

Natural and Unnatural supersymmetry

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Status of Supersymmetry

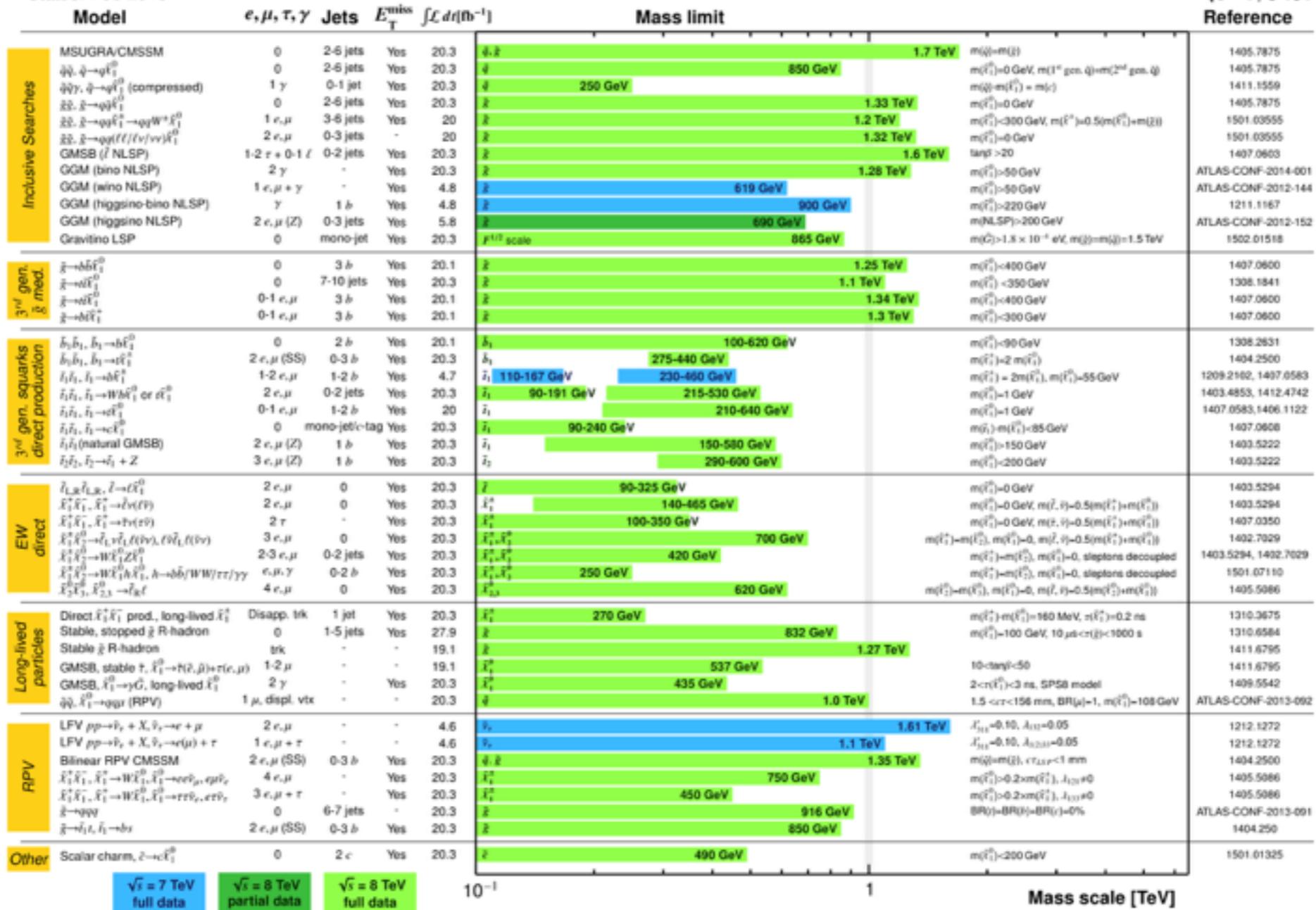
ATLAS SUSY Searches* - 95% CL Lower Limits

Status: Feb 2015

ATLAS Preliminary

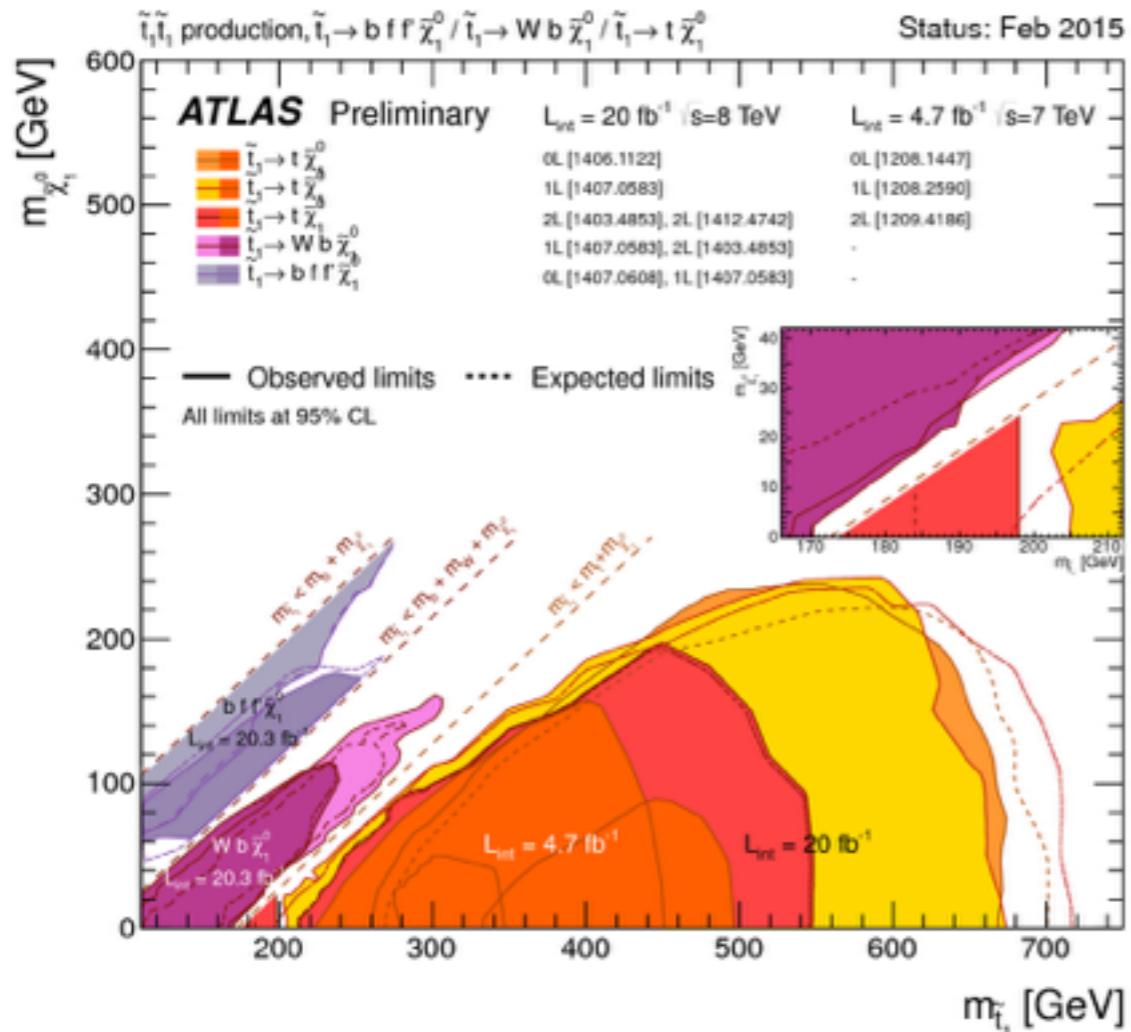
$\sqrt{s} = 7, 8 \text{ TeV}$

Reference

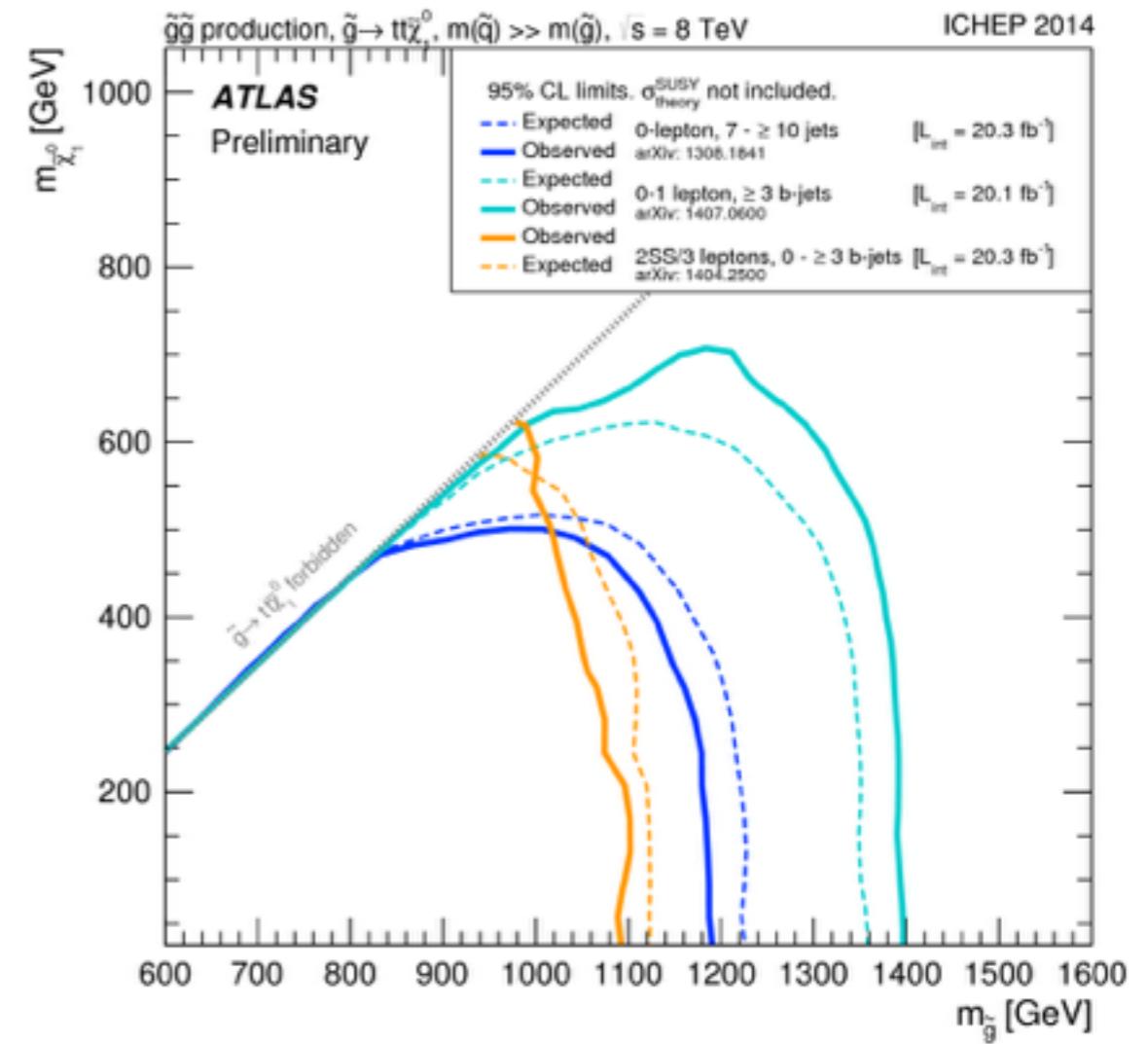


*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 for theoretical signal cross section uncertainty.

stop searches



gluino searches



$$m_{\tilde{t}} \gtrsim 700 \text{ GeV}$$

$$m_{\tilde{g}} \gtrsim 1.4 \text{ TeV}$$

What does it mean for naturalness?

‘natural SUSY’ (stop,gluino,higgsino)

Papucci, Ruderman,Weiler
‘11

stop

$$\delta m_h^2 = -\frac{3}{8\pi^2} y_t^2 \left(m_{Q_3}^2 + m_{u_3}^2 + |A_t|^2 \right) \log \left(\frac{\Lambda}{\text{TeV}} \right)$$

$$\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \lesssim 600 \text{GeV} \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

gluino

$$\delta m_h^2 = -\frac{2}{\pi^2} y_t^2 \left(\frac{\alpha_s}{\pi} \right) M_3^2 \log^2 \left(\frac{\Lambda}{\text{TeV}} \right)$$

$$M_3 \lesssim 900 \text{ GeV} \sin \beta \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

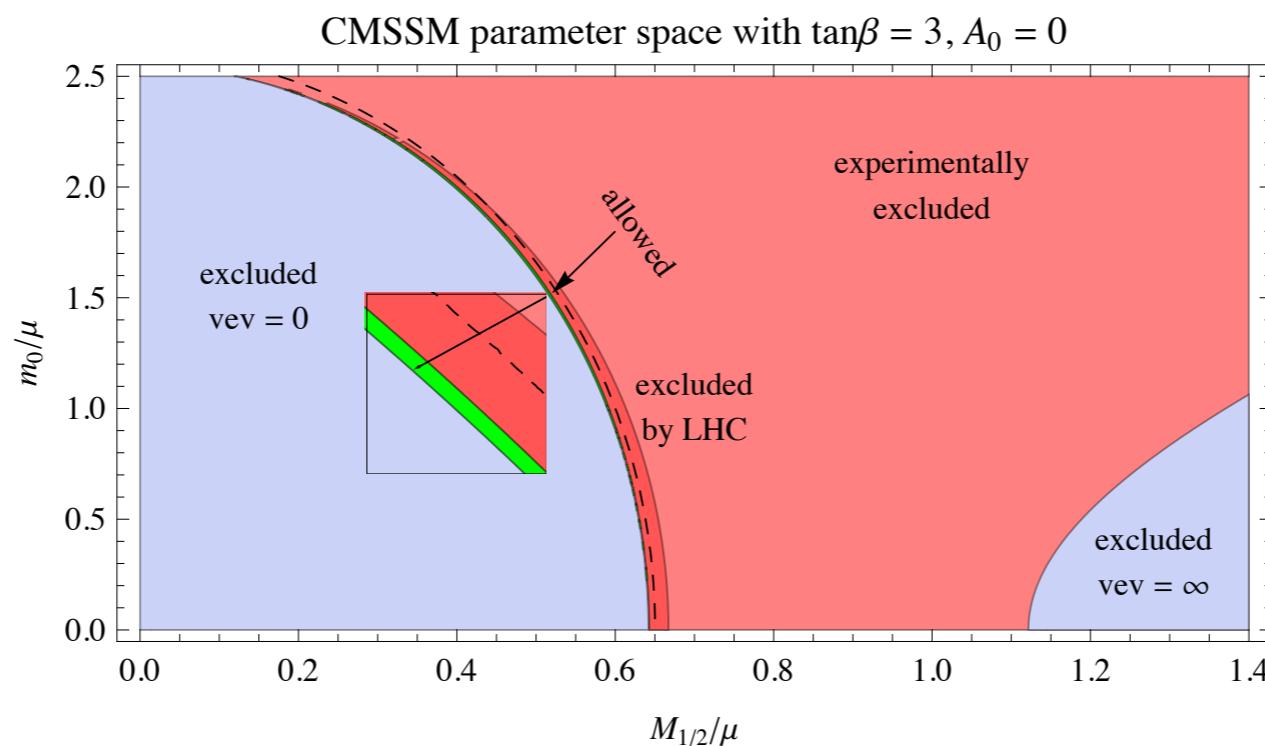
125 Gev Higgs : prefers heavy stops

$$\delta m_h^2 = \frac{3G_F}{\sqrt{2}\pi^2} m_t^4 \log \left(\frac{m_{\tilde{t}}^2}{m_t^2} \right)$$

~ 1-10 TeV stop depending on A-term

More complete models generally yield %-level fine-tuning

CMSSM



NMSSM, split generation, low-scale mediation and even R-parity breaking are generically tuned.

Arvanitaki, Baryakhtar,
Huang, Tiburg, Villadoro '13

More clever model building might save the day

Scherk-Schwartz SUSY breaking

Dimopoulos, March-Russell '14

Compressed spectrum

LeComte, Martin '11
Dimopoulos, March-Russell, Scoville '14

Stealth supersymmetry

Fan, Reece, Ruderman '11

Twin Higgs

Chacko, Goh, Harnik '05
Craig, Howe '13

Dirac gauginos

In the MSSM gauginos are Majorana

$$M\lambda\lambda$$

$$\int d^2\theta F_X \theta^2$$

A diagram showing a loop integral of a superfield X with gauge fields W_α and W^α . The integral is written as $\int d^2\theta F_X \theta^2$, where F_X is the superfield strength of X .

Can be Dirac if new superfields are added

$$W_\alpha^1, W_\alpha^2, W_\alpha^3$$

$$S, T, G$$

$$M_D \lambda \Psi$$

N=2 supersymmetry

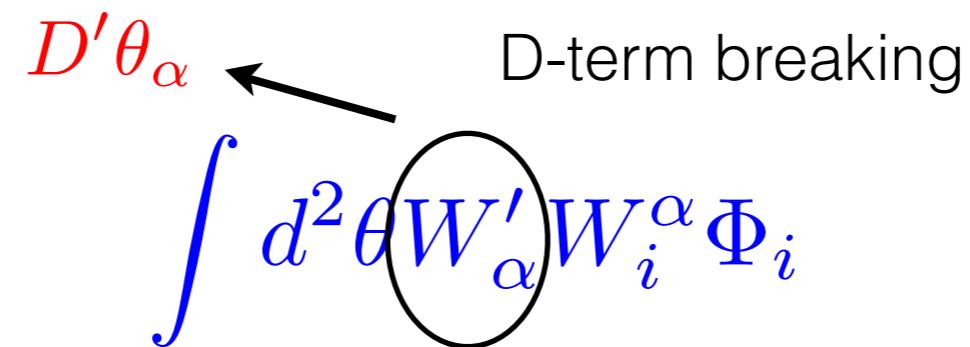
extra-dimension

Supersoft SUSY breaking

Fox, Nelson, Weiner '02

$$\int d^2\theta \textcolor{blue}{W'_\alpha} W^\alpha_i \Phi_i$$

$D'\theta_\alpha$ ← D-term breaking



Dirac gauginos do not feed into scalar masses through **renormalization**

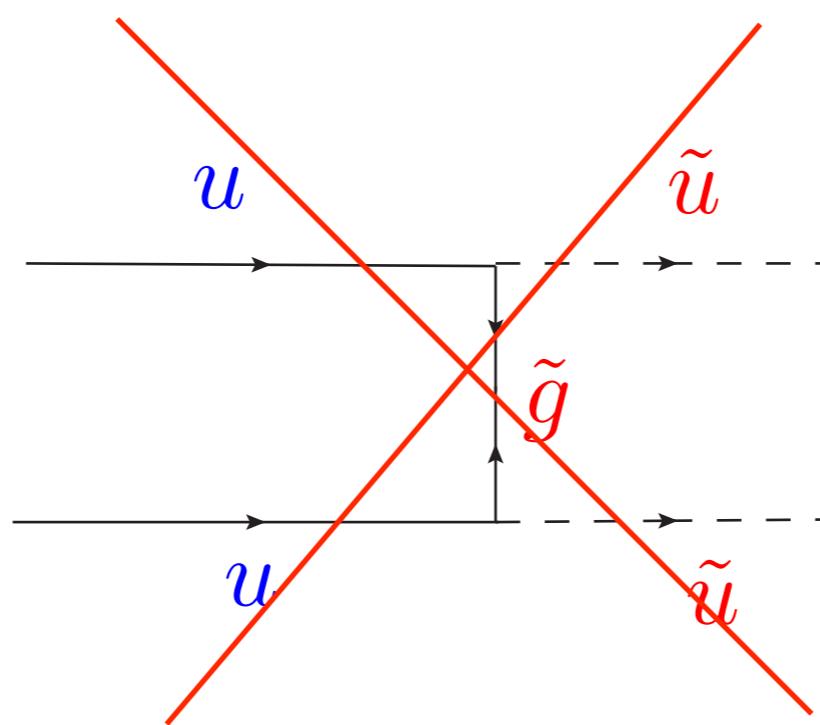
$$m^2 = \frac{C_i(r) \alpha_i m_i^2}{\pi} \log \left(\frac{\delta^2}{m_i^2} \right)$$

They can be naturally **heavier** than
scalars

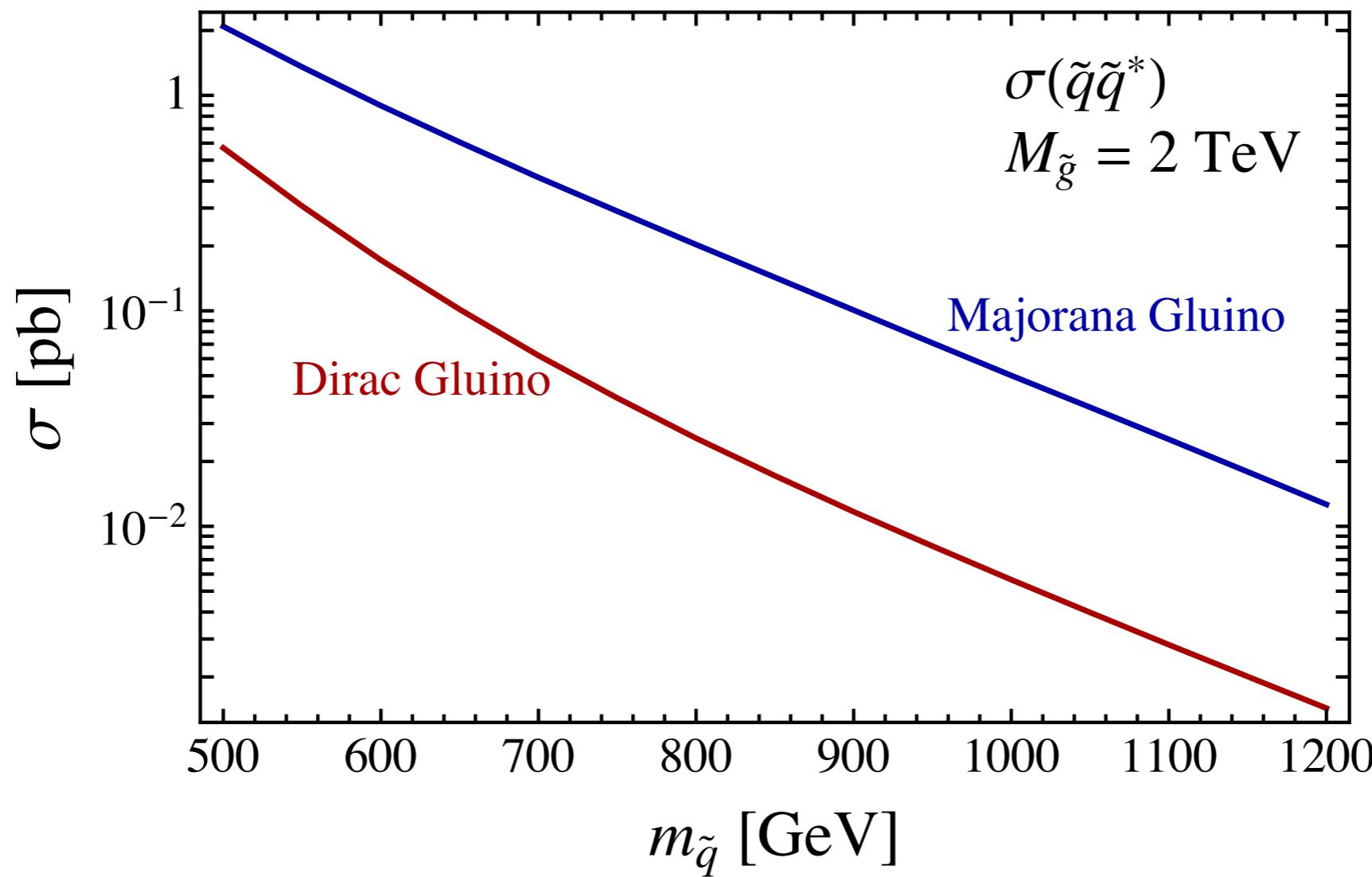
LHC will have a harder time seeing the
gluino...

M. Heikinheimo, M. Kellerstein, V. Sanz '12
Kribs, Martin '12

...and squarks



Squark production



Frugiuele, T.G., Kumar, Ponton

R-symmetry

With Dirac gaugino: possible to impose an
U(1) R-symmetry

$$M_D \lambda \Psi$$

Kribs,Poppitz,Weiner '02

- Bounds from FCNC are weaker: off diagonal m_{ij}

$$R[Q, U^c, D^c, L, E^c] = 1 \quad R[H_u, H_d] = 0$$



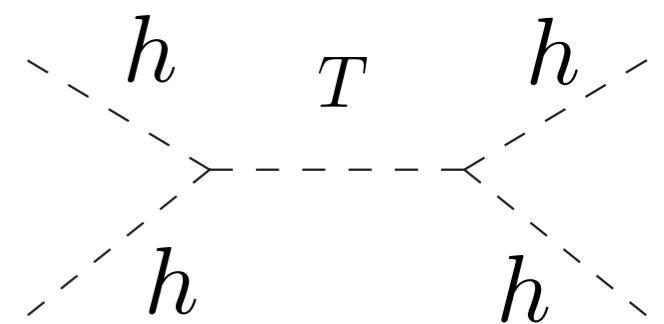
Higgs mass

Tree-level:

Reduced quartic, usual of Dirac gauginos

$$\int d^2\theta W'_\alpha W_i^\alpha \Phi_i \rightarrow D_2 = M_2 T^a + H_u^\dagger \sigma^a H_u + \dots$$

When the scalar T is integrated out:



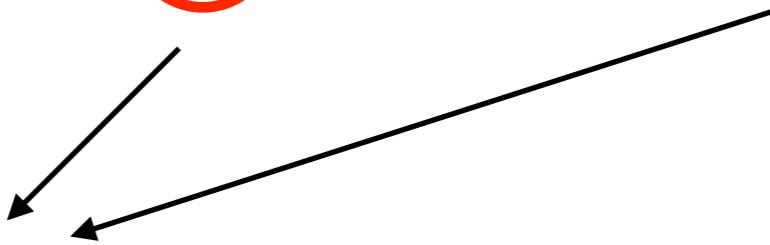
$$\lambda \rightarrow 0$$

Higgs quartic

If the mass of T is set by M_2 and $\lambda_T = 0$

No help (at tree-level) from

$$\lambda_T H_u T R_d + \lambda_S H_u S R_d$$



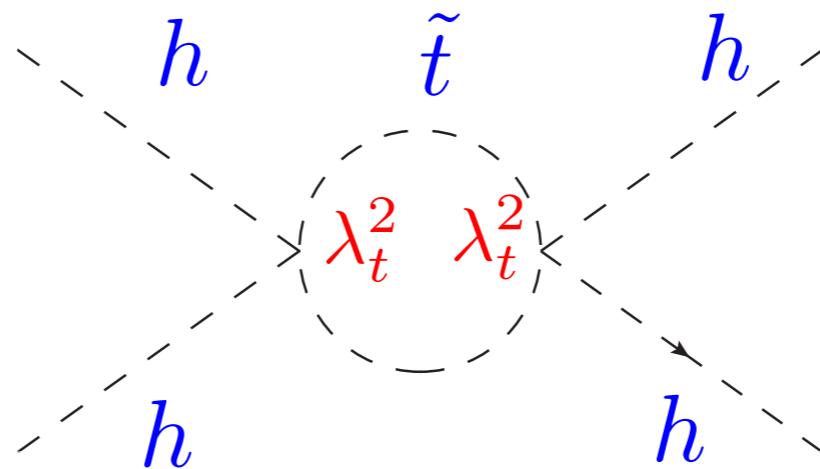
don't get a vev (In the limit of exact R-symmetry)

But do help in models without an R-symmetry

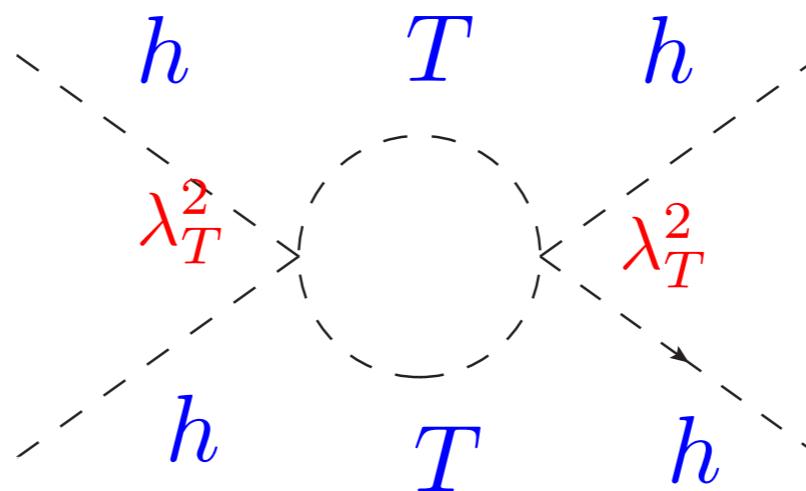
Benakli, Goodsell, Staub 1211.0552

Loop-level

Usual stop correction (but A-terms are 0)



Similar loop from the triplet



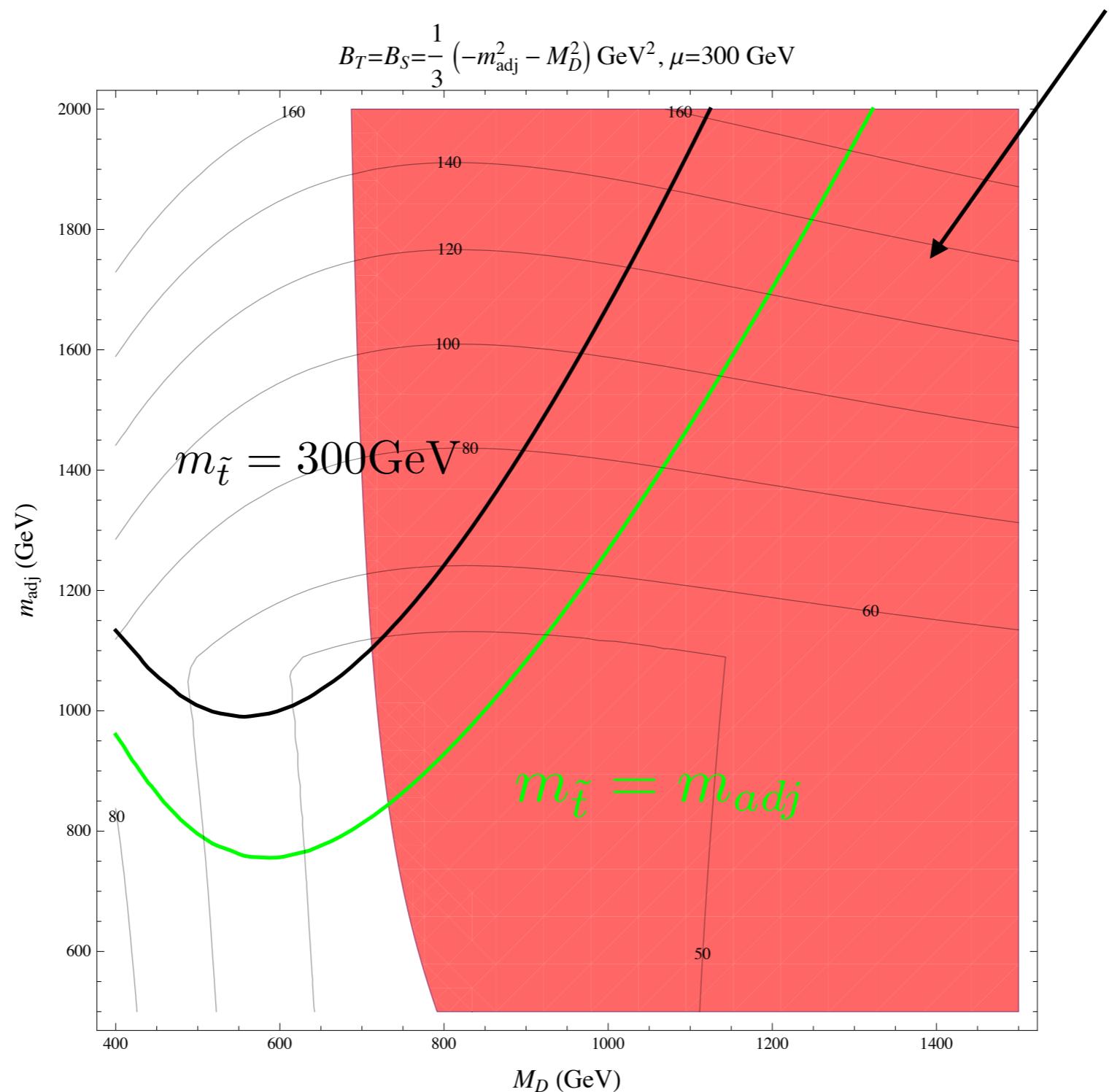
$$V_{\text{CW}} \sim \frac{1}{16\pi^2} \left(5\lambda_T^4 \log \frac{m_T^2}{M_2^2} + 3\lambda_t^4 \log \frac{m_{\tilde{t}}^2}{m_t^2} \right)$$



Very sensitive to λ_T

....but so are electroweak precision measurements

allowed by EWPT



Split Supersymmetry

Naturalness **might not be a good guide**

SUSY might still be relevant

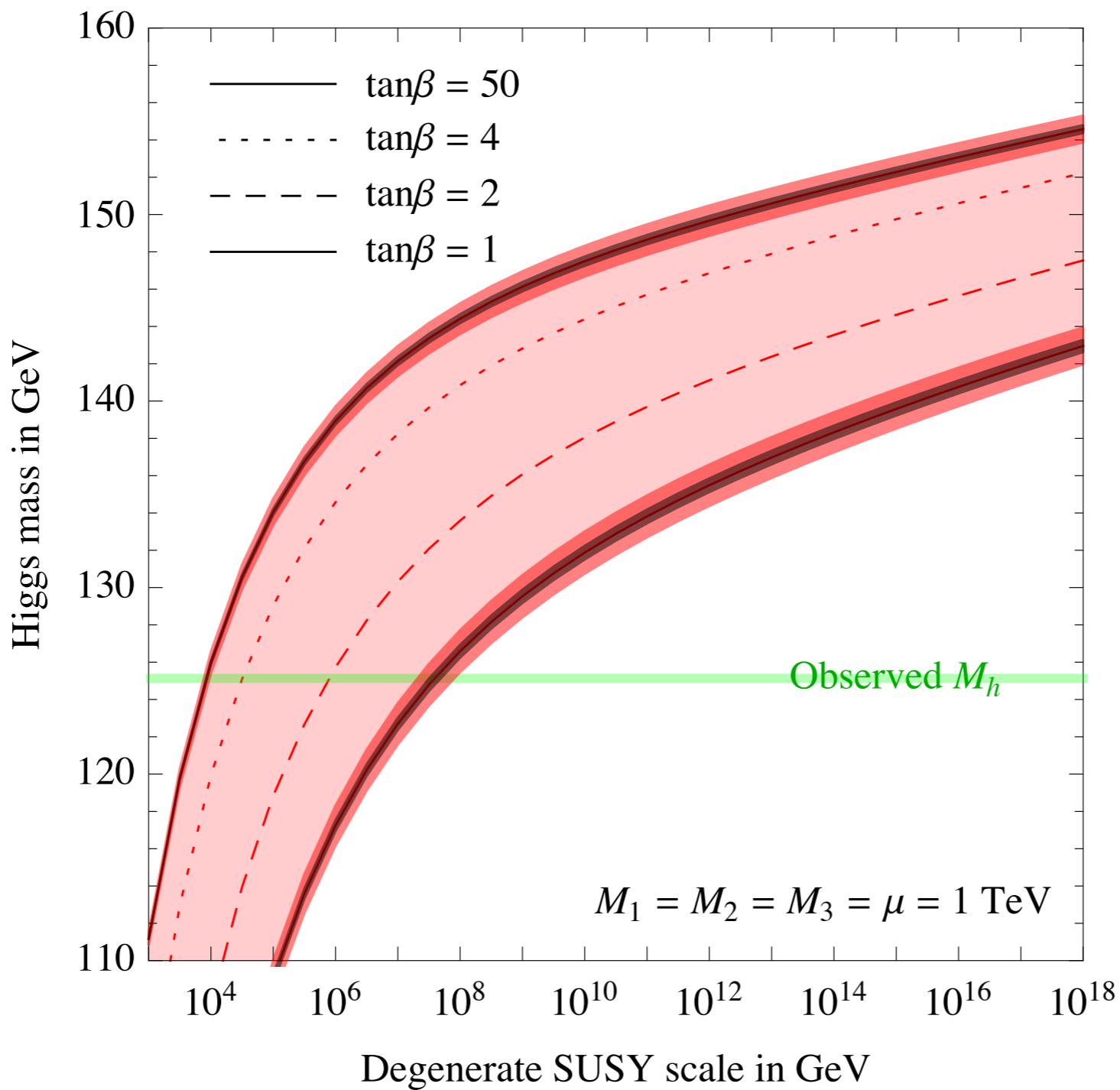
- Dark matter
- Gauge coupling unification
- ‘UV’ reasons

Gauginos and scalars might **not be at the same mass scale**

natural in for example anomaly mediation

Prediction for the Higgs mass

Split–SUSY



Bagnaschi,
Giudice,Slavich,Strumia
'14

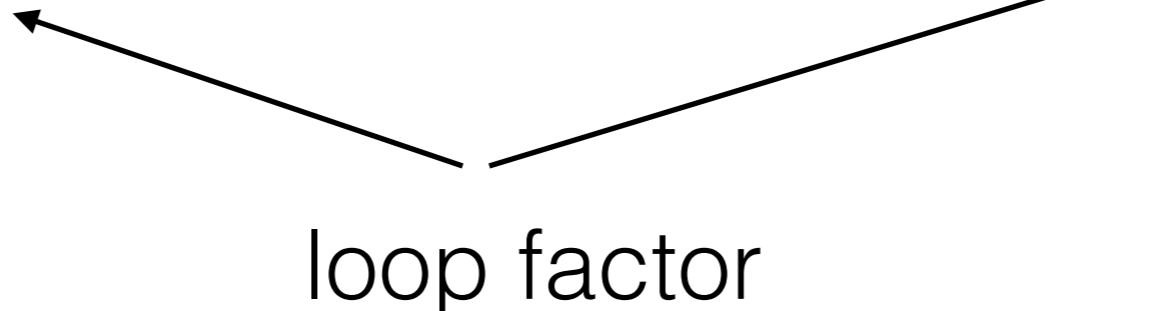
Mini-split in **anomaly mediation**

gaugino masses (AMSB):

$$M_i = \frac{b_i}{16\pi^2} g_i^2 m_{3/2}$$

scalar masses (gravity
mediation):

$$m^2 = m_{3/2}^2$$

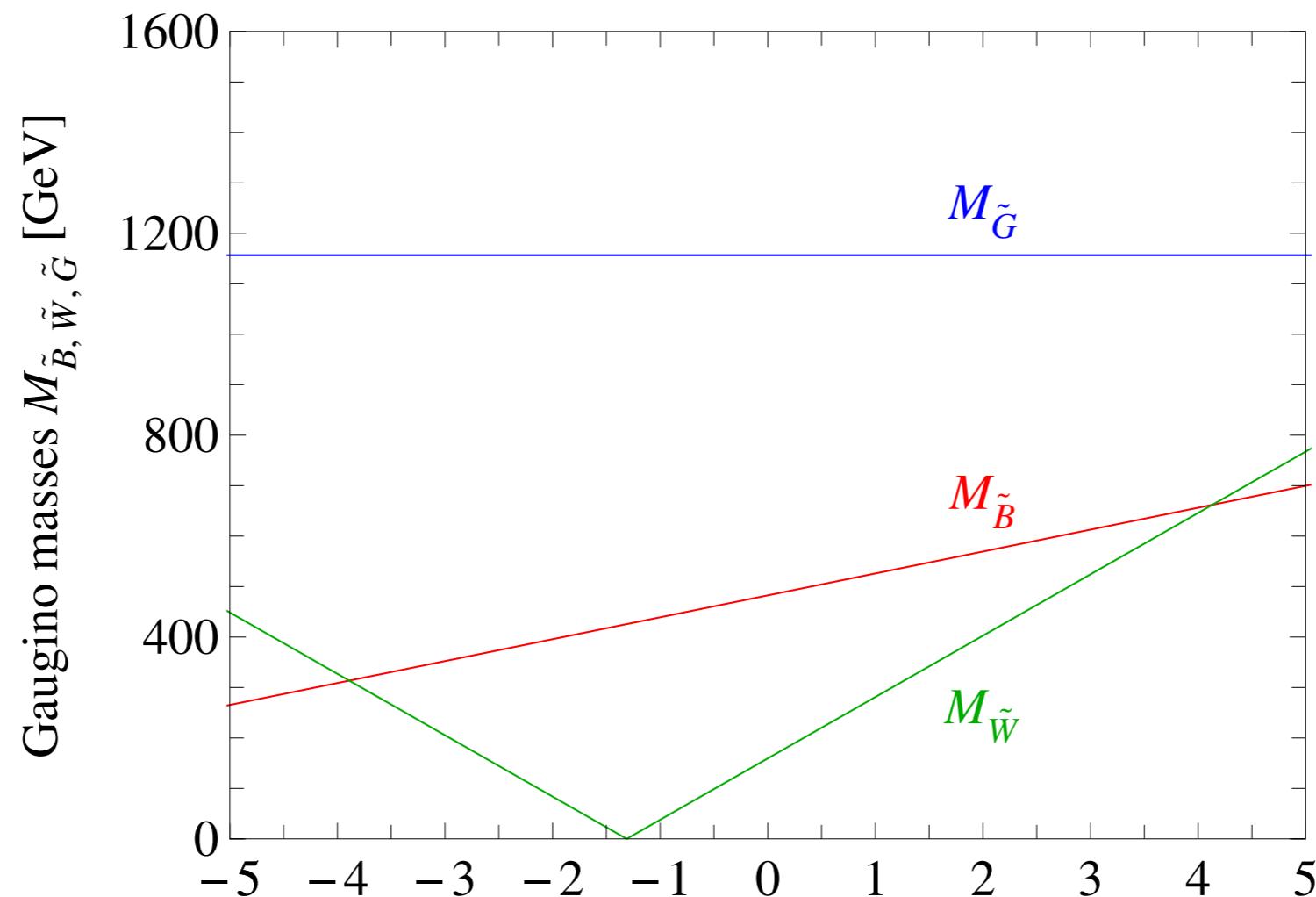


Higgsino (mu term): chosen
to be at $m_{3/2}$

Similar possibility in **gauge mediation**

Gaugino spectrum

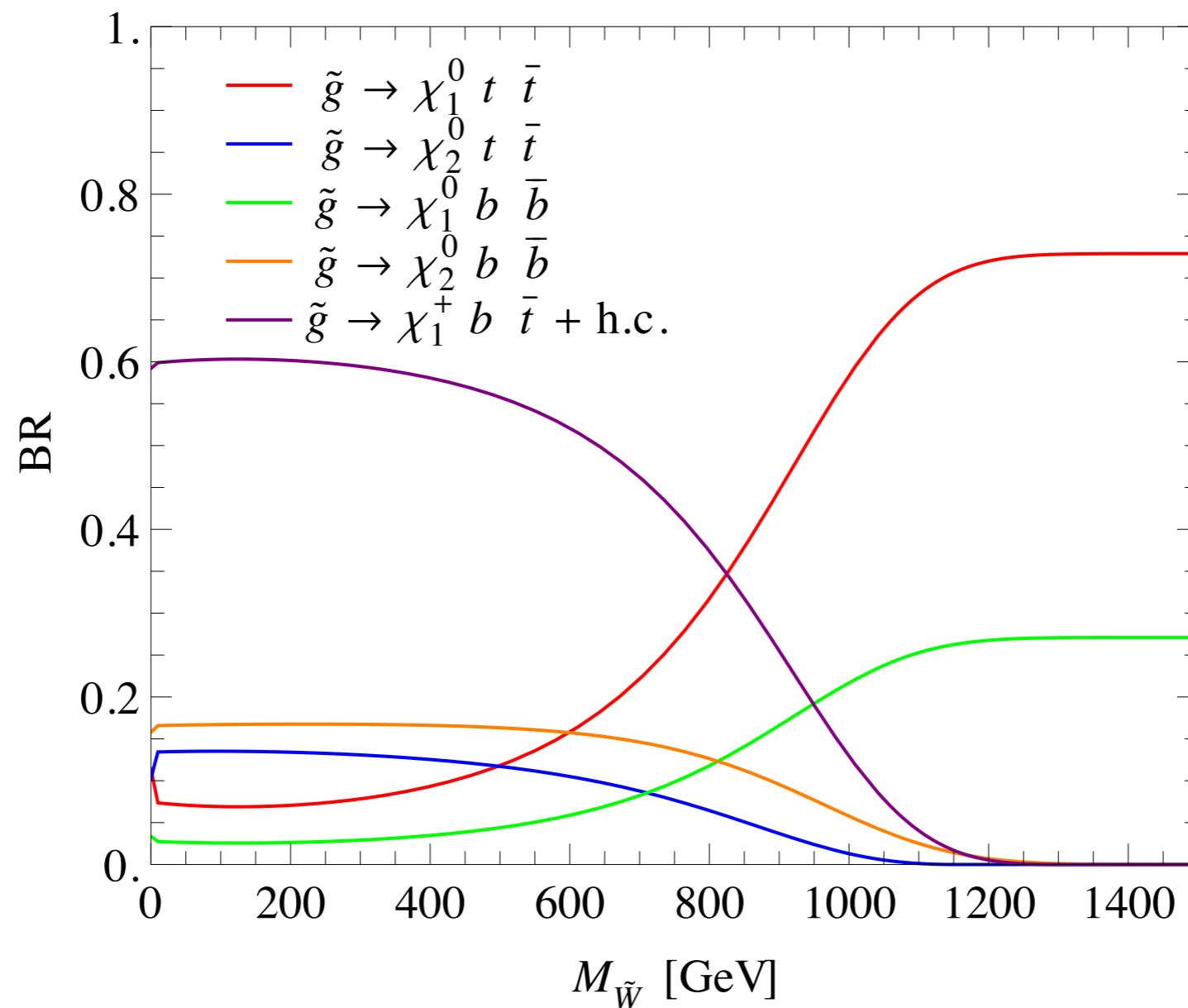
Beauchesne, Earl, T.G. '15



$$C_\mu = \frac{\mu}{m_{3/2}} \frac{m_A^2 \sin^2 \beta}{m_A^2 - \mu^2} \ln \frac{m_A^2}{\mu^2}$$

Parametrize deflection
from AMSB spectrum

Gluino decay (through 3rd generation)

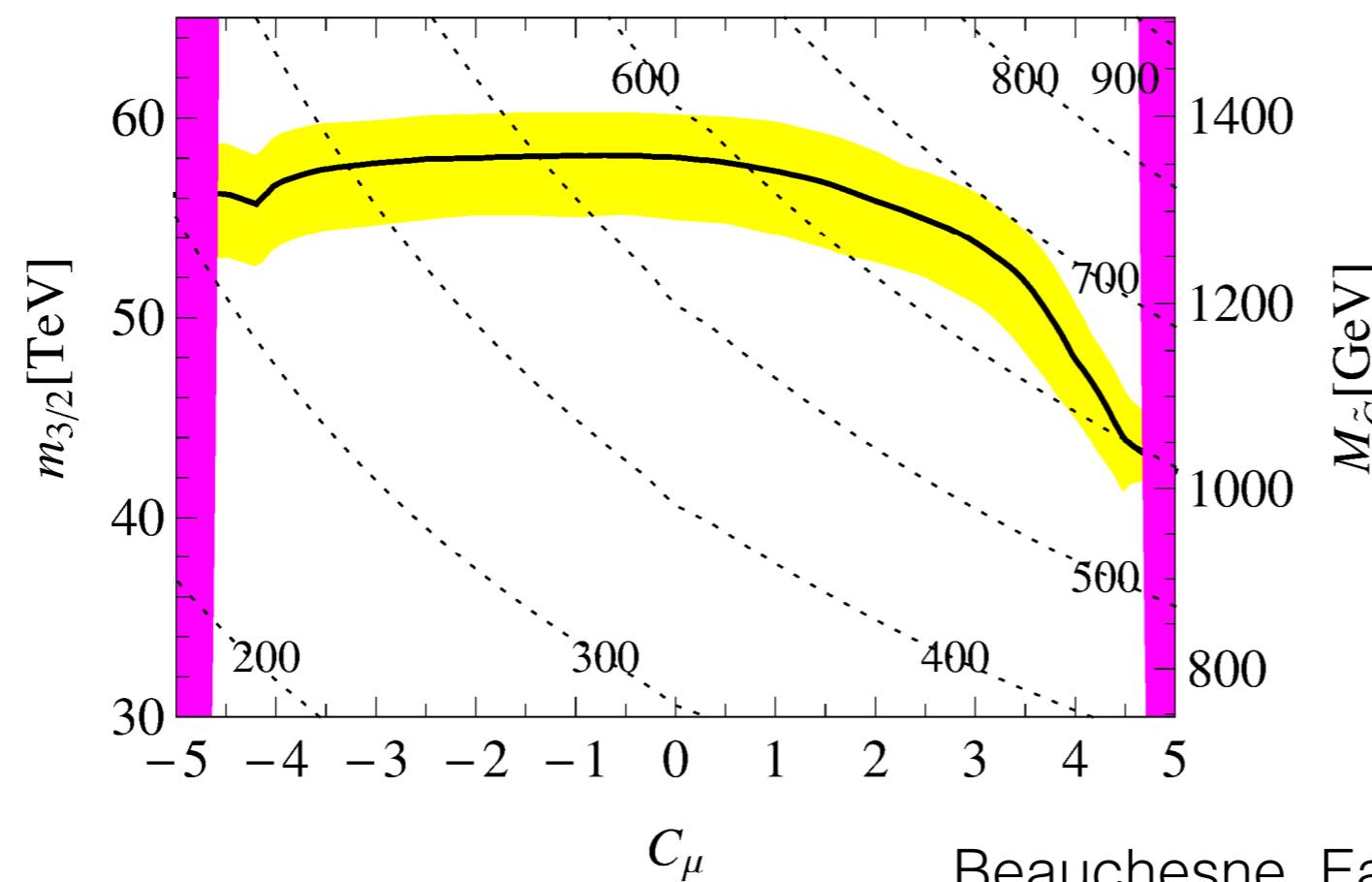


$$m_{\tilde{G}} = 1.5 \text{ TeV} \quad m_{\tilde{B}} = 0 \text{ GeV}$$

LHC limits

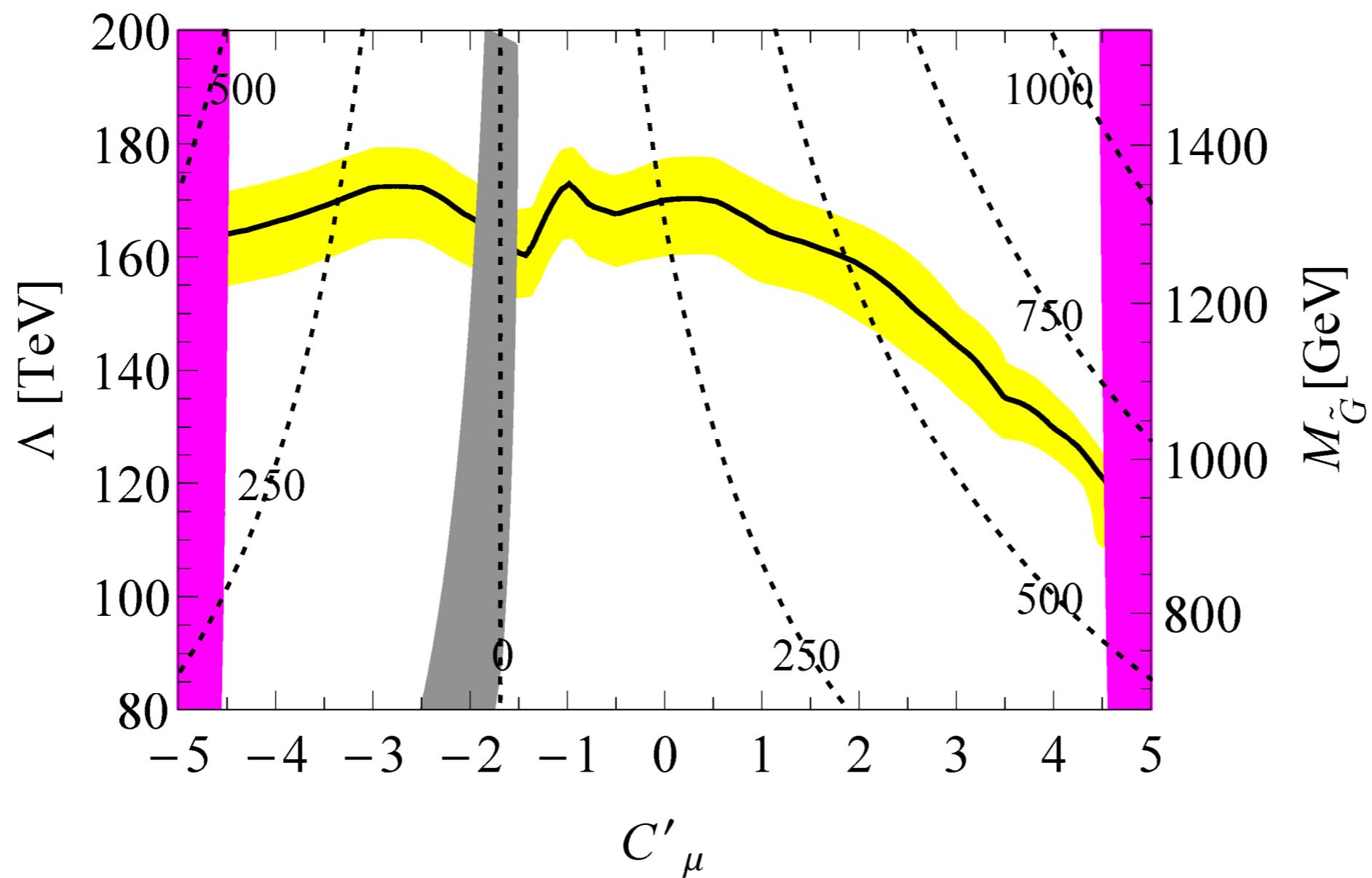
Gaugino spectrum and branching ratios are obtained as a function of $m_{3/2}$ and C_μ . The correct Higgs mass is obtained by fixing $\tan \beta$

Recast LHC searches



ATLAS multi-leptons+b-jets
ATLAS 0-1 lepton+b-jets
CMS high jets multiplicity
CMS 2 OS leptons+jets

Gauge mediation



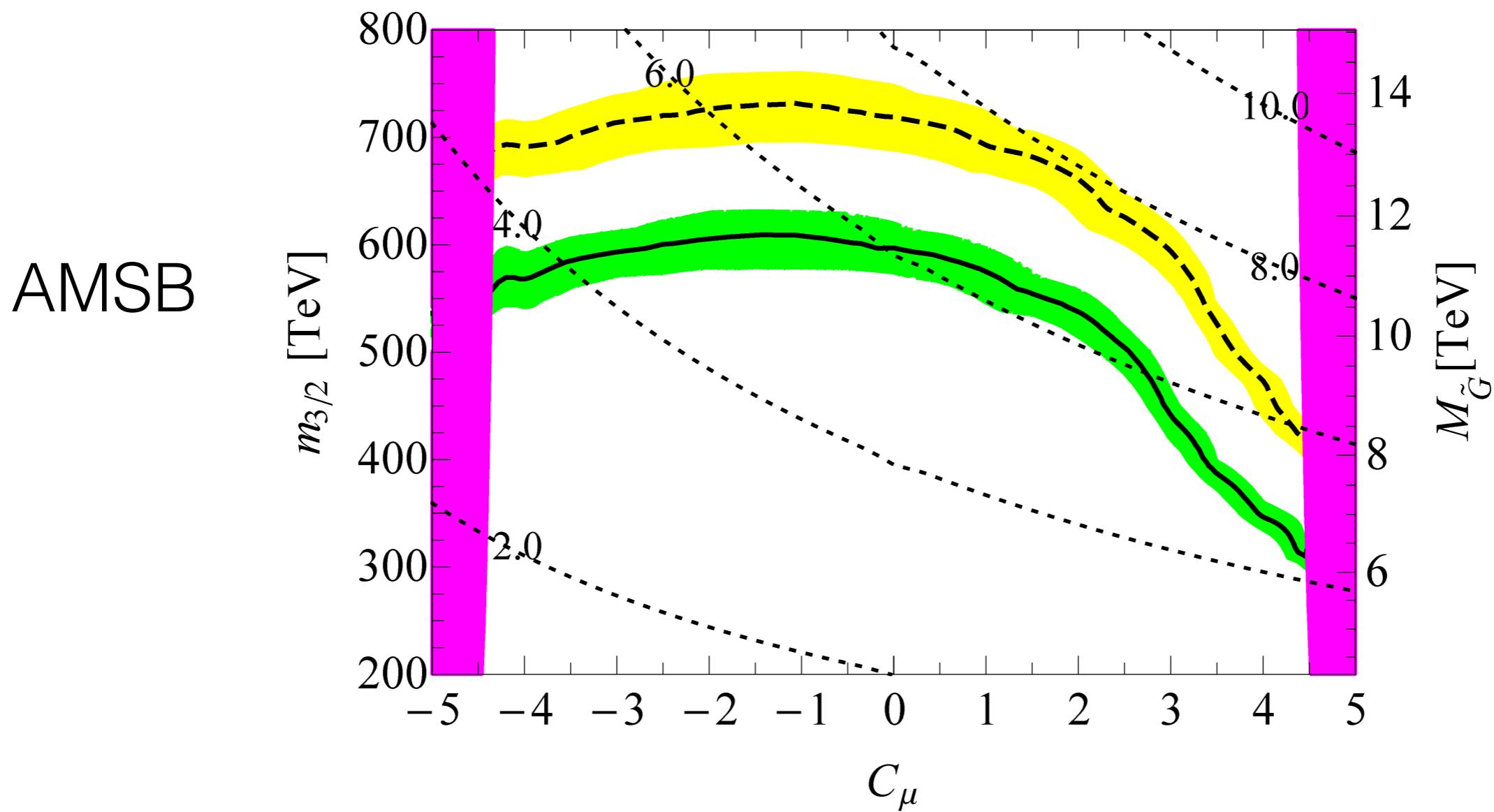
Prospect for a 100 TeV collider

0-lepton + jets +MET analysis

Jung, Wells '13

2 SS lepton analysis

Cohen et al. '14



GMSB

