

Higgs and BSM physics at CLIC

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on behalf of the CLICdp collaboration

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The Compact Linear Collider



 Compact Linear Collider

Future linear electron-positron collider providing high luminosity: $\text{few} \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$

CLIC would be implemented in energy stages, defined by physics and technical considerations

Staged construction gives access to three centre of mass energies:

Stage 1: 500 fb^{-1} at $\sqrt{s} = 350 \text{ GeV}$

▶ SM Higgs and top physics, top threshold scan

see “Top quark physics at a future linear collider” R. Poeschl

Stage 2: 1.5 ab^{-1} at $\sqrt{s} = 1.4 \text{ TeV}$

▶ Improved Higgs precision, top Yukawa coupling, first BSM searches

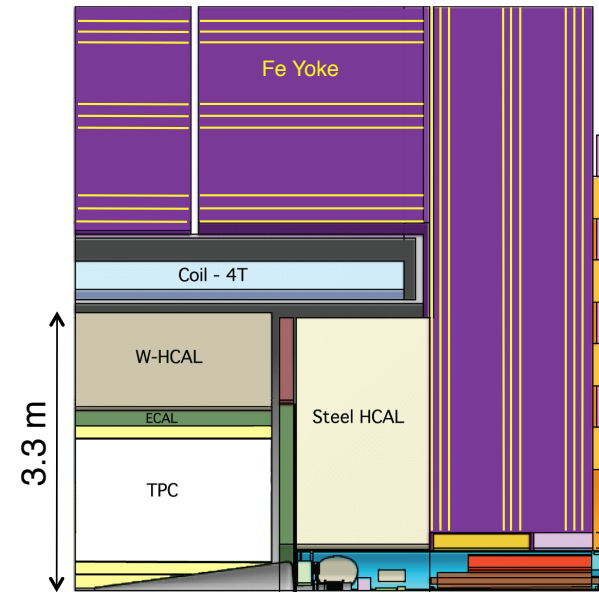
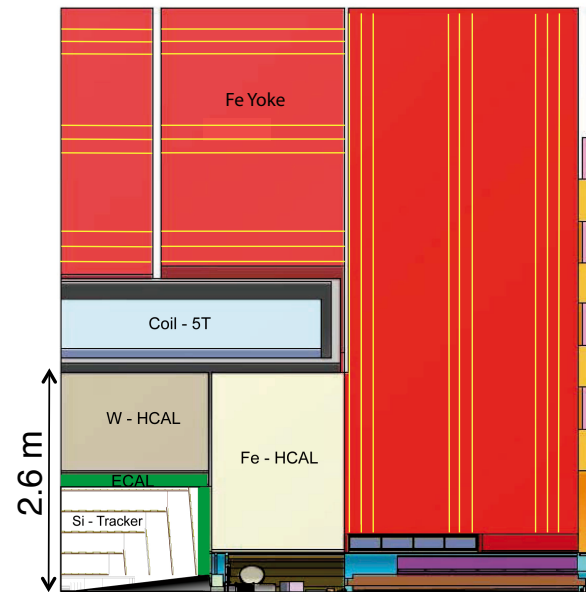
Stage 3: 2 ab^{-1} at $\sqrt{s} = 3 \text{ TeV}$

▶ Rarest Higgs processes, double Higgs production, best BSM sensitivity

Each stage corresponds to 4-5 years data taking

Detectors, beam conditions and backgrounds

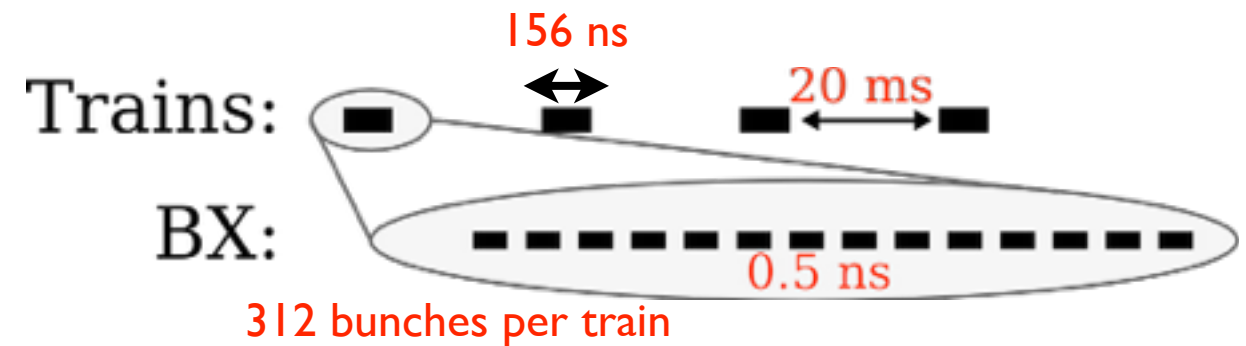
Physics studies in this talk use the **CLIC_SiD** and **CLIC_ILD** detector concepts



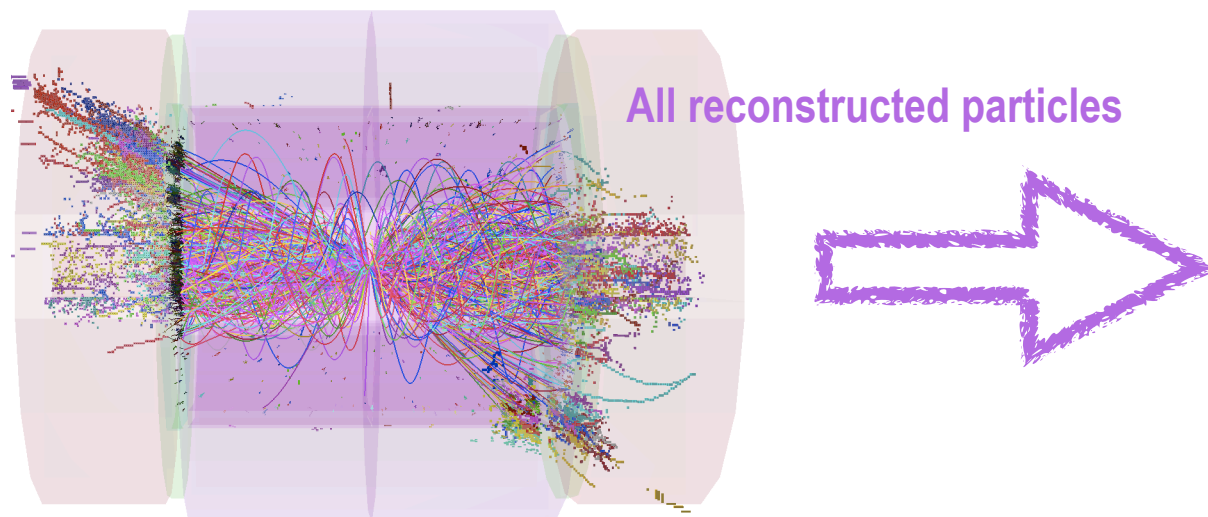
All studies use full physics and detector simulation

CLIC beams consist of short dense bunch trains

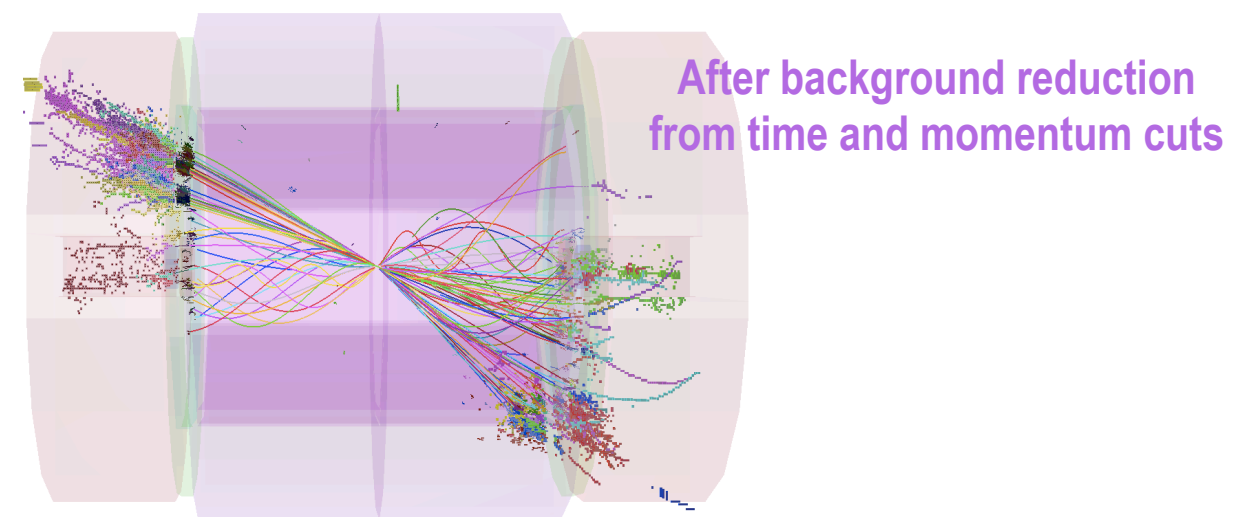
- ▶ enables power-pulsing, triggerless readout
- ▶ results in pile-up of beam-induced backgrounds



tt event plus 60 bunch crossings of beam-induced background



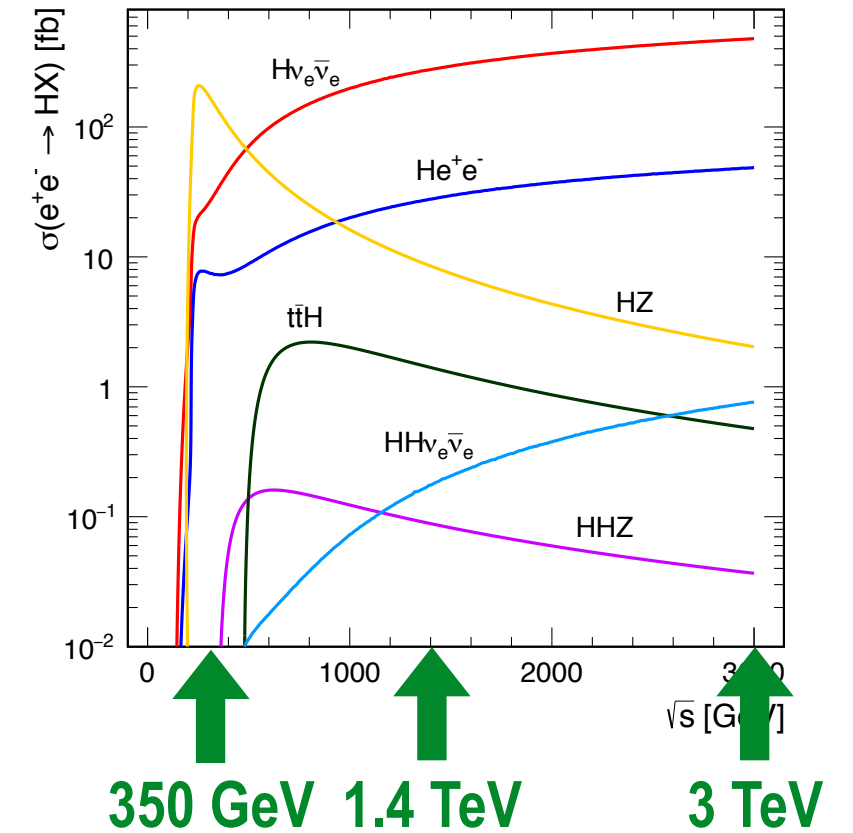
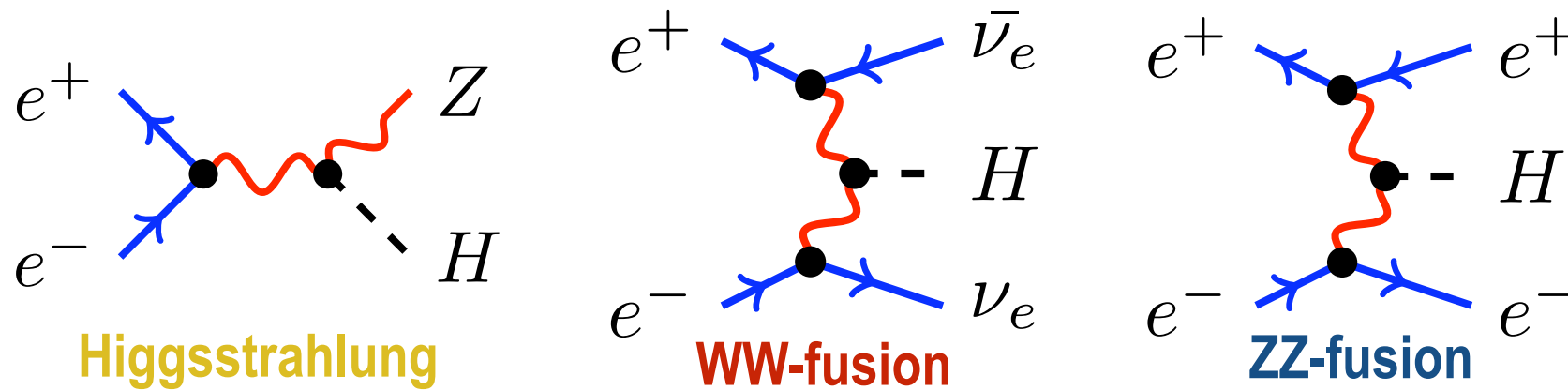
All reconstructed particles



After background reduction from time and momentum cuts

Single Higgs production at CLIC

Large samples of Higgs bosons will be produced at CLIC



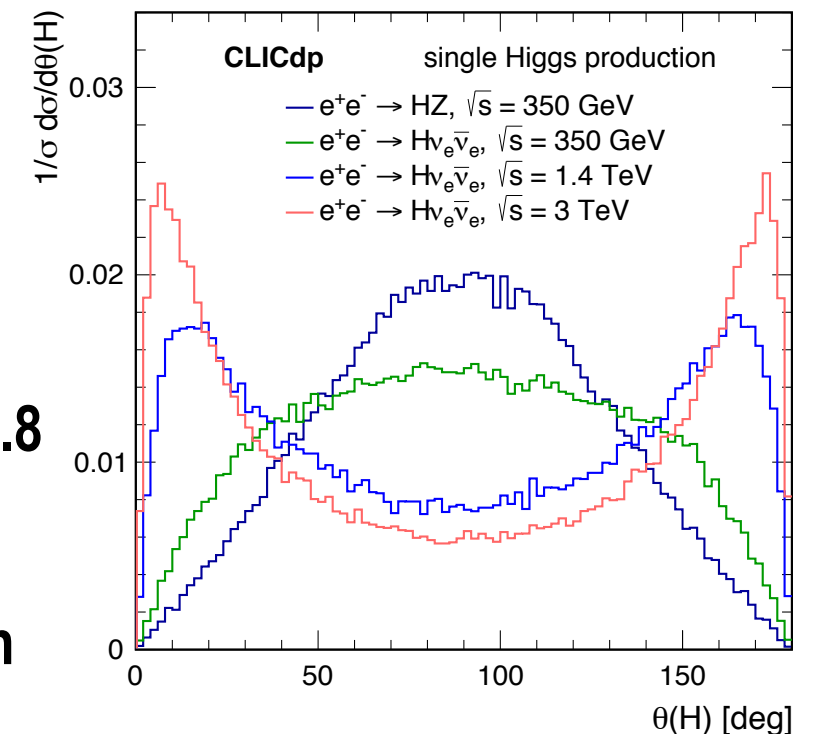
	350 GeV 500 fb ⁻¹	1.4 TeV 1.5 ab ⁻¹	3 TeV 2 ab ⁻¹
ZH	68,000	20,000	11,000
Hvν	17,000	370,000	830,000
Hee	3,700	37,000	84,000

unpolarised beams

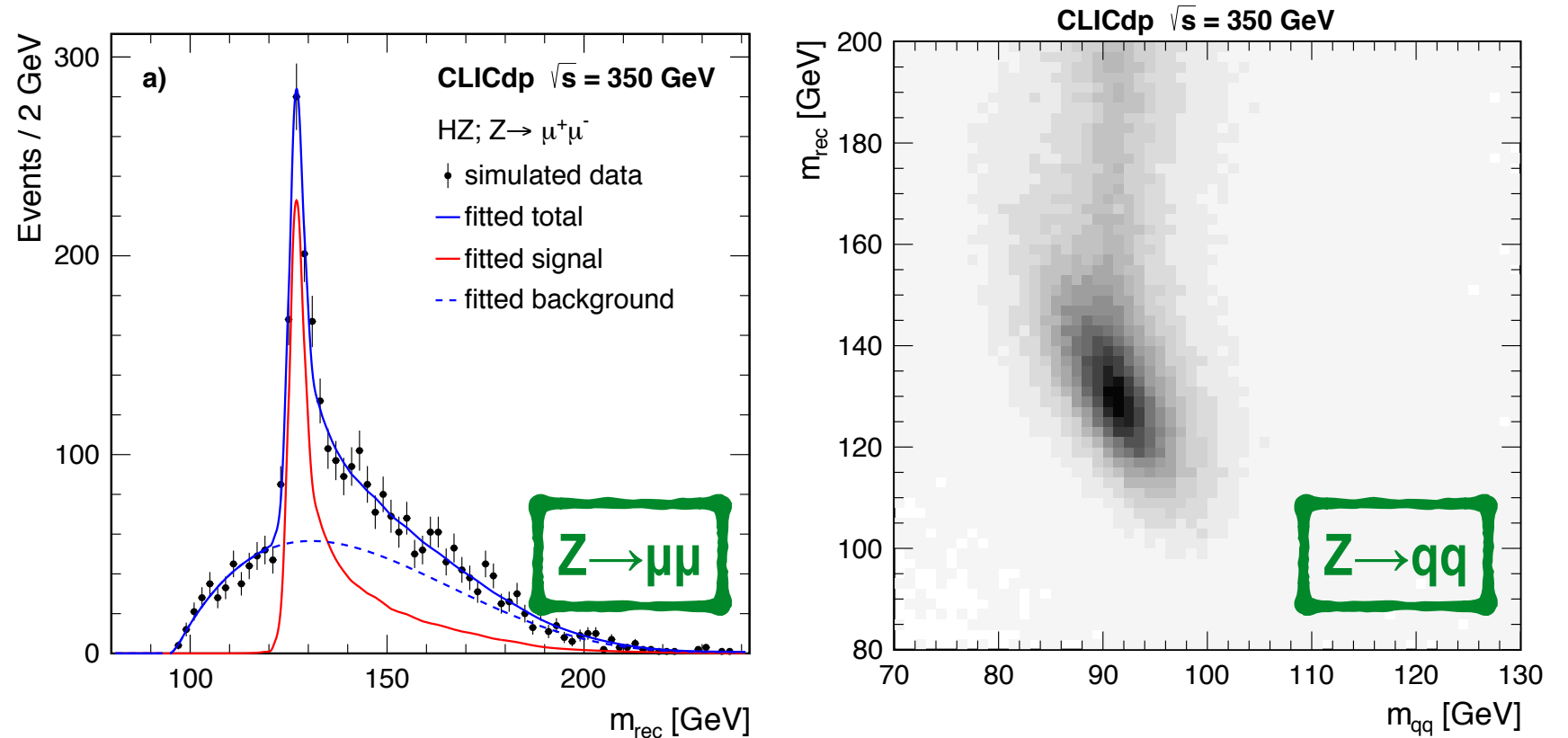
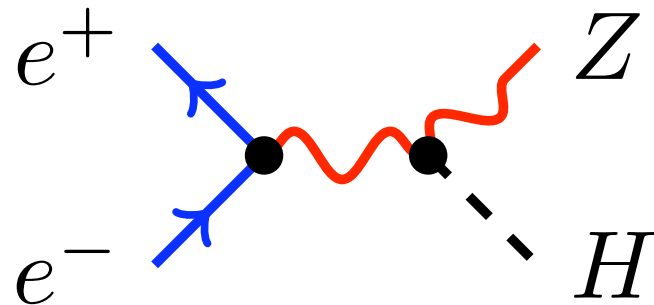
Polarising the electron beam (80%) enhances WW-fusion by a factor 1.8

► but benchmark studies assume unpolarised beams

Measurements at high \sqrt{s} require good detectors in the forward region



Higgsstrahlung at $\sqrt{s} = 350$ GeV



Identify HZ events from only the Z recoil mass: **model independent** measurement of g_{HZZ}

► very clean for $Z \rightarrow \mu\mu$, ee decays: 2% uncertainty on g_{HZZ}

► also possible for hadronic Z decays with minimal bias: 0.9% uncertainty on g_{HZZ}

Combined uncertainty 0.8% on model independent measurement of g_{HZZ}

Cross section x branching ratio measurements:

► precision at few % level

Constrain the invisible Higgs decay $< 1\%$ at 90% CL

	Stat precision
$\sigma(\text{ZH}) \times \text{BR}(\text{H} \rightarrow \tau\tau)$	6.2%
$\sigma(\text{ZH}) \times \text{BR}(\text{H} \rightarrow b\bar{b})$	1%
$\sigma(\text{ZH}) \times \text{BR}(\text{H} \rightarrow c\bar{c})$	5%
$\sigma(\text{ZH}) \times \text{BR}(\text{H} \rightarrow g\bar{g})$	6%
$\sigma(\text{ZH}) \times \text{BR}(\text{H} \rightarrow W\bar{W}^*)$	2%
$\sigma(\text{H}\nu\nu) \times \text{BR}(\text{H} \rightarrow b\bar{b})$	3%

unpolarised beams

Precision Higgs measurements at $\sqrt{s} > 1$ TeV

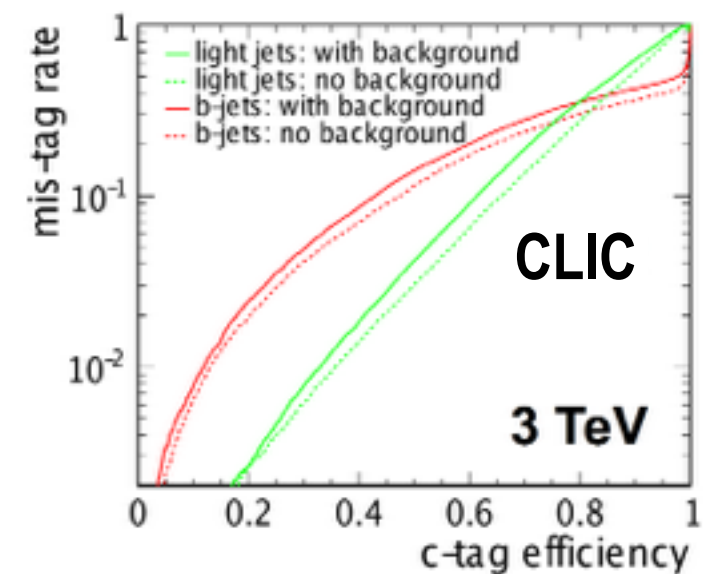
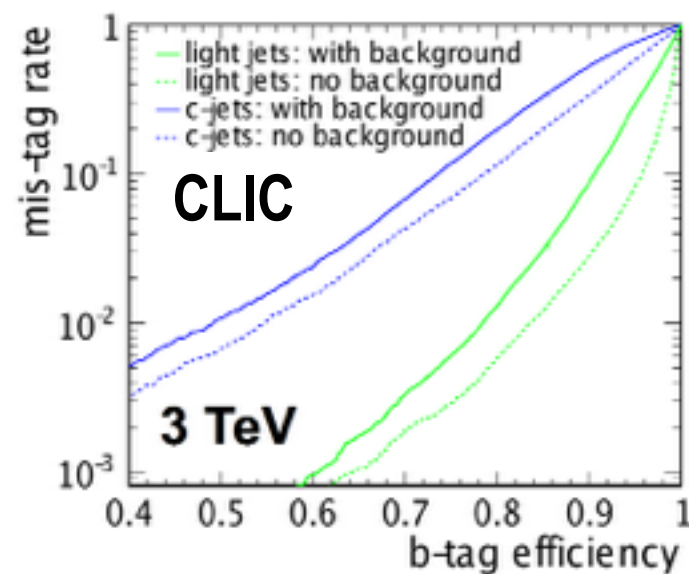
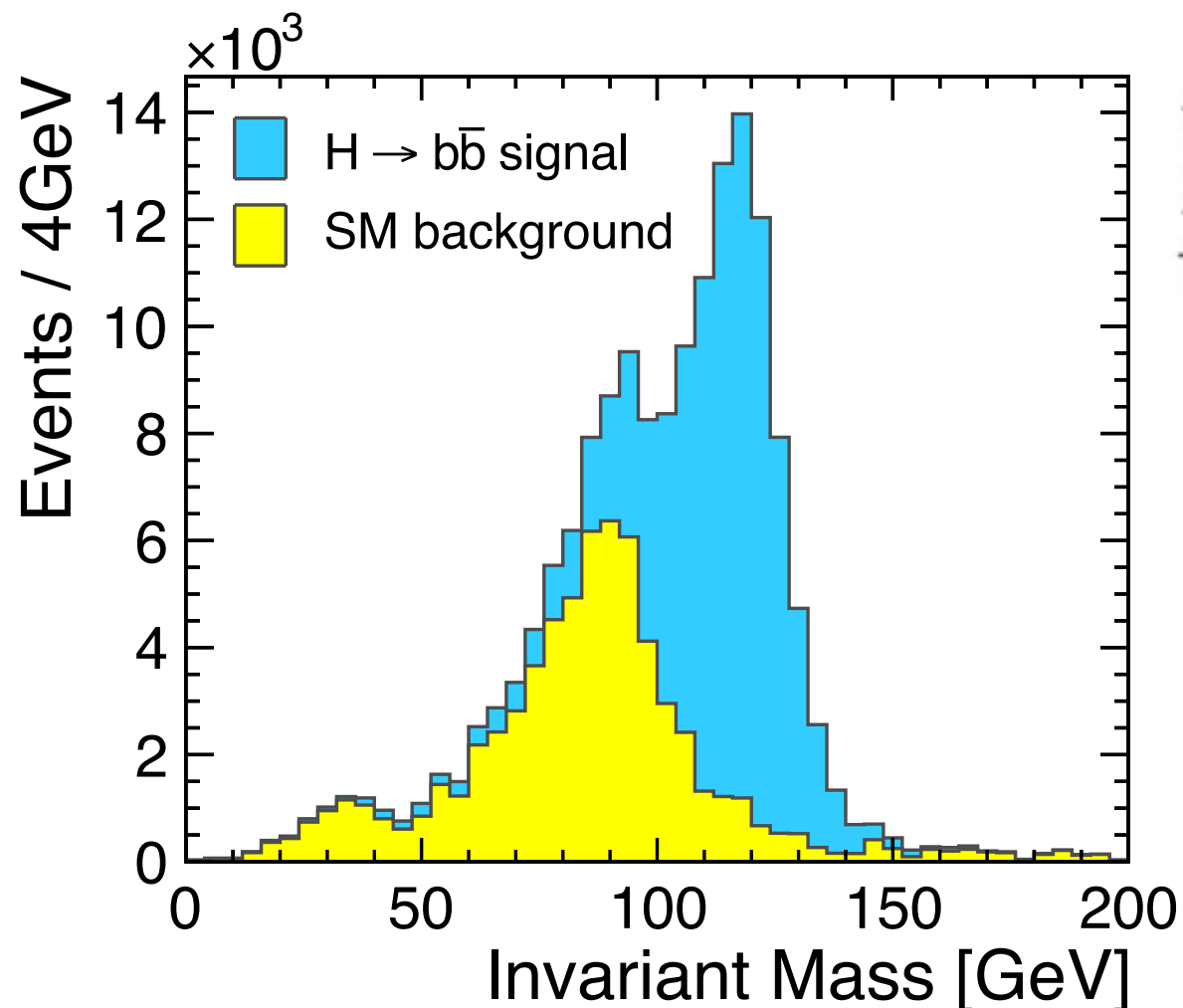
$H \rightarrow bb, cc, gg$

Separation of the different hadronic final states uses precise flavour tagging

► very challenging at the LHC

Higgs mass can be extracted from the $H \rightarrow bb$ invariant mass distribution: ± 33 MeV at 3 TeV

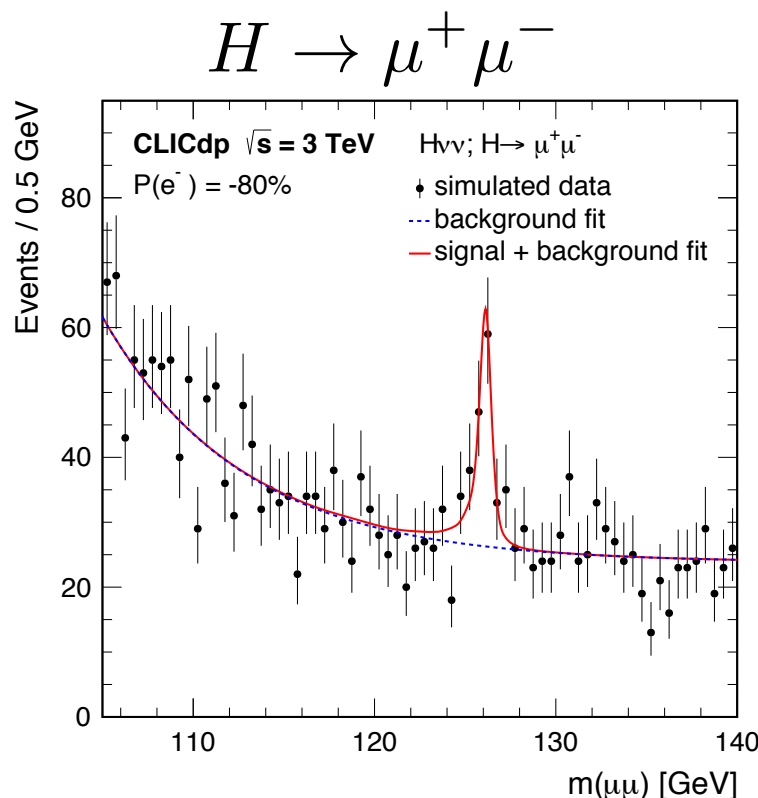
	3 TeV	unpolarised beams
$\sigma(H\nu\nu) \times BR(H \rightarrow bb)$	0.2%	
$\sigma(H\nu\nu) \times BR(H \rightarrow cc)$	2.7%	
$\sigma(H\nu\nu) \times BR(H \rightarrow gg)$	1.8%	



Rare Higgs decays at $\sqrt{s} > 1$ TeV

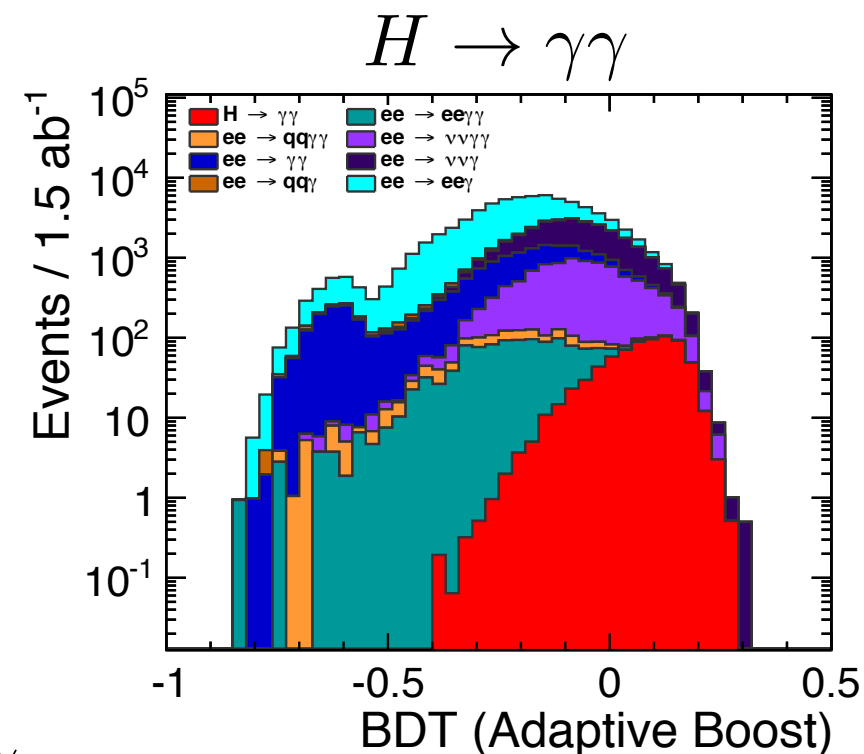
$\sigma(H\nu\nu) \times \text{BR}(H \rightarrow \mu\mu)$

- ▶ BR ~ 0.022%
- ▶ requires precise tracking
- ▶ 38% at 1.4 TeV, 16% at 3 TeV



$\sigma(H\nu\nu) \times \text{BR}(H \rightarrow \gamma\gamma)$

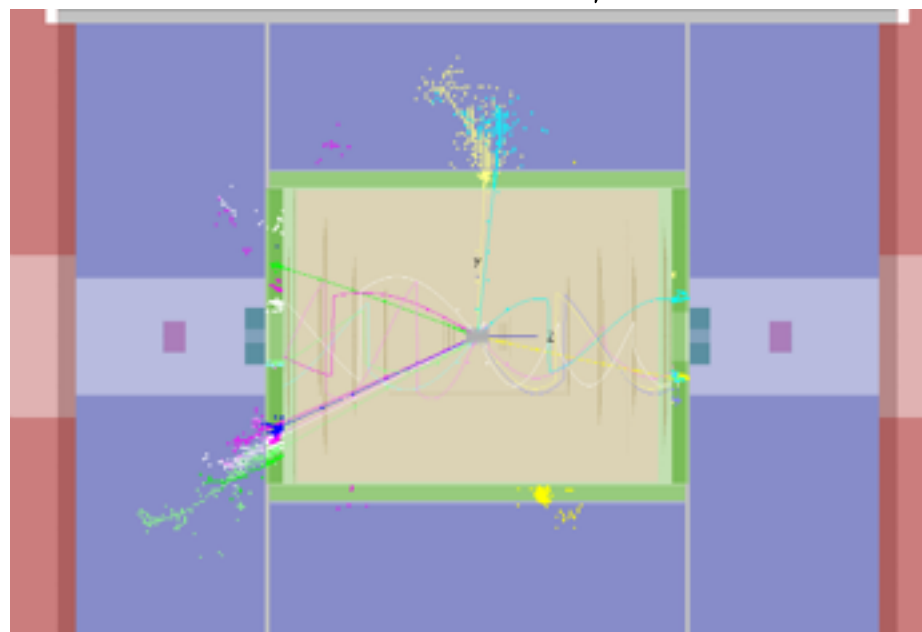
- ▶ BR ~ 0.23%
- ▶ reconstructed $\gamma\gamma$ mass resolution 3.3 GeV
- ▶ 15% at 1.4 TeV



$H \rightarrow Z\gamma$

$\sigma(H\nu\nu) \times \text{BR}(H \rightarrow Z\gamma)$

- ▶ BR ~ 0.16%
- ▶ hadronic Z decays used
- ▶ 42% at 1.4 TeV

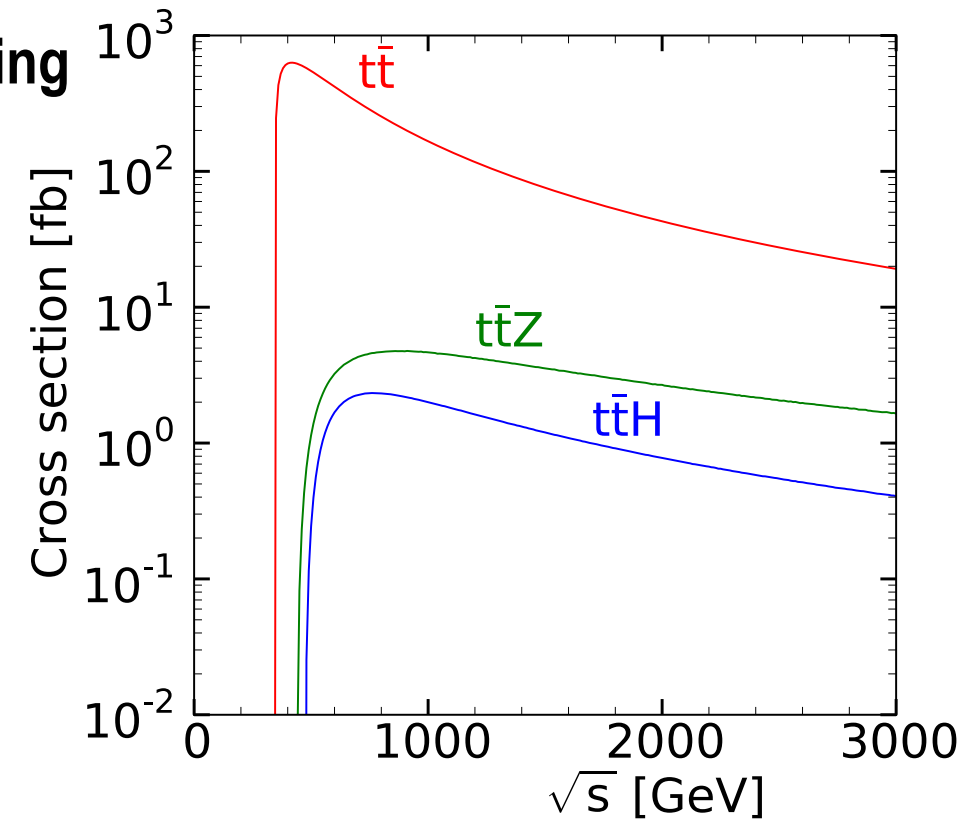
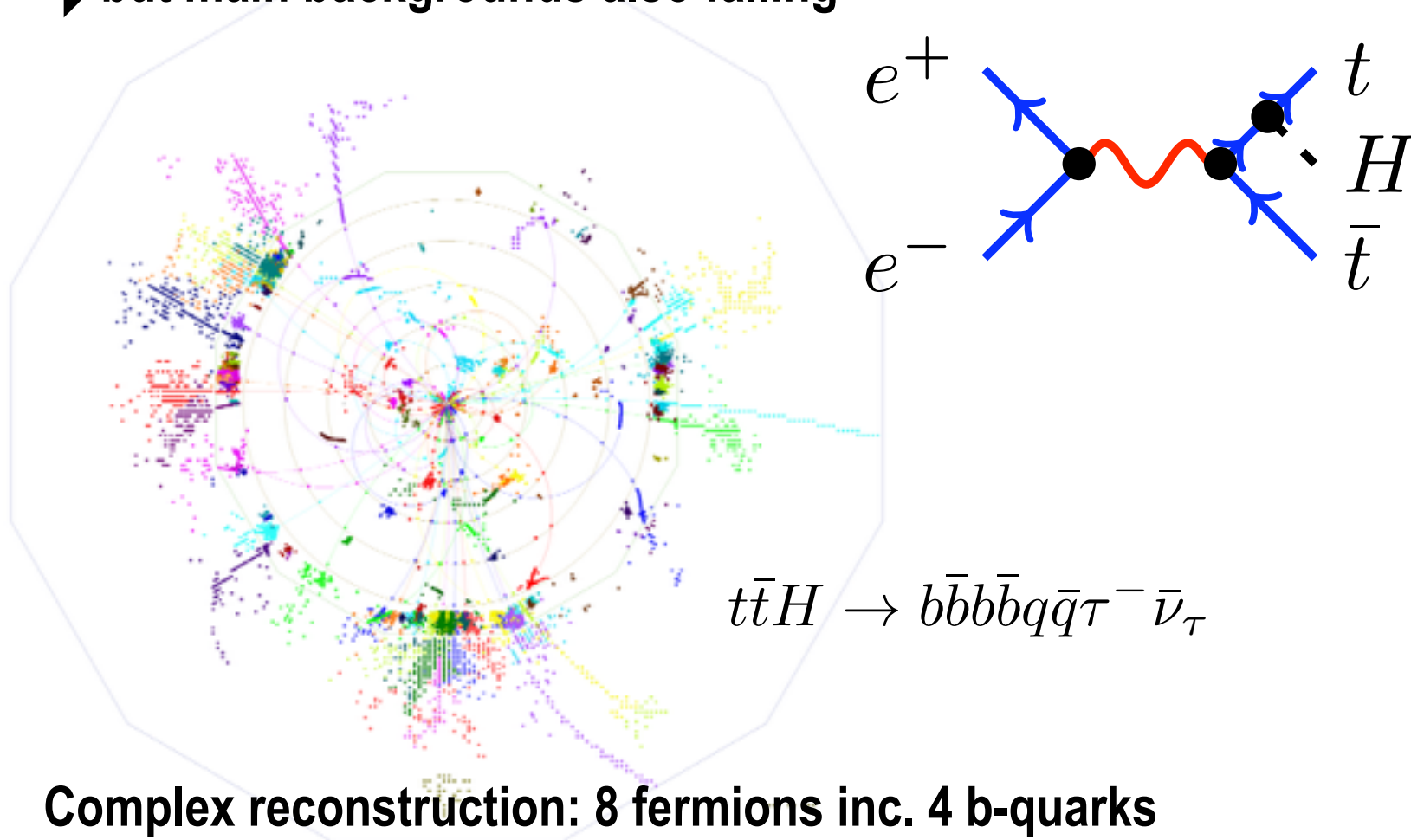


Top Yukawa coupling at $\sqrt{s} = 1.4$ TeV

The $t\bar{t}H$ cross section is directly sensitive to the top Yukawa coupling

$t\bar{t}H$ production peaks at ~ 800 GeV, measured at 1.4 TeV

► but main backgrounds also falling



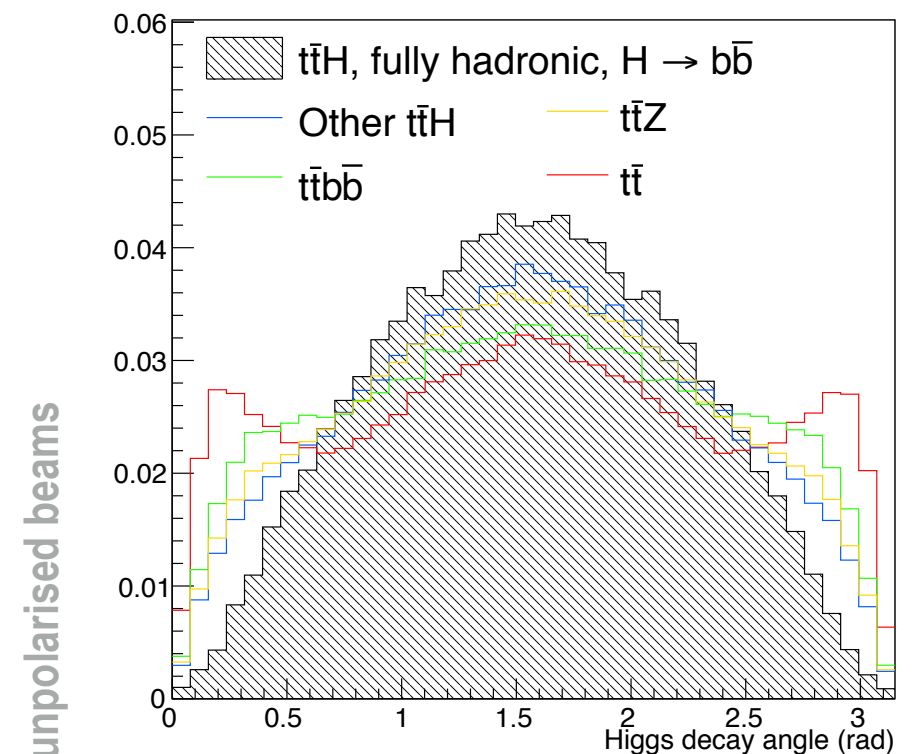
Complex reconstruction: 8 fermions inc. 4 b-quarks

► 6 jet + lepton + missing energy final state (semi-leptonic)

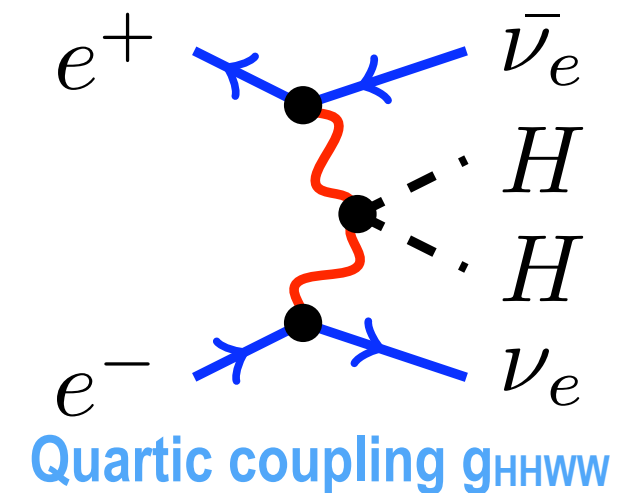
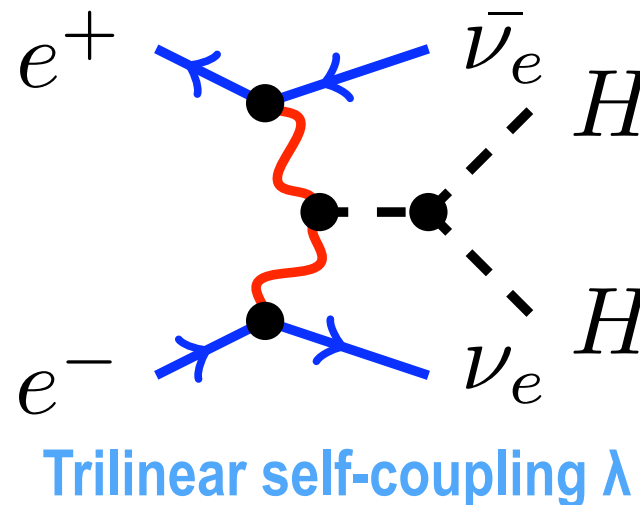
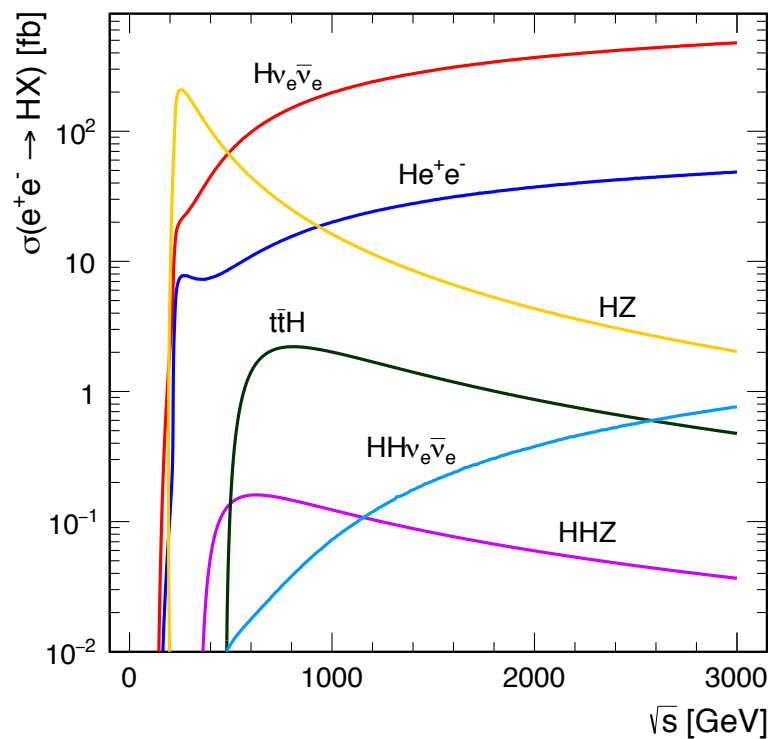
► 8 jet final state (fully hadronic)

Combined uncertainty on the top Yukawa coupling 4.5%

Comparable to other measurements at 1 TeV



Double Higgs production at $\sqrt{s} > 1$ TeV



The HH $\nu\nu$ cross section is sensitive to the Higgs self-coupling and the quartic coupling

Only 225 (1200) HH $\nu\nu$ events at $\sqrt{s} = 1.4$ (3) TeV

► high luminosity and high energy crucial

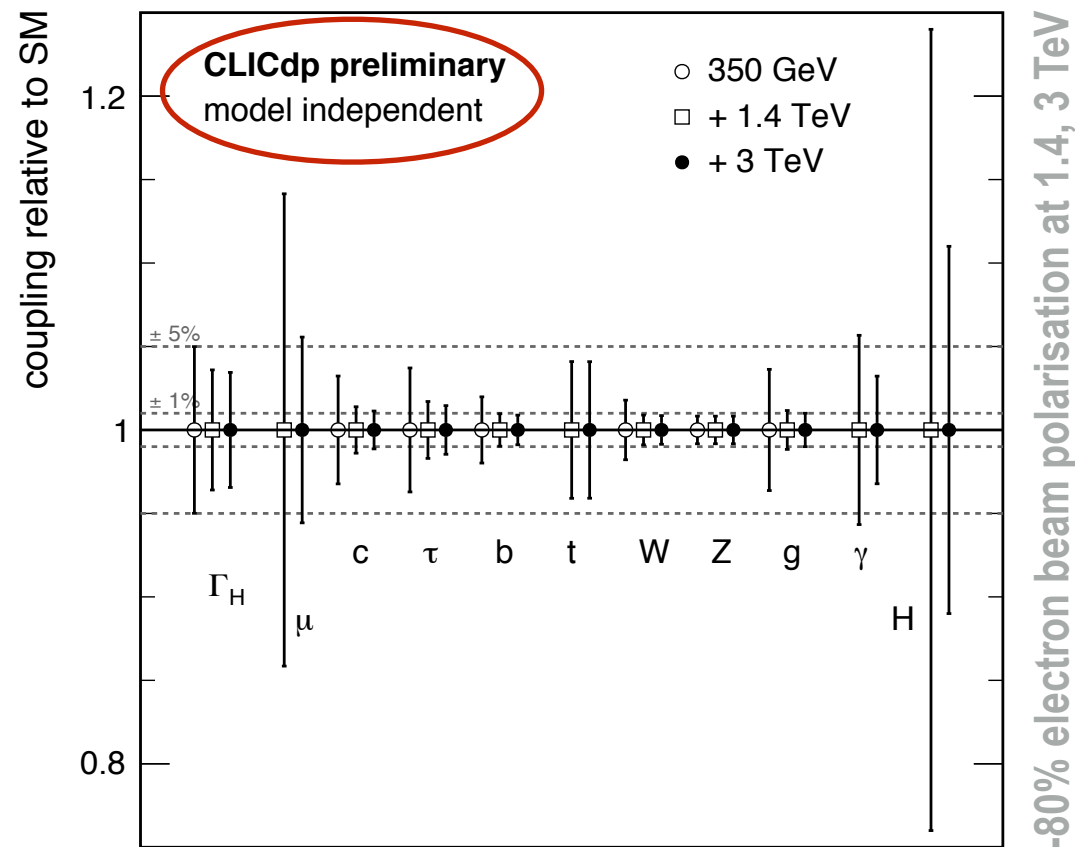
Events forced to 4-jet topology

Jets paired by hemisphere/to minimise mass χ^2

Template fit to extract sensitivity

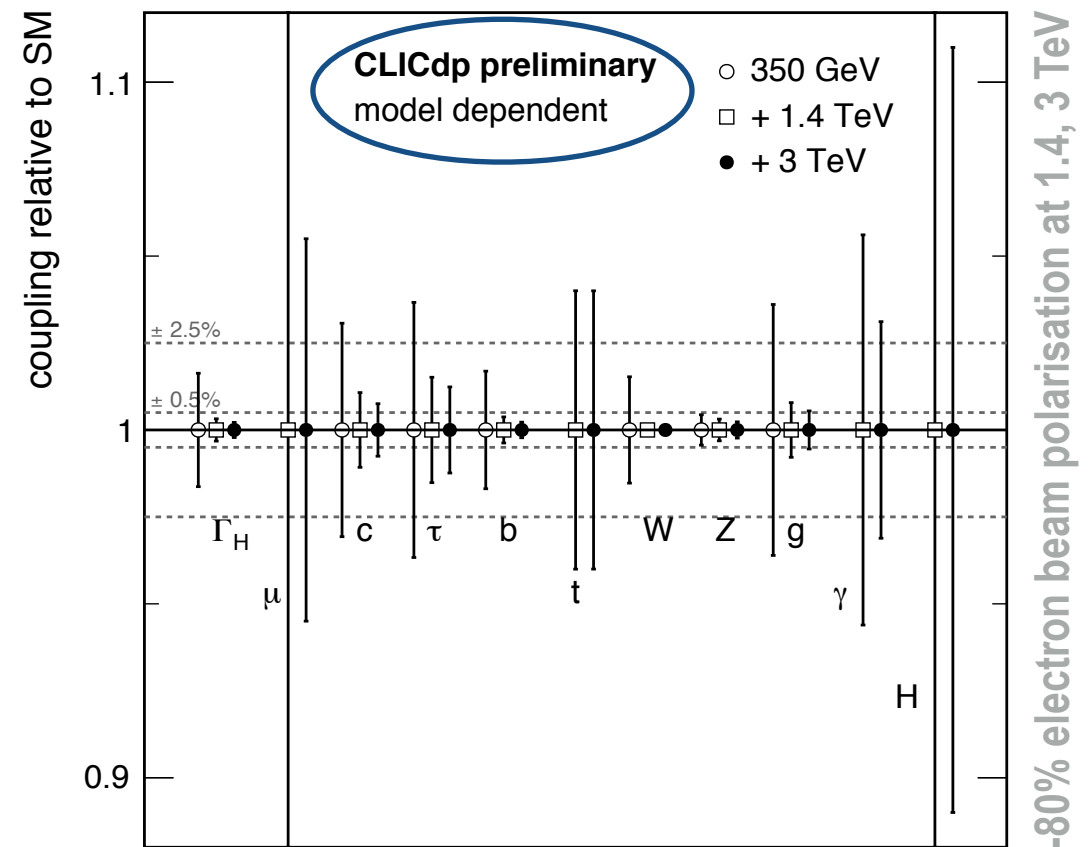
	1.4 TeV	3 TeV
$\Delta(g_{HHWW})$	7%	3%
$\Delta(\lambda)$	32%	16%
$\Delta(\lambda) P(e) = -80\%$	24%	12%

Model-independent fit to Higgs properties



Fully model independent fit

- ▶ only possible at lepton colliders
- ▶ stems from model independent measurement of g_{HZZ}
- ▶ Higgs width extracted with 5-3.5% precision



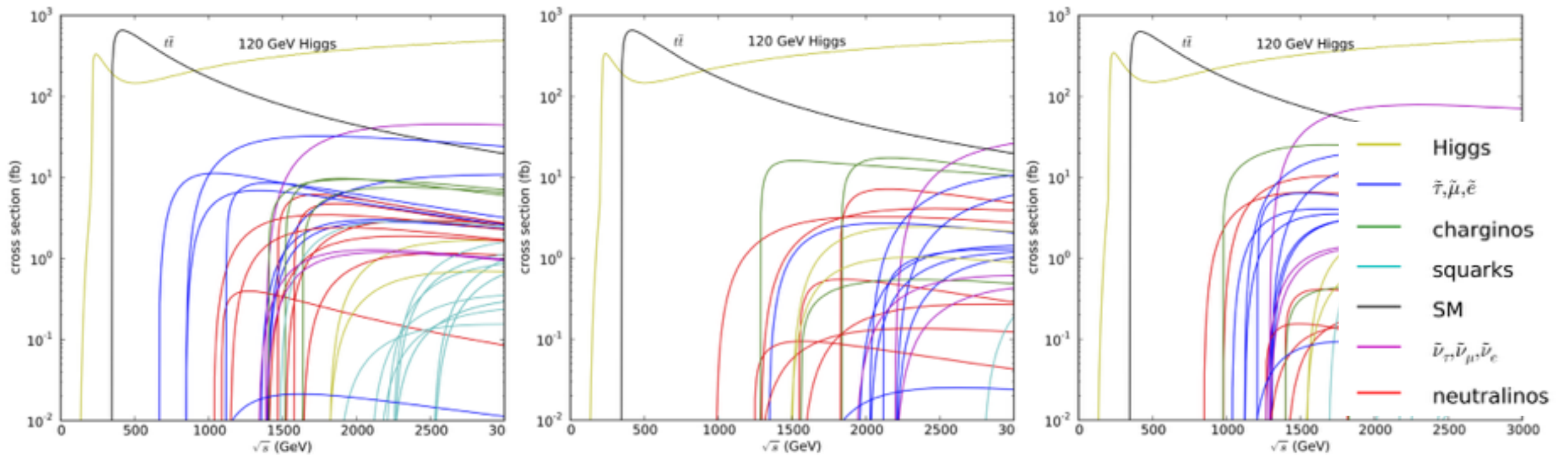
Model dependent fit

- ▶ as at LHC (no invisible H decays)
- ▶ sub-% precision at high energy
- ▶ Higgs width extracted with <1% precision
- ▶ but results depend strongly on fit assumptions

Prospects for BSM physics at CLIC

Direct searches via pair-production of new particles up to the kinematic limit $M < \sqrt{s} / 2$

► three SUSY models used to benchmark performance:



CDR model 1, $\sqrt{s} = 3 \text{ TeV}$

- squarks
- heavy Higgs

CDR model 2, $\sqrt{s} = 3 \text{ TeV}$

- smuons, selectrons
- gauginos

CDR model 3, $\sqrt{s} = 1.4 \text{ TeV}$

- smuons, selectrons, staus
- gauginos

Indirect searches via precision measurements of known variables, comparison with SM

► sensitive to, for example, Z' bosons and composite Higgs

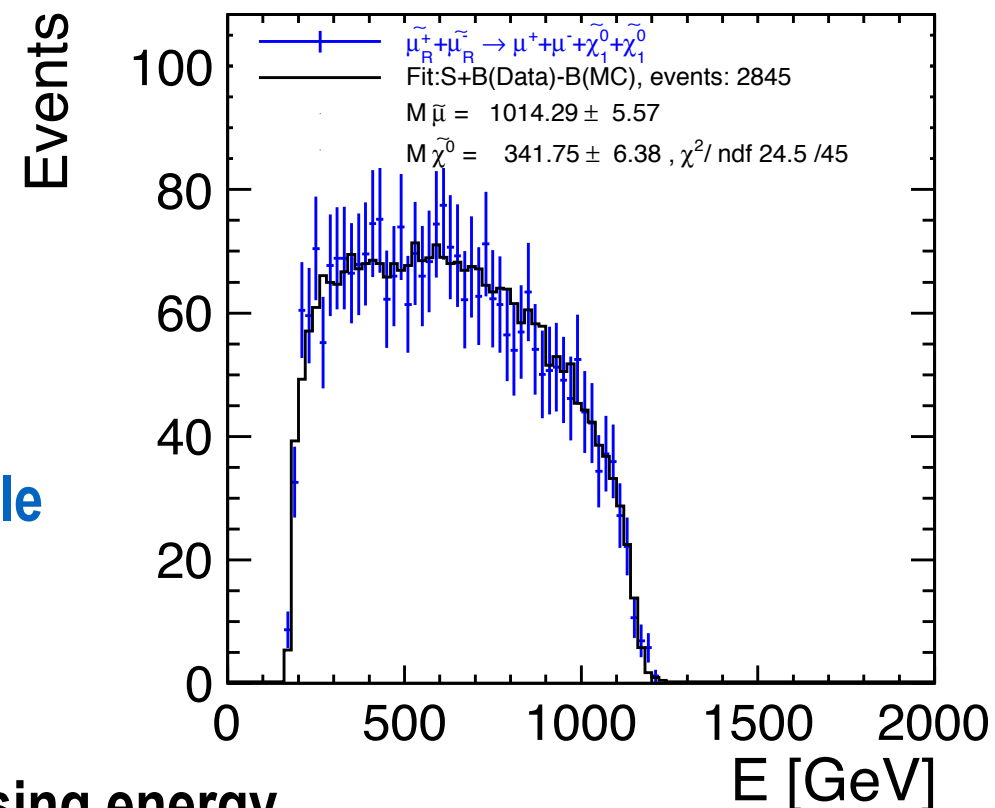
Sleptons and gauginos at $\sqrt{s} = 3$ TeV

Slepton signature very clean: leptons and missing energy

$$e^+ e^- \rightarrow \tilde{\mu}_R^+ \tilde{\mu}_R^- \rightarrow \mu^+ \mu^- \tilde{\chi}_1^0 \tilde{\chi}_1^0$$

Masses from end-points of energy spectra

► for slepton masses ~ 1 TeV, **precisions of $< 1\%$ achievable**



Chargino and neutralino pair production: four jets and missing energy

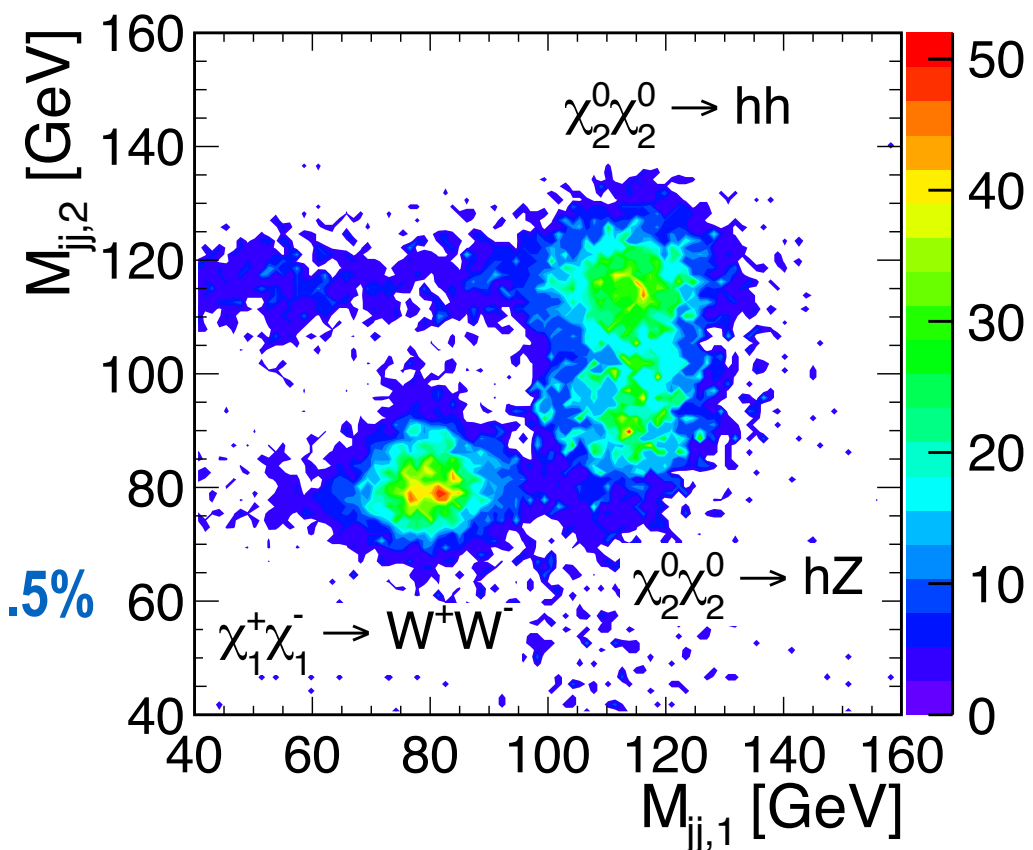
$$e^+ e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 W^+ W^-$$

$$e^+ e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 hh$$

$$e^+ e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 Zh$$

Masses from end-points of energy spectra

► for gaugino masses of few hundred GeV, **precisions of 1-1.5%**



Heavy Higgs bosons at $\sqrt{s} = 3$ TeV

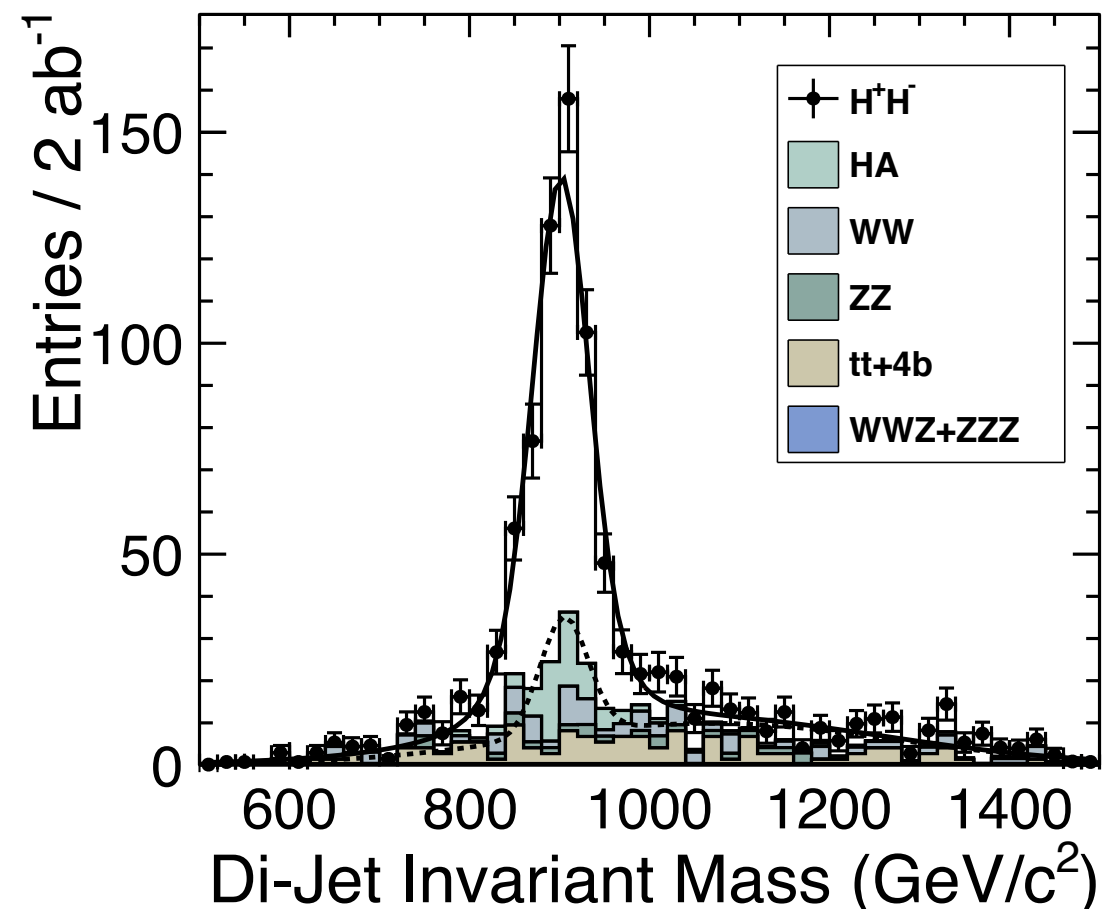
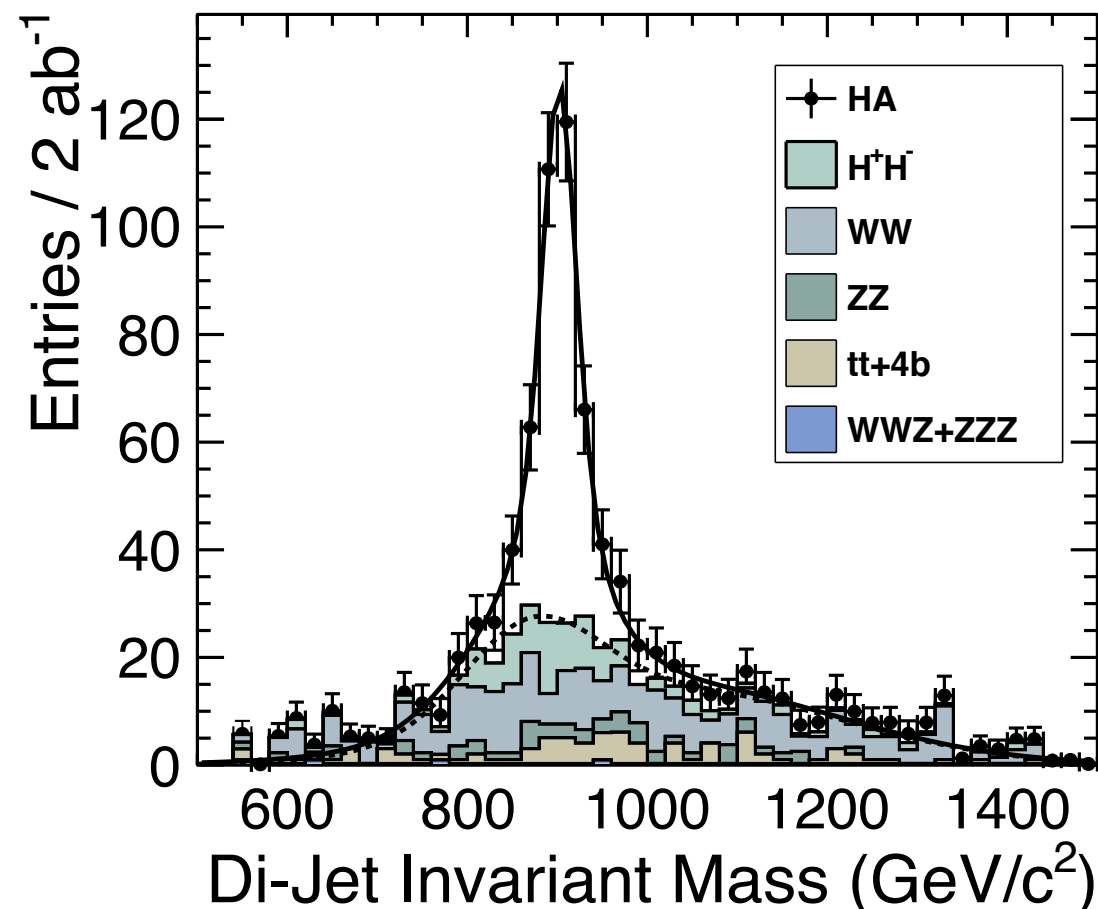
Heavy Higgs bosons (\sim degenerate in mass) would produce complex final states of heavy-flavour jets

$$e^+e^- \rightarrow HA \rightarrow b\bar{b}b\bar{b}$$

$$e^+e^- \rightarrow H^+H^- \rightarrow t\bar{b}b\bar{t}$$

Separation requires heavy-flavour tagging

► accuracy of Higgs mass measurements $\sim 0.3\%$

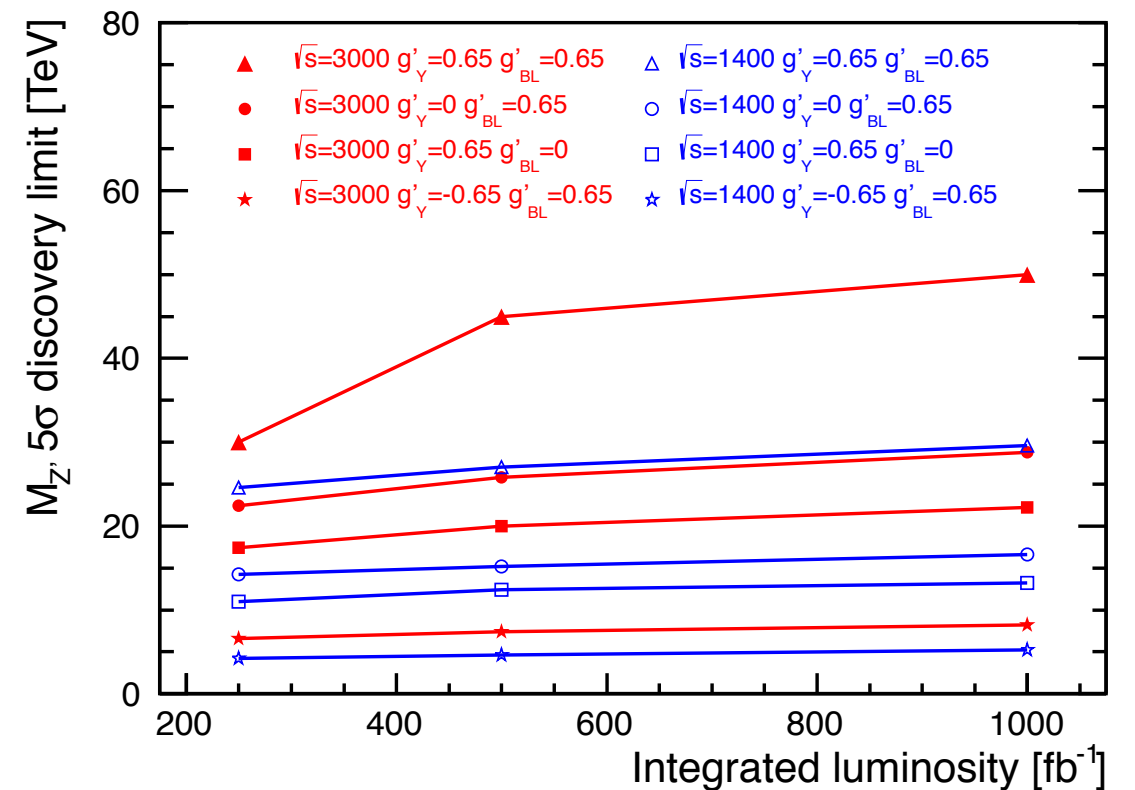
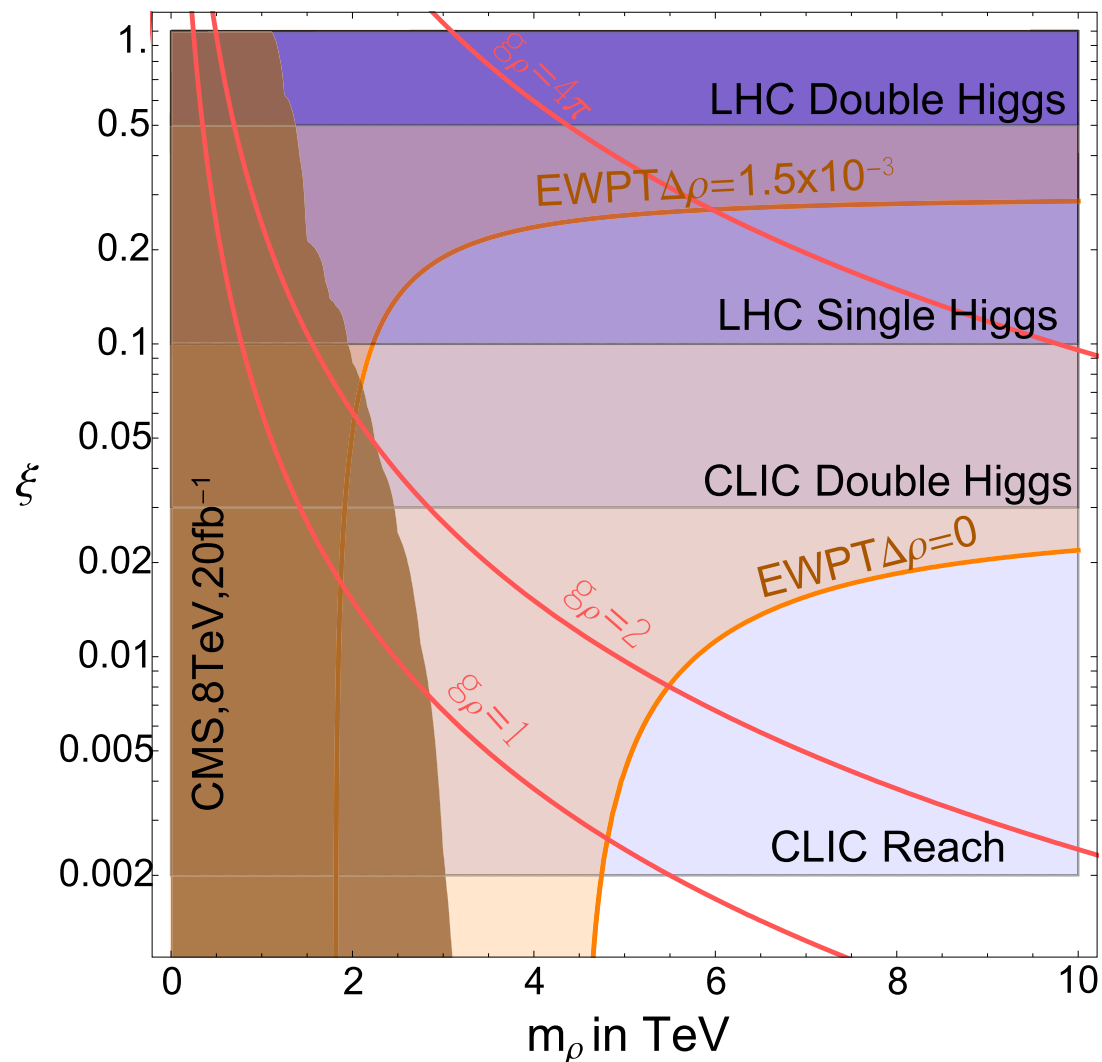


Indirect searches

Precision studies of di-muon pair production

- ▶ total cross section
- ▶ forward-backward asymmetry
- ▶ left-right asymmetry (requires beam polarisation)

Sensitive to Z' bosons, reach up to 10s of TeV



Composite Higgs theories (bound state of fermions)

- ▶ m_ρ : mass of the vector resonance
- ▶ ξ : strength of Higgs interactions

CLIC reach from single Higgs production provides an indirect probe of Higgs composite scale up to 70 TeV

Summary

CLIC offers a strong physics programme throughout its three energy stages

Higgs physics at CLIC

- ▶ model independent measurement of g_{HZZ}
- ▶ high statistics: precision measurements and rare decays
- ▶ top Yukawa coupling and Higgs self-coupling
- ▶ combined, model independent fit of Higgs parameters

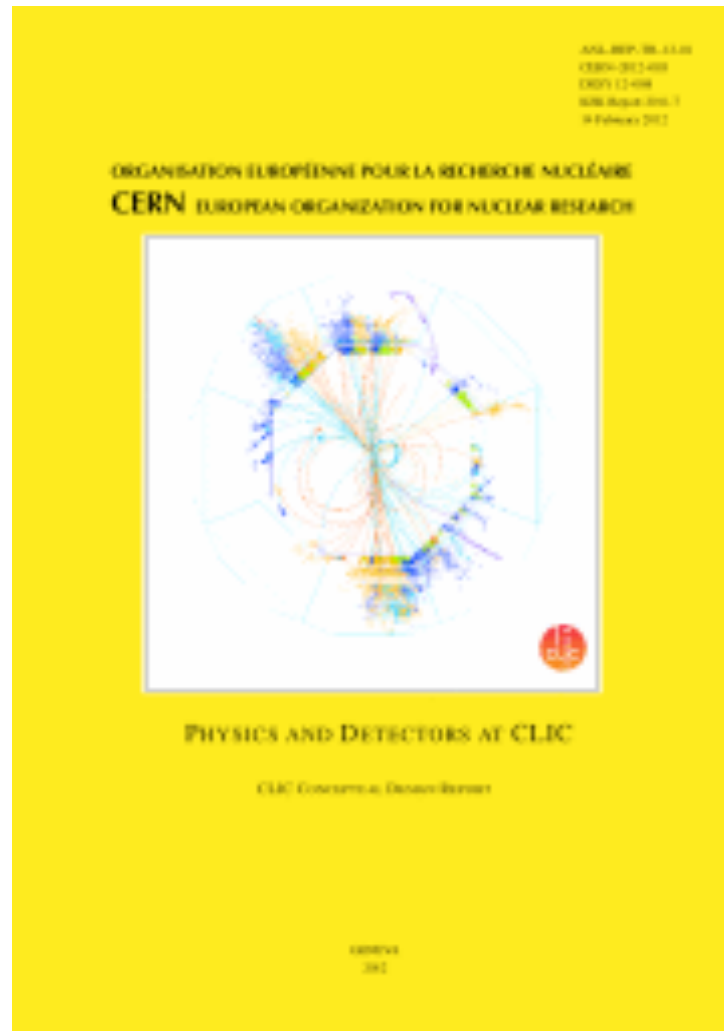
BSM physics at CLIC

- ▶ direct searches possible up to the kinematic limit
- ▶ mass measurements of SUSY particles with %-level precision
- ▶ indirect searches increase reach to tens of TeV

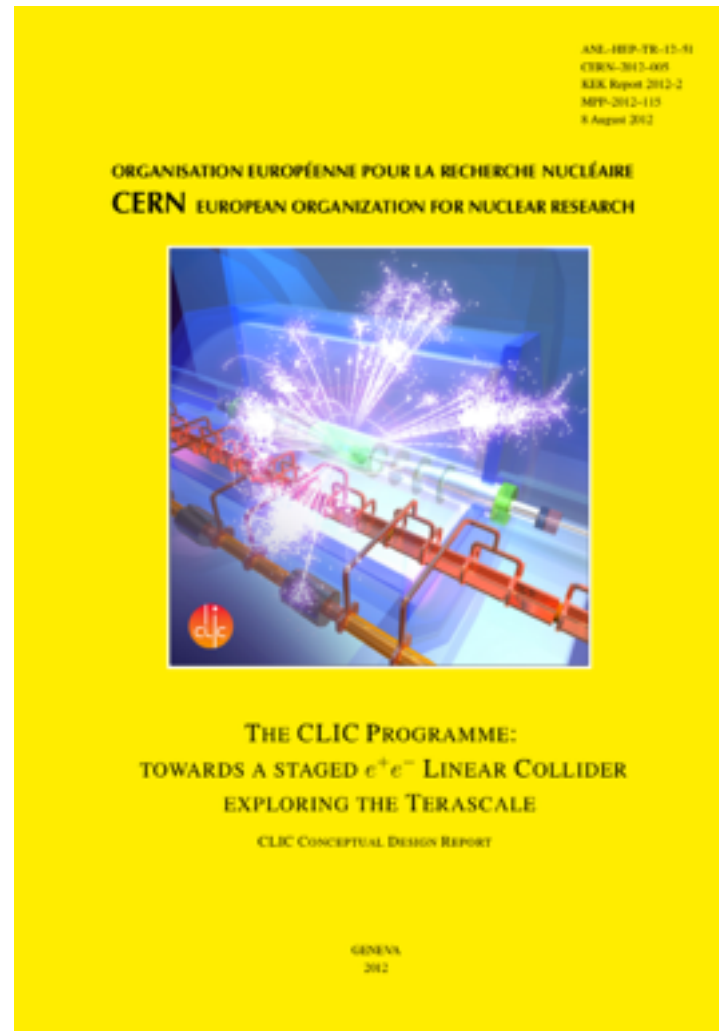


Thanks for your attention!

If you want to know more



CLIC CDR vol 2
Physics and Detectors
CERN-2012-003
arxiv: 1202.5940



CLIC CDR vol 3
The CLIC Programme
CERN-2012-005
arxiv: 1209.2543



CLIC Snowmass White Paper
Physics at CLIC
arxiv: 1307.5288

Optimising the first energy stage

Higgs mass measurement (Higgsstrahlung):

▶ $\sqrt{s} = 250 \text{ GeV}$: $\Delta m_H \sim 30 \text{ MeV}$

▶ $\sqrt{s} = 350 \text{ GeV}$: $\Delta m_H \sim 120 \text{ MeV}$

Alternatively: WW-fusion + $H \rightarrow bb$:

▶ $\sqrt{s} = 350 \text{ GeV}$: $\Delta m_H \sim 50 \text{ MeV}$ (tbc)

Higgs couplings:

▶ Requires access to Higgsstrahlung and WW-fusion: g_{HZZ} , g_{HWW} , Γ_H , couplings to fermions...

▶ $\sqrt{s} = 350 \text{ GeV}$ a good compromise

▶ But results may profit from higher energy&luminosity

$$\sigma(H\nu_e\bar{\nu}_e) \times BR(H \rightarrow WW^*) \propto \frac{g_{HWW}^4}{\Gamma_H}$$

Top physics:

▶ Threshold scan requires $\sqrt{s} = 360 \text{ GeV}$

CLIC with $\sqrt{s} > 1 \text{ TeV}$ will give best perspectives. But in some areas the first energy stage can already improve significantly on HL-LHC:

▶ Top A_{fb} and top couplings to Z, γ , W need $\sqrt{s} > 400 \text{ GeV}$

▶ It might be worth adapting the first stage to enable these studies

