

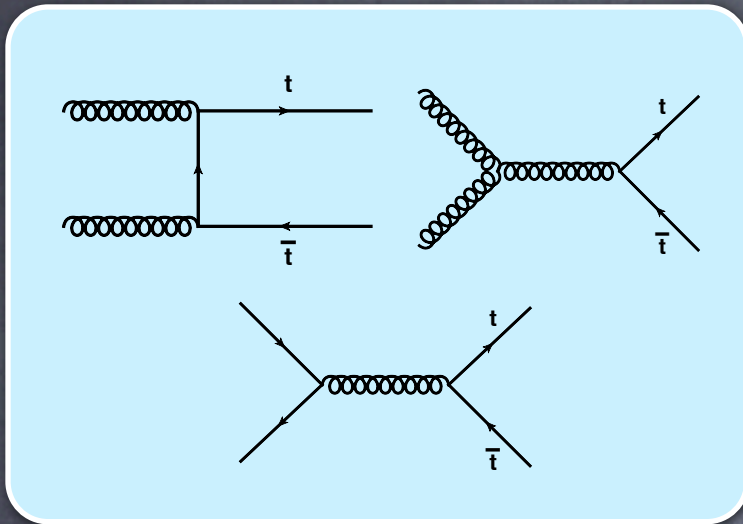
Phenomenological analysis using Diftop: impact of differential top quark-pair production on the gluon distribution.

M. Guzzi, K. Lipka*, S. Moch

JHEP 1501 (2015) 082

Probing QCD with top-quark pair production at the LHC

Top-quark pair production in pp collisions probes gluon distribution, top mass and α_s



with available NNLO theory for inclusive $t\bar{t}$ cross sections:

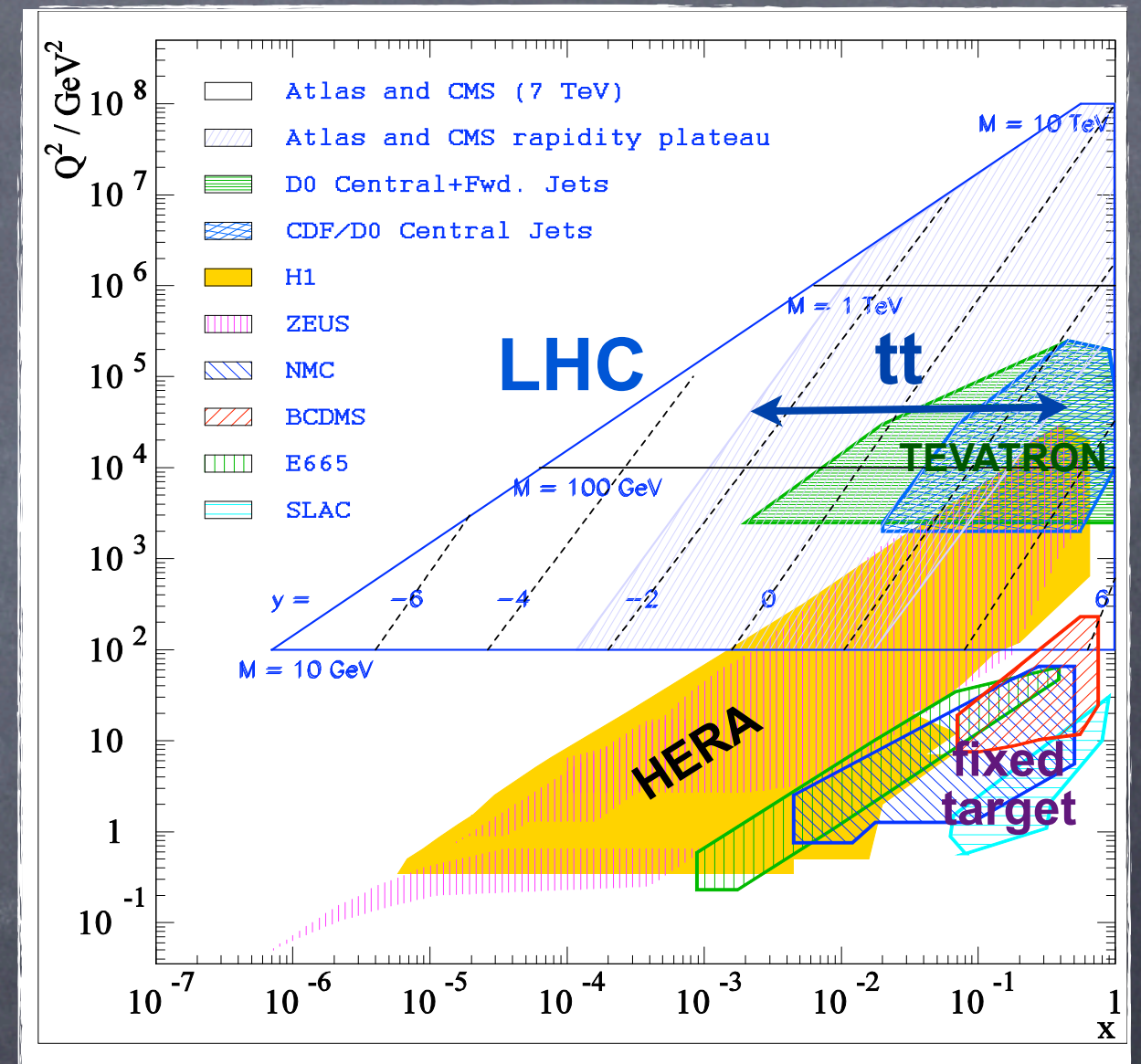
✓ correlations of $g(x)$, m_t , α_s studied

[CMS Collaboration, PLB 728 (2013) 496]

[S. Alekhin, J. Bluemlein, S. Moch PRD 89 (2014) 054028]

✓ impact on $g(x)$ is quantified

[Rojo et al., JHEP 1307 (2013) 167]

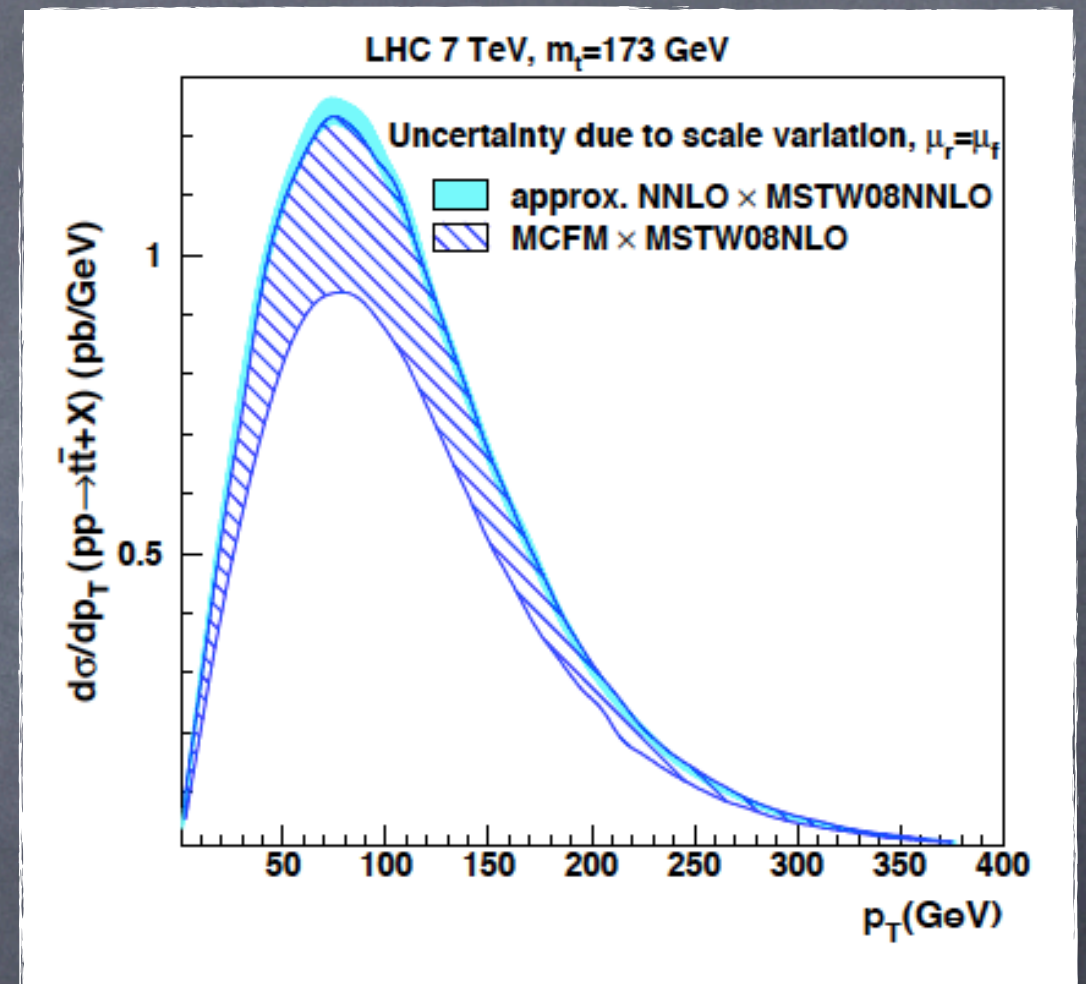
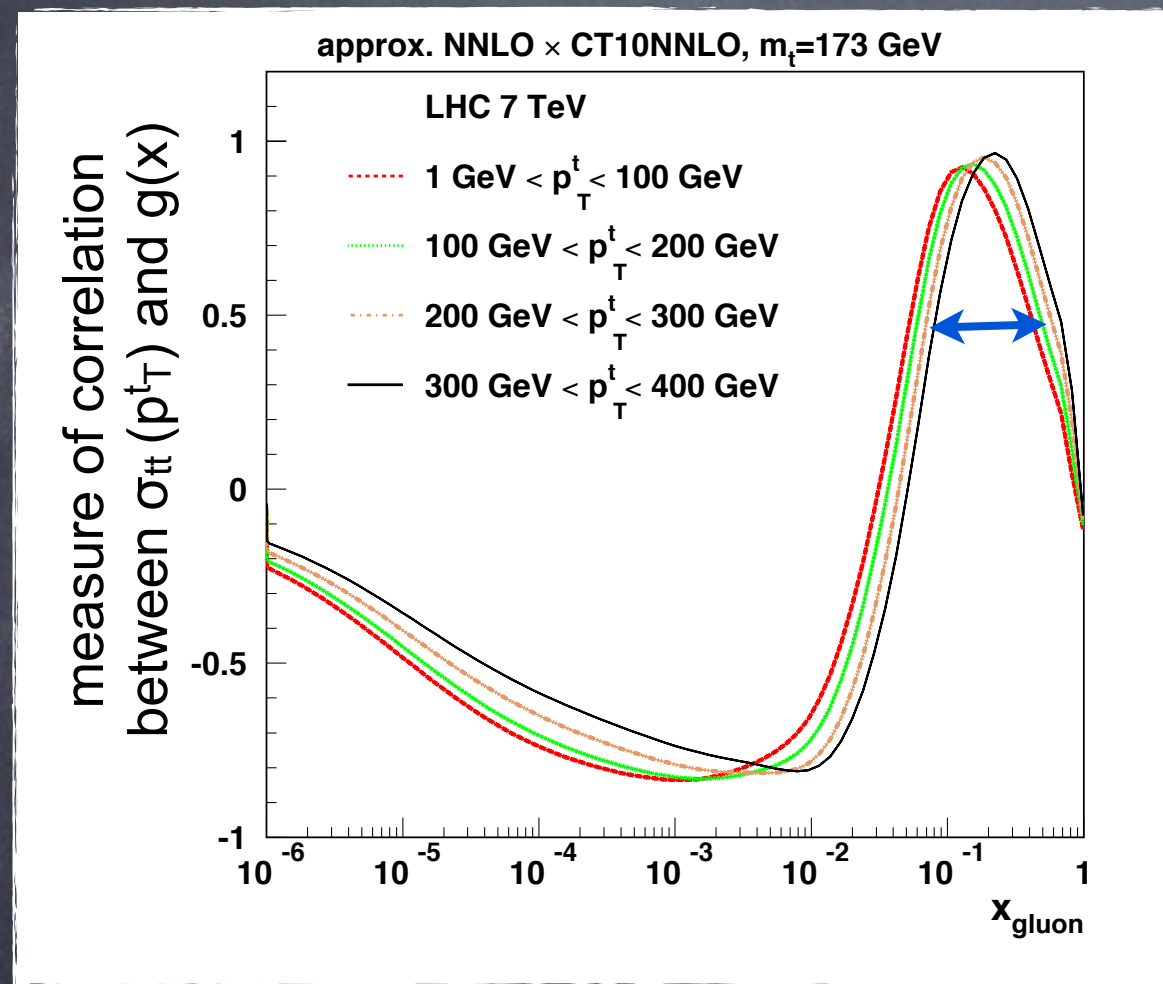


top-quark pair production reaches x-range, poorly known due to lack of experimental data

Probing gluon with $t\bar{t}$ differential cross sections

high sensitivity of $\sigma_{t\bar{t}}(p_T^t)$ to the gluon distribution is expected at high x

NLO calculation too imprecise...



Additional constraints from the shape of top-quark transverse momenta ?

Need a higher-order calculation suitable for a QCD analysis

Need effective description of the kinematic distributions capturing the main features of the full calculation

DiffTop: differential $t\bar{t}$ cross sections at aNNLO

Differential cross sections of heavy-quark pair production in proton- (anti)proton collisions calculated in pQCD at approximate NNLO (α_s^4) using methods of threshold resummation beyond the leading logarithmic accuracy.

Predictions for the single-particle inclusive kinematics (top-quark p_t , y_t)

Based on *N. Kidonakis, E. Laenen, S. Moch, R. Vogt Phys.Rev.D64 (2001)114001*

Details of the calculation: *M. Guzzi, KL, S.-O. Moch, JHEP 1501 (2015) 082*

Open-source code: <http://difftop.hepforge.org/>

DiffTop is hosted by Hepforge,

- Home
- Download Version 1.0.0
- User Manual
- Citations
- Contact

Welcome to DiffTop

DiffTop is the Fortran-based package, which allows the user to calculate the differential and total cross section for heavy-quark pair production at hadron colliders in One-particle inclusive (1PI) kinematics. The cross sections are calculated in perturbative QCD at approximate next-to-next-to-leading order (approx.NNLO) by using methods of threshold resummation beyond the leading logarithmic accuracy.

At present, only the simultaneous variation of the renormalisation and factorisation scales is allowed. The new version of the code will include additional terms, allowing for independent variation of the QCD scales.

The code is interfaced to the QCD analysis package **HERAFitter** via **fastNLOtoolkit**.

Marco Guzzi, Katerina Lipka, Sven-Olaf Moch [send mail to the authors](#) : difftop@projects.hepforge.org

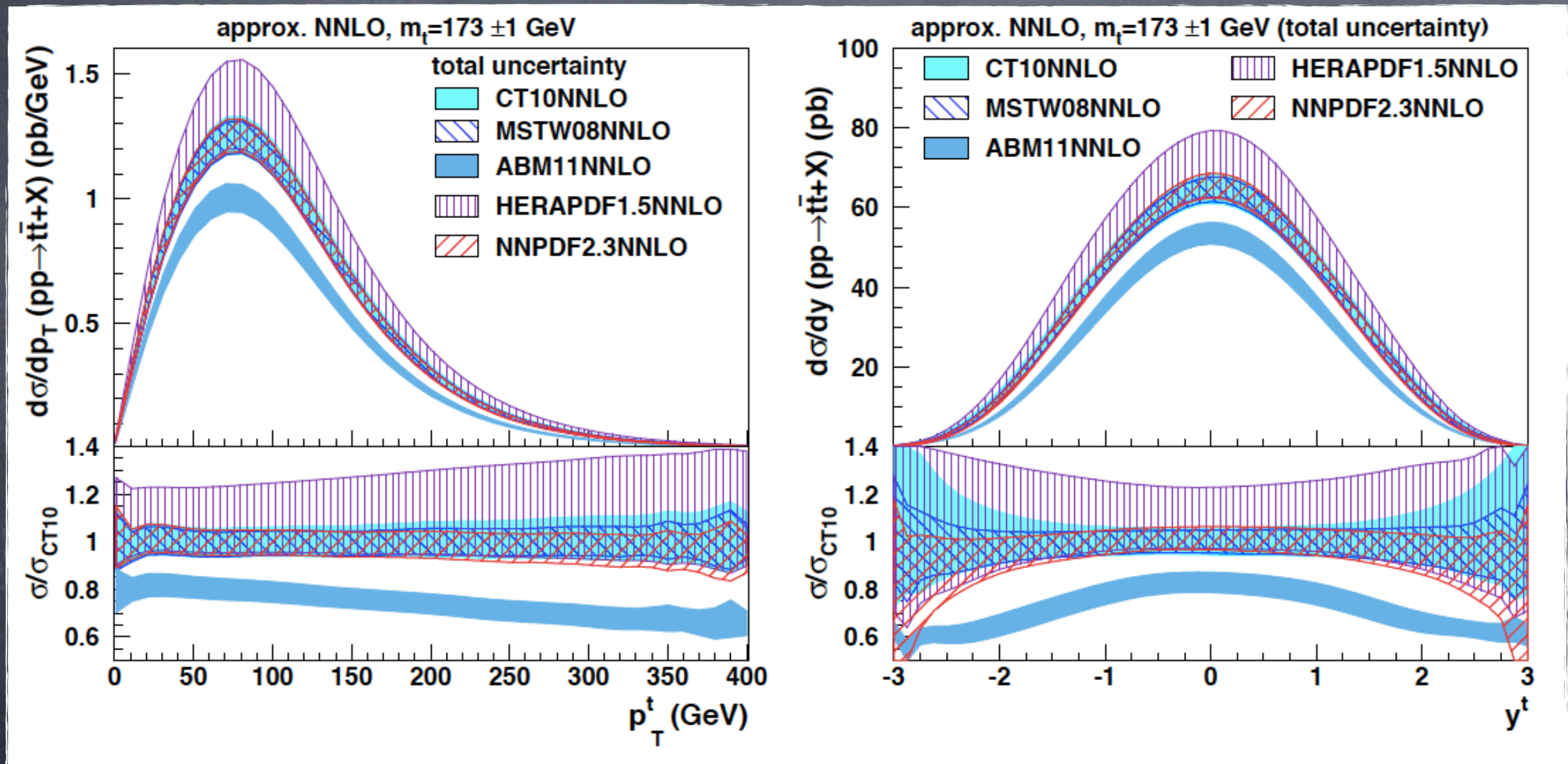
Last updated Sat 27 September 2014

Version 1.0.0: choose process, order, energy, PDFs, m_t , α_s , scales (NB: so far $\mu_r=\mu_f$)

Designed for fast QCD analyses: allows for various phenomenological studies

Phenomenology uncertainty on $t\bar{t}$ kinematics

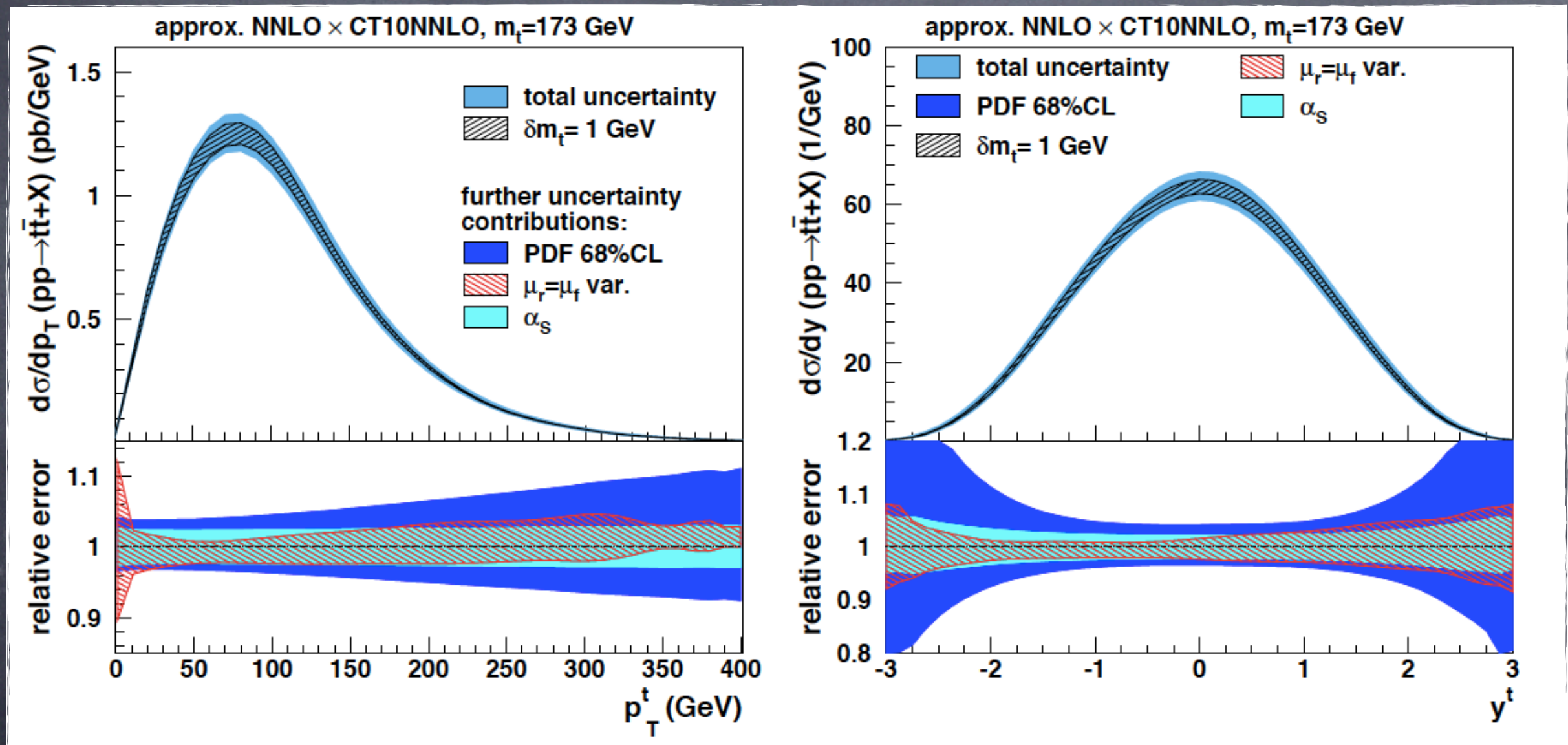
full uncertainty for the predicted $\sigma_{t\bar{t}}(p_T^t, y^t)$ **LHC 7 TeV** for different PDFs



PDF alters the normalization and shape of the predicted kinematics

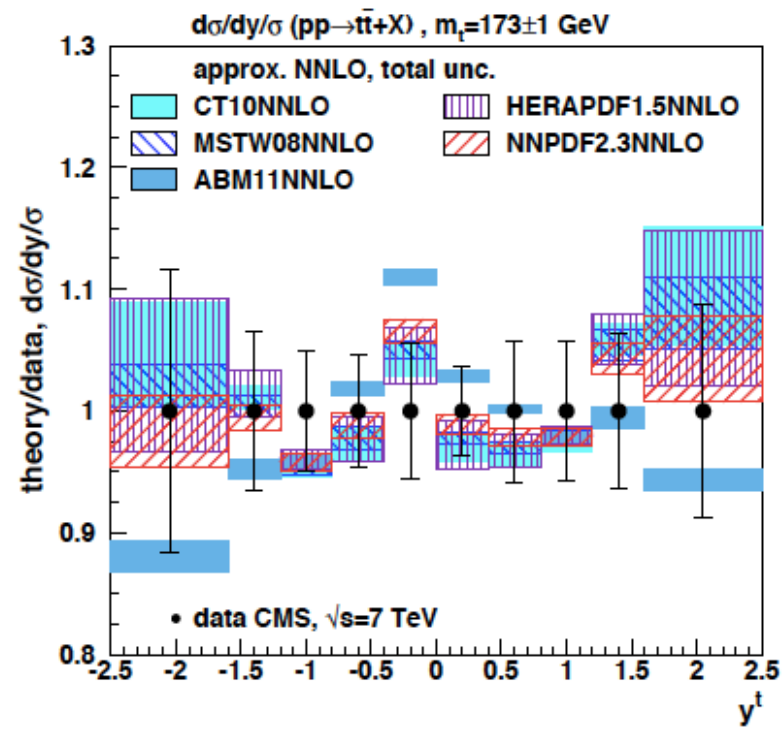
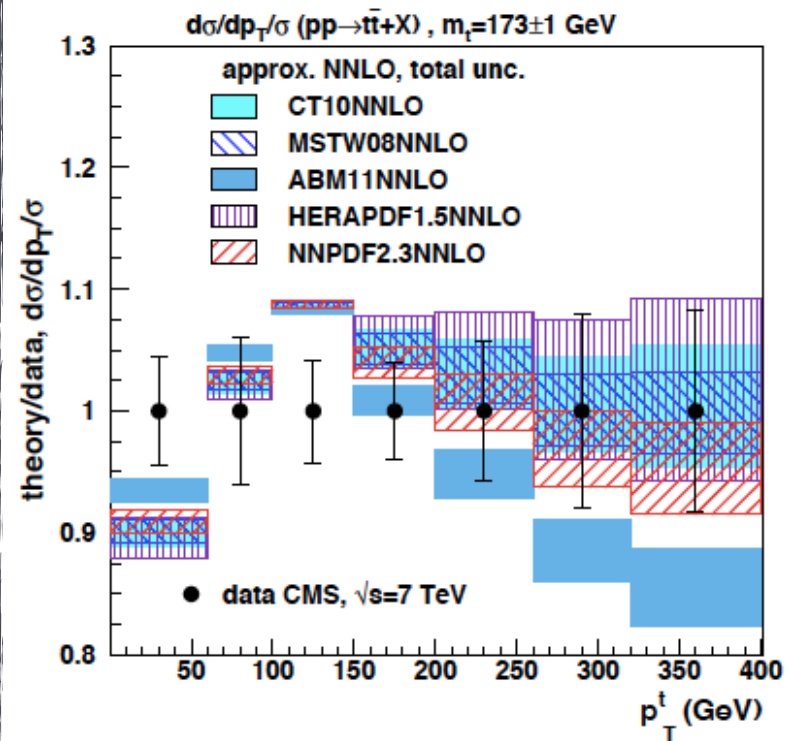
Phenomenology uncertainty on $t\bar{t}$ kinematics

breakdown of different contributions to the theory uncertainty



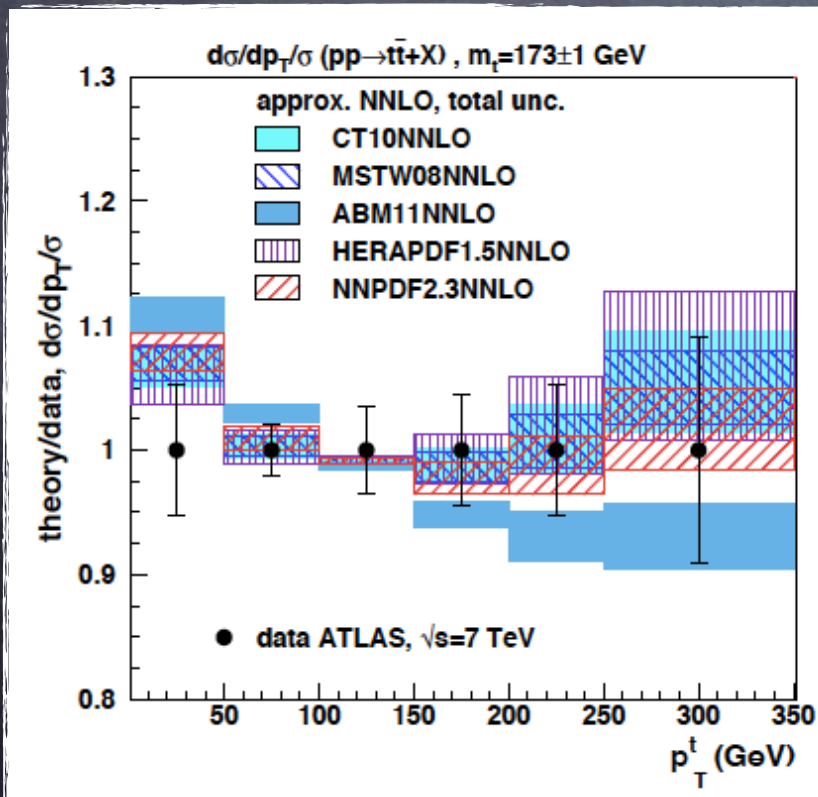
Dominant contribution to the theory uncertainty is due to PDF uncertainty

Comparison with LHC measurements: 7 TeV



CMS [EPJ C73 (2013) 2339]

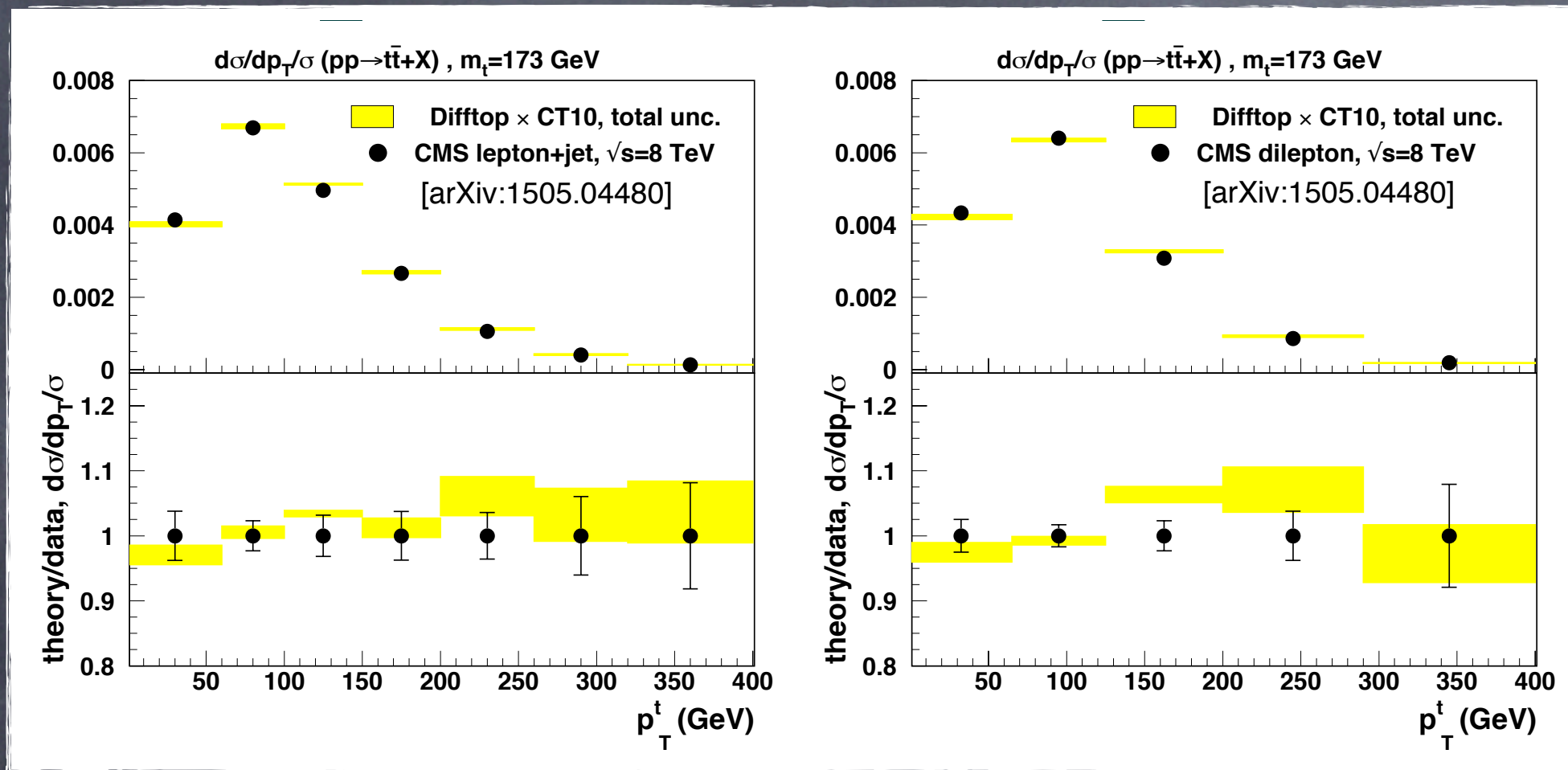
ATLAS [PRD 90 (2014) 7, 072004]



Diff-top predictions agree well with data,
agreement on different level for different PDFs

Comparison with LHC measurements: 8 TeV

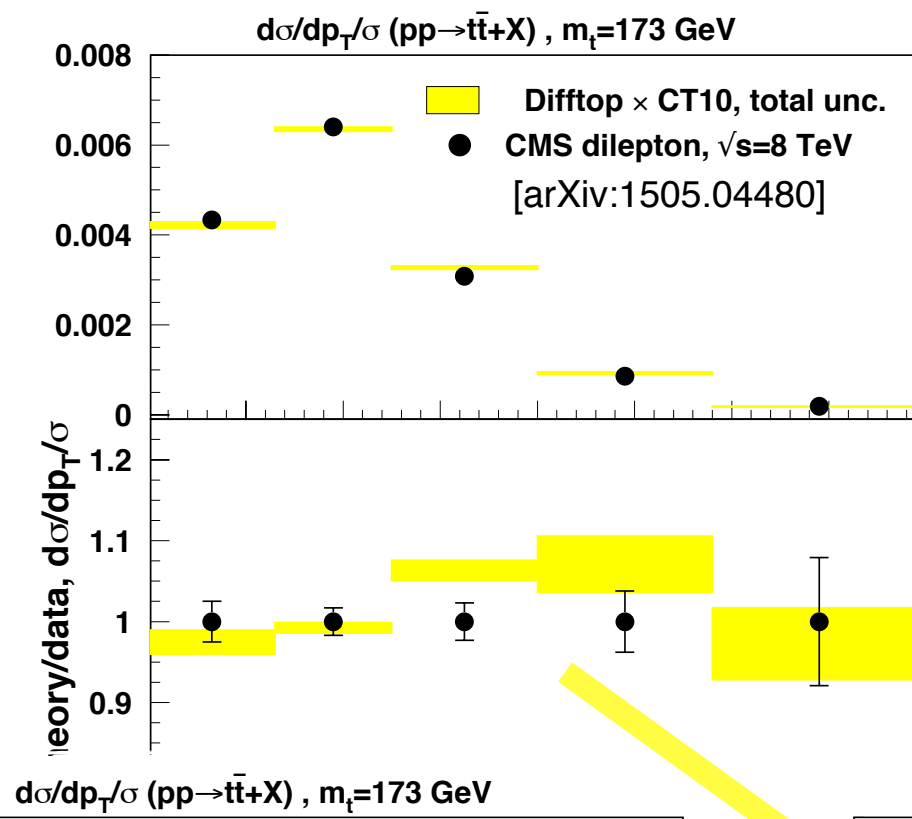
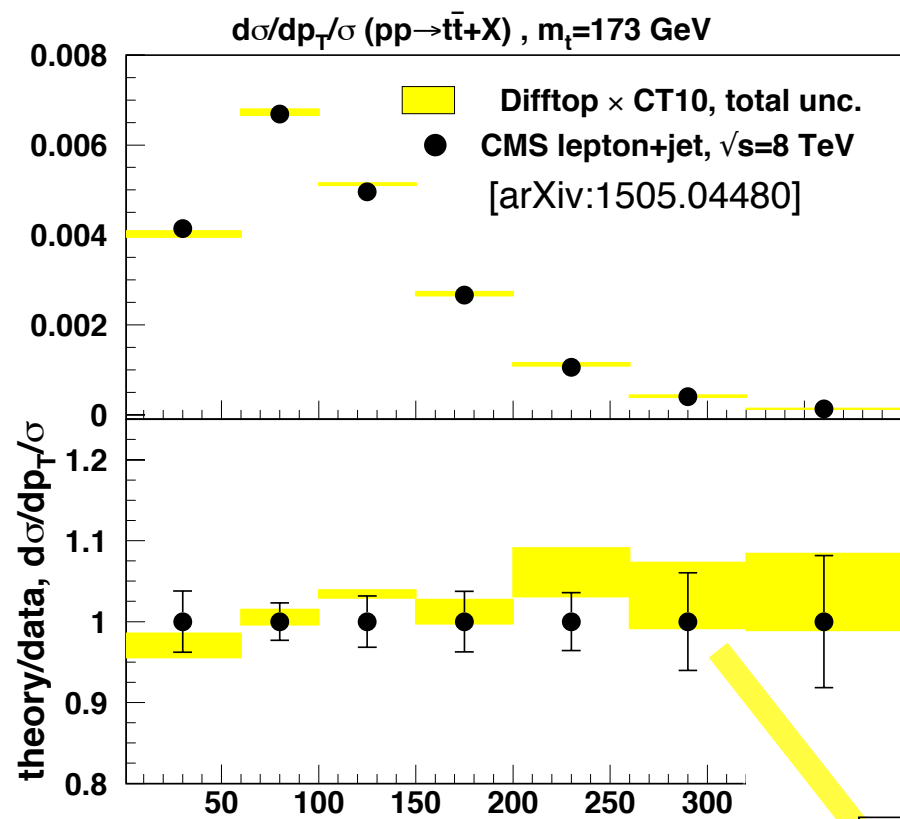
CMS Collaboration
[arXiv:1505.04480]



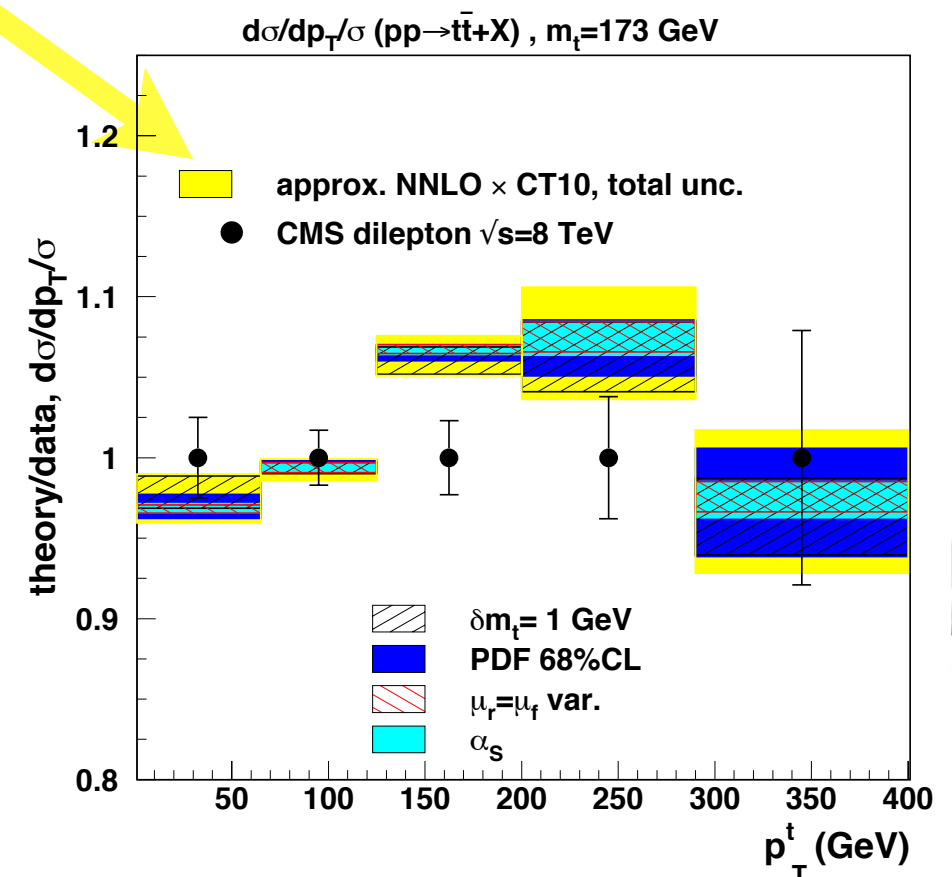
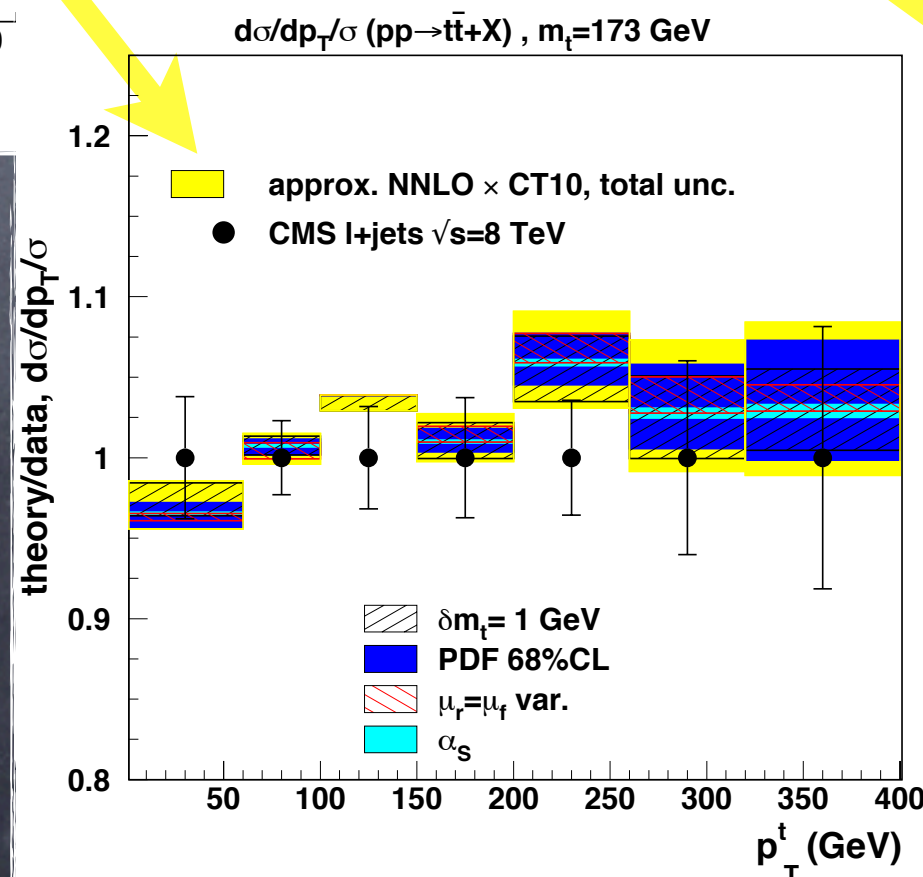
Comparison with LHC measurements: 8 TeV

CMS Collaboration
[arXiv:1505.04480]

+ break-down of
theory uncertainty:



dominant contributions:
top-quark mass variation
and PDF



QCD analysis: general strategy

Factorisation:
$$\sigma = \sum_k \sigma_k \times f_k$$

σ_k - Wilson coefficient

k - number of active flavours in the proton

PDF for flavor k : $f_k = f_k(x, Q^2)$

Q^2 - dependence predicted by QCD

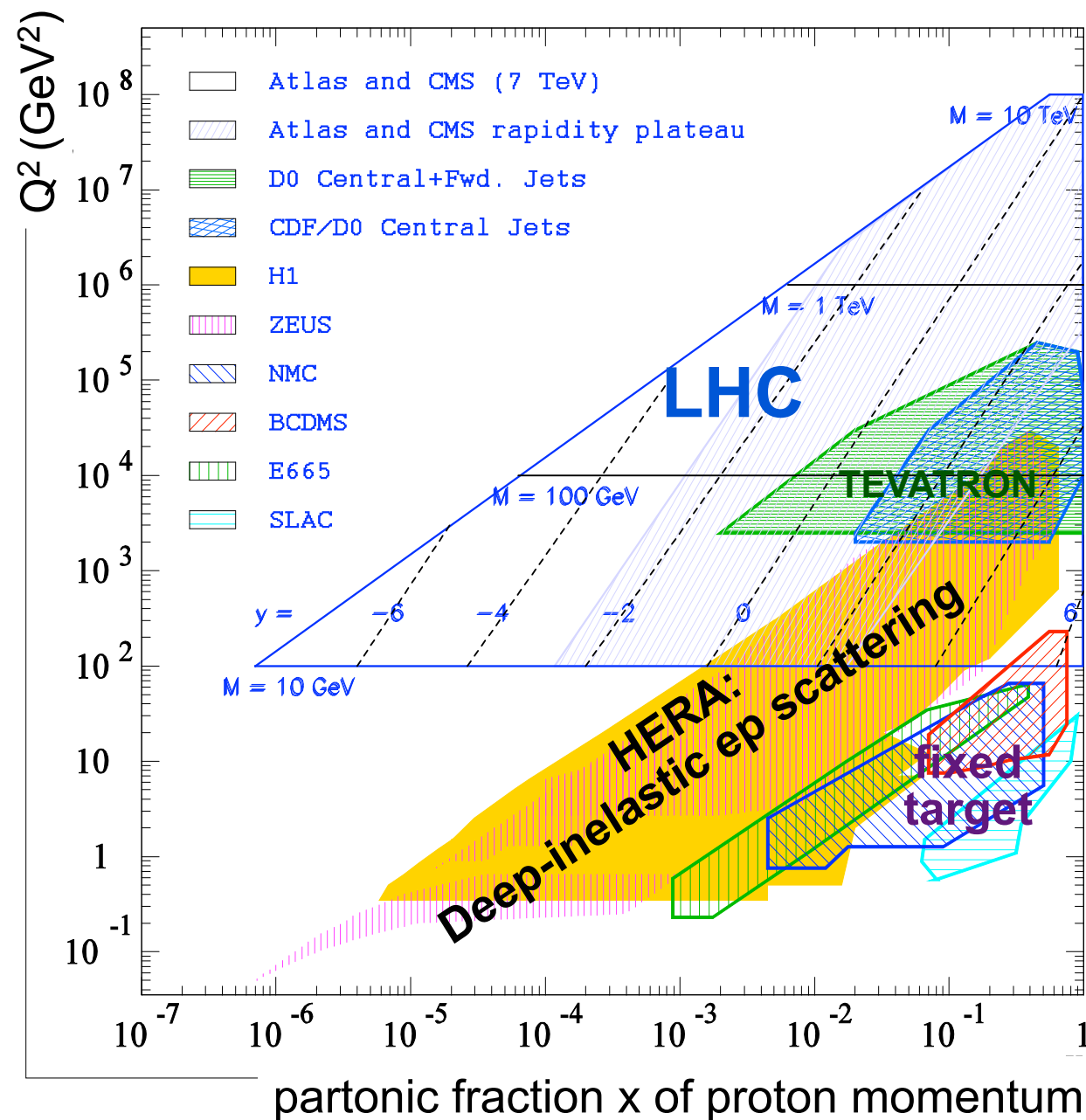
x - dependence determined from data

PDF determination

- parameterize PDFs at a starting scale Q^2_0 : $f(x) = Ax^B(1-x)^C(1+Dx+Ex^2)$
A: normalization, **B**: small- x behavior, **C**: $x \rightarrow 1$ shape
- evolve these PDFs to $Q^2 > Q^2_0$
- construct cross sections from PDFs and coefficients:
predictions for every data point in (x, Q^2) – plane
- χ^2 - fit to the experimental data

Experiments sensitive to PDFs

Measurements probing proton structure



- HERA DIS: quarks, gluon @ low, medium x

- HERA heavy-quarks: gluon, m_c , m_b

- Neutrino scattering: s-quark @ high x

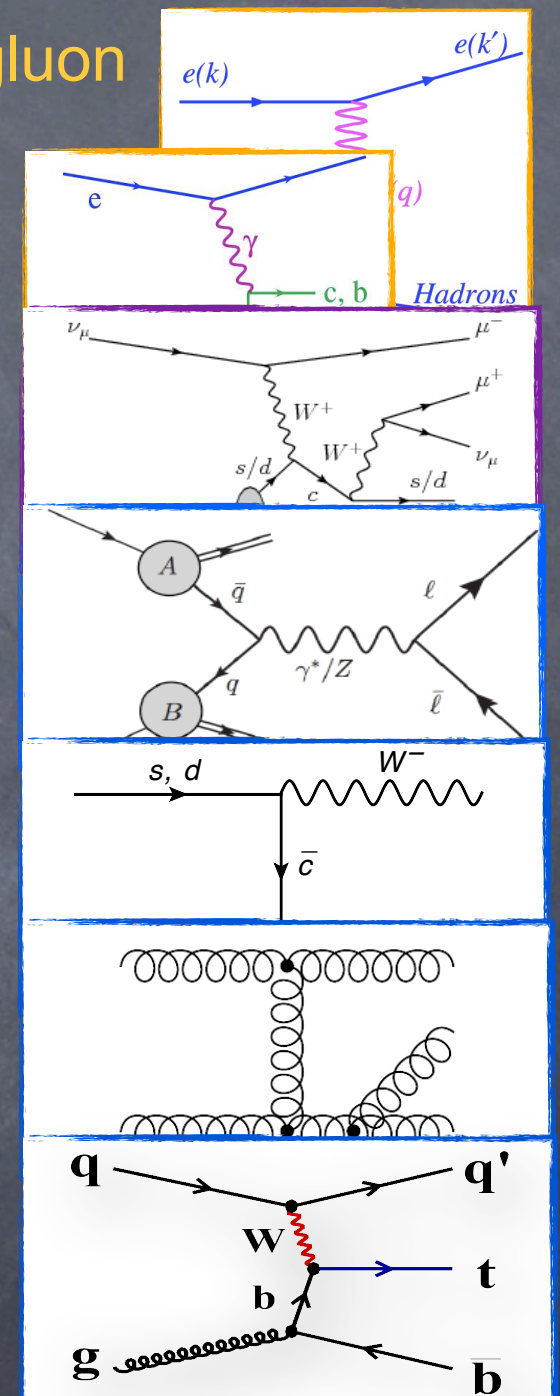
- LHC W,Z: light quarks at low and high x

- LHC W+c: s-quark medium x

- LHC jets: gluon at medium x

- LHC single top u, d and b quarks

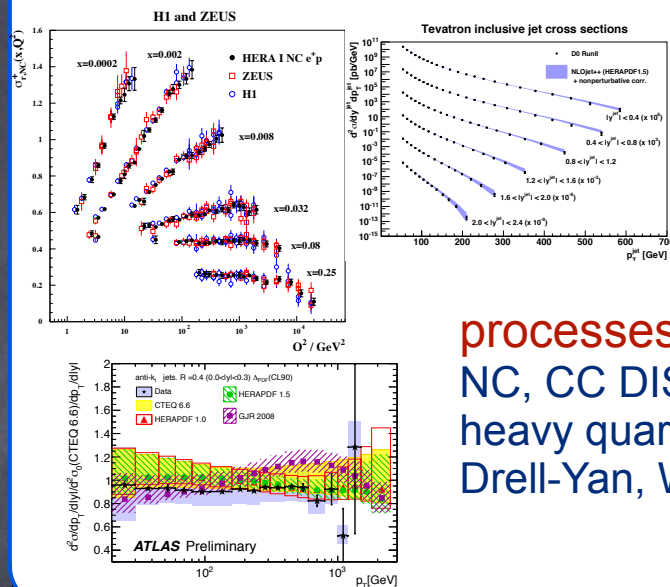
- LHC top-pairs: gluon at high x



HERAFitter QCD analysis framework

open-source code to test impact of the measurements on e.g. PDFs during data analysis

experimental input



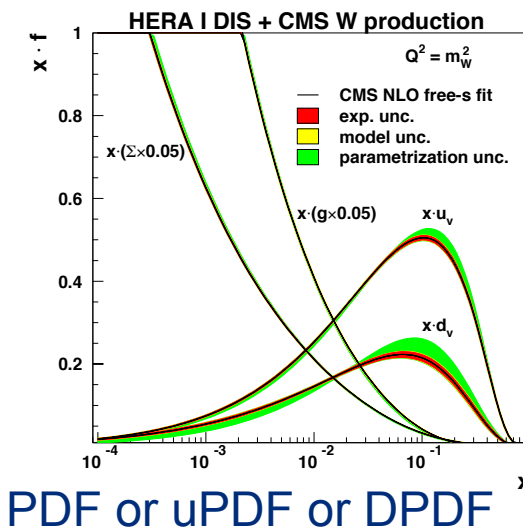
experiments:
HERA, Tevatron,
LHC, fixed target

processes:
NC, CC DIS, jets, diffraction,
heavy quarks (c,b,t)
Drell-Yan, W production

theoretical calculations/tools

Heavy quark schemes: MSTW, CTEQ, ABM
 Jets, W, Z production: fastNLO, Applgrid
 Top-quark production: NNLO (Hathor, Difftop)
 QCD Evolution: DGLAP (QCDNUM)
 Alternative tools: k_T factorisation
 Other models: NNPDF reweighting
 Dipole model
 + Different error treatment models
 + Tools for data combination (HERAaverager)

HERAFitter



$\alpha_S(M_Z), m_c, m_b, m_t, f_s, \dots$

Theory predictions

Benchmarking

Comparison of schemes

EPJ C75 (2015) 7, 304

<https://www.herafitter.org/HERAFitter>

Difftop interfaced to HERAFitter
via fastNLOtoolkit

Data in QCD analysis:

Deep Inelastic Scattering in ep collisions, combined HERA I *JHEP* 1001:109 (2010)

CMS muon charge asymmetry at $\sqrt{s}=7$ TeV ($L=4.7$ fb $^{-1}$) *Phys. Rev. D* 90 (2014) 032004

+ Top-quark pair production cross-sections:

total {

CMS Collaboration, *JHEP*11 (2012) 067; CMS-PAS-TOP-12-007

ATLAS Collaboration, *ATLAS-CONF-2012-024*, *ATLAS-CONF-2012-149*

CDF Collaboration, *CDF Conference Note* 9913 (2009)

differential {

$\sqrt{s}=7$ TeV: CMS Collaboration *Eur. Phys. J. C*73 (2013) 2339

ATLAS Collaboration *Phys.Rev. D*90 (2014) 7, 072004

NB: top-pair production at TEVATRON qq-dominated, need better constraint on light quarks
→ use CMS W asymmetry data

QCD analysis framework :

Use HERAFitter, DGLAP evolution at NNLO, parametrization at Q_0 :

$$\begin{aligned}
 xu_v(x) &= A_{u_v} x^{B_{u_v}} (1-x)^{C_{u_v}} (1 + D_{u_v}x + E_{u_v}x^2) & x \bar{U} &= x \bar{u} \\
 xd_v(x) &= A_{d_v} x^{B_{d_v}} (1-x)^{C_{d_v}} & x \bar{D} &= x \bar{d} + x \bar{s} \\
 x\bar{U}(x) &= A_{\bar{U}} x^{B_{\bar{U}}} (1-x)^{C_{\bar{U}}} & B_{\bar{U}} &= B_{\bar{D}} \\
 x\bar{D}(x) &= A_{\bar{D}} x^{B_{\bar{D}}} (1-x)^{C_{\bar{D}}} & A_{\bar{U}} &= A_{\bar{D}}(1 - f_s) \\
 xg(x) &= A_g x^{B_g} (1-x)^{C_g} + A'_g x^{B'_g} (1-x)^{C'_g} & f_s &= \bar{s}/(\bar{d} + \bar{s}) \equiv 0.31 \pm 0.08
 \end{aligned}$$

A_{u_v}, A_{d_v}, A_g are determined by QCD sum rules

Experimental (Hessian) uncertainties:

originate from uncertainties of the data considered, criterion $\Delta\chi^2=1$ is applied

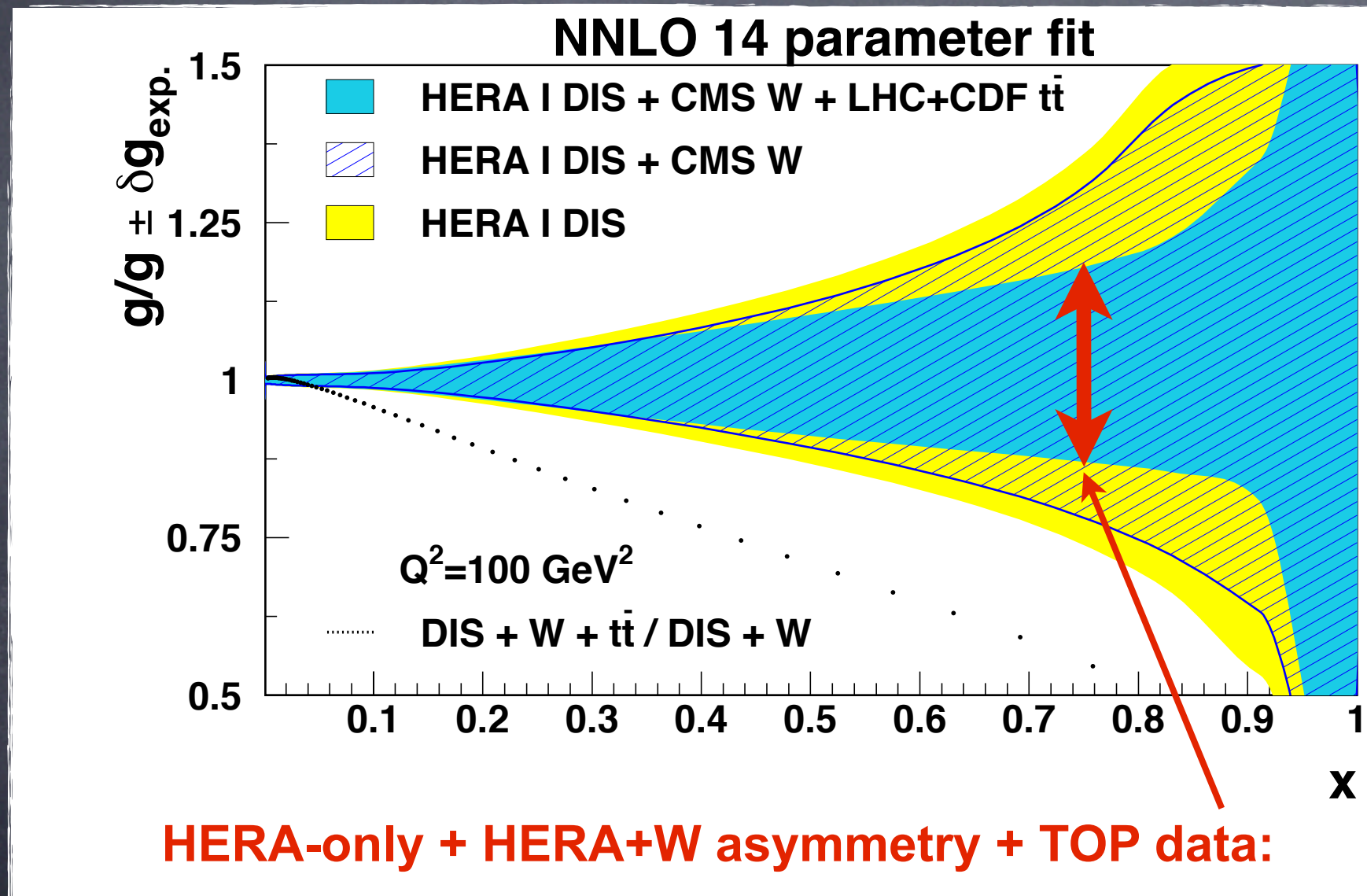
Details of the theory calculations and model input

- Calculation for W production via MCFM, interfaced via APPLGRID, K-factors by using FEWZ
- Calculation for top-quark pair production via fastNLO x DiffTop
- Starting scale of PDF evolution, $Q_0^2 = 1.9 \text{ GeV}^2$
- Heavy quark treatment: general mass variable flavor number (GMVFN) scheme by Thorne-Roberts (TR),
 $m_c = 1.4 \text{ GeV}, m_b = 4.75 \text{ GeV}.$

Impact on the gluon distribution

HERA-only vs HERA+W asymmetry:

- very similar effects on u and d-distributions as in CMS NLO fit *PRD 90 (2014) 032004*
- slight improvement in $g(x)$ only due to Sum Rules



NB:
only experimental
uncertainty shown

- Moderate improvement of the uncertainty on the gluon distribution for $x > 0.1$
- Significant change of the shape of the gluon distribution

Summary

Differential cross sections of top-pair production can be used in PDF fits @ NNLO

Data are yet not very precise, expect more effect on $g(x)$ using Run II measurements

The correlations between the total and differential cross sections are NOT used since those are NOT yet available

Needs from experimental measurements for PDF fits for RUN II data:

✓ absolute differential cross sections with list of sources of uncorrelated and correlated statistical and systematic uncertainties

or

✓ normalized differential cross sections with list of sources of uncorrelated and correlated statistical and systematic uncertainties
+ provided correlations with the total cross section

Back - up

Prediction: 1 PI kinematics, threshold approximation

Single-particle inclusive (1PI) kinematics

In our calculation, heavy-quark hadroproduction near the threshold is approximated by considering the partonic subprocesses

$$\begin{aligned} q(k_1) + \bar{q}(k_2) &\rightarrow t(p_1) + X[\bar{t}](p'_2), \\ g(k_1) + g(k_2) &\rightarrow t(p_1) + X[\bar{t}](p'_2) \quad p'_2 = \bar{p}_2 + k, \end{aligned} \quad (1)$$

where k is any additional radiation, and $s_4 = p'_2{}^2 - m_t^2 \rightarrow 0$ momentum at the threshold.

This kinematic is used to determine the p_T^t and rapidity y^t distribution of the final-state top-quark.

Hard scattering functions are expanded in terms of

$$\left[\frac{\ln^l(s_4/m_t^2)}{s_4} \right]_+ = \lim_{\Delta \rightarrow 0} \left\{ \frac{\ln^l(s_4/m_t^2)}{s_4} \theta(s_4 - \Delta) + \frac{1}{l+1} \ln^{l+1} \left(\frac{\Delta}{m_t^2} \right) \delta(s_4) \right\},$$

where corrections are denoted as leading-logarithmic (LL) if $l = 2i + 1$ at $\mathcal{O}(\alpha_s^{i+3})$ with $i = 0, 1, \dots$, as next-to-leading logarithm (NLL) if $l = 2i$, as next-to-next-to-leading logarithm (NNLL) if $l = 2i - 1$, and so on.

Prediction: 1 PI kinematics, threshold approximation

The hard scattering expansion

The factorized differential cross section is written as

$$S^2 \frac{d^2\sigma(S, T_1, U_1)}{dT_1 dU_1} = \sum_{i,j=q,\bar{q},g} \int_{x_1^-}^1 \frac{dx_1}{x_1} \int_{x_2^-}^1 \frac{dx_2}{x_2} f_{i/H_1}(x_1, \mu_F^2) f_{j/H_2}(x_2, \mu_F^2) \\ \times \omega_{ij}(s, t_1, u_1, m_t^2, \mu_F^2, \alpha_s(\mu_R^2)) + \mathcal{O}(\Lambda^2/m_t^2),$$

$$\omega_{ij}(s_4, s, t_1, u_1) = \omega_{ij}^{(0)} + \frac{\alpha_s}{\pi} \omega_{ij}^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2 \omega_{ij}^{(2)} + \dots$$

where $\omega_{ij}^{(2)}$ at parton level in 1PI is

$$\omega_{ij}^{(2)} = s^2 \frac{\hat{\sigma}_{ij}^{(2)}}{du_1 dt_1} \Big|_{1PI} = F_{ij}^{Born} \frac{\alpha_s^2(\mu_R^2)}{\pi^2} \left\{ D_{ij}^{(3)} \left[\frac{\ln^3(s_4/m_t^2)}{s_4} \right]_+ \right. \\ \left. + D_{ij}^{(2)} \left[\frac{\ln^2(s_4/m_t^2)}{s_4} \right]_+ + D_{ij}^{(1)} \left[\frac{\ln(s_4/m_t^2)}{s_4} \right]_+ + D_{ij}^{(0)} \left[\frac{1}{s_4} \right]_+ + R_{ij}^{(2)} \delta(s_4) \right\}.$$

Prediction: 1 PI kinematics, threshold approximation

Few more details...

- ▶ Hard and soft functions: $H_{ij} = H_{ij}^{(0)} + (\alpha_s/\pi)H_{ij}^{(1)} + \dots$ and $S_{ij} = S_{ij}^{(0)} + (\alpha_s/\pi)S_{ij}^{(1)} + \dots$, $H_{ij}^{(2)}$ and $S_{ij}^{(2)}$ are set to zero.
- ▶ Soft anomalous dimension matrices:
 $\Gamma_S = (\alpha_s/\pi)\Gamma_S^{(1)} + (\alpha_s/\pi)^2\Gamma_S^{(2)} + \dots$
In our calculation, $\Gamma_S^{(2)}$ at two-loop for the massive case is included. Becher (2009), Kidonakis (2009).
- ▶ Anomalous dimensions of the quantum fields $i = q, g$:
 $\gamma_i = (\alpha_s/\pi)\gamma_i^{(1)} + (\alpha_s/\pi)^2\gamma_i^{(2)} + \dots$
- ▶ The Coulomb interactions, due to gluon exchange between the final-state heavy quarks, are included at 1-loop level.

Prediction: 1 PI kinematics, threshold approximation

Matching

The matching conditions are determined by comparing the expansion in the Mellin moment space to the exact results for the partonic cross section.

Matching terms at NLO

$$Tr\{H^{(1)}S^{(0)} + H^{(0)}S^{(1)}\} \quad (2)$$

are included. Beenakker, Kuijf, Van Neerven, Smith (1989), Beenakker, Van Neerven, Meng, Schuler, Smith (1991), Mangano, Nason, Ridolfi (1992).

Matching terms at NNLO

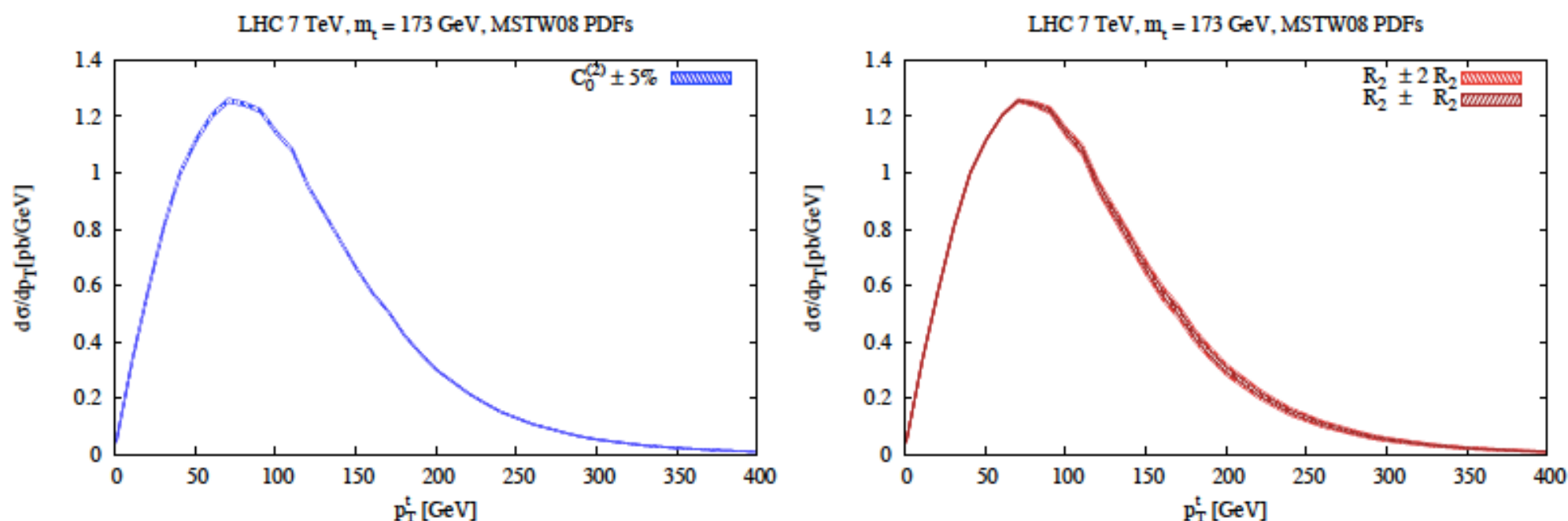
$$Tr\{H^{(1)}S^{(1)}\}, Tr\{H^{(0)}S^{(2)}\}, Tr\{H^{(2)}S^{(0)}\} \quad (3)$$

are set to zero.

Prediction: 1 PI kinematics, threshold approximation

Systematic uncertainties due to missing terms

The uncertainties due to missing contributions in $D_{ij}^{(0)}$ and R_2 are part of the systematic uncertainty associated to approximate calculations of this kind which are based on threshold expansions.



Left: The coefficient $C_0^{(2)}$ (scale ind. contribution in $D_{ij}^{(0)}$) is varied within its 5% while R_2 is kept fixed. Right: here the coefficient R_2 is varied by adding and subtracting $2R_2$ while $C_0^{(2)}$ is kept fixed.

Prediction: 1 PI kinematics, threshold approximation

QCD Threshold expansions: “pros and cons”

Approximate calculations based on threshold expansions are not perfect, but can be (easily) highly improved once the full NNLO calculation will be available.

- 😊 provide a local effective description of the p_T and y distributions that captures the main features of the full calculation.
- 😊 relatively easy interface to FASTNLO or APPLGRID.
- 😊 provide a fast tool for taking into account correlations (α_s , m_t , gluon); easy to implement different heavy-quark mass definitions. Dowling, Moch (2014)
- 😞 Very sensitive to the missing contribution in $D^{(0)}$ and R_2 .
- 😞 Scale uncertainty is also affected (at approx NNLO is underestimated at the moment. We'll improve on this)
- 😞 At the moment the description is valid near the threshold.