

# Discovery potential for $T' \rightarrow tZ$ in the trilepton channel at the LHC

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Based on LB, J. Andrea, JHEP **1502** (2015) 032  
arXiv:1411.7587 [hep-ph].

# Introduction

Vector-like quarks (equal LH and RH couplings) are common to BSM theories:  
Extra Dimensions, Little Higgs Models, Composite Higgs Models

At LHC, searched for in pair-production: QCD-like (model-independent)

Single production: model-*dependent*

- allows to access underlying model
- favoured when very heavy resonances

In this talk: LHC discovery potential of singly-produced  $T'$  (top-partner)

# Simplified model

M. Buchkremer et al. **Nucl.Phys. B876**, 376 (2013) [1305.4172]

$$\mathcal{L}_{T'} = g^* \left\{ \sqrt{\frac{R_L}{1+R_L}} \frac{g}{\sqrt{2}} [\bar{T'}_L W_\mu^+ \gamma^\mu d_L] + \sqrt{\frac{1}{1+R_L}} \frac{g}{\sqrt{2}} [\bar{T'}_L W_\mu^+ \gamma^\mu b_L] + \sqrt{\frac{R_L}{1+R_L}} \frac{g}{2 \cos \theta_W} [\bar{T'}_L Z_\mu \gamma^\mu u_L] + \sqrt{\frac{1}{1+R_L}} \frac{g}{2 \cos \theta_W} [\bar{T'}_L Z_\mu \gamma^\mu t_L] \right\} + h.c.$$

We allow for generic mixing to 1<sup>st</sup> generation quarks

Only 3 parameters:

- $M_{T'}$ , the vector-like mass of the top partner
- $g^*$ , the coupling strength to SM quarks, only relevant in single production.  
Rescaling:  $\sigma \propto (g^*)^2$
- $R_L$ , the mixing coupling to first generation quarks.  $R_L = 0$  corresponds to coupling to  $t/b$  only. Rescaling: by integrating  $1^{st} \propto \frac{R_L}{1+R_L}$  and  $3^{rd} \propto \frac{1}{1+R_L}$  gen. quark processes independently

# Monte Carlo simulation details

LO samples simulation with

- parton level: MG5\_aMC@NLO (CTEQ6L1)
- Hadronisation/showering: Pythia6 Tune Z2
- FastSim: Delphes3 ma5Tune
- Analysis: MadAnalysis5

## Signal:

5 benchmark points of  $T'$  mass in steps of 200 GeV:  $M_{T'} \in [800; 1600]$  GeV, with  $g^* = 0.1$  and  $R_L = 0.5$ . No  $k$ -factors

## Backgrounds (plus up to 2 jets):

- 3 prompt leptons:  $t\bar{t}W$ ,  $t\bar{t}Z$ ,  $tZj$ , and  $WZ$
- non-prompt leptons:  $t\bar{t}$  and  $Z/W + jets$

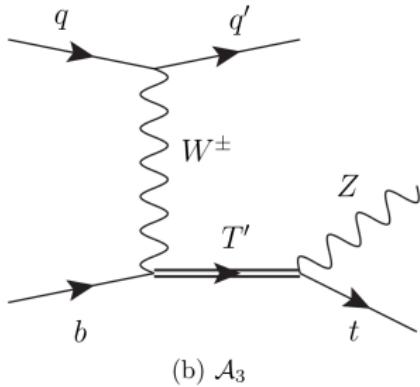
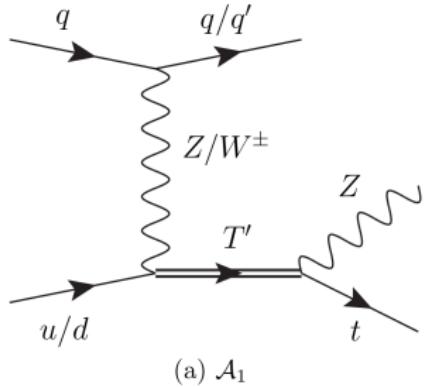
Samples normalised to NLO cross sections where available

CMS detector emulation

Anti- $k_T$  algorithm with  $R = 0.5$

b-tagging CSV medium working point:  $b$ -tag = 70%, mistag = 1%

# ANALYSES



Single production of  $T' \rightarrow tZ$  at the LHC at  $\sqrt{s} = 13$  TeV

Figures for  $\mathcal{L} = 100 \text{ fb}^{-1}$

Final state: trilepton channel,  $T' \rightarrow tZ \rightarrow (b\ell_W\nu)(\ell^+\ell^-)$

# Cut-and-count

## Objects identification

$$\begin{array}{ll} p_T(\ell) > 20 \text{ GeV}, & |\eta(e/\mu)| < 2.5/2.4, \\ p_T(j) > 40 \text{ GeV}, & \Delta R(\ell, j) > 0.4, \\ |\eta(j)| < 5, & |\eta(b)| < 2.4, \end{array}$$

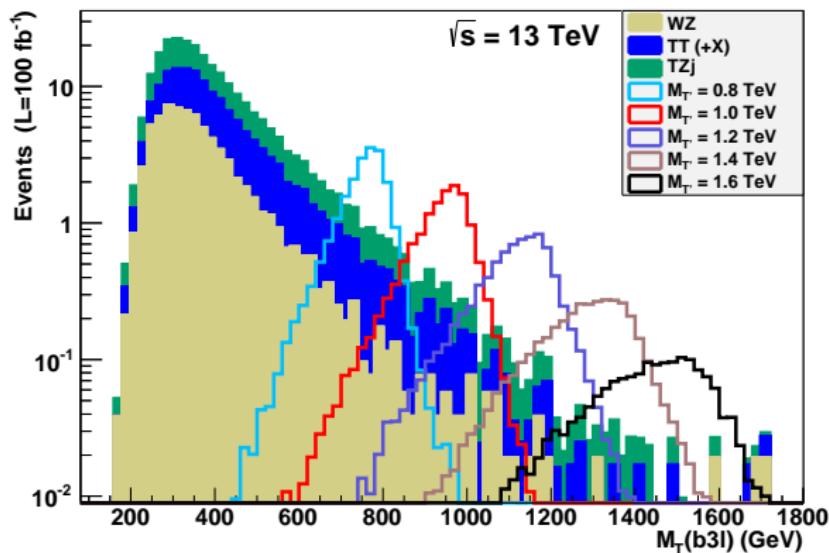
Cuts:

$n_\ell \equiv 3$	suppress $t\bar{t} + X \rightarrow \textcolor{red}{0.09\%}$
$1 < n_j < 3$	(remove pair-prod.)
$n_b \equiv 3$	suppress $WZ \rightarrow \textcolor{red}{4.2\%}$
$ M(\ell^+\ell^-)/\text{GeV} - M_Z  < 15$	$Z \rightarrow \ell^+\ell^- \text{ reco}$
$10 < M_T(\ell_W\nu)/\text{GeV} < 150$	$W \rightarrow \ell_W\nu \text{ reco}$
$0 < M_T(b\ell_W\nu)/\text{GeV} < 220$	$t \rightarrow bW \text{ reco}$

Cuts optimised to retain  $\geq 90\%$  of signal

# $M_T(b\,3\ell\,\nu)$

Signature:  $T' \rightarrow tZ \rightarrow (b\ell\nu)(\ell^+\ell^-)$



Signal clearly visible over background

Distribution in transverse mass, sharper peaks than invariant mass

Q: can we do better?

# Multi-Variate Analysis (MVA)

Cut-based strategy: suitable cuts on the most straightforward distributions  
Is it the best strategy?

Many additional variables to distinguish signal from background  
Recall the kinematics:  $T'$  very heavy,  $t - Z$  back-to-back and boosted

However, cutting on any of these variables unavoidably reduce also the signal

## Solution:

combine several variables using a *multivariate analysis* (MVA) to obtain the best signal/background discrimination

Here we used Boosted Decision Tree (BDT)

Variables drawn after  $Z$  mass cut:  $M_T(\ell_W)$ ,  $M_T(\ell_W b)$ ; MET,  $H_T$ ,  $S_T$ ;  $p_T$ ,  $\eta$ ;  $\Delta\eta$ ,  $\Delta\phi$ , angular correlations, ...

Some variables correlated, like  $p_T(Z)$  and  $p_T(\ell_1)$ : choose a reduced and uncorrelated set with still large sensitivity

# MVA: 13 variables

Variable	Importance	Variable	Importance
$M_T(b 3\ell)$	$2.60 \cdot 10^{-1}$	$\Delta R(b, \ell_W)$	$9.77 \cdot 10^{-2}$
$p_T(Z)/M_T(b 3\ell)$	$9.41 \cdot 10^{-2}$	$\Delta\varphi(t, Z)$	$8.17 \cdot 10^{-2}$
$\eta^{max}(j)$	$6.02 \cdot 10^{-2}$	$\Delta\varphi(\ell\ell _Z)$	$5.89 \cdot 10^{-2}$
$\Delta\varphi(Z, \not{p}_T)$	$5.37 \cdot 10^{-2}$	$p_T(j_1)/M_T(b 3\ell)$	$5.08 \cdot 10^{-2}$
$\Delta\eta(\ell\ell _Z)$	$5.05 \cdot 10^{-2}$	$\Delta\eta(b, \ell_W)$	$5.03 \cdot 10^{-2}$
$\eta(t)$	$4.99 \cdot 10^{-2}$	$\Delta\varphi(Z, \ell_W)$	$4.63 \cdot 10^{-2}$
$\eta(Z)$	$4.61 \cdot 10^{-2}$		

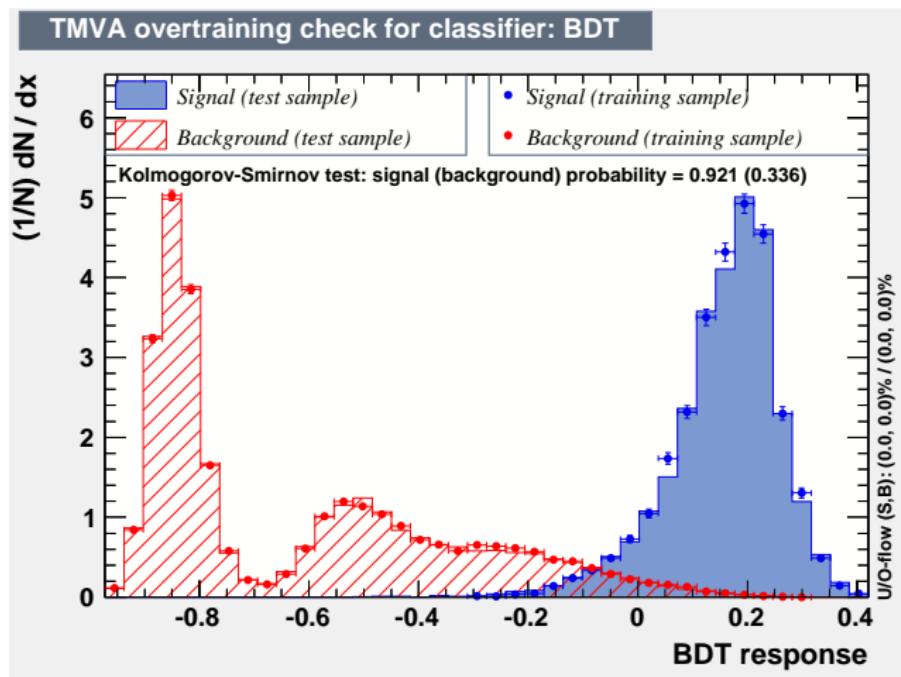
$(\ell\ell|_Z)$ : the pair of leptons reconstructing the  $Z$  boson

$\eta^{max}(j)$ : jet with largest rapidity (to account for associated jet)

$p_T(j_1)/M_T(b 3\ell)$  and  $p_T(Z)/M_T(b 3\ell)$  effectively decorrelated from  $M_T(b 3\ell)$

Angular variables from fully reconstructing the neutrino 4-momentum

# BDT output



Allows to check for “overtraining”: 2 random samples, one used for training and the other one for comparison, should get similar output

# Discovery power: benchmarks

Surviving events and significances for signal benchmark points  
( $g^* = 0.1$ ,  $R_L = 0.5$ )

- C&C: select a window around the peak in  $M_T(b3\ell)$
- MVA: perform a LH cut on BDT output

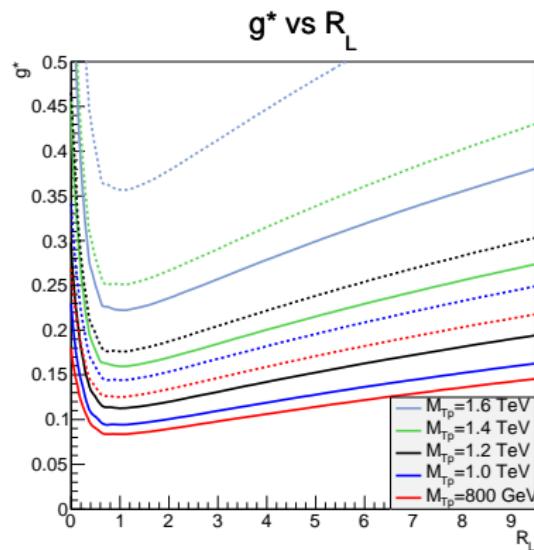
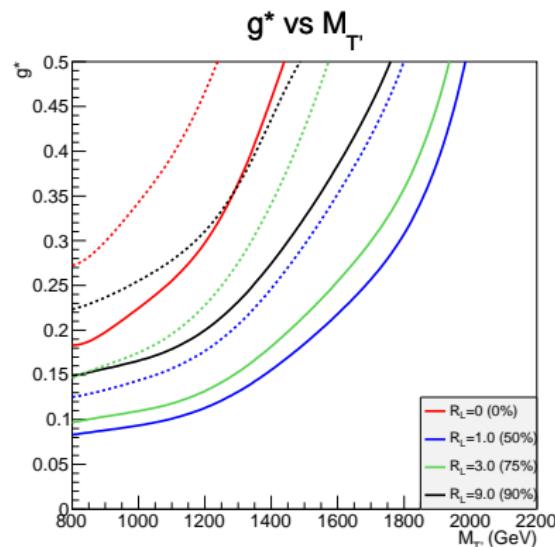
to maximise the significance:  $\sigma = S/\sqrt{S + B}$

Analysis	$M_{T'} = 0.8 \text{ TeV}$	$M_{T'} = 1.0 \text{ TeV}$	$M_{T'} = 1.2 \text{ TeV}$	$M_{T'} = 1.4 \text{ TeV}$	$M_{T'} = 1.6 \text{ TeV}$	
C&C	$M_T(b3\ell)$ cut (GeV)	[800 – 860]	[840 – 1200]	[1000 – 1340]	[1120 – 1640]	[1200 – 1800]
	S (ev.)	18.00	12.28	7.16	3.40	1.57
	B (ev.)	8.90	4.88	1.74	0.90	0.63
MVA	$\sigma$	3.47	2.96	2.40	1.64	1.06
	cut	0.07	0.08	0.11	0.12	0.12
	$\sigma$	3.64	3.10	2.50	1.62	1.15

**MVA: non-significant improvement (5%–8%)**

Significance depends on  $g^*$  and  $R_L$  per fixed  $T'$  mass

# Discovery power: parameter space

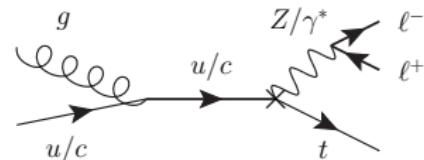
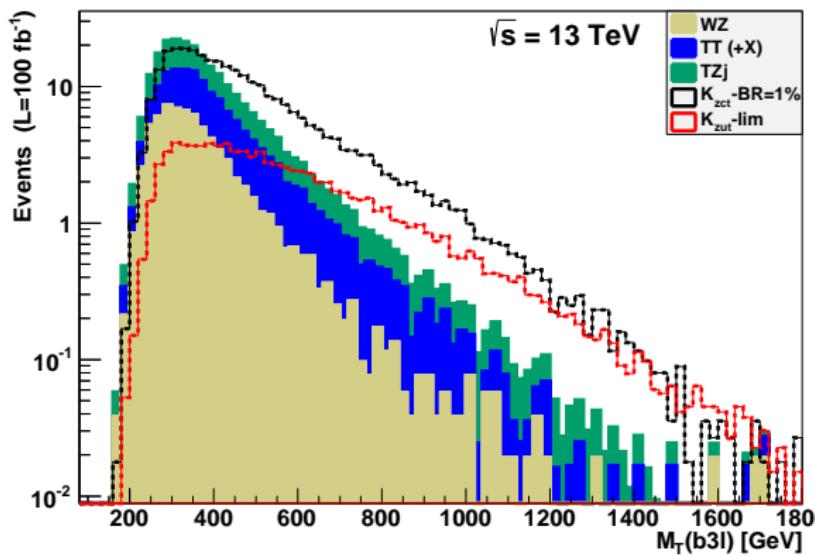


(dashed lines:  $5\sigma$ , solid lines:  $3\sigma$ )

$T'$  masses up to 2 TeV can be observed

Increased reach when  $R_L$  is non-vanishing (maximum for  $R_L \simeq 1$ , corresponding to 50%–50% mixing)

# Reinterpretation: top anomalous couplings



Cut	$\kappa_{tZu}$	$\kappa_{tZc}$
no cuts	2263(100%)	5360(100%)
$1 \leq n_j \leq 3$	1765(78.0%)	4452(83.0%)
$n_\ell \equiv 3$	191.8(10.9%)	623.3(14.0%)
$n_b \equiv 1$	113.8(59.3%)	381.0(61.1%)
Z-reco	103.2(90.7%)	342.7(90.0%)
W-reco	96.2(93.3%)	323.6(94.4%)
t-reco	91.1(94.7%)	304.7(94.1%)
$M_T(b3\ell)$	$> 400 \text{ GeV}$	$> 200 \text{ GeV}$
S	68.0	304.5
B	102.9	241.7
$\sigma$	5.2	13.0

Present limit:  $\text{BR}(t \rightarrow Zq) < 0.05\%$  (inclusive, from  $t\bar{t}$ )

MVA trained on  $T'$  signals: no improvements

In progress: training on the top anomalous signal

# Conclusions

Singlet top partners  $T'$  common to many BSM models trying to address the Higgs stability

Simplified model: only 3 parameters, simple rescaling to cover whole phase space

$T' \rightarrow tZ$ : study of the trilepton signature at  $\sqrt{s} = 13$  TeV in single production mode

$T'$  masses up to 2.0 TeV and couplings down to  $g^* = 0.1$  can be probed.  
Large gain if mixing with light generation is accounted for

Results from cut-based analysis: **simple and effective, no substantial improvements from MVA** → use cut-and-count

Reinterpretation to top anomalous couplings sharing similar signature

# Backup slides

# Examples of models

Very quick review: J. Reuter and M. Tonini, JHEP 1501 (2015) 088 [arXiv:1409.6962]

Composite Higgs models: Higgs boson is a composite state

Minimal case:  $SO(5)/SO(4)$   $\begin{cases} t_R \sim \mathbf{1}_4, \text{ complete rep. of } SO(4) \\ q_L \sim \text{incomplete rep. of } SO(5) \end{cases}$

New fermions:  $\Psi \begin{cases} \mathbf{1}_4 : T' \\ \mathbf{4}_4 : (T', B'), (X_{5/3}, X_{2/3}) \end{cases}$

Little Higgs models: Higgs is a pseudo-Goldstone boson

from a global spontaneous breaking of  $SU(5)/SO(5)$  (Littlest Higgs model)

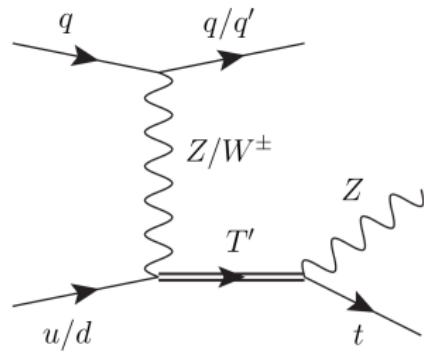
A vector-like heavy top is required to cancel loop quadratic divergences

Many models, many similarities  $\rightarrow$  simplified model

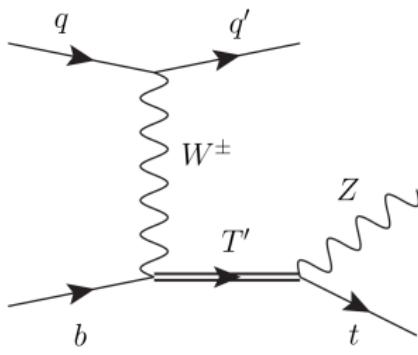
Here, **singlet top partner**:  $T'$

Typically,  $\text{BR}(T' \rightarrow qW^\pm) : \text{BR}(T' \rightarrow qZ) : \text{BR}(T' \rightarrow qh) \sim 2 : 1 : 1$

# Single production and $T' \rightarrow tZ$



(a)  $\mathcal{A}_1$



(b)  $\mathcal{A}_3$

$$\sigma_{pp \rightarrow T'}(M_{T'}, R_L) = \mathcal{A}_1(M_{T'}) \frac{R_L}{1 + R_L} + \mathcal{A}_3(M_{T'}) \frac{1}{1 + R_L}$$

$$BR_{T' \rightarrow tZ}(M_{T'}, R_L) = \mathcal{B}(M_{T'}) \frac{1}{1 + R_L}$$

$M_{T'} \text{ (GeV)}$	$\mathcal{A}_1(M_{T'}) \text{ (pb)}$	$\mathcal{A}_3(M_{T'}) \text{ (pb)}$	$\mathcal{B}(M_{T'}) \text{ (%)}$
800	1.2614	0.07242	22.4
1000	0.7752	0.03518	23.5
1200	0.5001	0.01826	24.0
1400	0.3331	0.00994	24.2
1600	0.2265	0.00561	24.4

# More simulation details

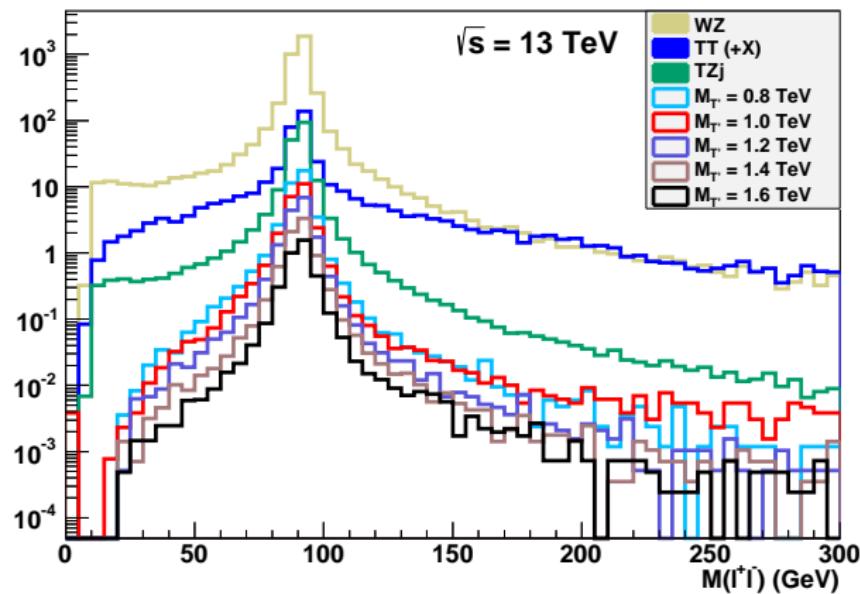
Massive background event generation to gather enough statistics:

Process	# Files	# Events	Process	# Files	# Events
SingleTop_W_madspin	189	18898481	SingleTop_s_madspin	188	18771372
SingleTop_t_5FS_madspin	83	8299246	TTdilep_WToLNu_madspin	1	64191
TTdilep_WToLLNuNu_madspin	1	99999	TTdilep_WZToLLNu_madspin	1	99991
TTdilep_ZToLL_madspin	1	99989	TTdilep_ZZToLLLL_madspin	1	99993
TTdilep_madspin	200	9427953	TTsemilep_WToLNu_madspin_1	1	59694
TTsemilep_WToLNu_madspin_2	1	59771	TTsemilep_WWToLLNuNu_madspin_1	1	99989
TTsemilep_WWToLLNuNu_madspin_2	1	99997	TTsemilep_WZToLLNu_madspin_1	2	199988
TTsemilep_ZToLL_madspin_1	1	99995	TTsemilep_ZToLL_madspin_2	1	99987
TTsemilep_ZZToLLLL_madspin_1	1	99993	TTsemilep_ZZToLLLL_madspin_2	1	99990
TTsemilep_madspin_1	172	8105465	TTsemilep_madspin_2	173	8156688
TZq2_W_trilep1	100	9999157	TZq2_W_trilep2	97	9672987
TZq2_s_trilep	94	9393276	TZq2_t5FS_trilep	97	9699081
WToLNu-0Jet_sm-no_masses	592	52785449	WToLNu-0Jet_sm-no_masses-run2	482	42972689
WToLNu-1Jet_sm-no_masses	586	32827404	WToLNu-2Jets_sm-no_masses	396	15769022
WToLNu-3Jets_sm-no_masses	488	12931463	WWToLLNuNu	194	11221071
WZToLLJJ	5	306339	WZToLLNu	120	7666801
WZToLNuNuNu	1	59147	WZToNuNuJJ	1	59420
ZToLL10-50-0Jet_sm-no_masses	1	97701	ZToLL10-50-1Jet_sm-no_masses	1	45361
ZToLL10-50-2Jets_sm-no_masses	1	38998	ZToLL10-50-3Jets_sm-no_masses	1	5690
ZToLL50-0Jet_sm-no_masses	9	784399	ZToLL50-1Jet_sm-no_masses	10	549567
ZToLL50-2Jets_sm-no_masses	9	350088	ZToLL50-3Jets_sm-no_masses_split	8	115396
ZToLL50-4Jets_sm-no_masses_split	1	2884	ZZTo4Nu	1	35808
ZZToLLLL	92	6222800	ZZToLLNuNu	1	64305

Monte Carlo errors below permil: neglected

# Cut-based analysis: optimisation

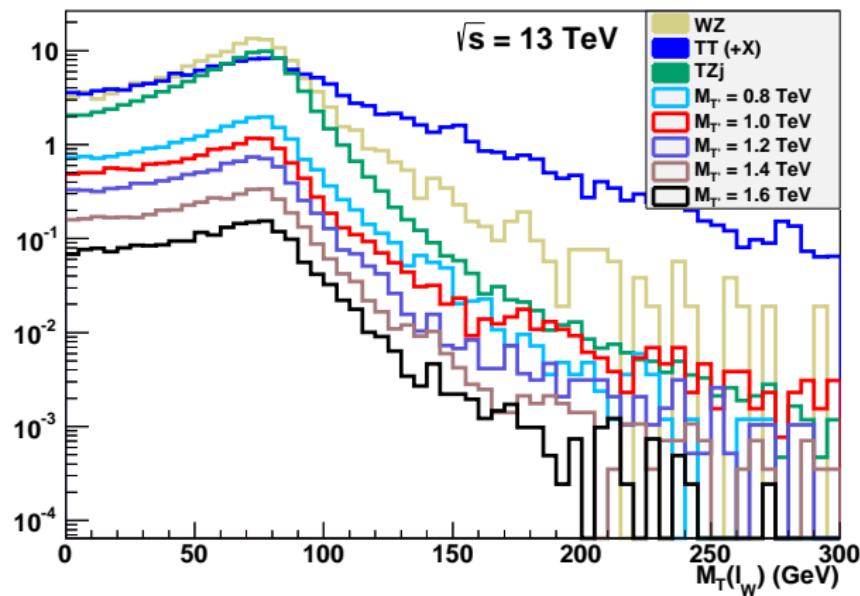
$Z$ -boson reco by minimising distance of OSSF leptons to  $M_Z$



$$|M(\ell^+ \ell^-)/\text{GeV} - M_Z| < 15$$

# Cut-based analysis: optimisation

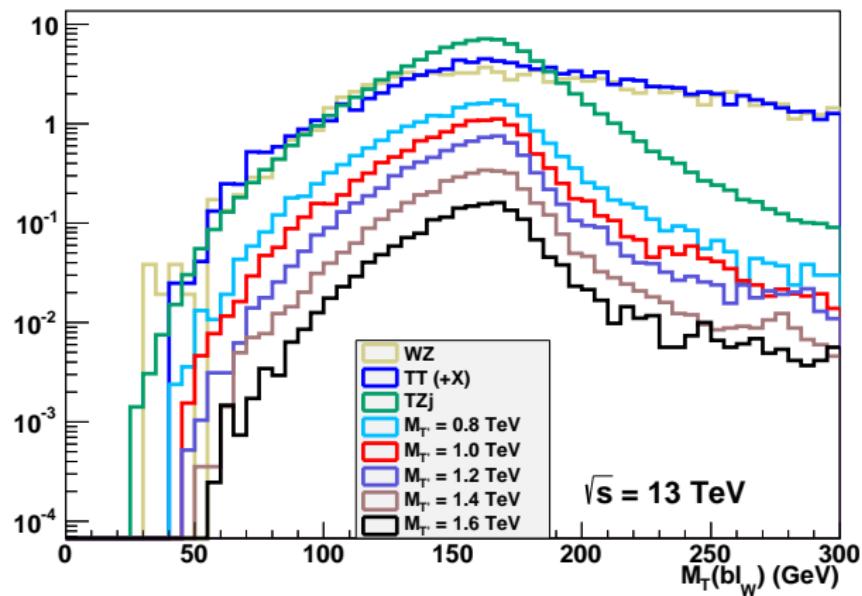
$W$  reco with remaining lepton



$$10 < M_T(\ell_W)/\text{GeV} < 150$$

# Cut-based analysis: optimisation

top reco with remaining lepton and  $b$ -tagged jet



$$0 < M_T(\ell_W b)/\text{GeV} < 220$$

# Objects selection

## Objects identification

$$p_T(\ell) > 20 \text{ GeV}, \quad |\eta(e/\mu)| < 2.5/2.4, \quad (1)$$

$$p_T(j) > 40 \text{ GeV}, \quad \Delta R(\ell, j) > 0.4, \quad (2)$$

$$|\eta(j)| < 5, \quad |\eta(b)| < 2.4, \quad (3)$$

Background	no cuts	$1 \leq n_j \leq 3$	$n_\ell \equiv 3$	$n_b \equiv 1$
$t\bar{t}(+X)$	$7.5 \cdot 10^6$ (100%)	$6.1 \cdot 10^6$ (81.2%)	<b>514.9 (0.09%)</b>	243.8 (47.3%)
$tZj$	3521 (100%)	2953 (83.9%)	290.6 (9.8%)	170.0 (58.5%)
$WZ$	$1.4 \cdot 10^5$ (100%)	$5.7 \cdot 10^4$ (41.9%)	3883 (6.9%)	164.3 (4.2%)
Total	$7.6 \cdot 10^6$ (100%)	$6.1 \cdot 10^6$ (80.5%)	4689 (0.08%)	578.0 (12.3%)
$M_{T'} \text{ (GeV)}$	no cuts	$1 \leq n_j \leq 3$	$n_\ell \equiv 3$	$n_b \equiv 1$
800	119.7 (100%)	105.0 (87.8%)	39.3 (37.4%)	25.5 (64.8%)
1000	77.1 (100%)	67.8 (87.9%)	26.0 (38.4%)	16.4 (63.2%)
1200	52.0 (100%)	45.3 (87.2%)	16.1 (35.6%)	10.1 (62.4%)
1400	35.3 (100%)	30.5 (86.6%)	8.0 (26.1%)	4.8 (60.1%)
1600	24.5 (100%)	21.1 (86.0%)	3.8 (18.0%)	2.2 (58.3%)

Signal generated without taus

# Cut-based analysis

## Selections

$$Z \rightarrow \ell^+ \ell^- \text{ reco} \quad |M(\ell^+ \ell^-)/\text{GeV} - M_Z| < 15, \quad (4)$$

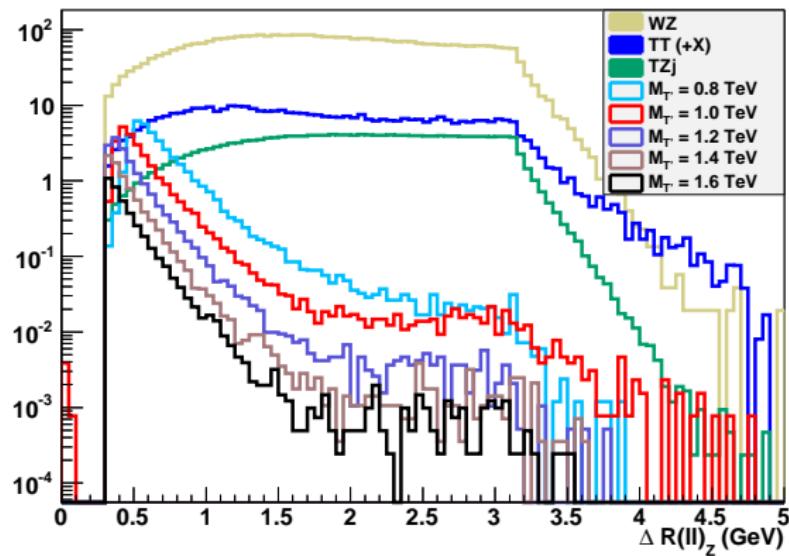
$$W \rightarrow \ell_W \nu \text{ reco} \quad 10 < M_T(\ell_W \nu)/\text{GeV} < 150, \quad (5)$$

$$t \rightarrow bW \text{ reco} \quad 0 < M_T(b \ell_W \nu)/\text{GeV} < 220. \quad (6)$$

Background	$n_b \equiv 1$	Z-reco	W-reco	t-reco
$t\bar{t}(+X)$	243.8 (47.3%)	154.8 (63.5%)	135.1 (87.3%)	83.0 (61.5%)
$tZj$	170.0 (58.5%)	155.6 (67.2%)	148.7 (95.6%)	139.8 (63.7%)
$WZ$	164.3 (4.2%)	146.9 (89.4%)	138.2 (94.1%)	71.5 ( <b>51.7%</b> )
Total	578.0 (12.3%)	457.2 (79.1%)	422.0 (92.3%)	294.3 (69.8%)
$M_{T'} \text{ (GeV)}$	$n_b \equiv 1$	Z-reco	W-reco	t-reco
800	25.5 (64.8%)	23.8 (93.6%)	22.2 (93.2%)	20.8 (93.6%)
1000	16.4 (63.2%)	15.4 (93.8%)	14.3 (92.4%)	13.4 (94.0%)
1200	10.1 (62.4%)	9.5 (94.2%)	8.7 (92.3%)	8.1 (92.3%)
1400	4.8 (60.1%)	4.5 (93.5%)	4.1 (92.1%)	3.8 (91.3%)
1600	2.2 (58.3%)	2.1 (93.3%)	1.9 (92.2%)	1.7 (90.0%)

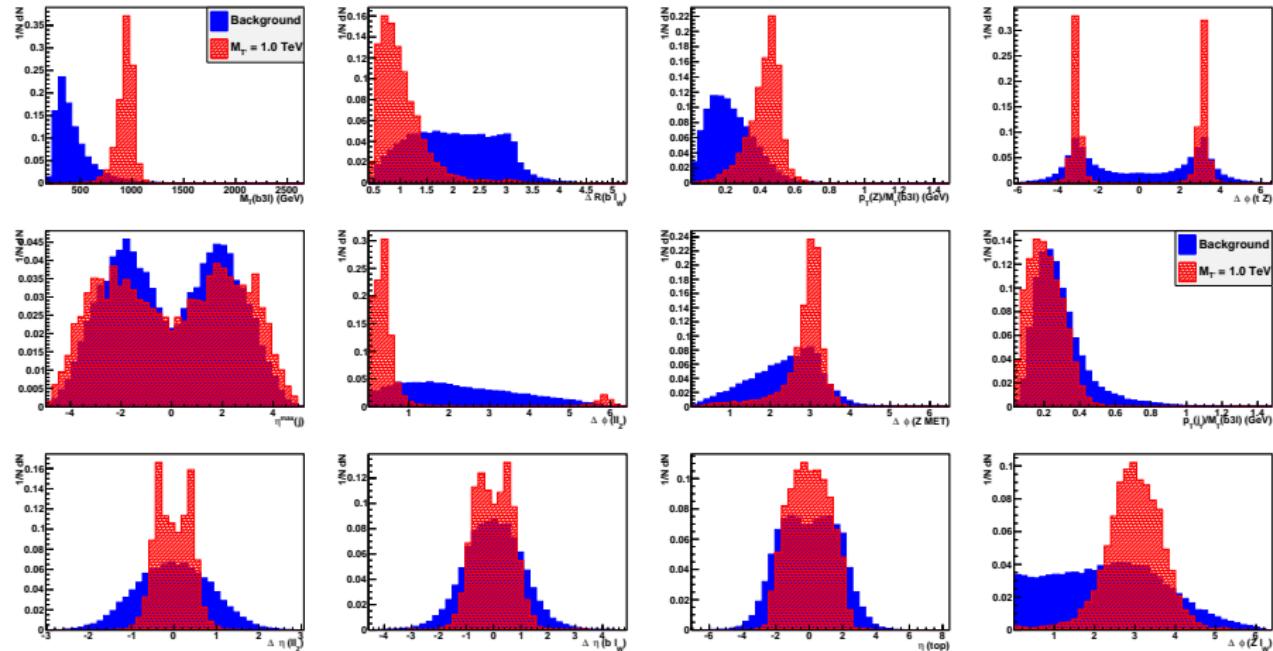
Cuts optimised to retain  $\geq 90\%$  of signal

# $\Delta R(\ell^+\ell^-)$ for $T'$ signals



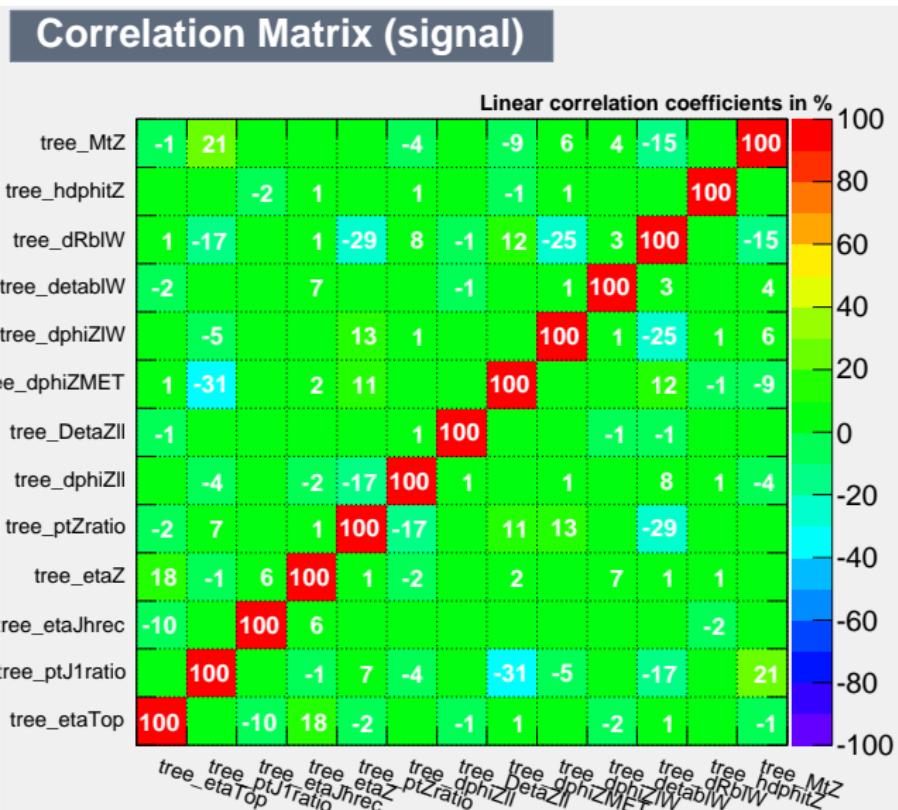
$T'$  is very massive, hence the decay products are boosted

# MVA variables



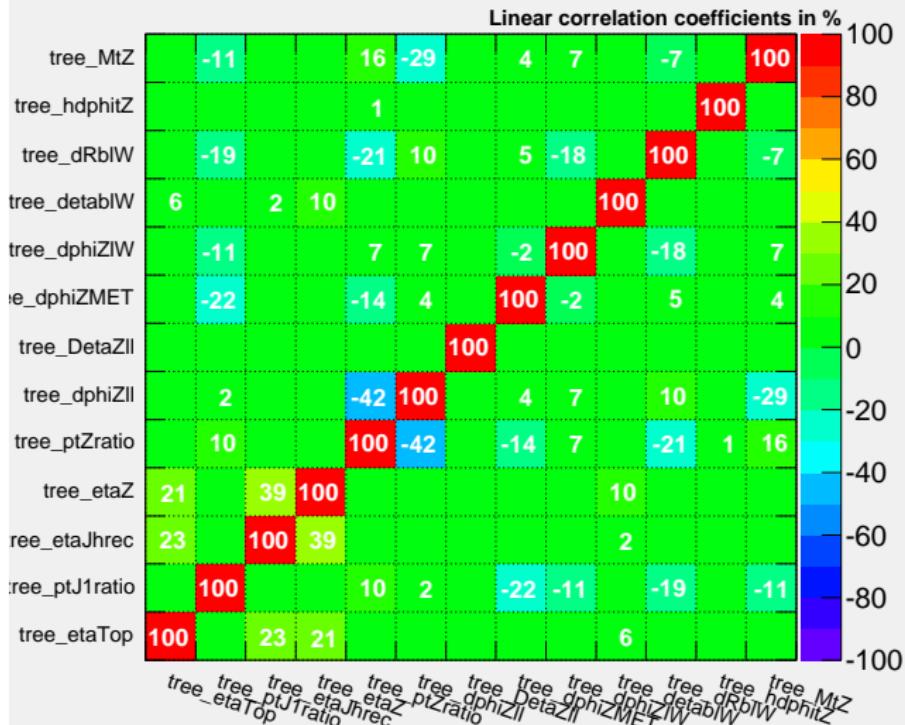
$p_T(Z)/M_T(b3\ell)$ ,  $p_T(j_1)/M_T(b3\ell)$ , and  $M_T(b3\ell)$  are decorrelated

# Correlations - $M_{T'} = 1 \text{ TeV}$



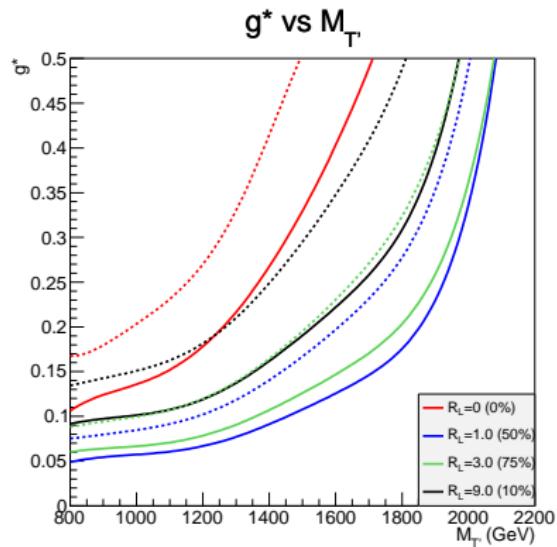
# Correlations - Background

Correlation Matrix (background)



# Comparison to dilepton channel

We set ourselves in similar conditions:  $\mathcal{L} = 300 \text{ fb}^{-1}$ ,  $\kappa_f = 1.14$ ,  $R_L = 0$

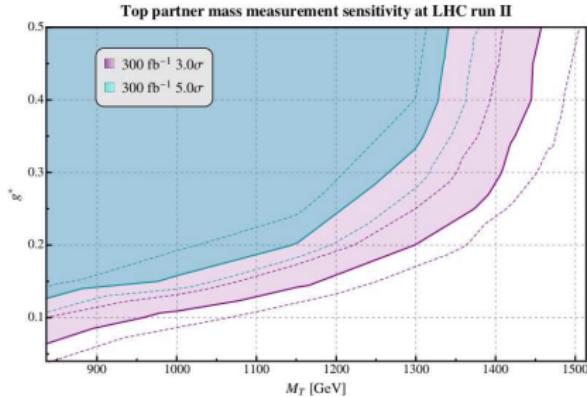


(dashed lines:  $5\sigma$ , solid lines:  $3\sigma$ )

Comparable reach at low  $T'$  masses (no pair-prod. here)

200  $\div$  300 GeV better sensitivity at high  $T'$  masses

J. Reuter and M. Tonini,  
JHEP 1501 (2015) 088 [arXiv:1409.6962]



# Reinterpretation: top anomalous couplings

The top-quark couplings can be parametrised in an effective field theory

The SM Lagrangian is extended by gauge-invariant (non-renormalisable) operators, obtained by integrating out heavy modes

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{C_i O_i}{\Lambda^2}$$

Here we consider only dimension 6 operators, the first non-vanishing terms in  $1/\Lambda$  expansion: total of 59 operators [W. Buchmuller, D. Wyler, Nucl.Phys. B268 \(1986\) 621](#)

Not all possible dim-6 operators that one can write are independent  
Redundant operators can be reduced by using equation of motions and other relations due to gauge invariance

[J. A. Aguilar-Saavedra, Nucl. Phys. B812, 181 \(2009\) \[0811.3842\]](#)

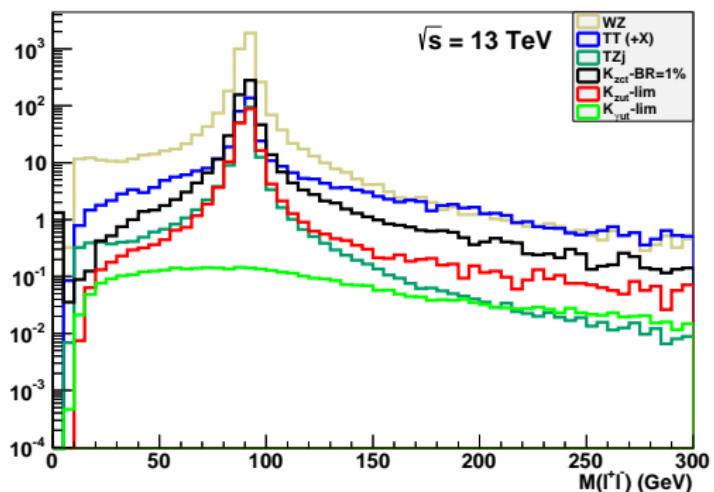
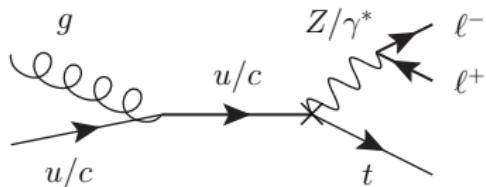
$$\mathcal{L} = \sum_{q=u,c} \frac{g}{\sqrt{2}c_W} \frac{\kappa_{tZq}}{\Lambda} \bar{t} \sigma^{\mu\nu} (f_{Zq}^L P_L + f_{Zq}^R P_R) q Z_{\mu\nu},$$

where  $\Lambda$  is the scale of new physics.

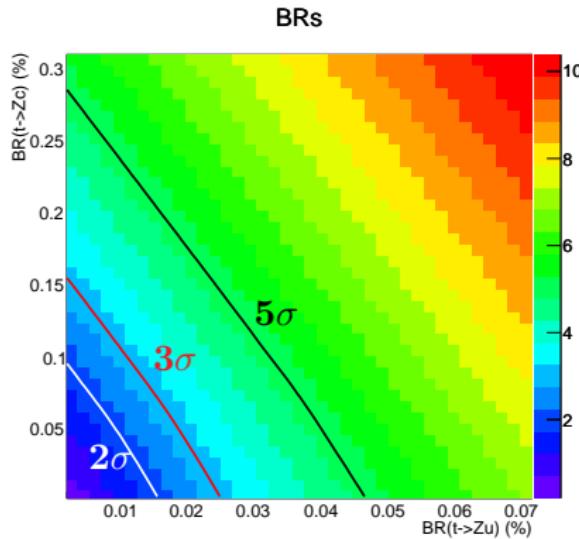
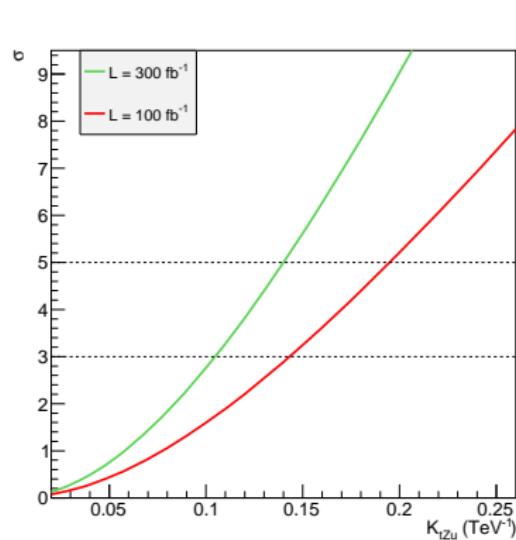
# Reinterpretation: optimisation

$tZq$  coupling gives similar final state as  $T' \rightarrow tZ \rightarrow t\ell^+\ell^-$

$t\gamma q$  coupling (with  $\gamma^*$ ) too. However, the cut around  $M_Z$  removes it



# Reinterpretation: parameter space



Actual limits:  $\text{BR}(t \rightarrow Zq) < 0.05\%$  (inclusive, from  $t\bar{t}$ )  $\Rightarrow \kappa_{tZu} < 0.2 \text{ TeV}^{-1}$

Otherwise from single top:  $\begin{cases} \text{BR}(t \rightarrow Zu) &< 0.51\% \\ \text{BR}(t \rightarrow Zc) &< 11.4\% \end{cases}$

See CMS-TOP-12-037 and CMS-TOP-12-021, respectively