

A Multiorbital DMFT Analysis of Electron-Hole Asymmetry in the Dynamic Hubbard Model



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The Hubbard Model

$$H = - \sum_{i,j,\sigma} t_{ij} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i c_{i\uparrow}^\dagger c_{i\uparrow} c_{i\downarrow}^\dagger c_{i\downarrow}$$

- restrict i and j to nearest neighbour lattice sites
- Pauli exclusion allows only two electrons per site
- U – double occupancy Coulomb repulsion
- t_{ij} – nearest-neighbour hopping

- single band model
- electron-hole symmetric

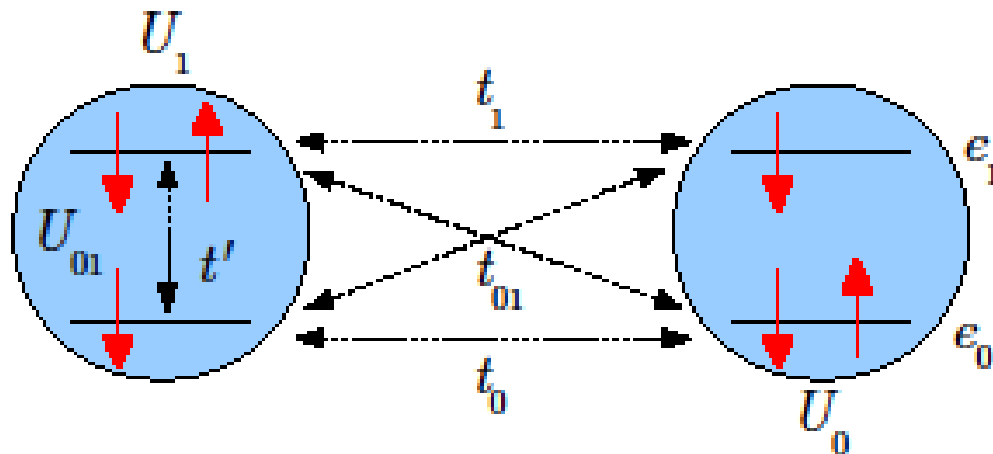
Double Occupancy and Orbital Relaxation

- The Hubbard model assumes a single orbital on each lattice site and an electron's state is static regardless of occupancy.
- J. E. Hirsch, Phys. Rev. B **65**, 184502 (2002):
The real electronic ground state includes higher-orbital contributions with weaker Coulomb repulsion which become especially important for strongly-correlated systems (large local Coulomb repulsion) at high filling
- Need to adjust the Hubbard model to capture the flexibility for electrons to change their state in response to changes in occupancy: dynamic Hubbard model

Dynamic Hubbard Model (DHM)

J. E. Hirsch, Phys. Rev. B **65**, 184502 (2002)

- Two non-degenerate orbitals: energies $\epsilon_0 < \epsilon_1$
- Three local Coulomb repulsions U_0, U_1, U_{01}
- Two intraband hopping parameters t_0, t_1
- Nonlocal hybridization (interband hopping) t_{01}
- Local interband hybridization t'



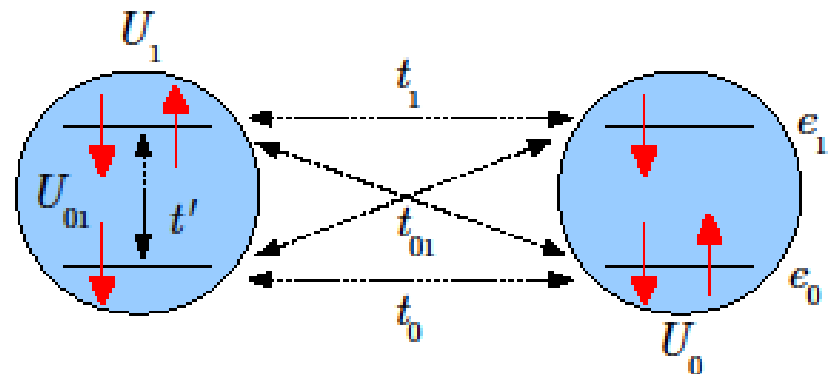
DHM Hamiltonian

$$H_{\text{DHMe}} = H_{\text{KE}} + H_{\text{local}} + H_{\text{hyb}}$$

$$H_{\text{KE}} = - \sum_{\langle ij \rangle \sigma} \left[t_0 (c_{0i\sigma}^\dagger c_{0j\sigma} + c_{0j\sigma}^\dagger c_{0i\sigma}) + t_1 (c_{1i\sigma}^\dagger c_{1j\sigma} + c_{1j\sigma}^\dagger c_{1i\sigma}) \right. \\ \left. + t_{01} (c_{1i\sigma}^\dagger c_{0j\sigma} + c_{0j\sigma}^\dagger c_{1i\sigma}) \right]$$

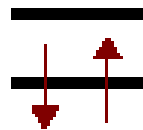
$$H_{\text{local}} = U_0 \sum_i n_{0i\uparrow} n_{0i\downarrow} + U_1 \sum_i n_{1i\uparrow} n_{1i\downarrow} + U_{01} \sum_{i\sigma\sigma'} n_{0i\sigma} n_{1i\sigma'} + (\epsilon_0 - \mu) \sum_{i\sigma} n_{0\sigma} \\ + (\epsilon_1 - \mu) \sum_{i\sigma} n_{1\sigma}$$

$$H_{\text{hyb}} = -t' \sum_{i\sigma} (c_{0i\sigma}^\dagger c_{1i\sigma} + c_{1i\sigma}^\dagger c_{0i\sigma})$$

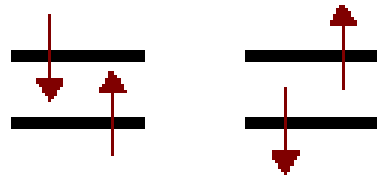


Orbital Relaxation in the DHM

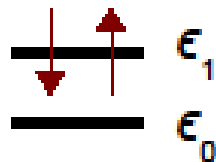
- Double occupancy energy-ordering conditions



$$E_2 = U_0 + 2\epsilon_0$$



$$E_1 = U_{01} + \epsilon_0 + \epsilon_1$$

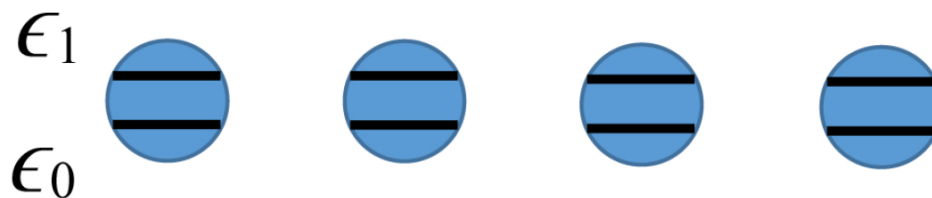


$$E_0 = U_1 + 2\epsilon_1$$

$$U_1 + 2\epsilon_1 < U_{01} + \epsilon_0 + \epsilon_1 < U_0 + 2\epsilon_0$$

Comparison: Four-Site Exact Diagonalization

- J. E. Hirsch, Phys. Rev. B **67**, 035103 (2003)

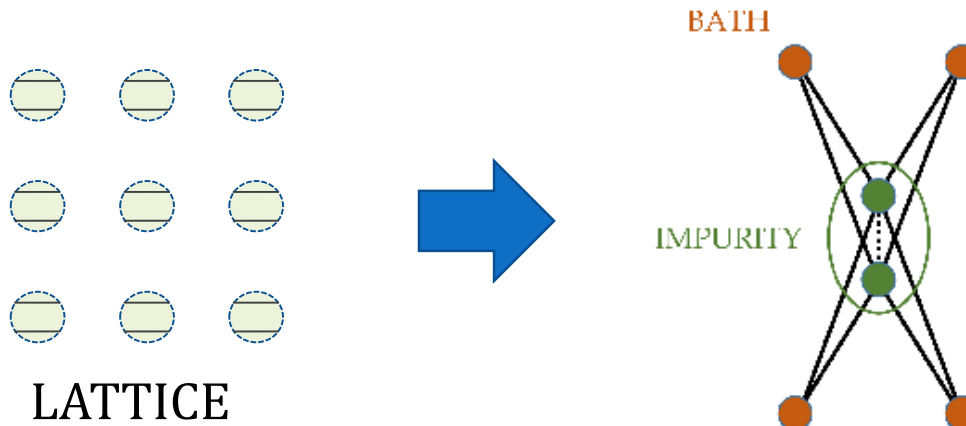


- Main result: electron-hole asymmetry in the Dynamic Hubbard Model
- fixed values of $t' = 0.2$, $t_{01} = 1.0 = t_0 = t_1$

Multiorbital Dynamical Mean Field Theory (MODMFT)

A. Georges, G. Kotliar, W. Krauth, and M. J. Rozenberg, Rev. Mod. Phys. **68**, 13 (1996)

- Maps an infinite-dimensional lattice model onto a local impurity model
- Solve the impurity model self-consistently for a set of effective mean field parameters which approximate the influence of the full lattice environment on a single site
- Retains the local dynamics of electronic occupancy of the impurity, yielding the Green's function and self energy of the system

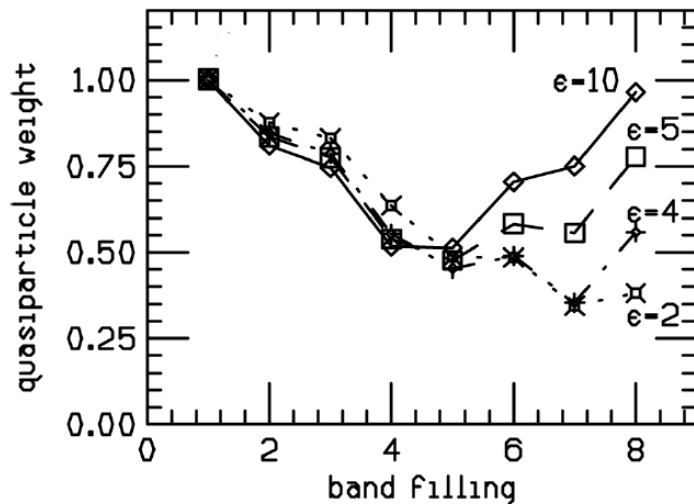


Result: Observed Asymmetry in Z

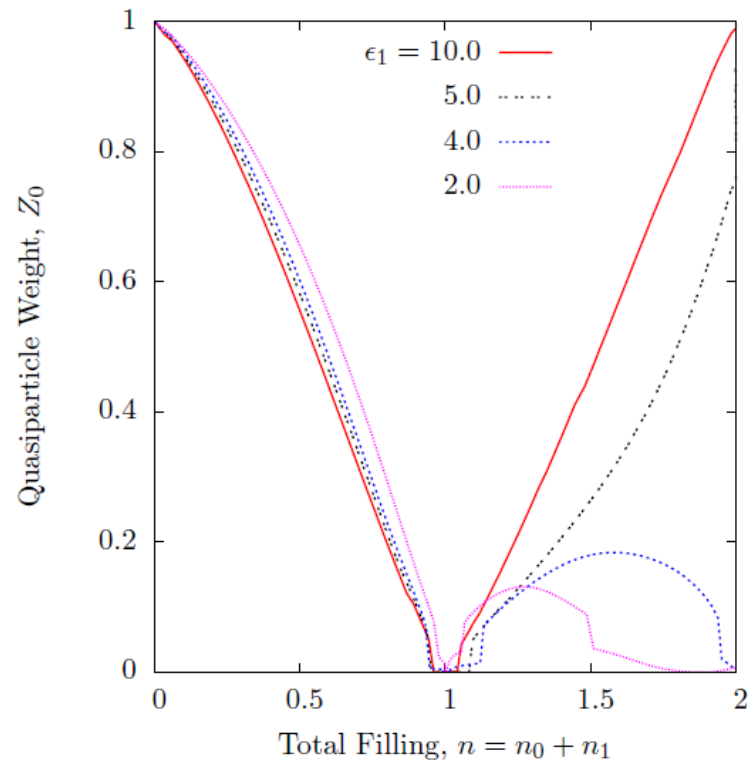
$$U_0 = 10.0, U_{01} = 6.0, U_1 = 5.0, t_0 = t_1 = t_{01} = 1.0, t' = 0.2$$

Quasiparticles become increasingly dressed with orbital relaxation.

Four Site ED



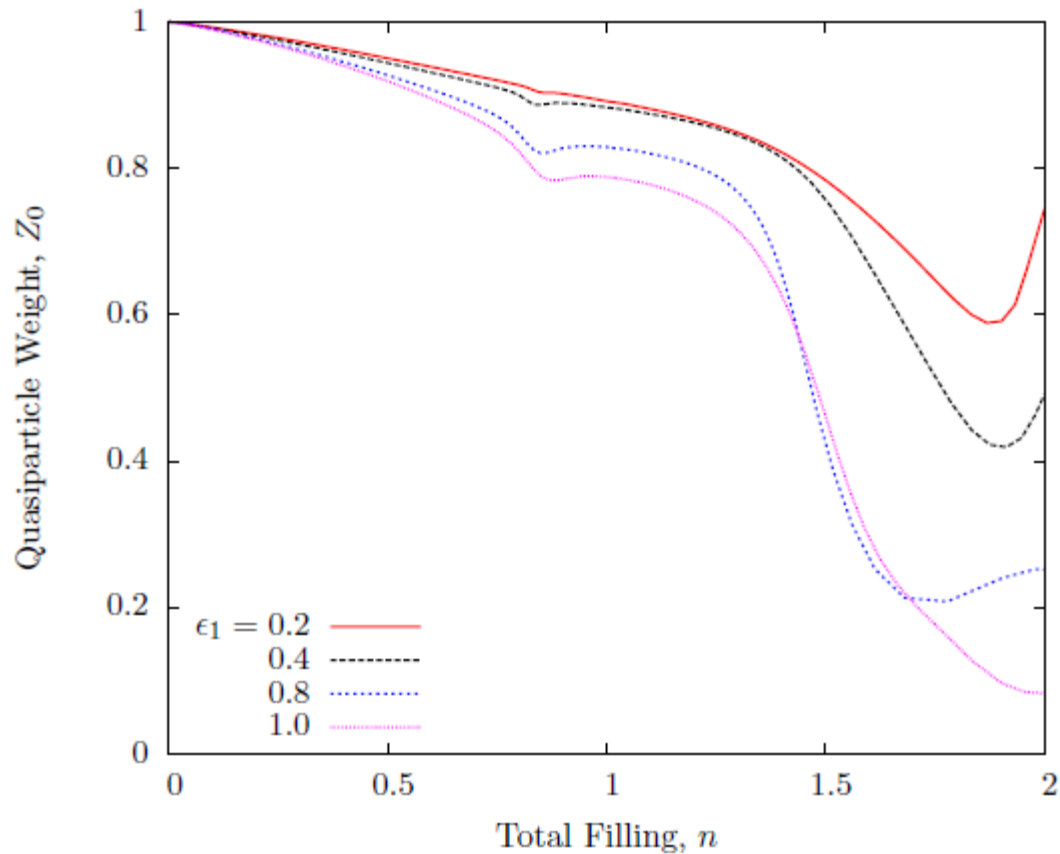
Bethe Lattice MODMFT



Or the Opposite Effect...

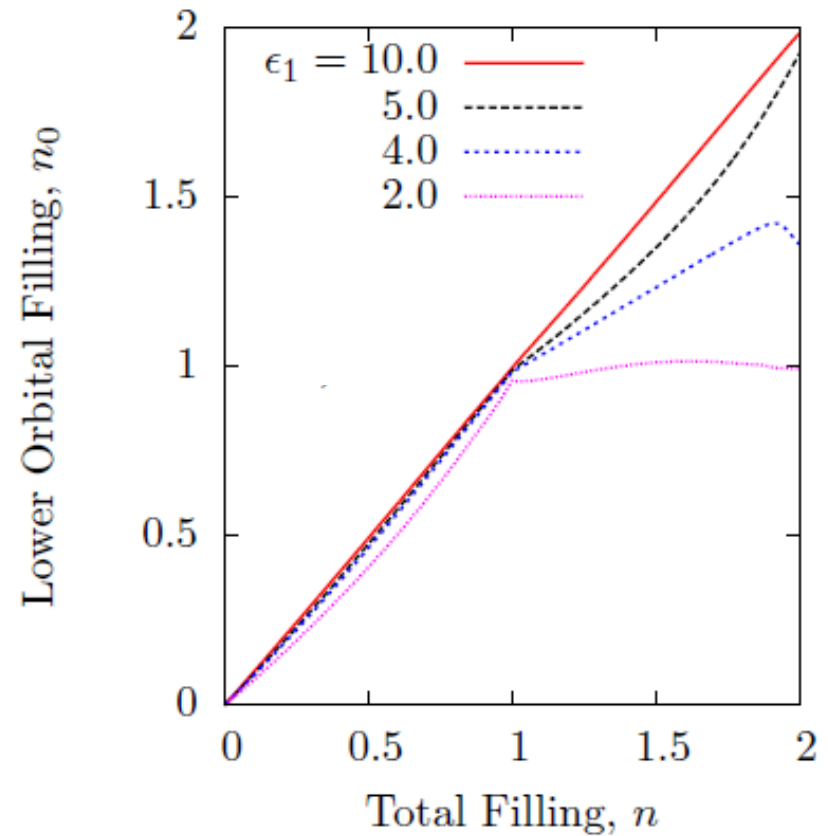
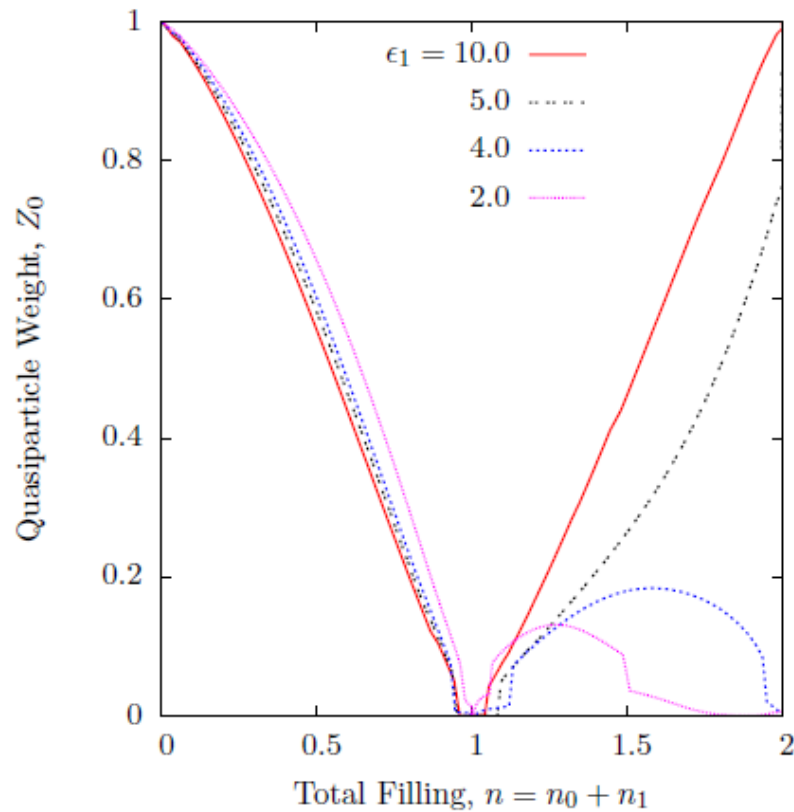
$$U_0 = 3.0, U_{01} = 2.0, U_1 = 1.0, t_0 = t_1 = t_{01} = 1.0, t' = 0.2$$

Quasiparticles can also *undress* with orbital relaxation.



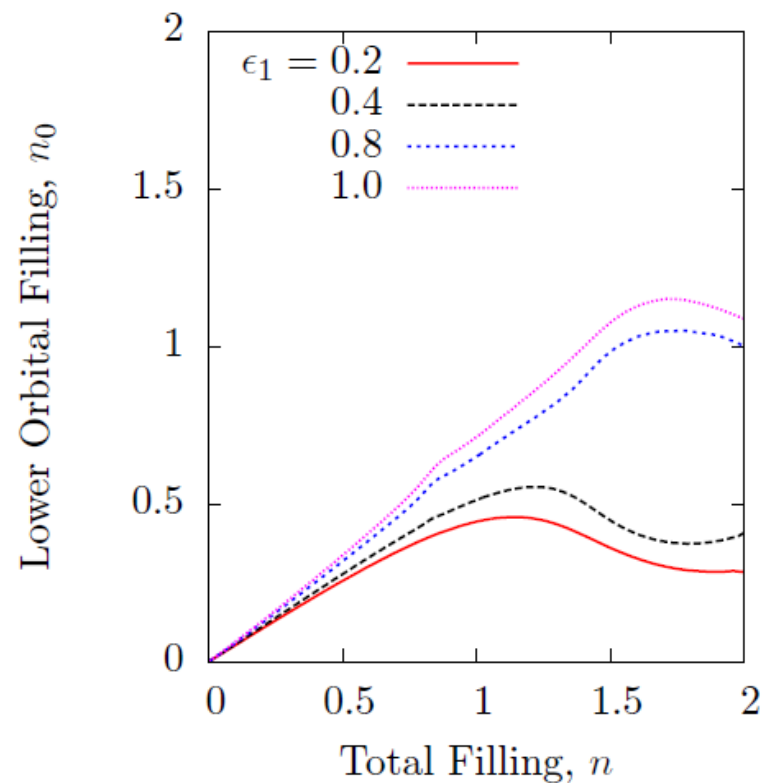
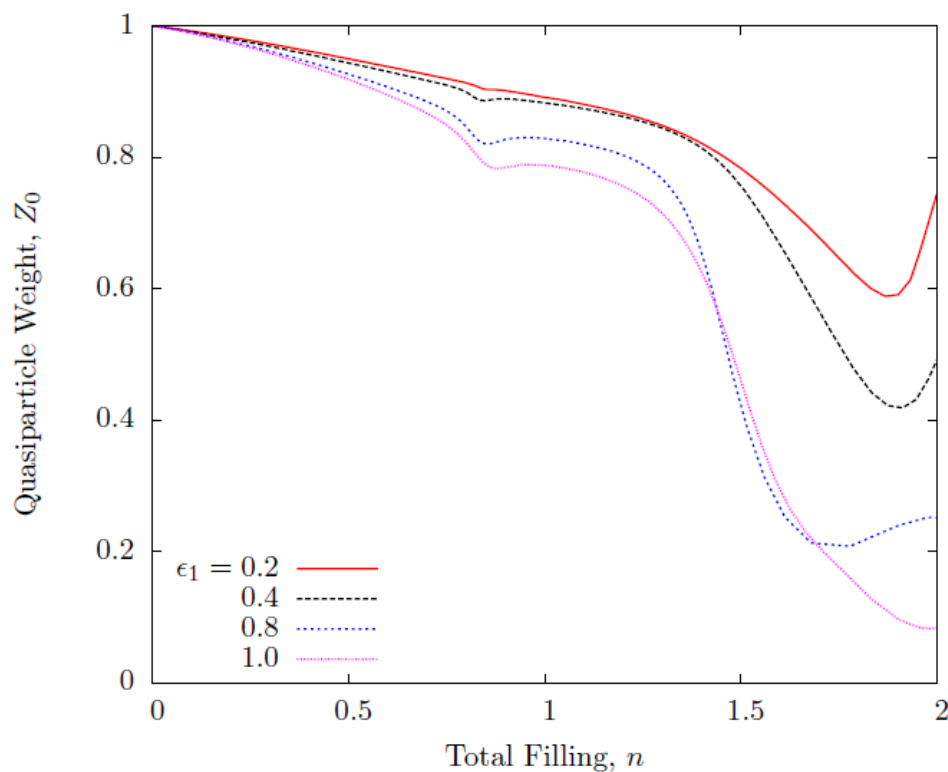
The Influence of Mott Physics on Dressing (Not Evaluated in Hirsch's ED Study)

$$U_0 = 10.0, U_{01} = 6.0, U_1 = 5.0, t_0 = t_1 = t_{01} = 1.0, t' = 0.2$$



The Influence of Mott Physics on Undressing

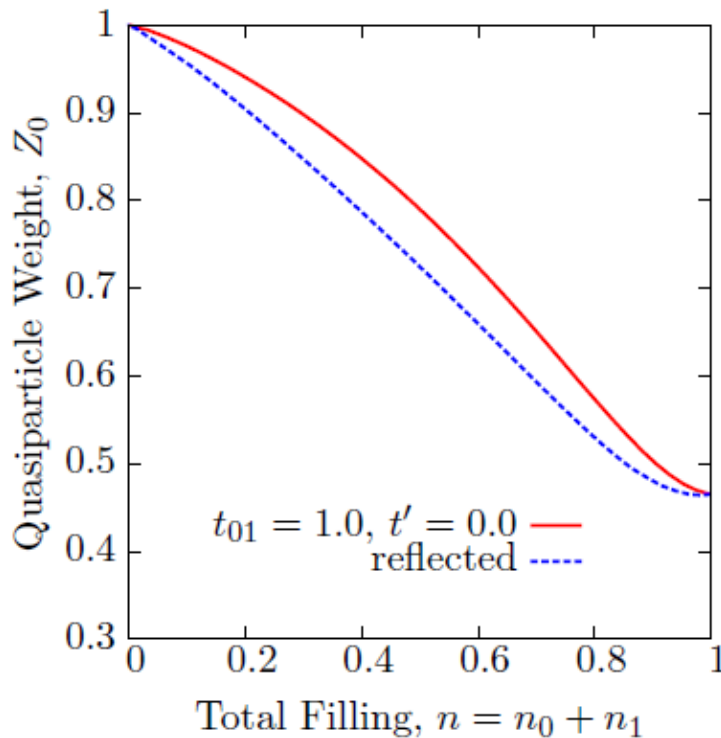
$$U_0 = 3.0, U_{01} = 2.0, U_1 = 1.0, t_0 = t_1 = t_{01} = 1.0, t' = 0.2$$



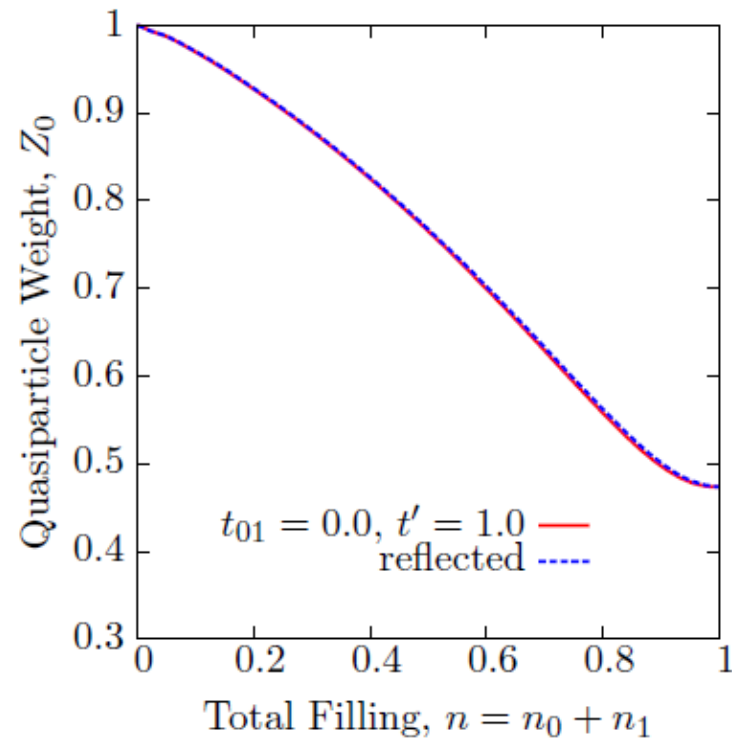
The Influence of Hybridization

- t_{01} is qualitatively more relevant to the physics of orbital relaxation than the (local) t' hybridization parameter. For example:

$$U_0 = 3.0, U_{01} = 2.0, U_1 = 1.0, t_0 = t_1 = t_{01} = 1.0, t' = 0.2, \epsilon_1 = 10.0$$



(a) Nonlocal hybridization.

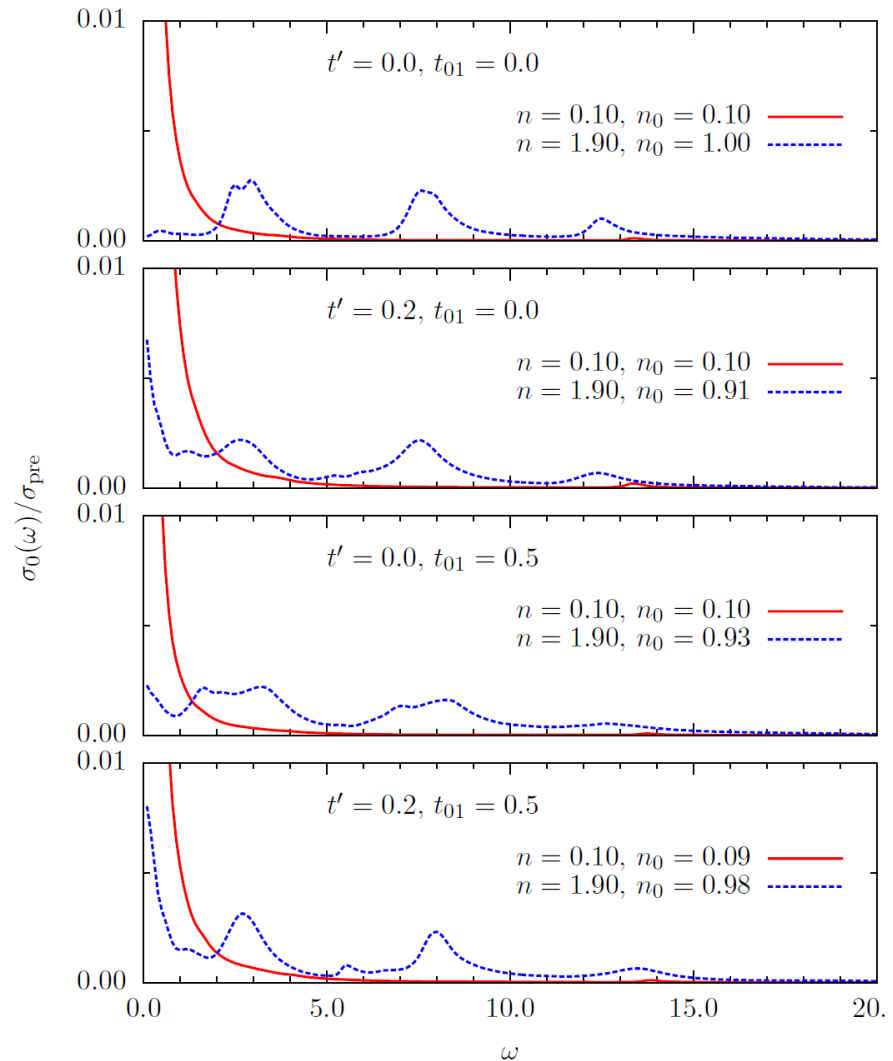


(b) Local hybridization.

Asymmetry Evidenced in Optical Conductivity Weight Transfer

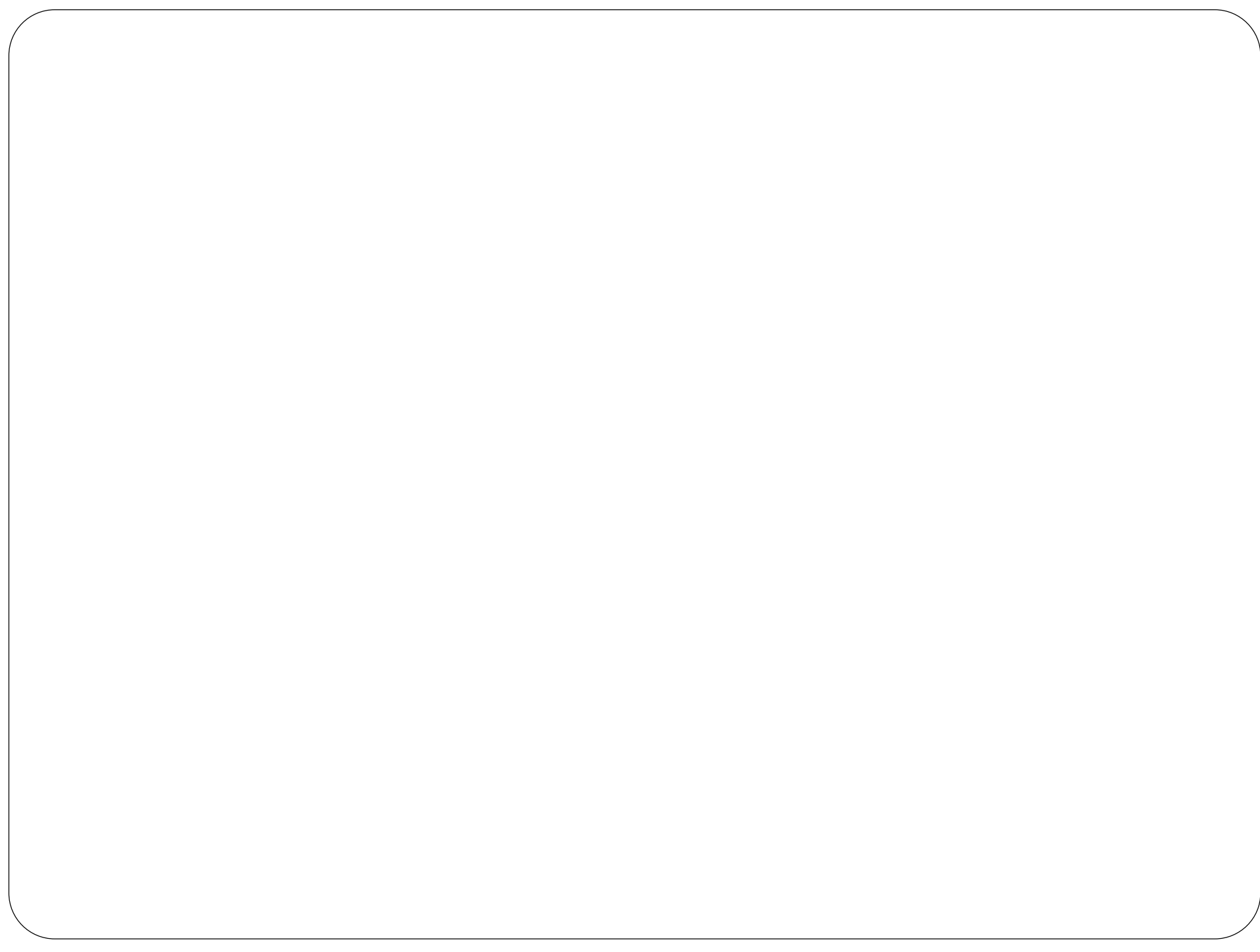
$$U_0 = 10.0, U_{01} = 5.0, U_1 = 0.5,$$
$$t_0 = t_1 = 1.0, \epsilon_1 = 4.0, \eta = 0.1$$

- Hole regime shows transfer of low energy to higher energy features: electron-hole asymmetry
- Significant effect of hybridization on the low energy Drude region



Conclusions

- Confirmed Hirsch's four-site ED observation of electron-hole asymmetry in the dynamic Hubbard model
 - in the quasiparticle weight
 - in optical conductivity weight transfer
- Nonlocal hybridization is qualitatively more important than local hybridization
- Complicated dependence of orbital relaxation on the energy gap, hybridization values and Mott physics in the DHM



MODMFT Background

- MODMFT has been in use since the earliest years of DMFT studies
- Q. Si and G. Kotliar, Phys. Rev. Lett. **70**, 3143 (1993)
- Q. Si and G. Kotliar, Phys, Rev. B **48**, 13881 (1993)

- Benchmark: A. Liebsch and H. Ishida, J. Phys.-Condens. Mat. **24**, 053201 (2012)

- Several studies of two-orbital systems with local hybridization t'
- Few with nonlocal hybridization t_{01}
- Focus has been on orbital selective Mott transitions with Hund's coupling; none appear to address the dynamic Hubbard model