

# $B \rightarrow K^{(*)} \nu \bar{\nu}$ decays in the Standard Model and beyond

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in collaboration with A. Buras, J. Girrbach-Noe and D. Straub

based on 1409.4557

EPS Vienna  
July 24, 2015

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→ No uncertainties from non-factorizable long-range photon exchange
- Recent theory progress allows to get much better SM predictions
- New measurements on  $B \rightarrow K^{(*)}\ell^+\ell^-$  allow to put strong constraints
- Hopefully, these decays will be accessible within the next few years

- 1 SM predictions
- 2 Model independent constraints
  - General remarks
  - Correlations with  $B \rightarrow K^{(*)} \ell^+ \ell^-$  decays
- 3 Conclusion



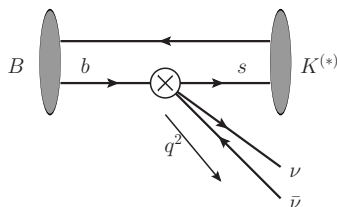
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## SM calculation

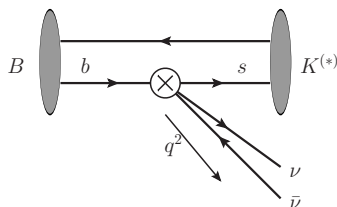


$$\mathcal{H}_{\text{eff}}^{\text{SM}} \propto C_L^{\text{SM}} \mathcal{O}_L + \text{h.c.}$$

$$\propto C_L^{\text{SM}} (\bar{s} \gamma_\mu P_L b) (\bar{\nu} \gamma^\mu P_L \nu) + \text{h.c.}$$

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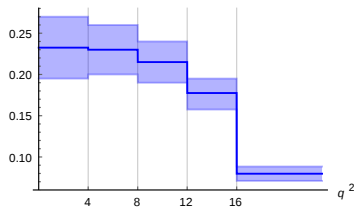
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We have to control:

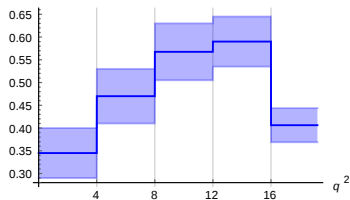
- Wilson coefficient  $C_L^{\text{SM}}$   
 $\rightarrow$  two-loop electroweak contributions (Brod,Gorbahn,Stamou 1009.0947)
- hadronic form factors  $\rho(q^2)$   
 $\rightarrow$  combined fit to LCSR and lattice results (Bharucha,Straub,Zwicky 1503.05534) ]

## updated SM predictions

$$10^6 \frac{d}{dq^2} \text{BR}(B^+ \rightarrow K^+ \nu \bar{\nu})$$



$$10^6 \frac{d}{dq^2} \text{BR}(B^0 \rightarrow K^{*0} \nu \bar{\nu})$$



$$\text{BR}(B^+ \rightarrow K^+ \nu \bar{\nu})_{\text{SM}} =$$

$$(3.98 \pm 0.43 \pm 0.19) \times 10^{-6}$$

$$< 1.7 \times 10^{-5} \text{ (BaBar)}$$

$$\text{BR}(B^0 \rightarrow K^{*0} \nu \bar{\nu})_{\text{SM}} =$$

$$(9.19 \pm 0.86 \pm 0.50) \times 10^{-6}$$

$$< 5.5 \times 10^{-5} \text{ (Belle)}$$

$$F_L^{\text{SM}} =$$

$$0.47 \pm 0.03$$

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## General remarks

Beyond the SM, a second eff. operator can contribute (right-handed currents!):

$$\mathcal{H}_{\text{eff}} \propto C_L \mathcal{O}_L + C_R \mathcal{O}_R + \text{h.c.},$$

$$\mathcal{O}_L \propto (\bar{s} \gamma_\mu P_L b) (\bar{\nu} \gamma^\mu P_L \nu)$$

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Reparametrize Wilson coefficients:

$$\epsilon = \frac{\sqrt{|C_L|^2 + |C_R|^2}}{|C_L^{\text{SM}}|}$$

$$\eta = \frac{-\text{Re}(C_L C_R^*)}{|C_L|^2 + |C_R|^2}$$

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Correlations in  $\mathcal{R}_K, \mathcal{R}_{K^*}, \mathcal{R}_{F_L}$

$\Rightarrow$  new invisible particles in final state?

# Correlations with $B \rightarrow K^{(*)} \ell^+ \ell^-$ decays

Idea: Use  $SU(2)_L$  symmetry to connect  $b \rightarrow s \nu \bar{\nu}$  decays to  $b \rightarrow s \ell^+ \ell^-$  decays, on which a lot of exp. data exists.

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Use most general  $\mathcal{G}_{SM}$ -invariant basis of dim6-operators. (Grzadkowski et al. 1008.4884)

$$Q_{Hq}^{(1)} = i(\bar{q}_L \gamma_\mu q_L) H^\dagger D^\mu H,$$

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$$\begin{aligned}
 O_L &\propto (\bar{s} \gamma_\mu P_L b) (\bar{\nu} \gamma^\mu P_L \nu) & O_9 &\propto (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \ell) \\
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So, one finds a dictionary:

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 C_L &= C_L^{\text{SM}} + \tilde{c}_{ql}^{(1)} - \tilde{c}_{ql}^{(3)} + \tilde{c}_Z, & C_R &= \tilde{c}_{dl} + \tilde{c}'_Z, \\
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with  $\tilde{c}_Z = \frac{1}{2}(\tilde{c}_{Hq}^{(1)} + \tilde{c}_{Hq}^{(3)})$ ,  $\tilde{c}'_Z = \frac{1}{2}\tilde{c}_{Hd}$ .

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- Now, use  $b \rightarrow s \ell^+ \ell^-$  data to constraint the Wilson coefficients and see how large effects in  $b \rightarrow s \nu \bar{\nu}$  can still get. (Altmannshofer, Straub 1411.3161)



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- Now, use  $b \rightarrow s \ell^+ \ell^-$  data to constraint the Wilson coefficients and see how large effects in  $b \rightarrow s \nu \bar{\nu}$  can still get. (Altmannshofer, Straub 1411.3161)
- Consider only certain scenarios of NP where only a subset of operators is active.

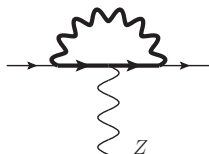
# Two different scenarios

## Two different scenarios

NP dominated by:

- Modified (flavour changing)  $Z$ -couplings [e.g. MSSM, partial compositeness]

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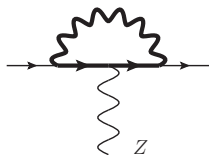


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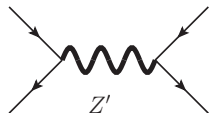
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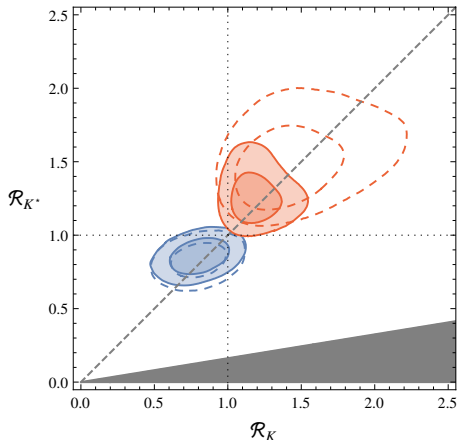
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- 4-Fermion-Operators [e.g. exchange of heavy  $Z'$  boson]

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blue: modified Z couplings

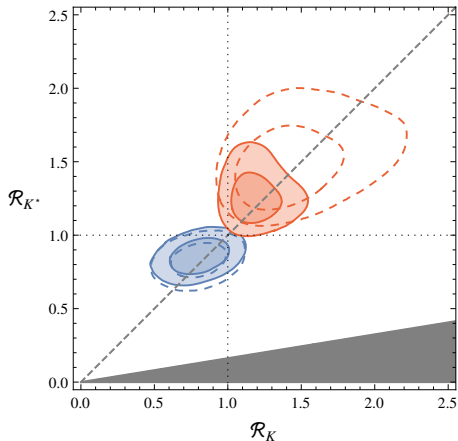
red: 4-Fermion operators

solid: real

dashed: complex

Current  $b \rightarrow s \ell^+ \ell^-$  data  
favour:

- Suppression of  $\mathcal{R}_{K^{(*)}}$  if NP mainly in modified Z couplings
- Enhancement of  $\mathcal{R}_{K^{(*)}}$  if NP mainly in 4-Fermion operators



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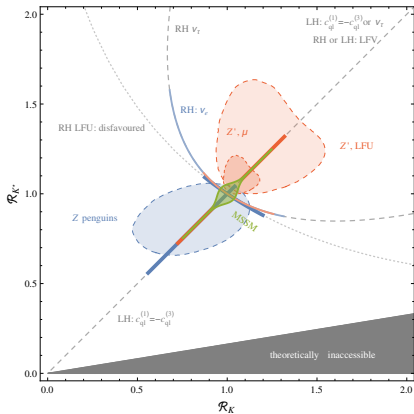
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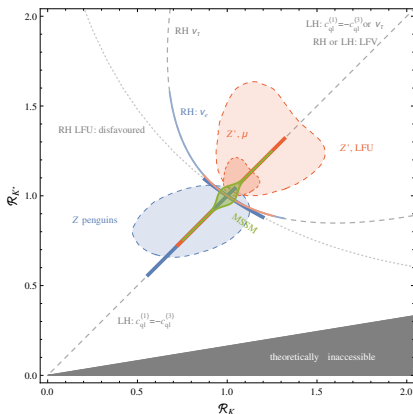
Correlations between  $\mathcal{R}_K$  and  $\mathcal{R}_{K^*}$  allow to disentangle both scenarios.

Due to tensions in  $b \rightarrow s \ell^+ \ell^-$ .

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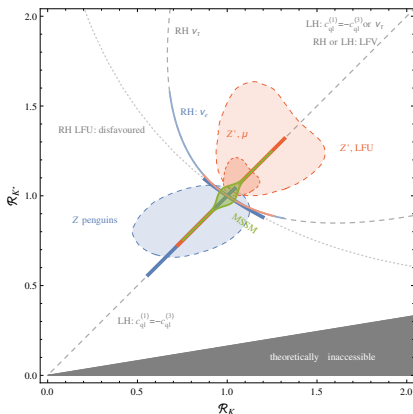




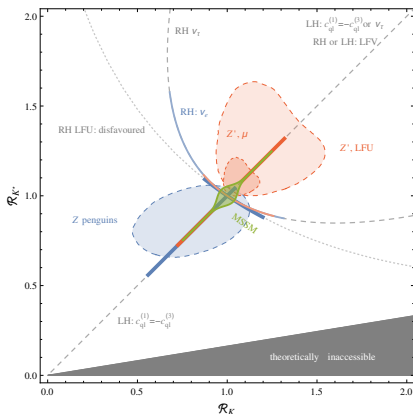


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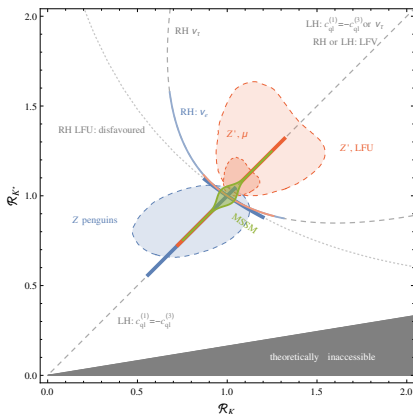
- MFV
- LH  $Z'$  couplings
- certain PC scenarios
- MSSM
- $SU(2)_L$  singlet or triplet LQ



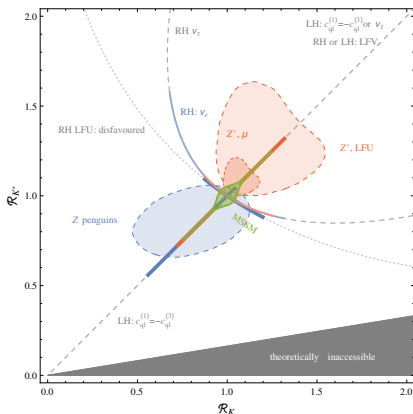
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- otherwise effects max.  $\begin{cases} \pm 60\% \text{ LFU} \\ \pm 20\% \text{ only } \mu \end{cases}$

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- Correlations between  $B \rightarrow K\nu\bar{\nu}$  and  $B \rightarrow K^*\nu\bar{\nu}$  (and also  $B \rightarrow K^{(*)}\ell\ell$  and  $B_s \rightarrow \ell\ell$ ) can help to identify possible NP scenarios