▲ロト ▲帰 ト ▲ ヨ ト ▲ ヨ ト ・ ヨ ・ の Q ()

# Dark Matter Self-Interactions via Collisionless Shocks in Cluster Mergers

Christian Spethmann

**EPS-HEP** Conference,

Vienna, 23.07.2015

based on arXiv: 1504.04371 M. Heikinheimo, M. Raidal, C.S. and H. Veermäe (accepted for publication by PLB)

#### Contents

#### Observations

Plasma dynamics

The model

▲□▶ ▲圖▶ ★園▶ ★園▶ - 園 - のへで

# Observational Evidence, Exhibit A: Bullet Cluster



- The two dark matter halos move through each other.
- The visible gas gets shocked and stays behind.

 $\Rightarrow$  (Most of) dark matter collisionless!

# Observational Evidence, Exhibit B: Abell 520



- The visible gas gets shocked and stays behind.
- Microlensing: Excess of dark matter on top of the gas.
- $\Rightarrow$  (some component of) dark matter collisional after all?

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

## Observational Evidence, Exhibit C: Abell 3827



- Microlensing: dark matter halos stay behind visible stars.
- $\Rightarrow$  Drag force on dark matter from intracluster medium?

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

### Preliminary Conclusions from Evidence

The observations can be explained if:

- Most of dark matter is collisionless (Bullet cluster)
- A subcomponent (10 30 %) of dark matter imitates the visible gas (Abell 520)
- The interacting component is slowed down and stays behind (Abell 3827)
- $\Rightarrow$  The interacting dark matter is a plasma like the visible gas

 $\Rightarrow$  <u>Collisionless Shocks</u>

#### Contents

#### Observations

Plasma dynamics

The model

## Plasma Physics 101

Some characteristic properties of (astrophysical) plasmas:

- Mix of charge carriers, interacting via long range forces.
- Debye shielding length  $\lambda_D = \sqrt{rac{ au}{4\pi lpha n}}$  (here  $\sim$  10 km).
- Physical size > λ<sub>D</sub> (bulk interactions dominate over surface effects).
- Collective effects dominate dynamics if  $\Lambda = \frac{4\pi}{3} \lambda_D^3 n \gg 1$ (here  $\Lambda = 10^{19}$  to  $10^{20}$ ).
- Electrostatic interactions dominate over direct  $2 \rightarrow 2$  scattering
- Plasma instabilities important (see next slides).

# Collisionless Shocks (12min talk cartoon version)



Counter-streaming plasma  $\Rightarrow$  plasma instabilities

- $\Rightarrow$  large EM fieds  $\Rightarrow$  saturation regime (non-linear)
- $\Rightarrow$  shock waves  $\Rightarrow$  energy dissipation

Typical time scale:  $\sim 100 \ \omega_p^{-1} = 100 \left(\frac{4\pi\alpha n}{m}\right)^{-1/2}$ [1502.00626 [physics.plasm-ph]].

# Are Collisionless Shocks Real?

- Observations of visible astrophysical plasmas:
  - Earth's bow shock
  - Expansion of supernova remnants
  - Galaxy collisions and cluster mergers
- Numerical Studies:
  - Particle in cell (PIC) simulations
- Experimental Studies:
  - electron-positron plasmas
  - ionized gases produced with laser pulses
- Dedicated numerical studies of dark plamas (in progress).

 $\Rightarrow$  YES, and they will affect dark matter dynamics if massless dark force carriers are present.

#### Contents

#### Observations

Plasma dynamics

The model

▲□▶ ▲圖▶ ▲≣▶ ▲≣▶ = 差 = のへで

# A Concrete Model of Dark Plasma

Details of the model:

- 70% of DM generic WIMP, 30% dark plasma.
- One Dirac fermion charged under an unbroken *U*(1) gauge group:

$$\mathcal{L} = \frac{1}{4} F_{D\mu\nu} F_D^{\mu\nu} + \bar{\chi} \left( i \not D - m_D \right) \chi.$$

- No kinetic mixing term  $F_{D\mu\nu}F^{\mu\nu}$  (highly constrained).
- Dark matter abundance is produced as thermal relic from annihilation into dark photons,  $\bar{\chi}\chi \rightarrow \gamma_D \gamma_D$ .

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

# Observational Constraints (1), BBN

- Big Bang Nucleosynthesis:  $N_{\rm eff} < 3.38$ 
  - $\Rightarrow$  Constrains temperature of dark photons during BBN.
- Dark photon temperature

$$T_D = T_\gamma \left( rac{g_{*s,\gamma}(T_\gamma)g_{*s,D}(T_*)}{g_{*s,D}(T_D)g_{*s,\gamma}(T_*)} 
ight)^{1/3},$$

(thermal equilibrium assumed at  $T_*$ ).

- Constrains number of fermions in the dark sector:  $N_D < 2.35$ 

# Observational Constraints (2), CMB/SSS

• <u>Structure formation</u> suppressed until kinetic decoupling of the dark matter and dark radiation, which occurs at

$$T_{\rm kin} = \left(rac{4\pi}{45}g_*
ight)^{rac{1}{4}} \sqrt{rac{135}{64\pi^3}} rac{m_D^{rac{3}{2}}}{\sqrt{m_P}lpha_D}$$

- If  $T_{\rm kin} > 640$  eV, only multipoles above l > 2500 are affected in the <u>CMB</u>, and thus temperatures above this limit are unconstrained by PLANCK.
- For  $T_{\rm kin} \approx 500$  eV the small scale structure is suppressed for structures below the size of  $\sim 10^9 M_{\odot}$ , alleviating the missing satellites problem.

イロト 不得 トイヨト イヨト

э

## Exclusion plot: DM mass vs. coupling constant



< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

### Outlook: Where to go from here?

- Conflicting observations of different cluster mergers, low resolution of weak lensing mass reconstruction
  - $\Rightarrow$  More <u>observations</u> required!
- Effects of dark plasma dynamics on galactic and cluster halos, structure formation etc.

 $\Rightarrow$  More <u>simulations</u> required!

- Natural explanation for the partially interacting dark matter scenario, *e.g.* partially ionized dark atoms?
  - $\Rightarrow$  More modelbuilding required!

# The Minimal Take Home Message

- 1. Collisionless shocks dictate the dynamics of rare astrophysical plasmas.
- 2. <u>Dark matter</u> coupled to <u>dark radiation</u> behaves the same way.
- 3. Considering only 2  $\rightarrow$  2 scattering is in this case not adequate.

For the full story: 1504.04371