

DIFFERENCES OF MASSES AND WIDTHS OF THE CHARGED AND NEUTRAL $\rho(770)$, $\rho(1450)$, $\rho(1700)$ MESONS FROM DATA ON ELECTRO-WEAK PROCESSES E. Bartoš^{1,*}, S. Dubnička¹, A.-Z. Dubničková², H. Hayashii³

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Abstract

The $\rho(770)$, $\rho(1450)$, $\rho(1700)$ mesons exist in three charged states ρ^0 , ρ^+ and ρ^- , whereby masses of positively charged mesons are identical with masses of negatively charged mesons, due to the CPT theorem. However, there is no reason for the identity of charged meson masses with neutral meson masses. For determination of differences of masses and decay widths of charged and neutral $\rho(770)$, $\rho(1450)$, $\rho(1700)$ mesons are employed the data on $e^+e^- \rightarrow \pi^+\pi^-$ an $\tau^- \rightarrow \pi^-\pi^0\nu_{\tau}$ processes to be analyzed by the Unitary and Analytic models of the electromagnetic and weak pion form factors, respectively.

Parameters of $\rho(770)$ $\Gamma_{ ho^0} - \Gamma_{ ho^\pm} \, [{ m MeV}]$ Processes $m_{ ho^0}-m_{ ho^\pm} \, [{ m MeV}]$ Processes 2012: PDG average $3.6 \pm 1.8 \pm 1.7$ $e^+e^- \to \pi^+\pi^-\pi^0$ 2.4 ± 2.1 $\pi^{\pm} p \rightarrow \rho N$ Our result $au^- ightarrow \pi^- \pi^0 u_{ au}$ -0.2 ± 1.0 -5 ± 5 $\overline{p}p \to \rho\pi$ 2005 Scha 2003: Aloisid $\pi^- p \rightarrow \rho N$ -4 ± 4



Weak and EM FF

Conserved Vector Current hypothesis [1]: relation of spectral functions for systems $\pi^-\pi^0$ and $\pi^+\pi^-$

$$\sigma_{\overline{\nu}ee^{-} \to \pi^{-}\pi^{0}}(t) = \frac{\pi\alpha^{2}}{6t} \left(1 - \frac{4m_{\pi}^{2}}{t}\right)^{3/2} \left|F_{\pi}^{W}(s)\right|^{2}, \qquad F_{\pi}^{W}(s) = \sqrt{2} F_{\pi}^{E,I=1}(s)$$
$$\sigma_{e^{+}e^{-} \to \pi^{+}\pi^{-}}(t) = \frac{\pi\alpha^{2}}{3t} \left(1 - \frac{4m_{\pi}^{2}}{t}\right)^{3/2} \left|F_{\pi}^{E,I=1}(s) + Re^{i\phi} \frac{m_{\omega}^{2}}{m_{\omega}^{2} - t - im_{\omega}\Gamma_{\omega}}\right|^{2}$$

Comparison of fits

U&Amodel of pion

Main properties of Unitary & Analytic model of π

F_π(t) – analytic function in the whole complex t-plane besides the cut 4m²_π → ∞
|F_π(t)|_{t→∞} ~ t^{n_q-1} ~ t⁻¹ – asymptotic behavior of FF (n_q – number of quarks)
F_π(0) = 1, normalization in t = 0, next unitarity condition

 $\frac{1}{2\mathrm{i}} \left\{ \left\langle \pi^{+} \pi^{-} | J_{\mu}^{\mathrm{EM}}(0) | 0 \right\rangle \left\langle 0 | J_{\mu}^{\mathrm{EM}}(0) | \pi^{+} \pi^{-} \right\rangle^{\star} \right\} = \sum_{n} \left\langle \pi^{+} \pi^{-} | T^{+} | n \right\rangle \left\langle n | J_{\mu}^{\mathrm{EM}}(0) | 0 \right\rangle$ with reality condition $(\mathrm{F}_{\pi}(t))^{\star} = \mathrm{F}_{\pi}(t^{\star})$ leads to elastic unitarity condition for form factor: $\mathrm{Im} \, \mathrm{F}_{\pi}(t) = (A_{1}^{1}(t))^{\star} \mathrm{F}_{\pi}(t) + \sigma(t)$



 $A_1^1(t)$ – P wave isovector elastic $\pi\pi$ scattering amplitude, $\sigma(t)$ – higher contributions





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$m_{ ho(770)}$	758.23 ± 0.46	761.60 ± 0.95	-3.37 ± 1.06
$m_{ ho(1450)}$	1342.31 ± 46.62	1373.83 ± 11.37	-31.53 ± 47.99
$m_{ ho(1700)}$	1718.50 ± 65.44	1766.80 ± 52.36	-48.30 ± 83.81
$\Gamma_{ ho(770)}$	144.56 ± 0.80	139.90 ± 0.46	4.66 ± 0.85
$\Gamma_{ ho(1450)}$	492.17 ± 138.38	340.87 ± 23.84	151.30 ± 140.42
$\Gamma_{ ho(1700)}$	489.58 ± 16.95	414.71 ± 119.48	74.87 ± 120.67

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