Probing the Nature of Inflation

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CMB History

Inflation and the CMB

Inflation today

Inflation tomorrow

CMB History

The CMB is todays workhorse of cosmology



What we observe is the light the leftover after neutral hydrogen formed, 378 000 years after the big bang

This year marks the 50th anniversary of the CMB



A MEASUREMENT OF EXCESS ANTENNA TEMPERATURE AT 4080 Mc/s

Measurements of the effective zenith noise temperature of the 20-foot horn-reflector antenna (Crawford, Hogg, and Hunt 1961) at the Crawford Hill Laboratory, Holmdel, New Jersey, at 4080 Mc/s have yielded a value about 3.5° K higher than expected. This excess temperature is, within the limits of our observations, isotropic, unpolarized, and

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Note added in proof.—The highest frequency at which the background temperature of the sky had been measured previously was 404 Mc/s (Pauliny-Toth and Shakeshaft 1962), where a minimum temperature of 16° K was observed. Combining this value with our result, we find that the average spectrum of the background radiation over this frequency range can be no steeper than $\lambda^{0.7}$. This clearly eliminates the possibility that the radiation we observe is due to radio sources of types known to exist, since in this event, the spectrum would have to be very much steeper.

A. A. PENZIAS R. W. WILSON

May 13, 1965 Bell Telephone Laboratories, Inc Crawford Hill, Holmdel, New Jersey

CMB was first observed by Penzias & Wilson (1965)



Consistent with a uniform temperature $\ T\sim 3\,{ m K}$

CMB dipole was first observed by Conklin (1969)



Temperature fluctuation of $\Delta T \sim 3 \,\mathrm{mK}$

CMB anisotropy observed by COBE (1992)



First detection of primordial density fluctuations

First acoustic peak observed (1993–2001)



Consistent with spatial flat universe

First acoustic peak(s) observed (1993-2001)



Consistent with spatial flat universe

The CMB today

WMAP (2003-2012)



Cosmic variance limited to $\ell = 548$

The CMB today

Planck (2013-present)



Cosmic variance limited to $\ell \sim 1500$

The CMB today

Data (today):



Precision measurement cosmological parameters

Inflation and the CMB

What we see is a snap-shot of the sound waves



Acoustic peaks show that they are in-phase



Phase coherence is a stringent requirement of CMB

Any local source will have arbitrary phases

$$\frac{\delta T}{T} \sim A_k \cos kr_s + B_k \sin kr_s$$

Observed power spectrum requires $B_k = 0$

However, if mode existed outside the "horizon"

$$\frac{k}{aH} \ll 1 \to B_k \propto a^{-3} \to 0$$

Phase coherence "prove" modes are super-horizon

In a decelerating universe, this is hard to achieve

$$\partial_t(aH) = \ddot{a} < 0$$

Physical wavelengths only decrease w.r.t Hubble

Two options:

(1) Non-local production of fluctuations

(2) Change the matter content of the universe

A definition:

1. A period of quasi-dS expansion Guth



During inflation, fluctuations stretched

$$a \sim a_0 e^{Ht} \qquad \frac{k}{aH} \to 0$$

Long wavelengths evolve from short wavelengths

Production of fluctuations can be local



Even a cosmological constant has this effect

Inflation also requires that the phase ends



We must get the hot "big bang" eventually

A definition:

- 2. A physical clock Linde; Albrecht & Steinhardt
- "End of inflation" needs a physical definition
- Inflation must end everywhere at the same "time"
- Different regions synched their clocks in the past

Slow-roll Inflation



Slow-roll Inflation



End of inflation defined by value of field Afterwards, energy converted to radiation No clock is perfect (uncertainty principle)

The amount of inflation will vary from place to place.

$$\zeta(x) \sim \frac{\delta a(x)}{a} \sim \frac{\dot{a}\delta t(x)}{a} \equiv H\delta t$$

RMS fluctuations of the clock $\sqrt{\langle (\delta t)^2 \rangle} \sim \frac{H}{f_-^2}$

Time between "ticks" defines an energy scale f_{π}

For slow-roll inflation
$$\delta t \sim \frac{\delta \phi}{\dot{\phi}}$$
 and $f_{\pi}^2 = \dot{\phi}$

These are the adiabatic fluctuations Determine CMB temperature fluctuations

$$\frac{\delta T(\mathbf{n})}{T} \sim \int d^3 k F(\mathbf{k} \cdot \mathbf{n}, k) \zeta_{\mathbf{k}}$$

Think of

$$\langle \zeta(x_1)..\zeta(x_n) \rangle \to$$



Conditions have been formalized: EFT of Inflation Creminelli et al.; Cheung et al.

Write a theory directly for the clock

Different models controlled by a few parameters

E.g. We may choose any H(t) as long as

 $|\dot{H}| \ll H(t)^2$

"Predictions" of (single-field) Inflation

On large scales:

- Homogeneous & Isotropic
- Universe is spatially flat

Density fluctuations are:

- Adiabatic (i.e. uniform in all energy densities)
- Nearly scale-invariant
- Gaussian

All of these 'predictions' are now supported by data

Raises the question: what was the clock?

We have lots of ways to make clocks

Slow-roll inflation is easiest to understand (dynamics are very simple)

But, does the data prefer slow-roll inflation?

Raises the question: what was the clock?

Most "scalar fields" are not fundamental particles (e.g. temperature on earth)

I.e. Could the clock be an emergent phenomena, not at fundamental scalar field?

How would we tell?

Little non-gaussanity in (single-field) slow-roll Creminelli

Non-linearity controlled by a new scale: Λ

Slow-roll is very linear at energy $\dot{\phi}$



Leads to a qualitatively picture:



This level of linearity is not necessary

E.g. K-inflation, DBI inflation,...

Non-linear kinetic energy

Armendáriz-Picón et al., Silverstein & Tong; Alishahiha et al.; ...

K.E. =
$$F(\dot{\phi}^2) = \frac{1}{2}\dot{\phi}^2 + \frac{1}{\Lambda^4}\dot{\phi}^4 + \dots$$

Not even approximately linear for $\Lambda^2 < \dot{\phi} = f_\pi^2$

Closely related to the theory of superfluids

Leads to two qualitatively different pictures: Baumann & DG



Leads to two qualitatively different pictures: Baumann & DG



Leads to two qualitatively different pictures: Baumann & DG



Natural boundary between the pictures Baumann, DG & Porto



What do we know today?

For this picture, 2 numbers matter:

1. Amplitude of the power spectrum

$$\left\langle \left(\frac{\delta T}{T}\right)^2 \right\rangle = 2.2 \times 10^{-9} \sim \frac{H^4}{f_\pi^4}$$

Very well measured from CMB

This fixes one ratio $f_{\pi} = 58 \ H$

For this picture, 2 numbers matter.

1. Amplitude of the power spectrum



For this picture, 2 numbers matter:

2. Amount of non-gaussanity

Typically reported in terms of

$$f_{
m NL}\sim rac{f_\pi^2}{\Lambda^2}$$
 (or $f_{
m NL}\zeta\sim rac{H^2}{\Lambda^2}$)

This is the amplitude for a bispectrum

However, there is no unique "shape" of bispectrum

Planck reports limits on 3 templates.



 $f_{\rm NL}^{\rm local} = 0.8 \pm 5.0$ (68% C.I.)

Planck reports limits on 3 templates.



Planck reports limits on 3 templates.



Can be translated into lower bounds on non-linearity

For single-field, there is no "local" shape Maldacena; Creminelli & Zaldarriaga

Equilateral and Orthogonal are related to two scales.

Planck (68%) $\Lambda_1 > 4.5 H$ $\Lambda_2 > 1.6 H$

Consistent with gaussianity at 10^{-3} level

Still a weak result in terms of scales

For this picture, 2 numbers matter:

2. Amount of non-gaussianity



For this picture, 2 numbers matter:

2. Amount of non-gaussianity



Tensors play a special role

Fixes the overall scale of Hubble



Tensors have major model building implications



Already a strong constraint on many popular ideas

How will we do better?

One goal is to test the full non-slow-roll region



One goal is to test the full non-slow-roll region

In terms of measurable parameters we need

$$f_{\rm NL}^{\rm equilateral} < 1 \ (2\sigma)$$

Best limit today is $\Delta f_{\rm NL}^{\rm equilateral} = 84 \ (2\sigma)$

WMAP to Planck (2015) was a factor of 4 improvement

The brute force approach is to find more "modes"

When each bin is cosmic variance limited

$$\Delta f_{\rm NL} \sim \frac{10^5}{\sqrt{N_{\rm modes}}}$$

E.g. From Planck we get roughly

$$N_{\rm modes, Planck} \sim \ell_{\rm max}^2 \sim 2 \times 10^6$$

To improve by 10^2 we will need 10^{10} modes!

There aren't many more modes in the CMB



Small scales (high ℓ) dominated by foregrounds

Future will be dominated by Large Scale Structure



Future will be dominated by Large Scale Structure

Basic advantage is that there are a lot more modes!

$\underline{Reason \ 1}: LSS \ is \ 3d \ versus \ 2d \ CMB$

For same range of scales $N_{
m modes} \sim 10^9$

<u>Reason 2</u> : Large range of scales

Total number of linear modes $N_{
m modes} \sim 10^{18}$

Large Scale Structure



LSS faces many new challenges

To take advantage of 3d modes we need:

- Very accurate redshifts
- Good model for galaxy formation
- Control of many many new systematics

For these (and other) reasons, no one has actually performed the analysis that will be needed

Summary

Inflation covers a lot more than slow-roll

Ultimately data should decide the correct picture

CMB data today is inconclusive

Large Scale Structure surveys are poised to overtake the CMB in raw sensitivity

CMB will remain vital through search for tensors (and as a probe of the LSS between us and the CMB)

Thank you