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Thermal Transport of the Solar Captured Dark Matter and its Impact on the Indirect Dark Matter Search

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Outline

- ❖ Motivations
- ❖ How does thermal energy transport?
- ❖ Dark matter (DM) evolution and thermal transport
- ❖ Physical implications and its impact on indirect search
- ❖ Summary

Motivations

- ❖ Recent studies suggest that DM is not collisionless
- ❖ The DM temperature, T_χ , is not necessarily to be the same as the core temperature of the Sun, T_c
- ❖ Thermal energy can exchange between DM-DM and DM-nucleus or even dissipates through annihilation
- ❖ For much weaker $\sigma_{\chi p}$, the thermal energy exchange between DM-nucleus is less efficient
- ❖ In such case, T_χ can be distinct from T_c eventually
- ❖ As a consequence, it may alter the DM annihilation rate

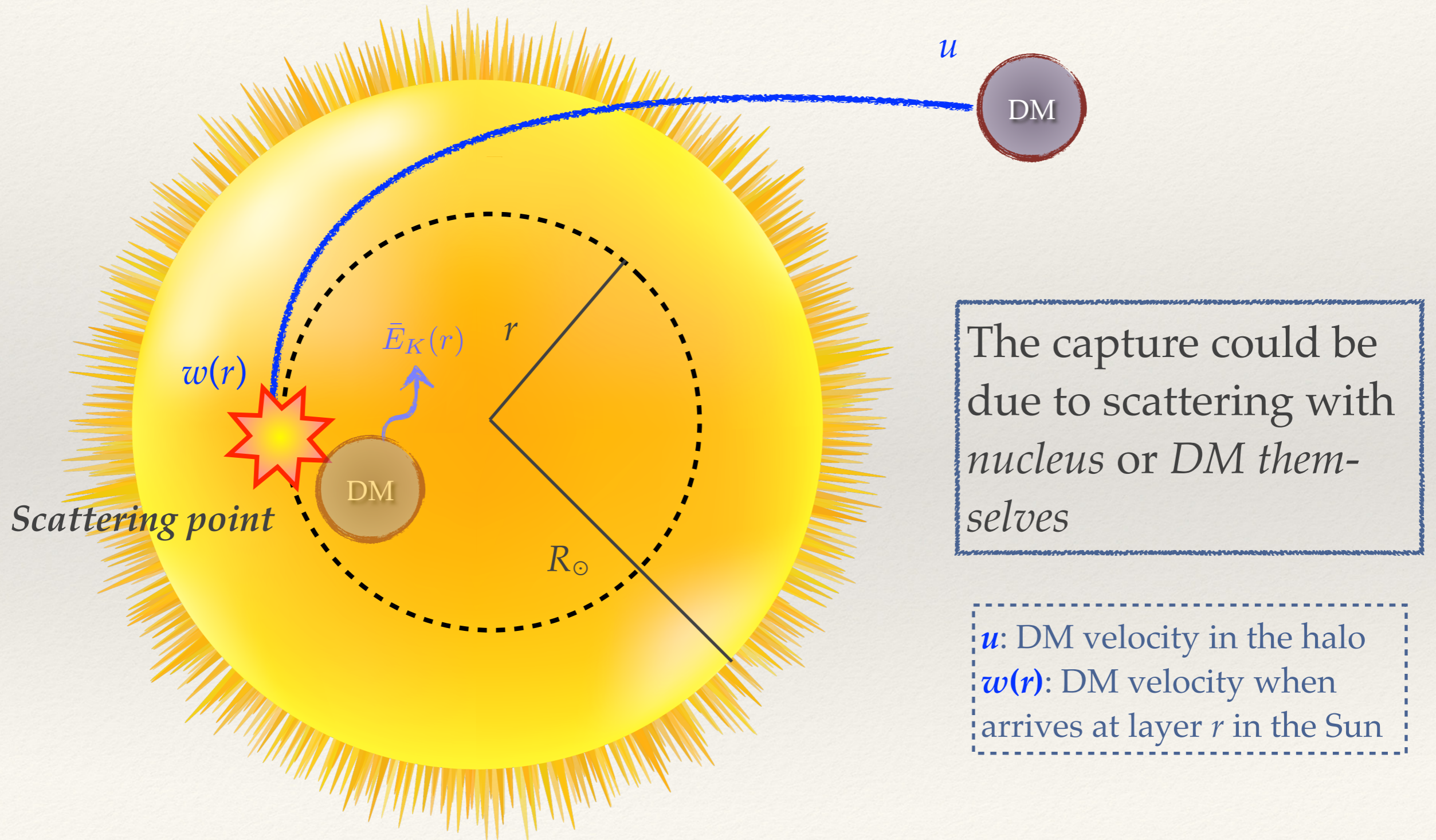
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How does thermal energy transport?



Energy flow via capture



DM kinetics after scattering

- ❖ When DM reaches layer r in the Sun, it carries velocity $w = \sqrt{u^2 + v_{\text{esc}}^2(r)}$

- ❖ Thus, the average DM kinetic energy after scattering with nucleus is given by

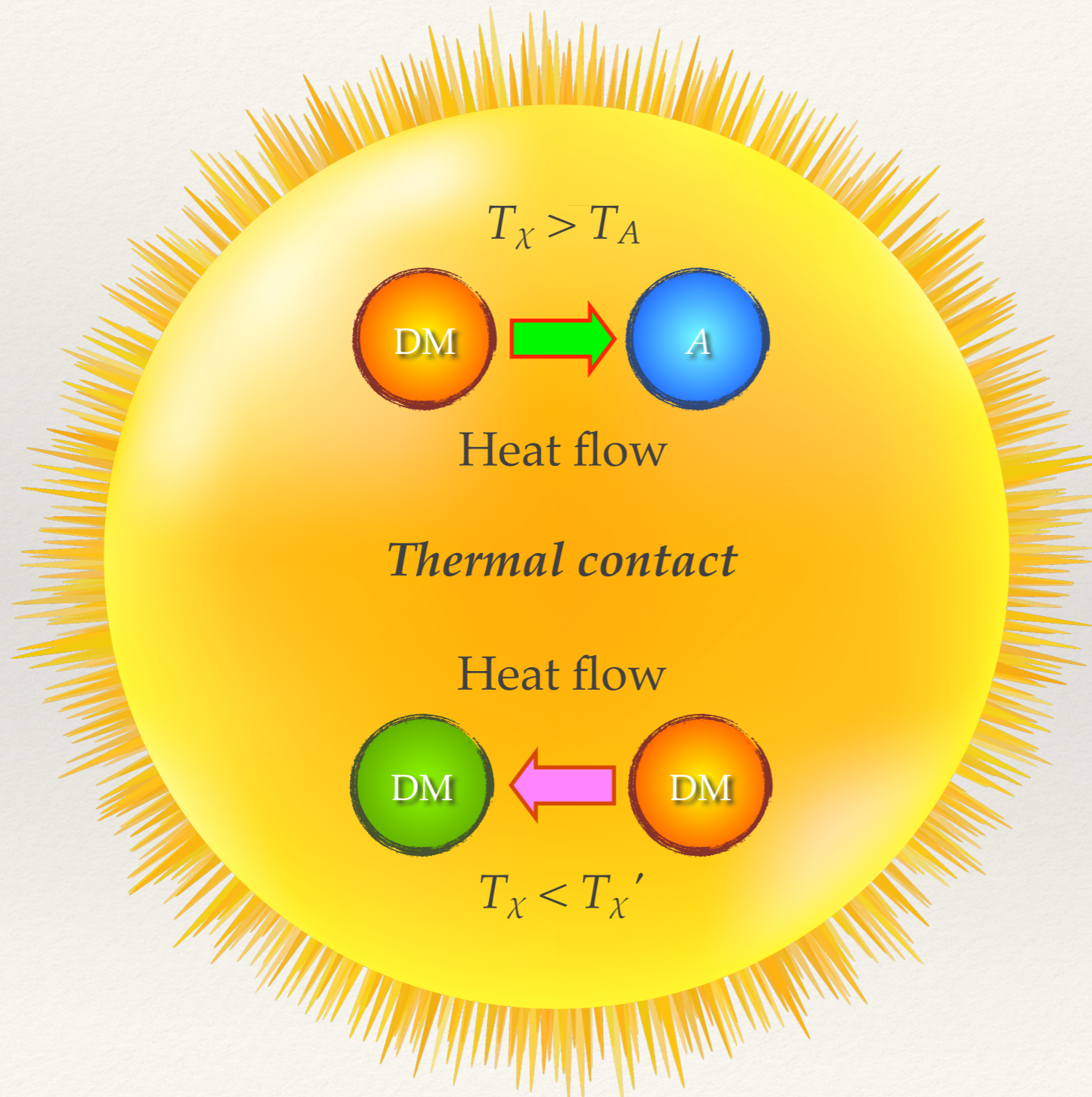
$$\bar{E}_K(r) = \frac{m_\chi}{4} \left(\frac{m_\chi - m_A}{m_\chi + m_A} \right)^2 u^2 + \frac{m_\chi}{2} \frac{m_\chi^2 + m_A^2}{(m_\chi + m_A)^2} v_{\text{esc}}^2(r)$$

- ❖ However, the scattering among DM themselves gives

$$\bar{E}_K(r) = \frac{1}{4} m_\chi v_{\text{esc}}^2(r)$$

- ❖ m_A is the nucleus mass of element A in the Sun

Energy flow via thermal contact



Thermal contact via *DM*
and nucleus *A*

Thermal contact via *DM*
themselves

Mean collision time

- ❖ The mean collision time is used to determine which one reaches thermal equilibrium earlier, DM-DM or DM-nucleus
- ❖ Thereupon, by $\tau \sim 1 / (n\sigma v)$ we have

$$\begin{cases} \tau_{\chi\chi}(t) \simeq \frac{V_{\odot}}{N_{\chi}(t)\sigma_{\chi\chi}v}, & \text{for DM-DM} \\ \tau_{\chi\odot} \simeq \frac{V_{\odot}}{\sum_i N_i \sigma_{\chi A_i} v}, & \text{for DM-nucleus} \end{cases}$$

- ❖ Suppose the time scale for DM to reach thermal equilibrium is $t = \tau_{\chi}^{\text{eq}}$, we have

$$\tau_{\chi}^{\text{eq}} \simeq \tau_{\chi\chi}(\tau_{\chi}^{\text{eq}}) = \frac{V_{\odot}}{N_{\chi}(\tau_{\chi}^{\text{eq}})\sigma_{\chi\chi}v}$$

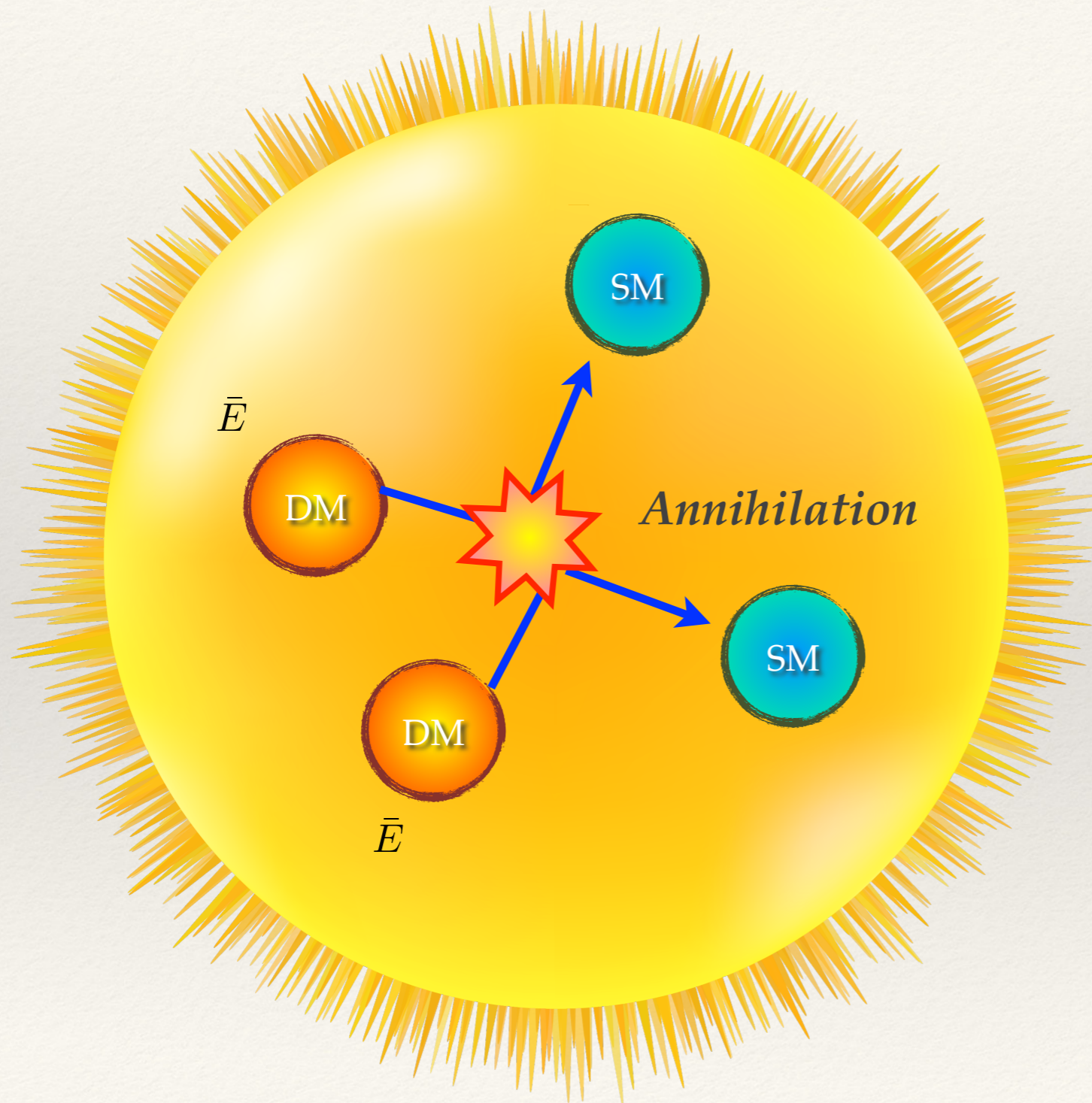
Thermal equilibrium time scale

- ❖ Assuming when DM reaches its own thermal equilibrium, N_χ is still its early stage of accumulation
- ❖ In other words, the capture by DM-nucleus scattering is dominant, $N_\chi \approx C_c t$
- ❖ By such assumption, we can define the ratio r

$$r \equiv \frac{\tau_\chi^{\text{eq}}}{\tau_{\chi\odot}} \simeq 10^9 \sqrt{\frac{\sigma_{\chi p}}{\sigma_{\chi\chi}}}$$

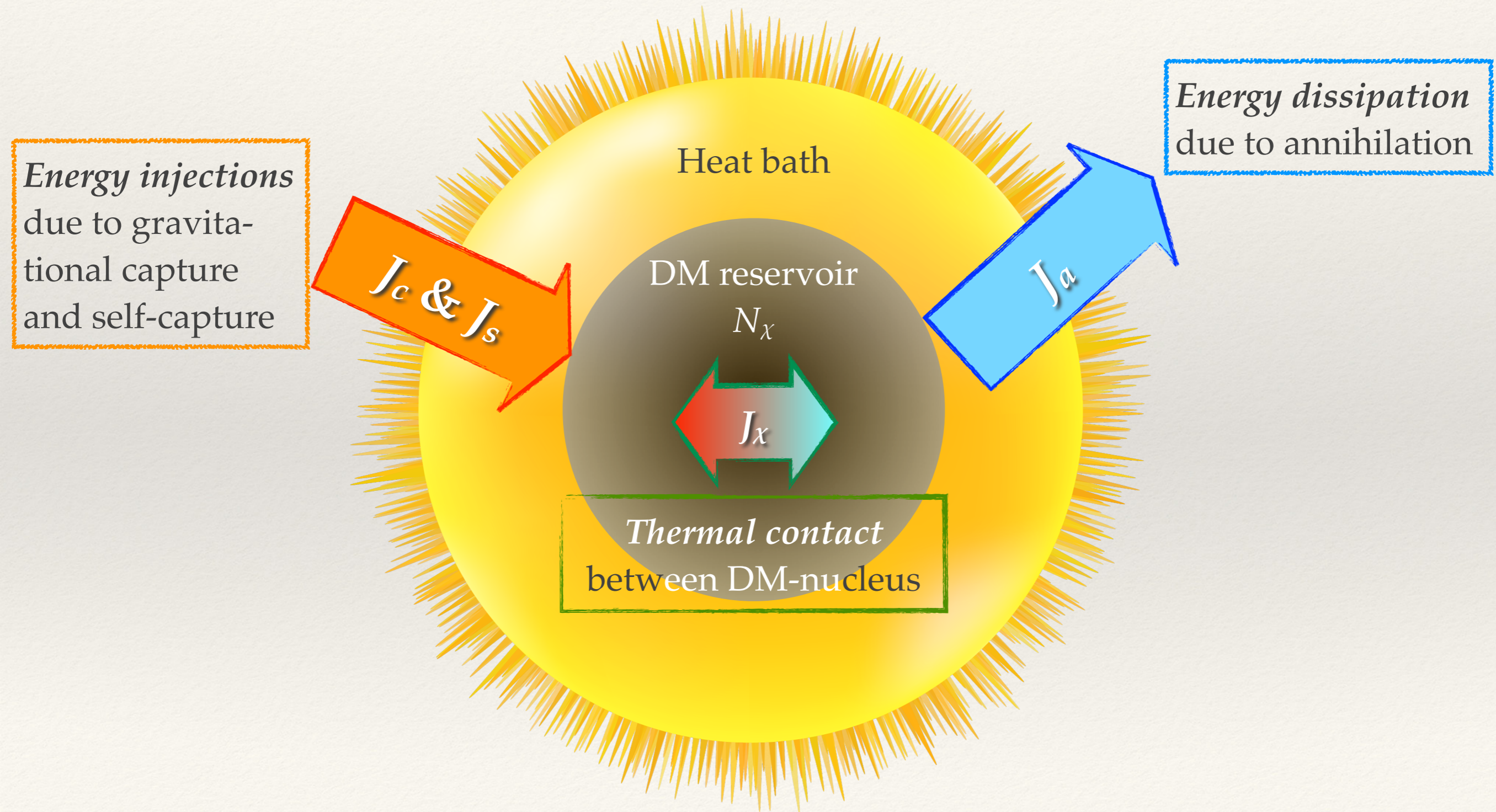
- ❖ For $r < 1$, DMs will reach their own thermal equilibrium quicker than with the solar medium
- ❖ Thus it is sufficient to just consider the thermal contact between DM and nucleus after $t = \tau_\chi^{\text{eq}}$

Energy dissipation via annihilation



Annihilation takes DM away and transports them into the final SM products

Defining physical quantities



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DM evolution process and thermal exchange



DM evolution in the Sun

- ❖ The common DM evolution equation:

$$\frac{dN_\chi}{dt} = C_c + C_s N_\chi - C_a N_\chi^2$$

- ❖ Additionally, N_χ contributes to the total energy evolving as well. Thus, we have the energy evolution equation:

$$\frac{d(N_\chi \bar{E})}{dt} = J_c + (J_\chi + J_s) N_\chi - J_a N_\chi^2$$

$$\bar{E} = \frac{s}{2} k_B T_\chi(t)$$

- ❖ $T_\chi(t)$ is the average DM temperature and s denotes the d.o.f. of DM. We take $s = 3$ for subsequent discussions

T_χ -dependent quantity in dN_χ/dt

- ❖ The two differential equations are mutual dependent since C_a is T_χ -dependent^[2]

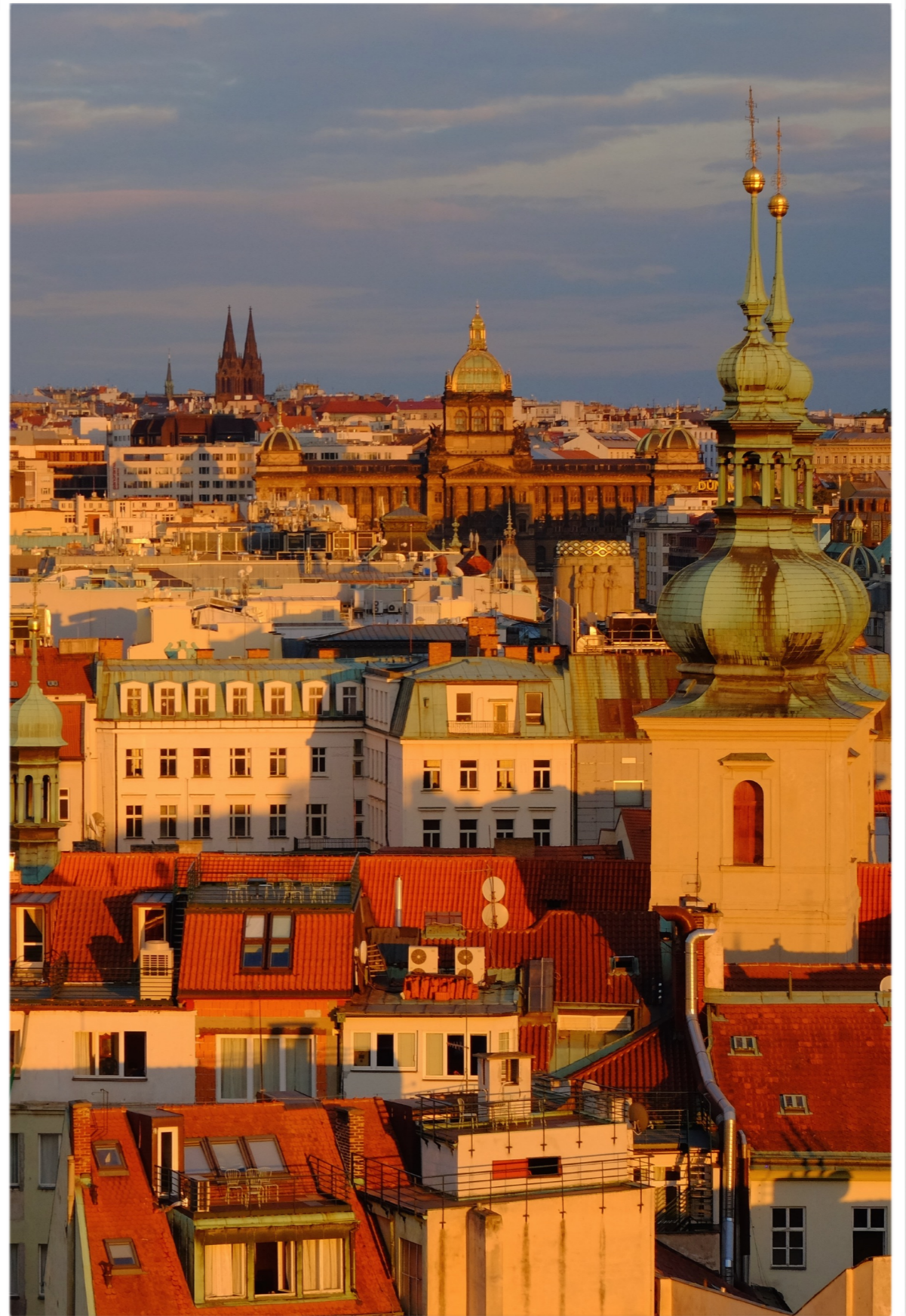
$$C_a = \frac{\langle \sigma v \rangle V_2}{V_1^2} \quad V_j \propto \left(\frac{T_\chi}{T_c} \frac{10 \text{ GeV}}{j m_\chi} \right)^{3/2} \quad j = 1, 2$$

- ❖ How T_χ evolving will affect N_χ accumulation and then feedback to T_χ itself
- ❖ The factors $J_{c,s,\chi,a}$ are all T_χ -dependent

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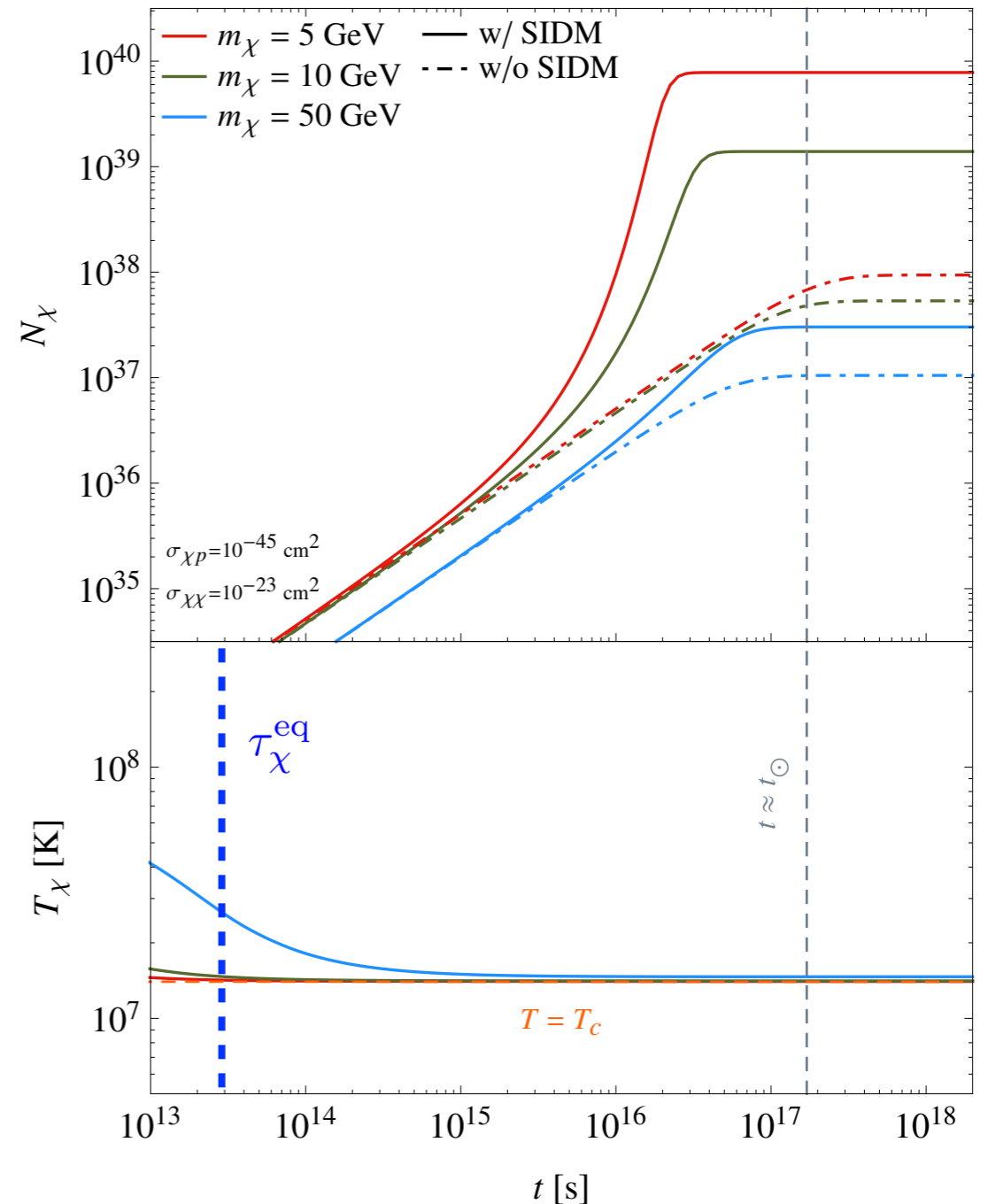


Physical implications



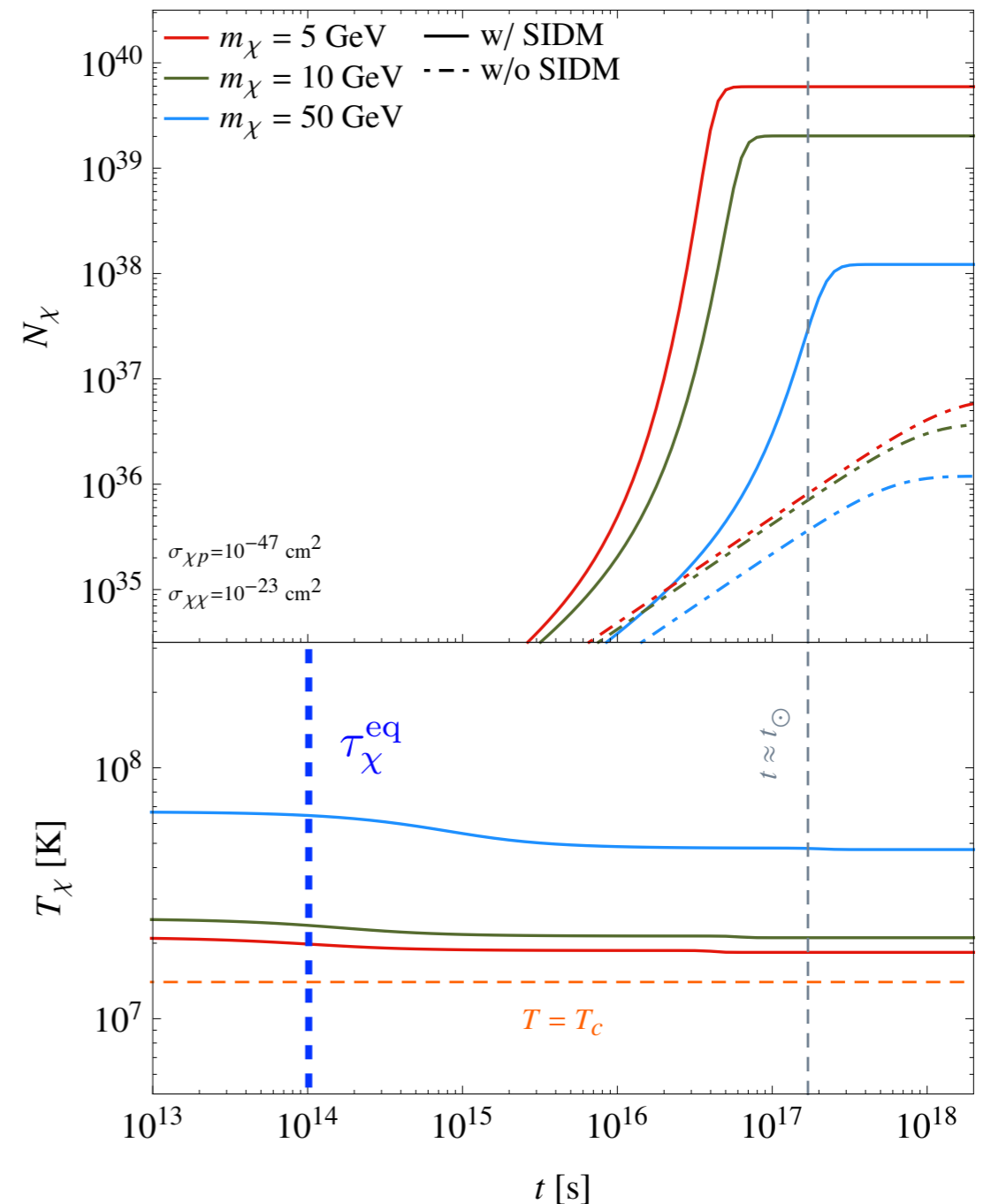
N_χ and T_χ evolutions: Stronger $\sigma_{\chi p}$

- ❖ $r \sim 0.01$ for $(\sigma_{\chi p}, \sigma_{\chi\chi}) = (10^{-45}, 10^{-23})$ cm^2
- ❖ Although J_c and $J_s N_\chi$ transport energy into the DM reservoir, the thermal contact $J_\chi N_\chi$ makes the heat out of the reservoir due to larger $\sigma_{\chi p}$ as well as the $J_a N_\chi^2$
- ❖ Eventually T_χ will be balanced by the Sun with the core temperature T_c



N_χ and T_χ evolutions: Weaker $\sigma_{\chi p}$

- ❖ $r \sim 10^{-4}$ for $(\sigma_{\chi p}, \sigma_{\chi\chi}) = (10^{-47}, 10^{-23})$ cm^2
- ❖ In such case, the energy injection via J_c and J_s can overcome the dissipation due to J_χ and J_a
- ❖ Moreover, energy injects to the DM reservoir constantly and makes them form their own thermal system
- ❖ The T_χ is distinct from T_c in this case

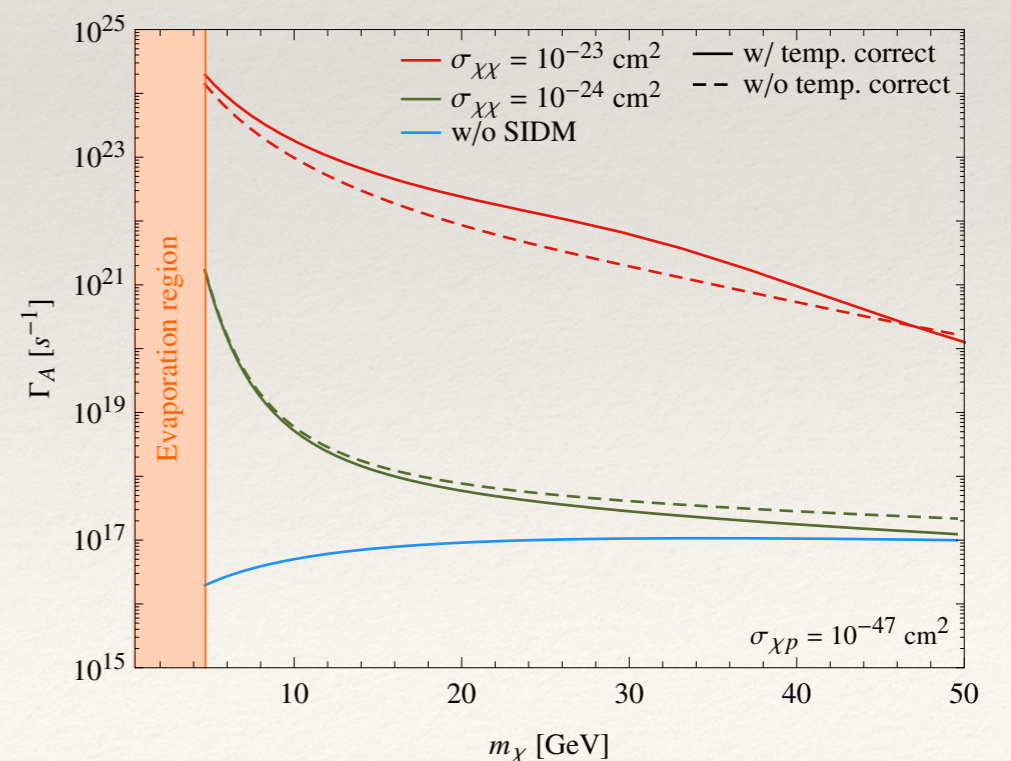
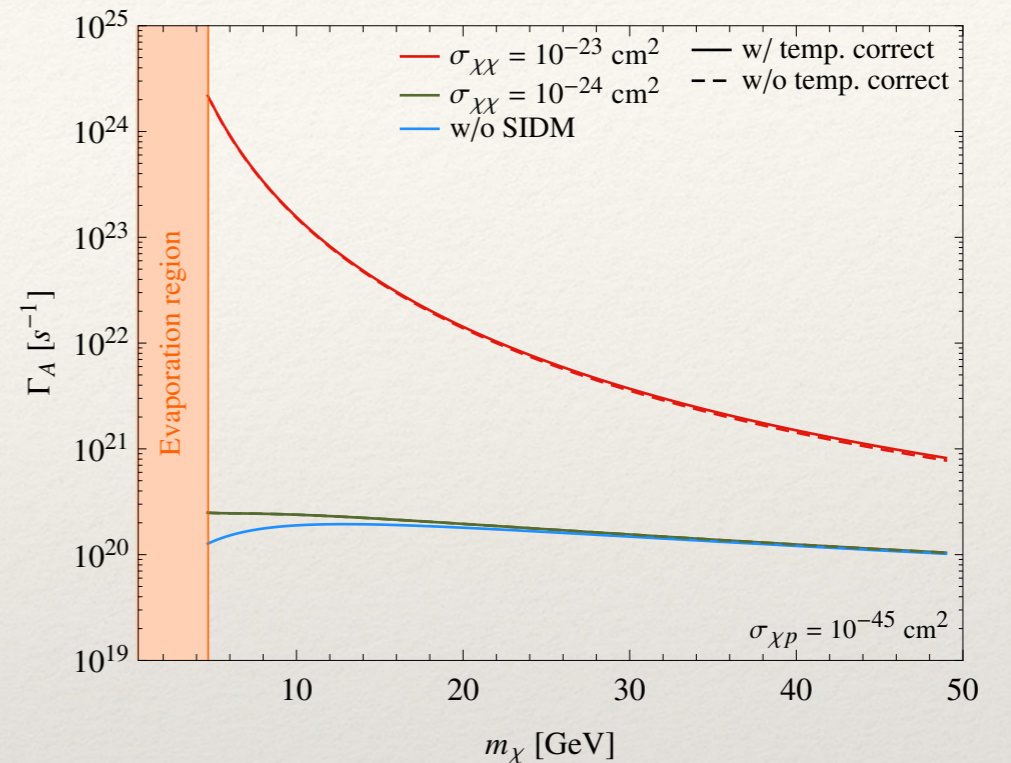


The total annihilation rate Γ_A

- ❖ The total annihilation rate

$$\Gamma_A = \frac{1}{2} C_a N_\chi^2$$

- ❖ Since $C_a \propto T_\chi^{-1.5}$, larger T_χ makes smaller C_a .
- ❖ However, smaller C_a leads to more N_χ accumulation once $dN_\chi/dt = 0$ is attained^[3]
- ❖ When N_χ^2 increment can overcome the C_a suppression, we have larger Γ_A compares to those without temperature correction
- ❖ For some cases, $dN_\chi/dt = 0$ does not attain at current epoch. It makes the corresponding Γ_A smaller than those without temperature correction



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Summary



To summarize so far...

- ❖ The DM temperature does not necessarily equal to the Sun's core temperature for collisional DM
- ❖ As a consequence of $\tau_{\chi}^{\text{eq}} < \tau_{\chi\odot}$, DM reaches thermal equilibrium before it starts thermal exchange with the solar nucleus efficiently
- ❖ If $\sigma_{\chi p}$ is small enough, eventually DM will have its own temperature inside the Sun regardless of T_c
- ❖ For more accurate prediction on the DM signal, the temperature correction factor should be considered