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Horizontal gauge symmetries and the LHC flavour anomalies

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Outline:

- Introduction: LHC flavor anomalies and $L_\mu - L_\tau$
- 1HDM with gauged $L_\mu - L_\tau$ and vector-quarks
 - $B \rightarrow K^* \mu^+ \mu^-$
 - $B \rightarrow K \mu^+ \mu^- / B \rightarrow K e^+ e^-$
- 2HDM with gauged $L_\mu - L_\tau$ and vector- quarks
 - $h \rightarrow \tau \mu$ and $\tau \rightarrow \mu \mu \mu$
- 2HDM and 3HDM with gauged horizontal charged (no vector-quarks)
- Conclusions

Based on: Andreas Crivellin, Giancarlo D'Ambrosio and Julian Heeck

Explaining $B \rightarrow K^* \mu \mu$, $B \rightarrow K \mu \mu / B \rightarrow K e e$ and $h \rightarrow \tau \mu$ in a two-Higgs-doublet model with gauged $L_\mu - L_\tau$
arXiv:1501.00993

Addressing the LHC flavour anomalies with horizontal gauge symmetries (in preparation)

The LHC flavour anomalies

$B \rightarrow K^* \mu \mu$

- 2-3 σ deviation from the SM mostly in C_5'

- Can be explained by $O_9 = \bar{s} \gamma^\mu P_L b \bar{\ell} \gamma_\mu \ell$

Descotes-Genon et al. 1307.5683, Altmannshofer and DS 1308.1501, Beaujean et al. 1310.2478

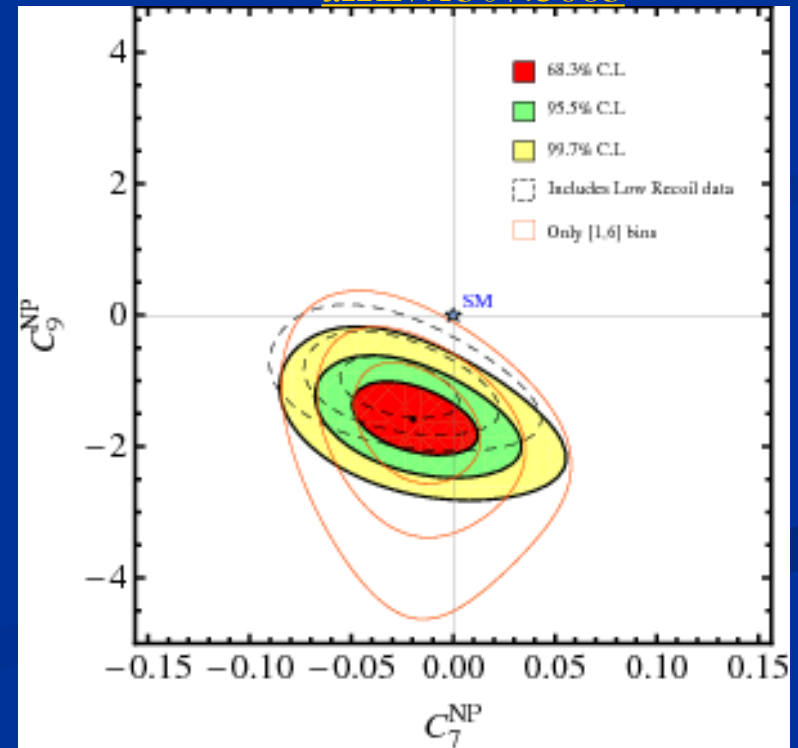
- New physics explanation is not easy [arXiv:1307.5683](https://arxiv.org/abs/1307.5683)

(MSSM, 2HDM, extra dimensions not possible).

- Most natural explanation: Z'

Gauld et al. 1310.1082,
Buras et al. 1311.6729, ...

- Subleading hadronic effects might be larger than expected...



$R(K) = \mathcal{B}(B \rightarrow K_{\mu\mu}) / \mathcal{B}(B \rightarrow Ke e)$

- Lepton flavour universality violation
- 2.6 σ deviation from the theoretically rather clean SM expectation

$$R_K^{\text{SM}} = 1.0003 \pm 0.0001$$

C. Bobeth, G. Hiller, and G. Piranishvili, 0709.4174.

$$R_K^{\text{exp}} = 0.745^{+0.090}_{-0.074} \pm 0.036$$

LHCb 1406.6482

■ Explanation:

- Leptoquarks
- flavour non-universal Z'

e.g. G. Hiller, M. Schmaltz, 1408.1627

W. Altmannshofer, et al. 1403.1269

S. L. Glashow et al. 1411.0565.



Also LFV in B decays?

Talk of Lars Hofer

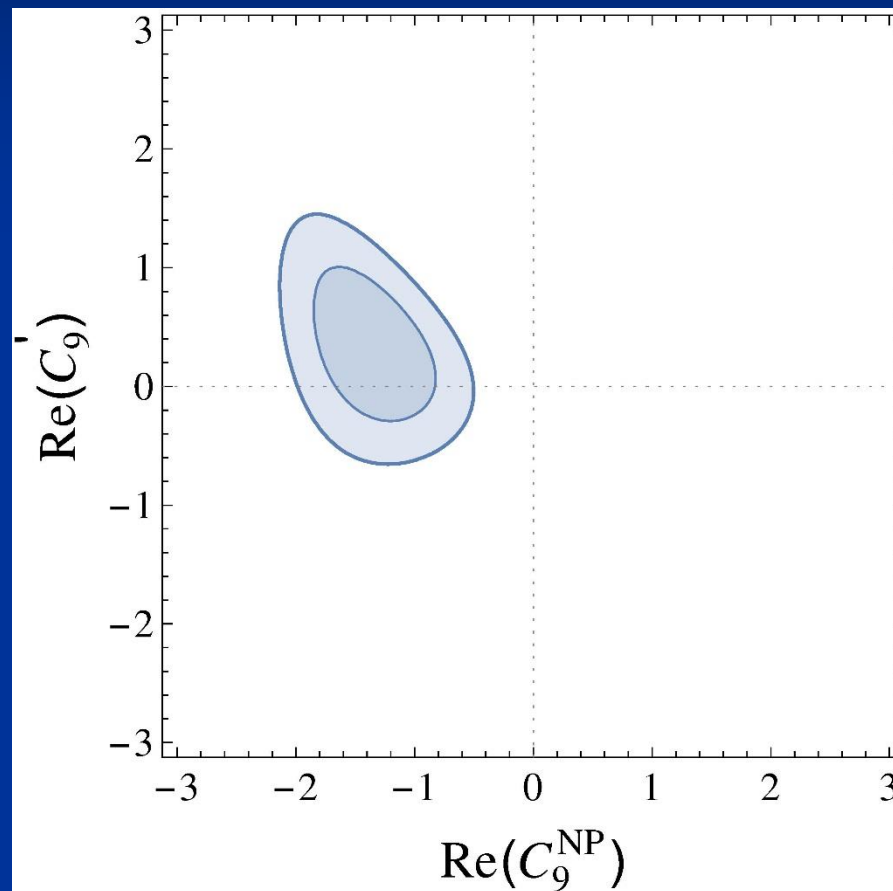
$B \rightarrow K^* \mu\mu$ and $R(K)$

- Contribution to $C_9^{\mu\mu}$ but not C_9^{ee} gives simultaneously a good fit

W. Altmannshofer and D. M. Straub, 1411.3161.
T. Hurth, F. Mahmoudi, and S. Neshatpour, 1410.4545.

(see talk of Tobias Hurth)

- $C_9^{\mu\mu} = -C_{10}^{\mu\mu}$ possible but less good fit.



arXiv:1411.3161

$$h \rightarrow \tau\mu$$

- 2.4 σ difference from zero

$$\text{Br}[h \rightarrow \mu\tau] = \left(0.89^{+0.40}_{-0.37}\right) \quad \text{CMS-PAS-HIG-14-005}$$

- Can be explained in the effective field theory approach by

$$\mathcal{Q}_{e\phi}^{fi} = \ell_f \phi e_i \phi^\dagger \phi$$

R. Harnik, J. Kopp, and J. Zupan, 1209.1397.
G. Blankenburg, J. Ellis, and G. Isidori, 1202.5704.
S. Davidson and P. Verdier, 1211.1248.

- No dominant contribution from vector-like fermions

A. Falkowski, D. M. Straub, and A. Vicente, 1312.5329

A. Dery, et. al., 1408.1371.
M. D. Campos, et. al., 1408.1652.
A. Celis, et. al., 1409.4439.
D. Aristizabal Sierra and A. Vicente, 1409.7690.
C.-J. Lee and J. Tandean, 1410.6803.
J. Heeck, et. al., 1412.3671.



Extended
Higgs sector

Gauged $L_\mu - L_\tau$

- Vectorial U(1) gauge group:
 $Q(e)=0, Q(\mu)=1, Q(\tau)=-1$
- Anomaly free X. He, G. C. et al., Phys.Rev. **D43**, 22 (1991).
R. Foot, Mod.Phys.Lett. **A6**, 527 (1991).
- Good zero order approximation to the PMNS matrix:
 - maximal atmospheric and
 - vanishing reactor neutrino mixing angle

$$M_\nu = \begin{pmatrix} X & 0 & 0 \\ 0 & 0 & Y \\ 0 & Y & 0 \end{pmatrix}$$

P. Binetruy, et al., hep-ph/9610481.
N. F. Bell and R. R. Volkas, hep-ph/0008177.
S. Choubey and W. Rodejohann, hep-ph/0411190.
J. Heeck and W. Rodejohann, 1107.5238

➔ Breaking necessary for a realistic neutrino sector

1HDM with vector-quarks

The Model

W. Altmannshofer, S. Gori, M. Pospelov, and I. Yavin, 1403.1269.

- Gauged $L_\mu - L_\tau$: Z' boson with

$$-ig' \bar{l}_f \gamma^\mu Z'_\mu l_i \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}_{fi}$$

- Vector-like quarks charged under $L_\mu - L_\tau$

$$m_Q \bar{Q}_L \tilde{Q}_R + m_D \bar{D}_L \tilde{D}_R + m_U \bar{U}_L U_R + \text{h.c.}$$

- Effective Z' quark couplings

P. Langacker, 0801.1345., A. J. Buras, F. De Fazio, and J. Girrbach, 1211.1896.

$$ig' \gamma^\mu d_f \left(\Gamma_{fi}^L P_L + \Gamma_{fi}^R P_R \right) d_i Z'_\mu$$

$$\Gamma_{ij}^{dR} \square - \frac{V_\Phi^2}{2m_D^2} (Y_i^D Y_j^{D*}), \quad \Gamma_{ij}^{dL} \square \frac{V_\Phi^2}{2m_Q^2} (Y_i^Q Y_j^{Q*})$$

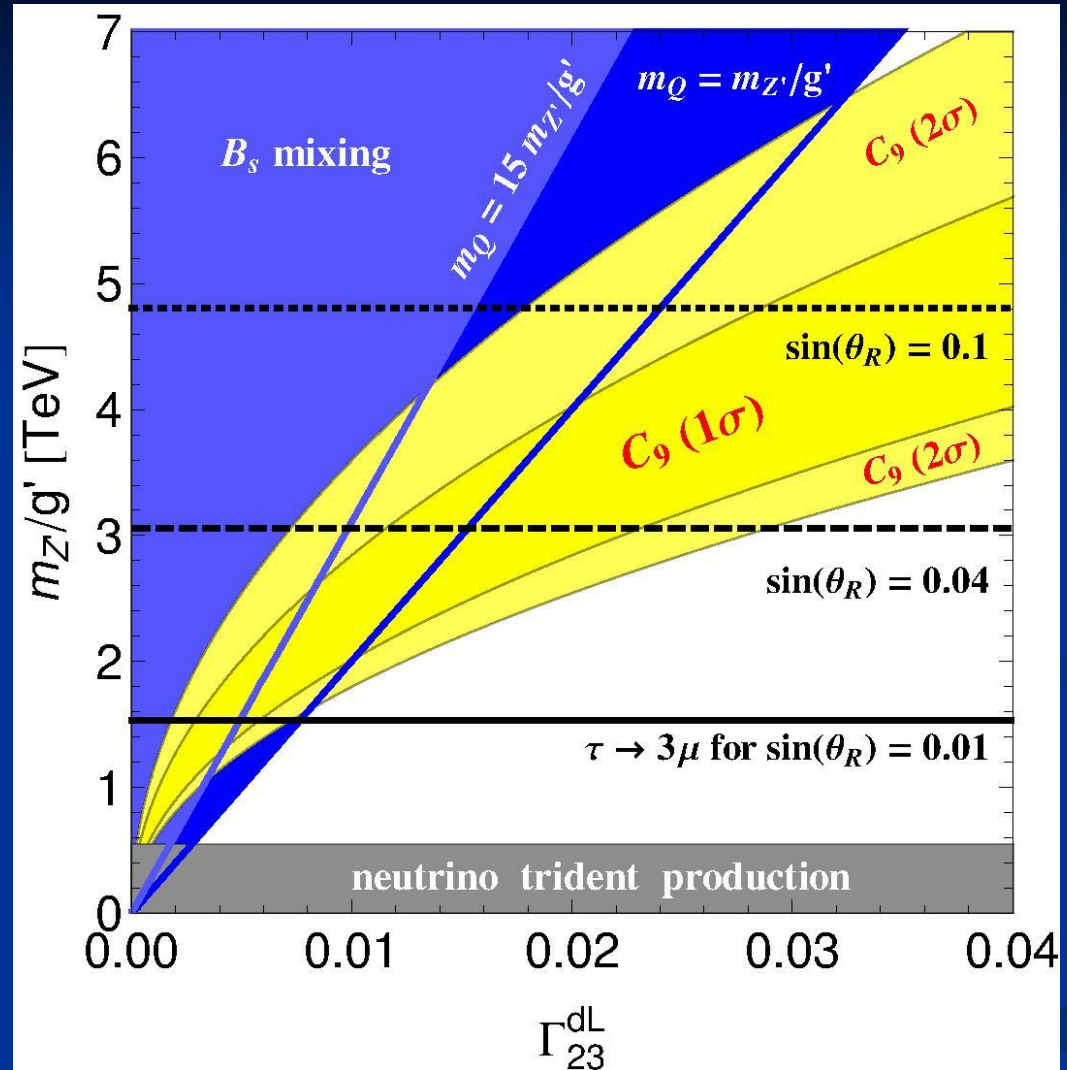
$$m_D^2 \rightarrow \infty$$

$B \rightarrow K^* \mu \mu, R(K)$

$$C_9^{\mu\mu} \square \Gamma_{23}^{dL} \frac{g'^2}{m_{Z'}^2}$$

B_s mixing

$$\frac{\Delta M_{12}}{M_{12}^{\text{SM}}} \square \Gamma_{23}^{dL} \frac{g'^2}{m_{Z'}^2}$$



allowed regions

2HDM with vector-quarks

Explaining $B \rightarrow K^* \mu \mu$, $B \rightarrow K \mu \mu$ / $B \rightarrow K e e$ and $h \rightarrow \tau \mu$ in a two-Higgs-doublet model with gauged $L_\mu - L_\tau$
arXiv:1501.00993

2nd Doublet breaks $L_\mu - L_\tau$

J. Heeck, M. Holthausen, W. Rodejohann and Y. Shimizu, 1412.3671

■ Two Higgs doublets

$$Q_{L_\mu - L_\tau}(\Psi_2) = 0 \quad Q_{L_\mu - L_\tau}(\Psi_1) = 2$$

■ Yukawa couplings

$$L_Y \supset -\bar{\ell}_f Y_i^\ell \delta_{fi} \Psi_2 e_i - \xi_{\tau\mu} \bar{\ell}_3 \Psi_1 e_2 - \bar{Q}_f Y_{fi}^u \tilde{\Psi}_2 u_i - \bar{Q}_f Y_{fi}^d \Psi_2 d_i + \text{h.c.}$$

■ Flavour changing SM-like Higgs coupling

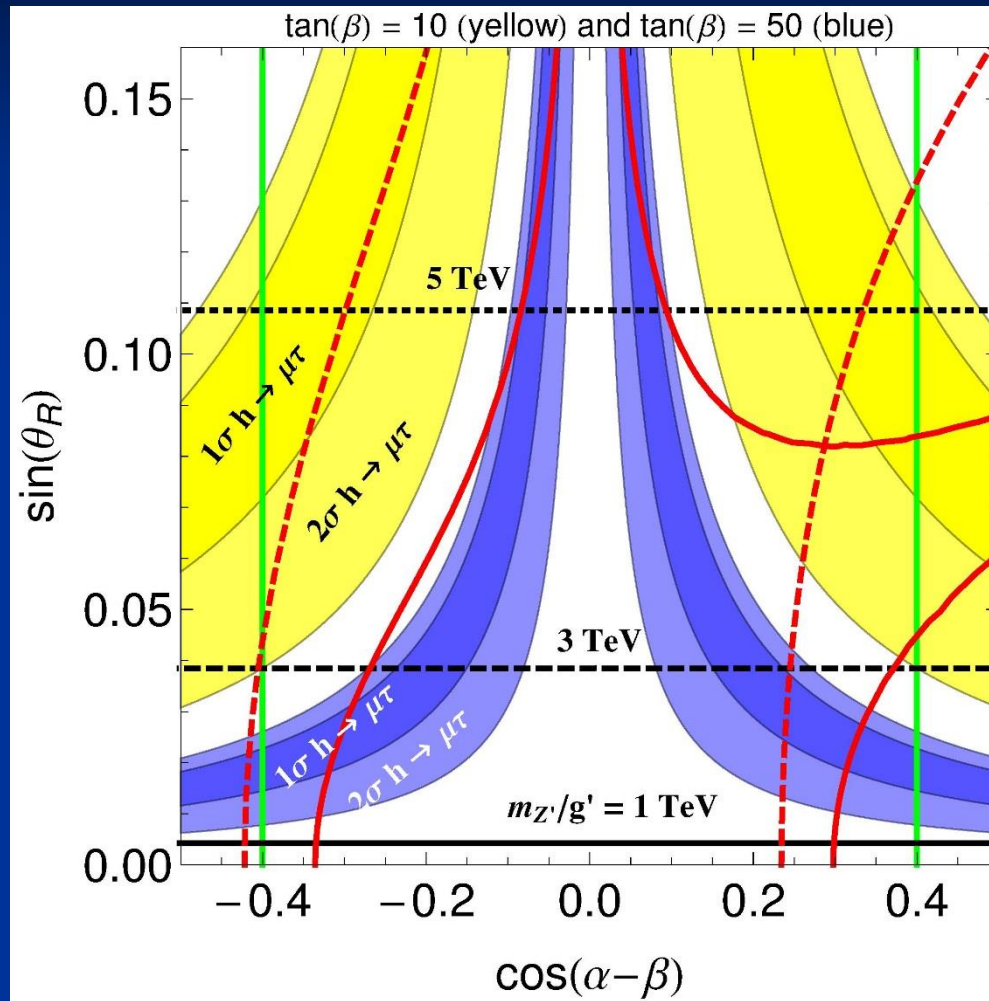
$$\Gamma_{\tau\mu}^h \bar{\tau} P_R \mu h^0 \approx \frac{m_\tau}{v} \frac{\cos(\alpha - \beta)}{\cos(\beta) \sin(\beta)} \theta_R \bar{\tau} P_R \mu h^0 \quad \sin \theta_R \square \frac{v}{\sqrt{2} m_\tau} \xi_{\tau\mu} \cos \beta$$

$$\sin \theta_L \square 0$$

■ Lepton flavour violating Z' couplings

$$g' Z'(\bar{\mu}, \bar{\tau}) \begin{pmatrix} \cos 2\theta_R & \sin 2\theta_R \\ \sin 2\theta_R & -\cos 2\theta_R \end{pmatrix} \gamma^\nu P_R \begin{pmatrix} \mu \\ \tau \end{pmatrix}$$

$h \rightarrow \mu\tau$ and $\tau \rightarrow \mu\mu\mu$



$h \rightarrow \gamma\gamma$ etc.

$h \rightarrow \tau\tau$



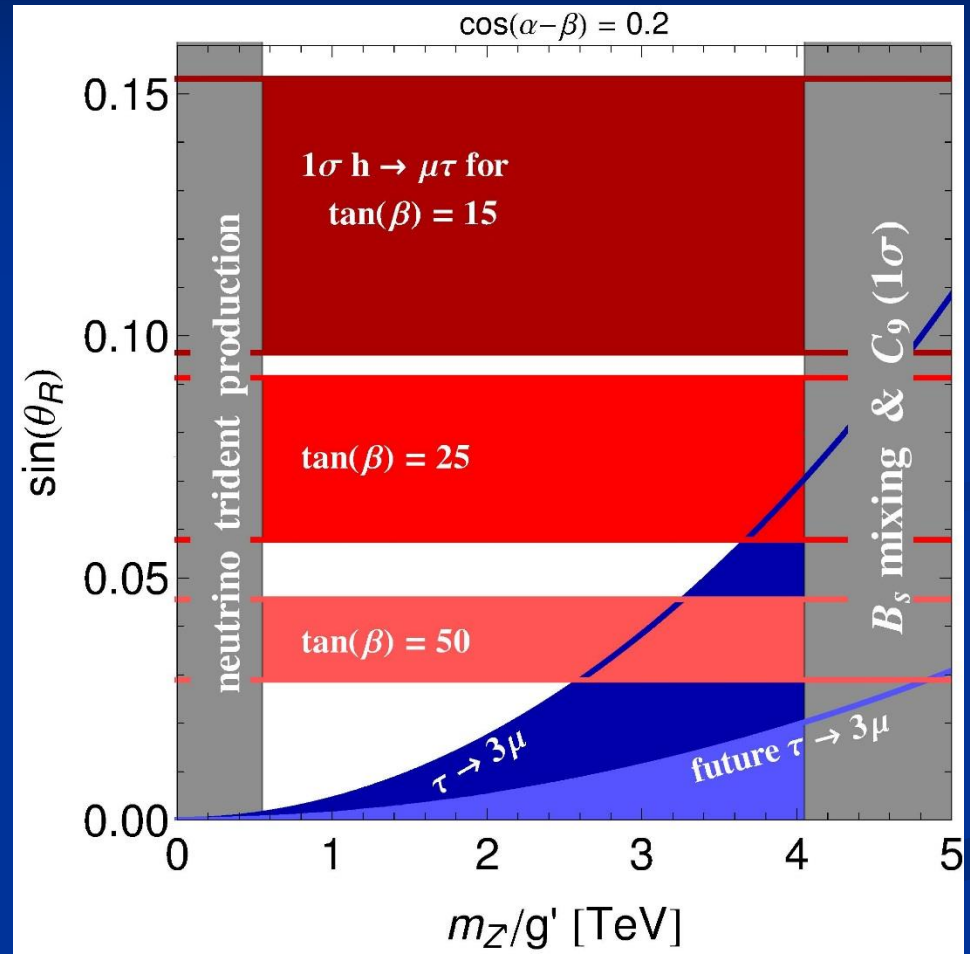
$h \rightarrow \mu\tau$ ($\tan\beta = 50$)



$h \rightarrow \mu\tau$ ($\tan\beta = 10$)

$\tau \rightarrow \mu\mu\mu$ and $h \rightarrow \mu\tau$

- excluded
- allowed by $h \rightarrow \tau\mu$
- allowed by $\tau \rightarrow \mu\mu\mu$



Horizontal charges

Andreas Crivellin, Giancarlo D'Ambrosio and Julian Heeck

Addressing the LHC flavour anomalies with horizontal gauge symmetries (in preparation)

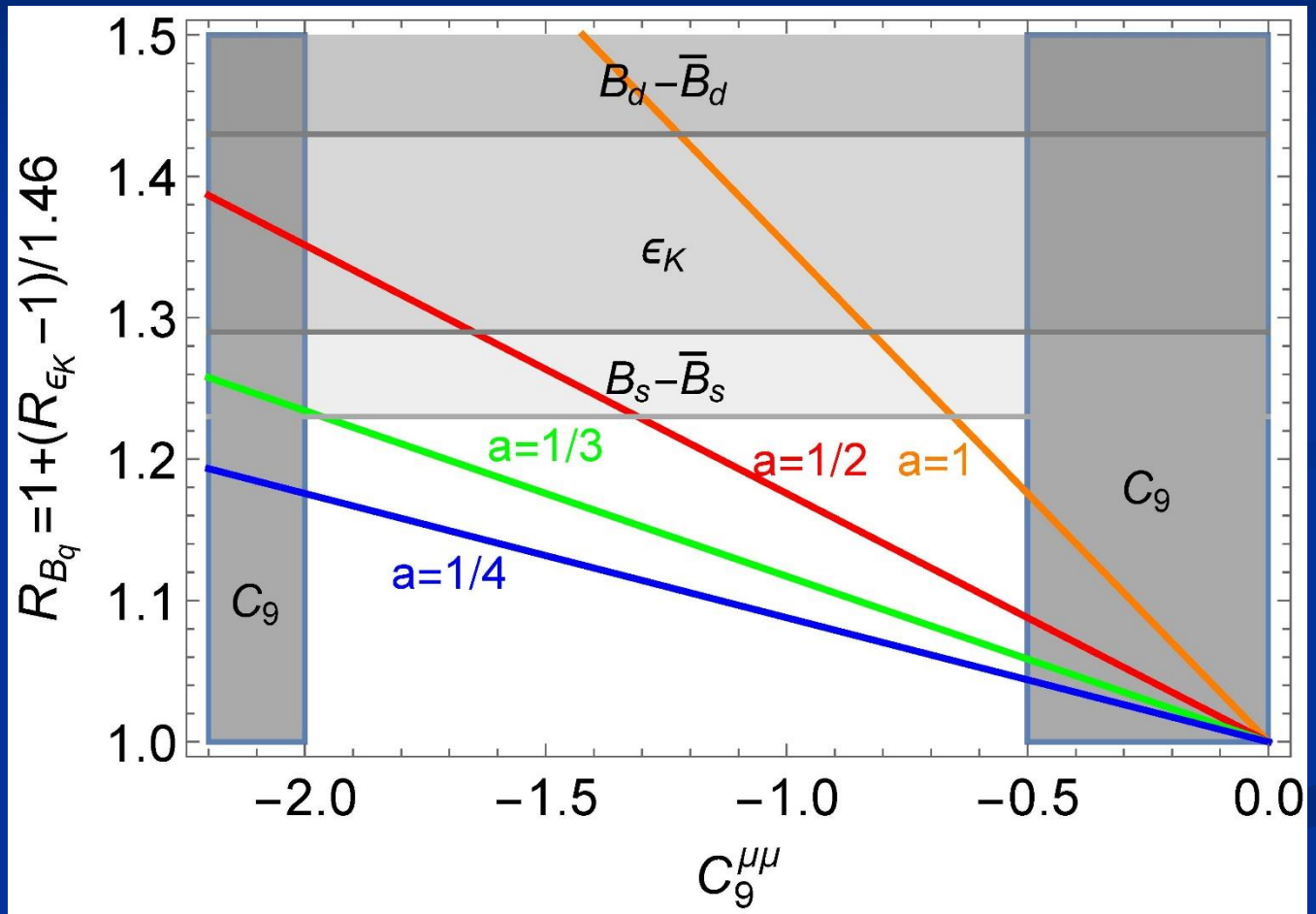
Charge assignmet

- Avoid vector-like quarks by assigning charges to baryons as well
 - ➔ same mechanism in the quark and lepton sector
- Use $L_\mu - L_\tau$ in the lepton sector
 - ➔ good symmetry for the PMNS matrix
 - ➔ effect in $C_9^{\mu\mu}$ but not C_9^{ee}
- First two quark generations must have the same charges because the large Cabibbo angle would lead to huge effect in Kaon mixing
- Anomaly free
 - ➔ $Q(B) = (-a, -a, 2a)$

$\Delta F=2$: Z' contribution

$$R_{B_q} \equiv \frac{\Delta m_{B_q}}{\Delta m_{B_q}^{SM}}$$

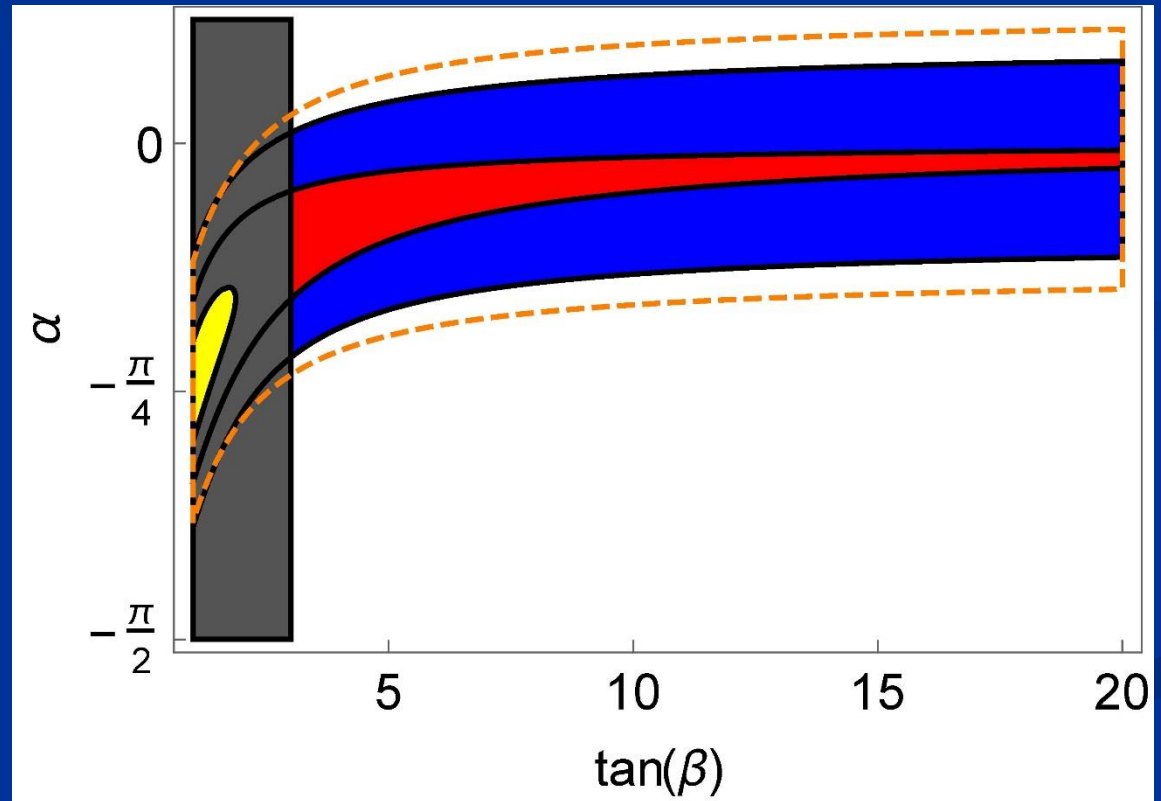
$$R_{\epsilon_K} \equiv \frac{\epsilon_K}{\epsilon_K^{SM}}$$



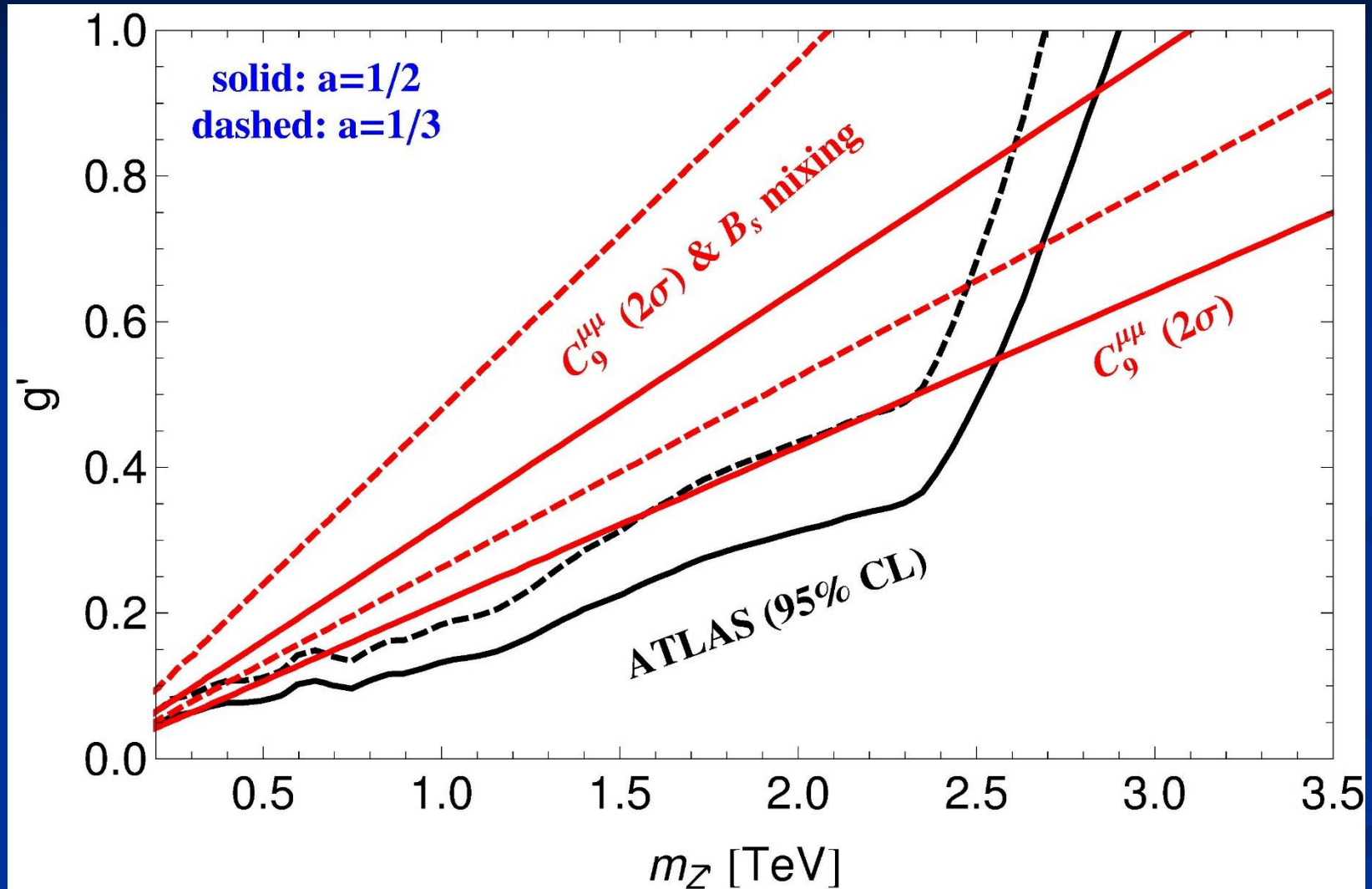
$\Delta F=2$: Higgs contributions

$$m_H = 300 \text{ GeV}, \quad C_9^{\mu\mu} = -1.3$$

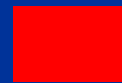
- $m_A = 250 \text{ GeV}$
- $m_A = 300 \text{ GeV}$
- $m_A = 350 \text{ GeV}$



LHC limits



ATLAS



$C_9^{\mu\mu} \& B_s - \bar{B}_s$

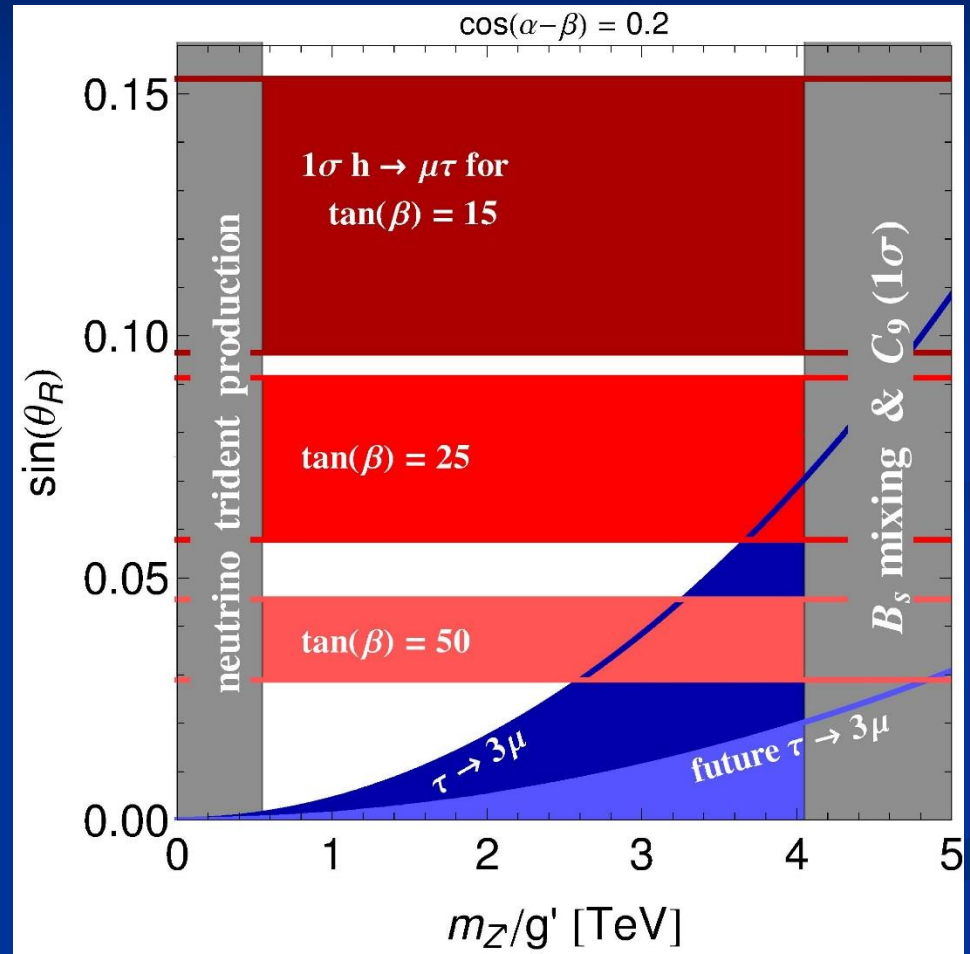
3HDM

- Same effect in

$$\tau \rightarrow \mu\mu\mu$$

$$h \rightarrow \mu\tau$$

provided that the mixing among the doublets is small



Conclusions

- The LHC found three anomalies in the flavour sector
 - $h \rightarrow \tau\mu$
 - $B \rightarrow K^* \mu^+ \mu^-$
 - $B \rightarrow K \mu^+ \mu^- / B \rightarrow K e^+ e^-$
- All three anomalies can be explained in a model with gauged $L_\mu - L_\tau$
 - 2HDM with vector-quarks
 - 3HDM with gauged flavour dependent B-L charges