2.5	-2.0	β_3	0.42	
I_n^π	E_n (MeV)	$0s_{1/2}$	$0p_{3/2}$	$0p_{1/2}$	$0d_{5/2}$
0^+_1	0.0	10^6	10^6	10^6	0.0
0^+_2	6.049	10^6	10^6	0.0	0.0
3^-_1	6.13	10^6	10^6	5.0	0.0
2^+_1	6.92	10^6	10^6	0.0	0.0
1^-_1	7.12	10^6	10^6	5.0	1.0

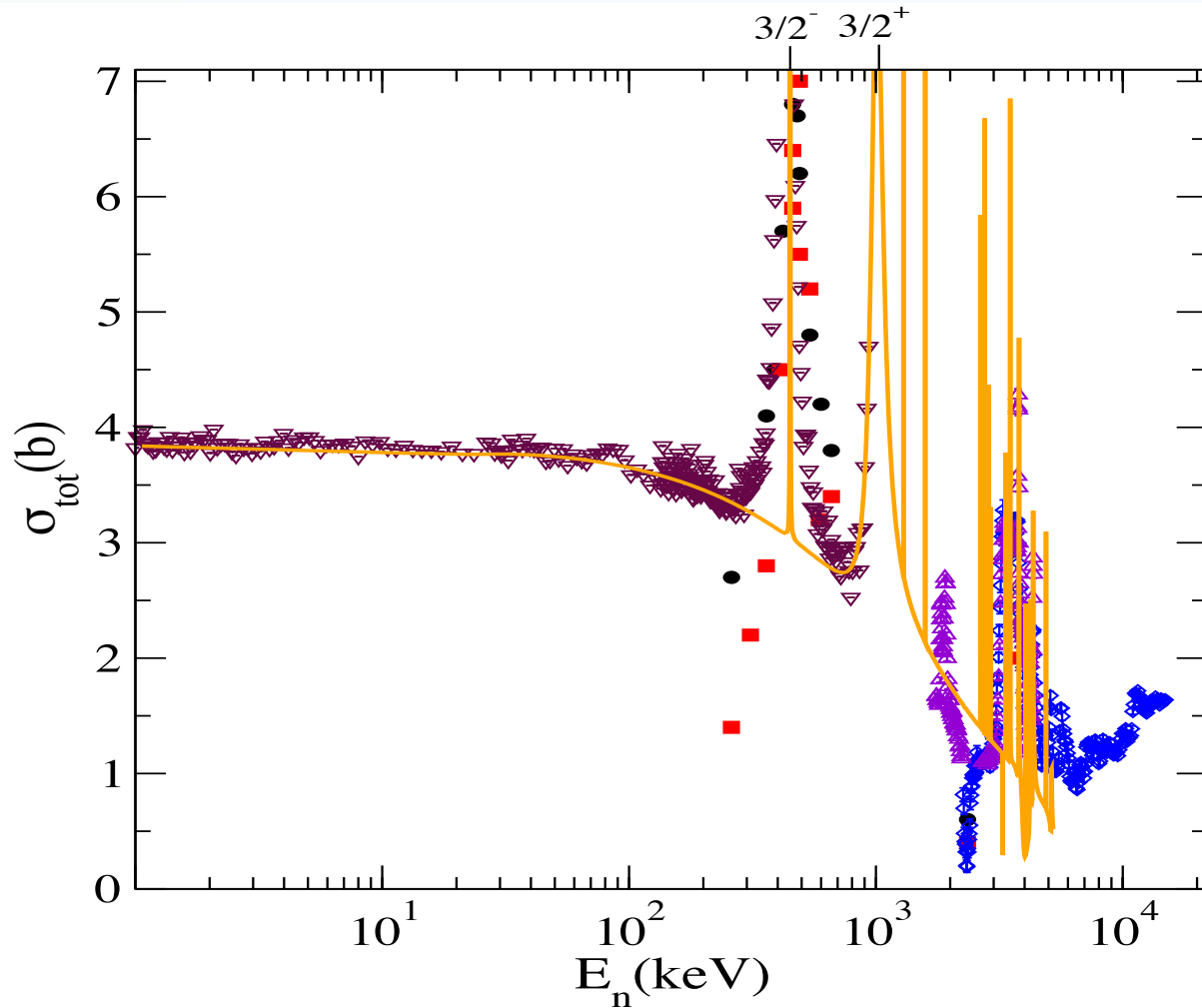
Spectra of ^{17}O and ^{17}F



The 10 lowest-E states

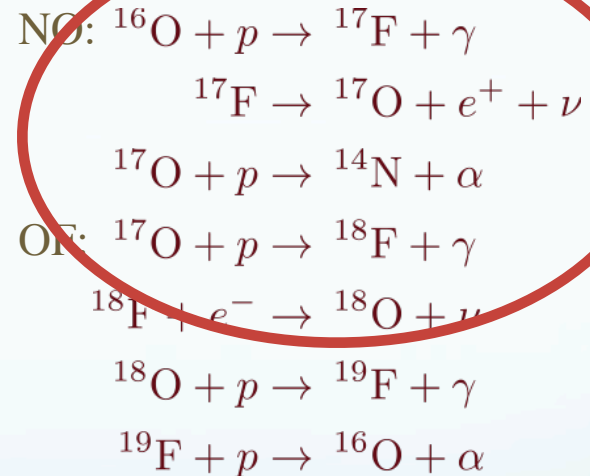
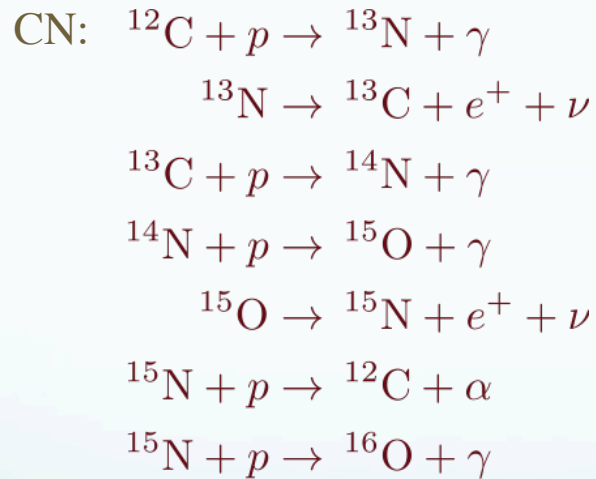
J^π	$^{17}\text{O}_{\text{exp}}$	Γ_{exp}	E_{mcas}	Γ_{mcas}	$^{17}\text{F}_{\text{exp}}$	Γ_{exp}	E_{mcas}	Γ_{mcas}
(5/2)+	-4.1436	—	-4.1432	—	-0.6005	—	-0.8079	—
(1/2)+	-3.2729	—	-3.4426	—	-0.1052	—	-0.3927	—
(1/2)-	-1.0882	—	-0.7781	—	2.5035	19	2.8874	$5.58(10)^{-5}$
(5/2)-	-0.3008	—	-0.4732	—	3.2565	1.5	2.5644	$9.80(10)^{-6}$
(3/2)-	0.4102	40	0.42264	1.277	4.0395	225	3.2104	0.00552
(3/2)+	0.9412	96	0.9534	129	4.3995	1.530	3.9557	0.906
(9/2)-	1.0722	< 0.1	2.1528	$1.08(10)^{-7}$	4.6195	—	5.3930	$1.26(10)^{-9}$
(3/2)-	1.2356	28	2.7332	0.2923	4.8875	68	5.826	$6.8(10)^{-5}$
(7/2)-	1.5537	3.4	1.2185	0.1615	5.0715	40	4.3679	$1.95(10)^{-3}$
(5/2)-	1.5892	<1	3.1504	0.1982	5.0815	< 0.6	6.3027	$6.8(10)^{-4}$

$n+^{16}\text{O}$ total scattering cross section

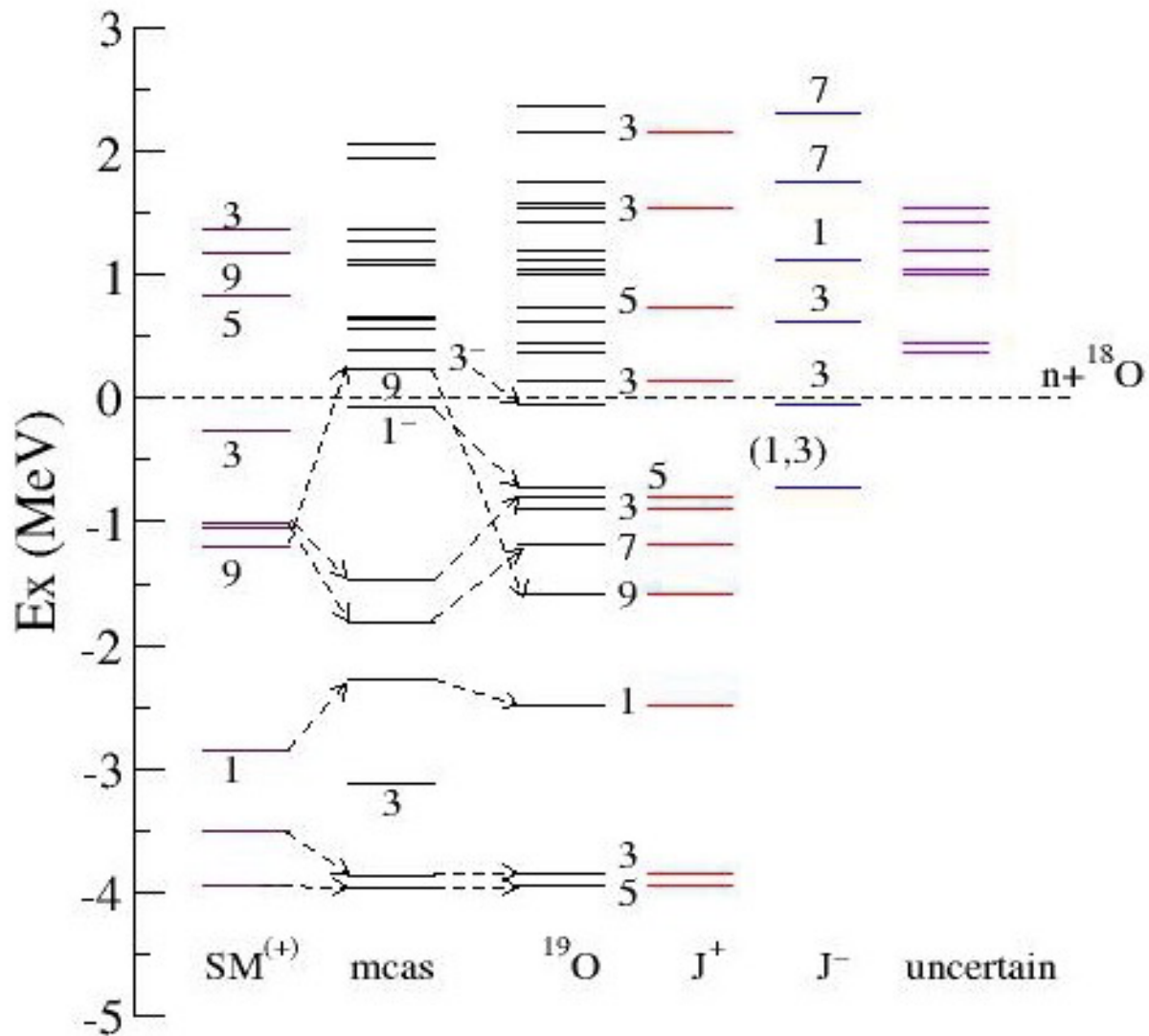


Significance of the mass 17-19 systems

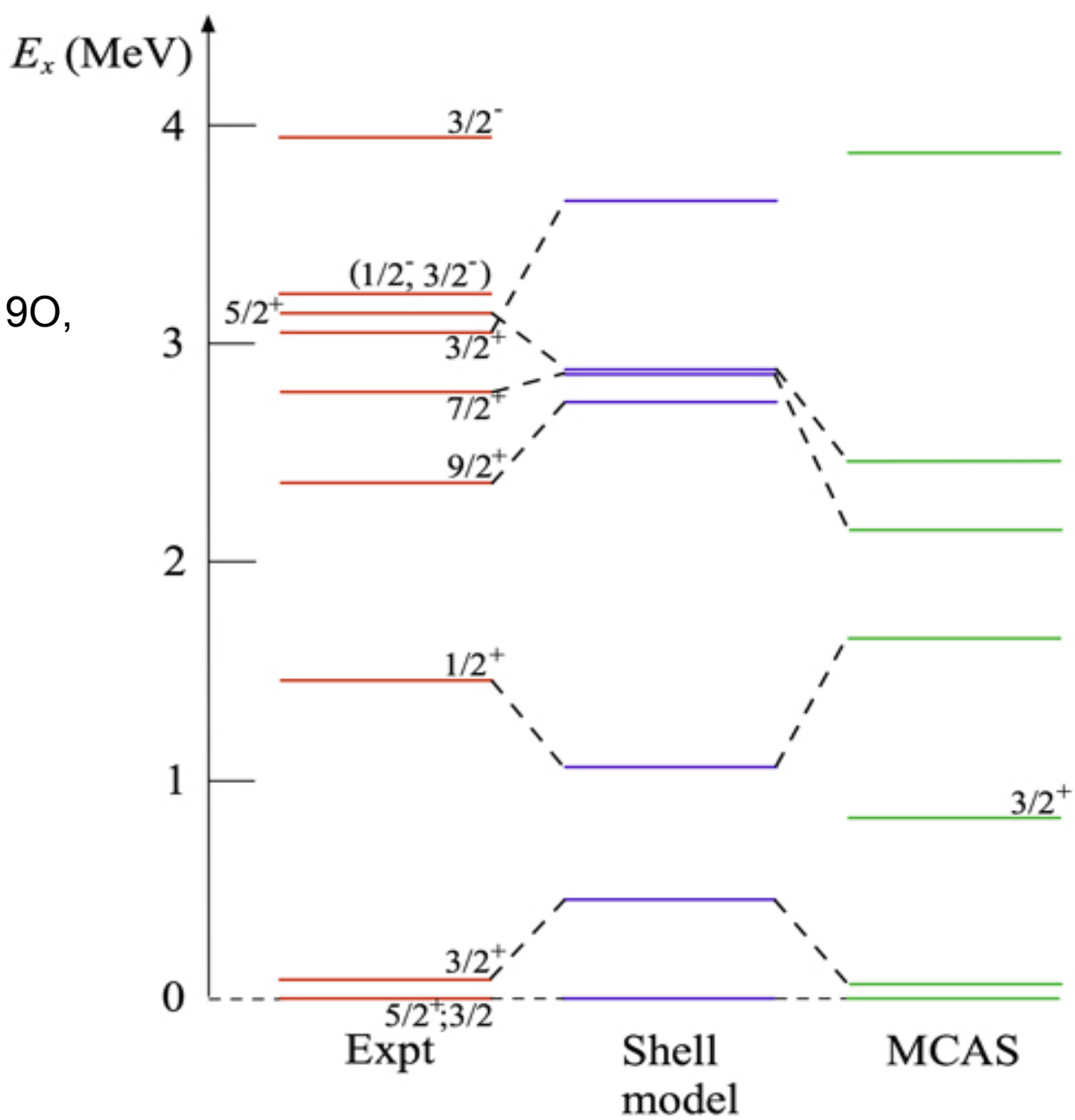
The structures of ^{17}O and ^{17}F are critical in the synthesis of elements, beyond carbon, within the stellar environment. The CNO cycle:



Mass 18 and 19

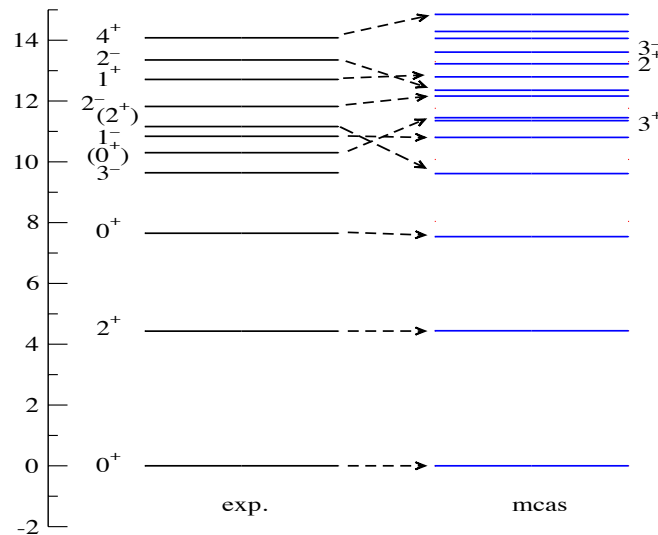


Spectrum of ^{19}O ,
below $n+18\text{O}$
threshold

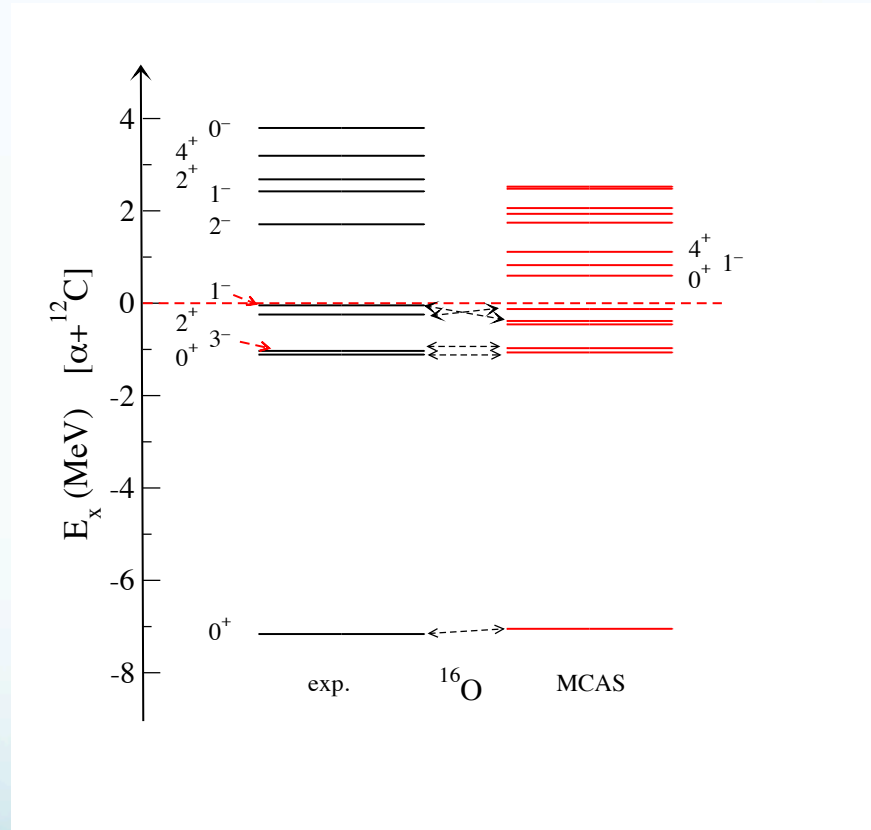


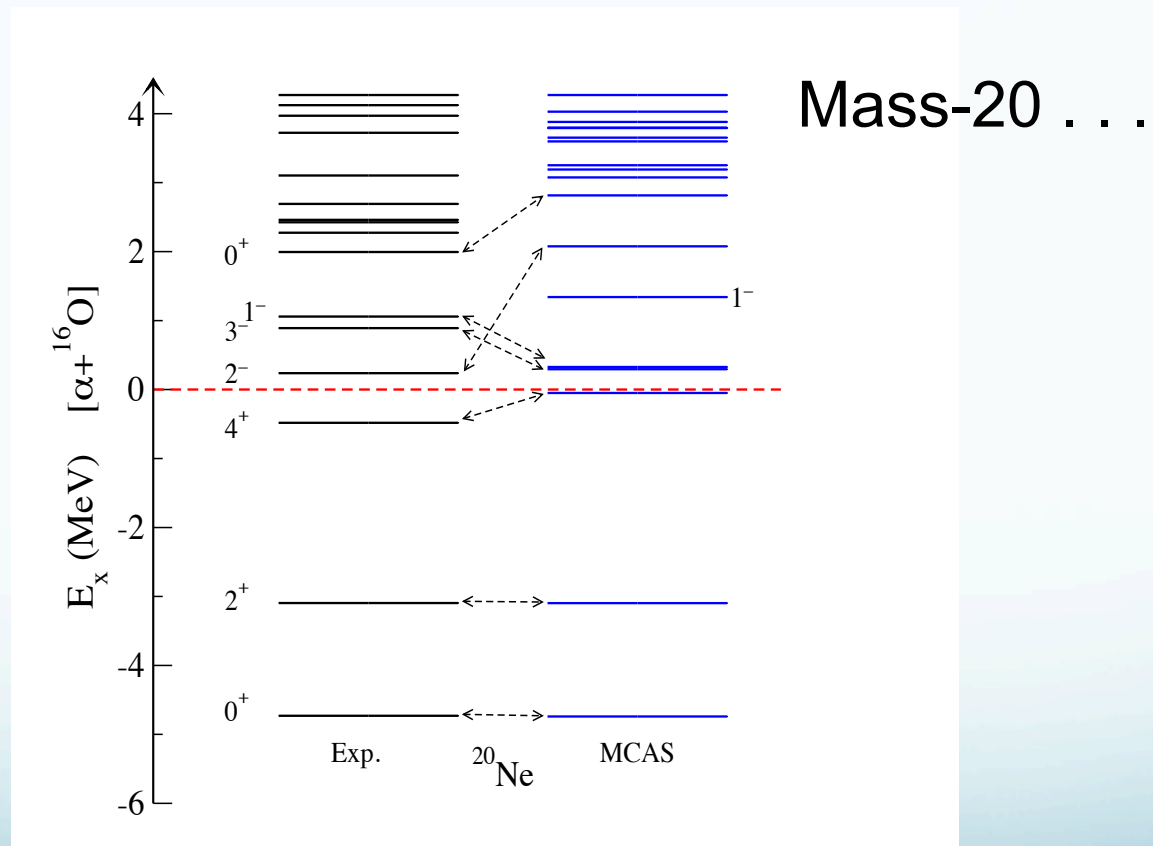
Results for α scattering from even-even nuclei.

Here it is assumed an α particle scattering from the unstable ${}^8\text{Be}$ nucleus fuses to form ${}^{12}\text{C}$.
The levels shown are the lowest levels in ${}^{12}\text{C}$.



Alpha + $^{12}\text{C} \rightarrow ^{16}\text{O}$





... and beyond: mass 22-23

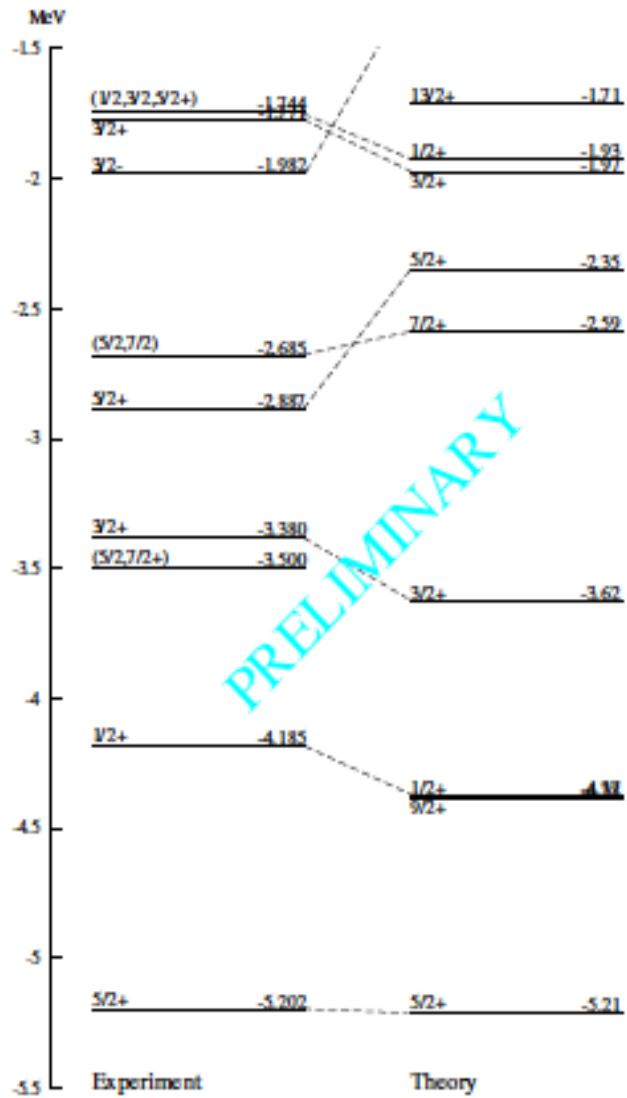


Figure 1. Experimental and MCAS ^{23}Ne spectrum.

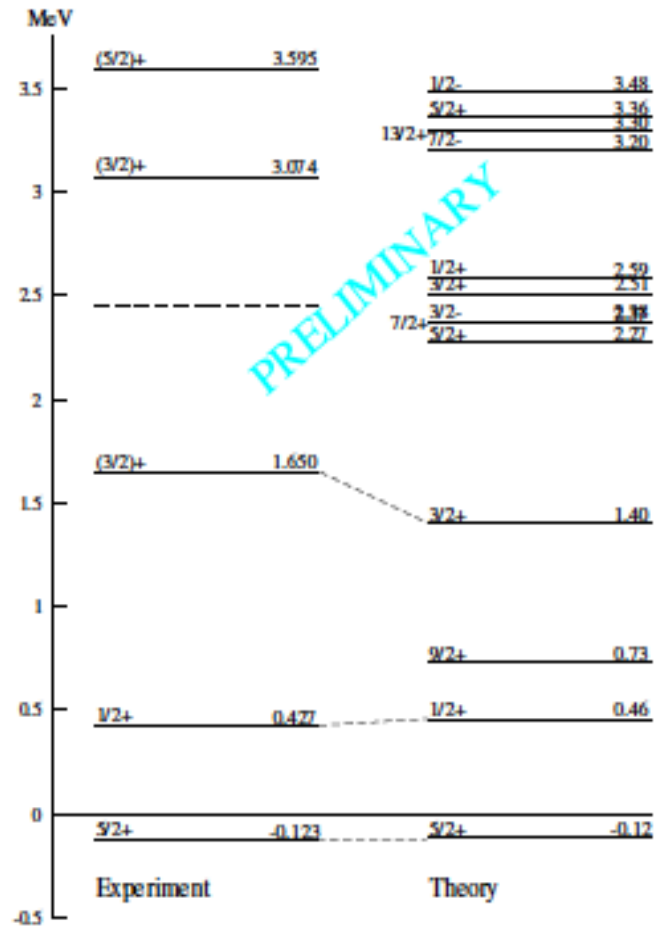
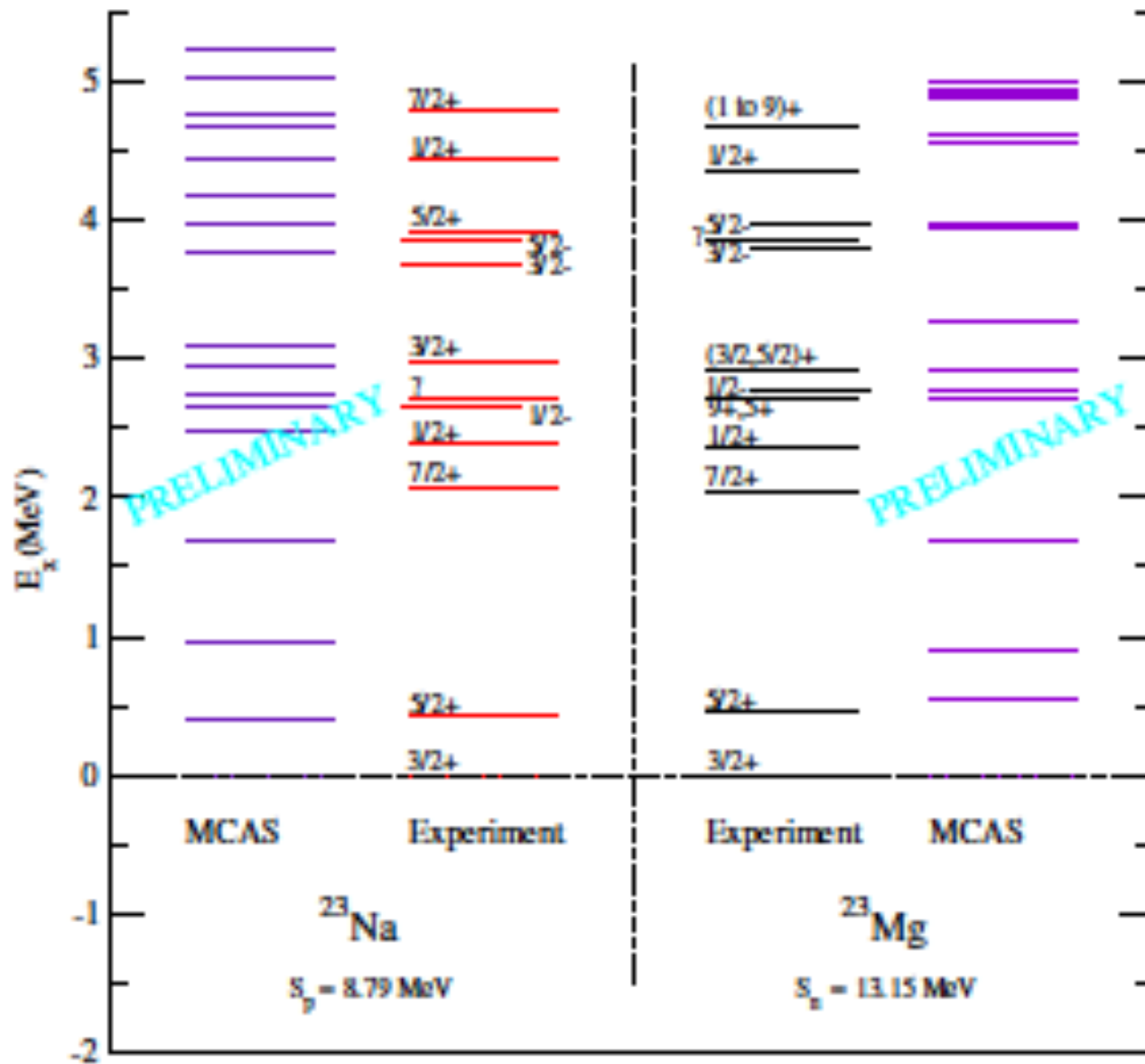


Figure 2. Experimental and MCAS ^{23}Al spectrum.

. . . to mass 23



Concluding remarks

Work on neutron- ^{16}O scattering is near completion and results are in good agreement with experimental data.

Proton- ^{16}O scattering calculations are also well advanced, and a spectrum of ^{17}F was shown.

Current work involves the mass 18-19 system, and alpha-nucleus scattering.

Work on mass-20 and beyond has commenced; preliminary results are promising.

MCAS Collaboration

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