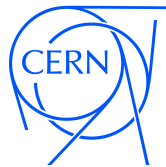


27th Rencontres de Blois

Constraining top quark Z-boson
couplings at the LHC

Markus Schulze



in collaboration with R. Röntsch

based on
JHEP 1407 (2014) 091; arXiv:1501.05939

Twenty years after the top quark discovery



- Our understanding of the top quark as an elementary particle and its dynamics in QCD is very solid.
- Many of its properties were explored at the Tevatron.
- What about interactions with the electroweak sector of the SM?

The Top Quark and Electroweak Theory

→ **What do we know about top quark interactions with γ , Z , W or H ?**

- This question immediately motivates the study of the processes

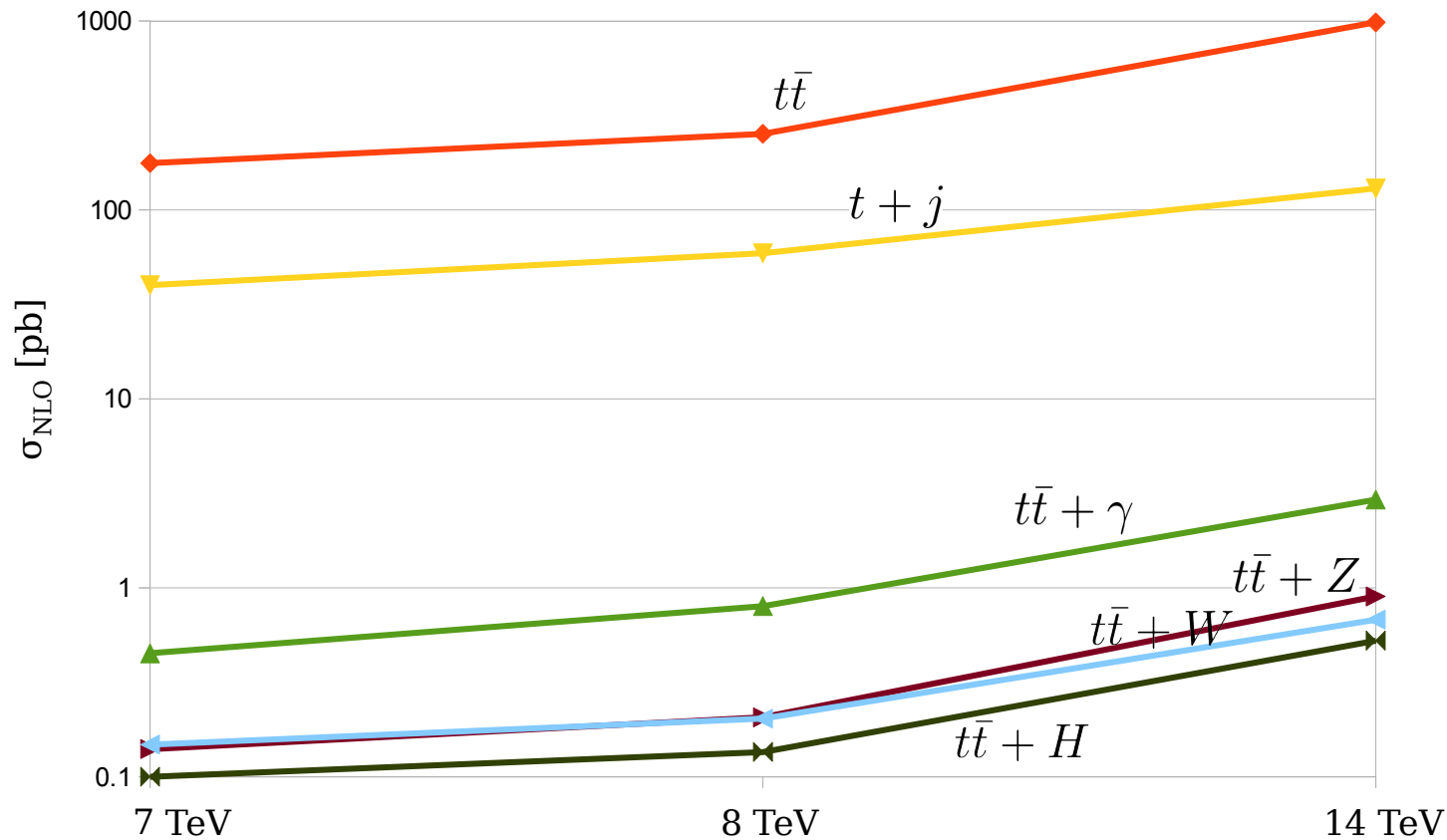
$$t\bar{t} + V \quad t\bar{t} + H$$

- CDF finds 30 $t\bar{t}b + \gamma$ events. All other processes have never been observed at the Tevatron (high threshold + penalties from branching fractions).

→ Study of these processes at the LHC will open a new era in top quark physics

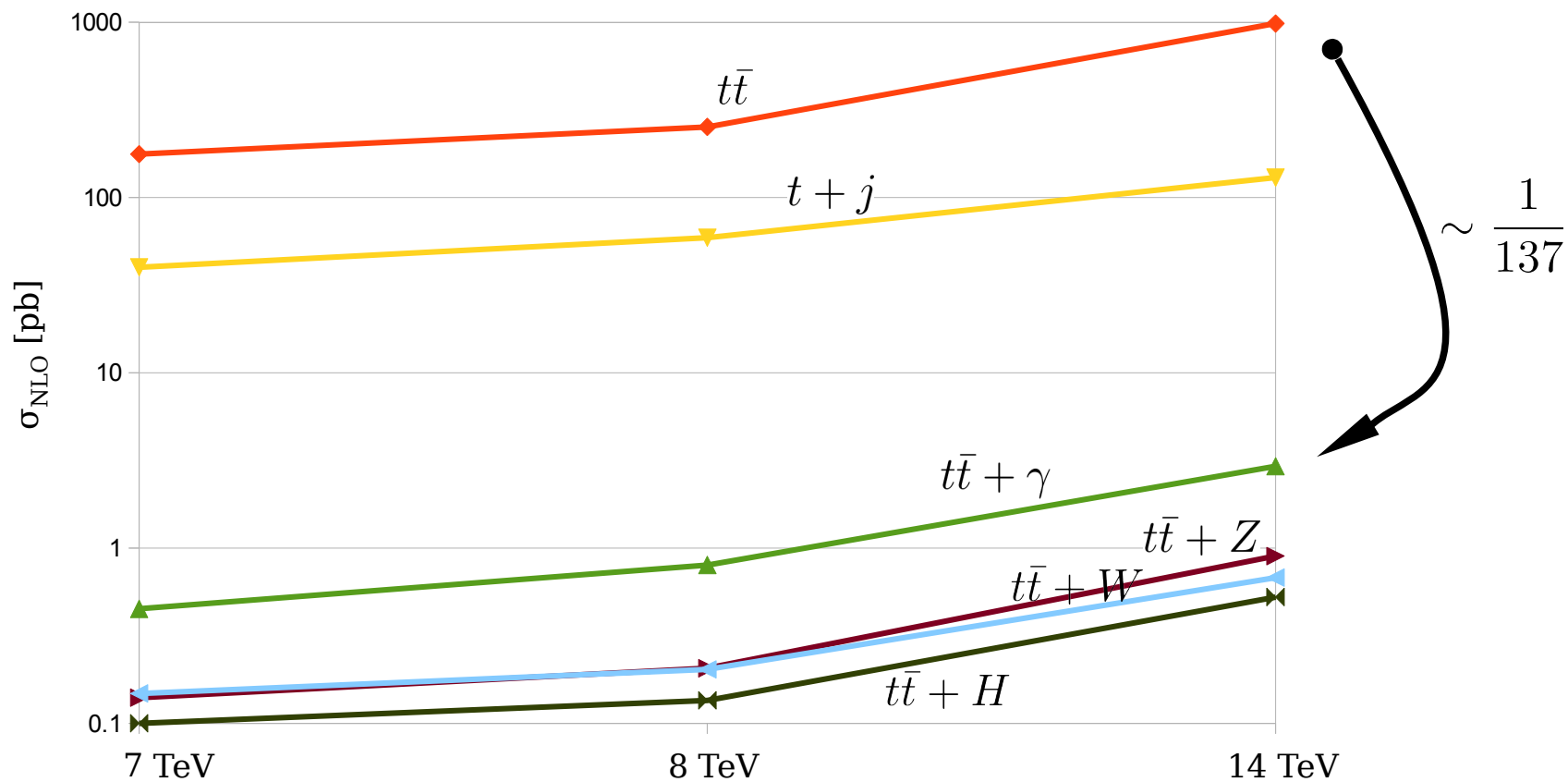
→ What precision can be reached?

Spectrum of cross sections



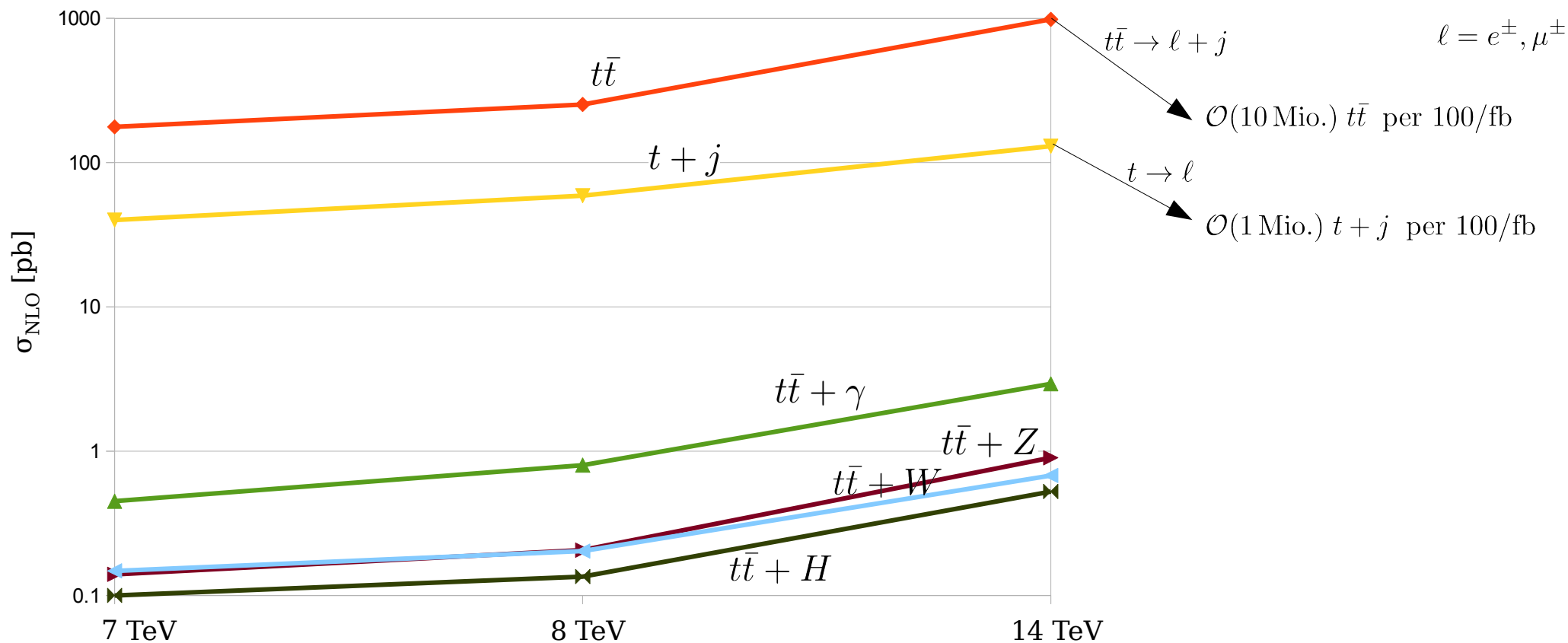
[Baernreuther, Czakon, Fiedler, Mitov], [Campbell, Ellis, Röntsch],
[Melnikov, Scharf, M.S.], [Garzelli, Kardos, Papadopoulos, Trocsanyi],
[Demartin, Maltoni, Mawatari, Page, Zaro]

Spectrum of cross sections



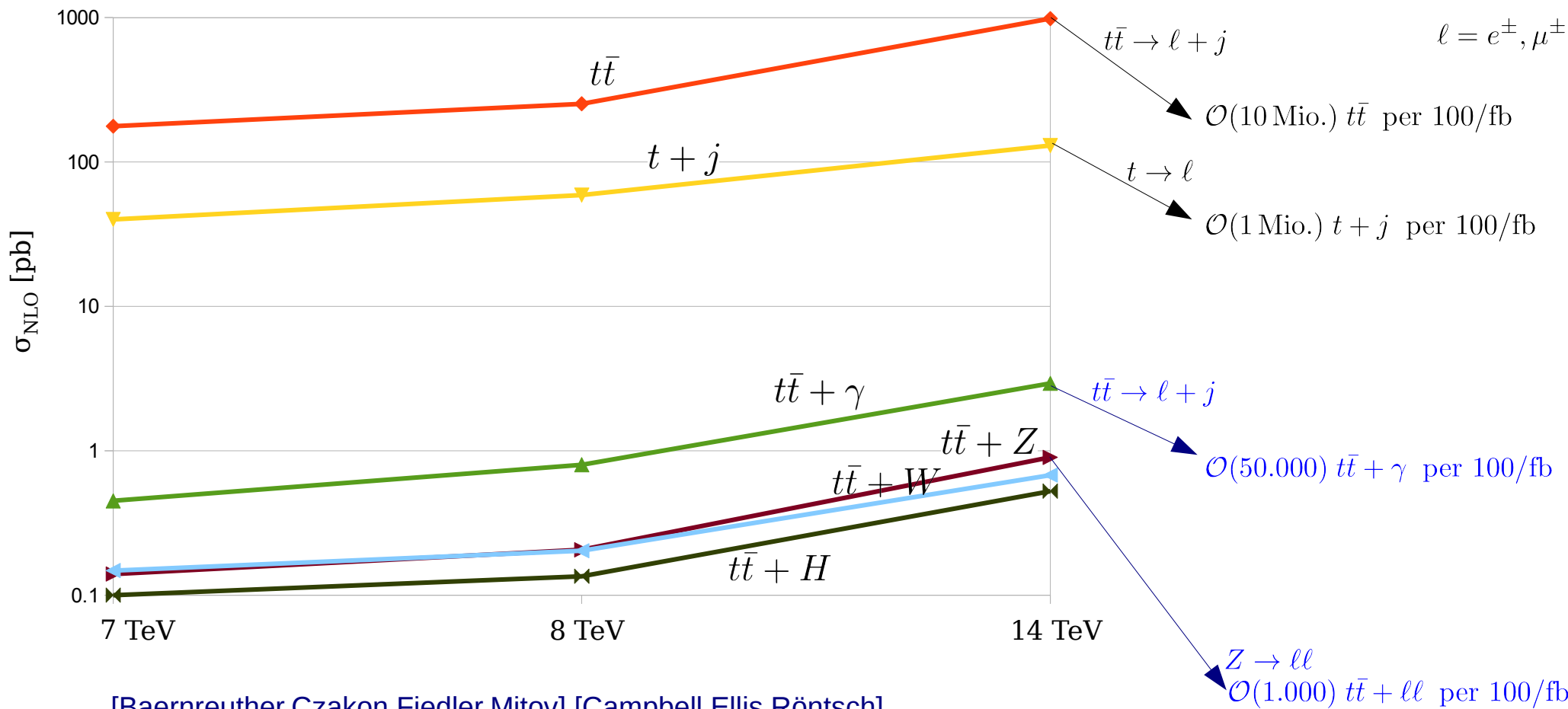
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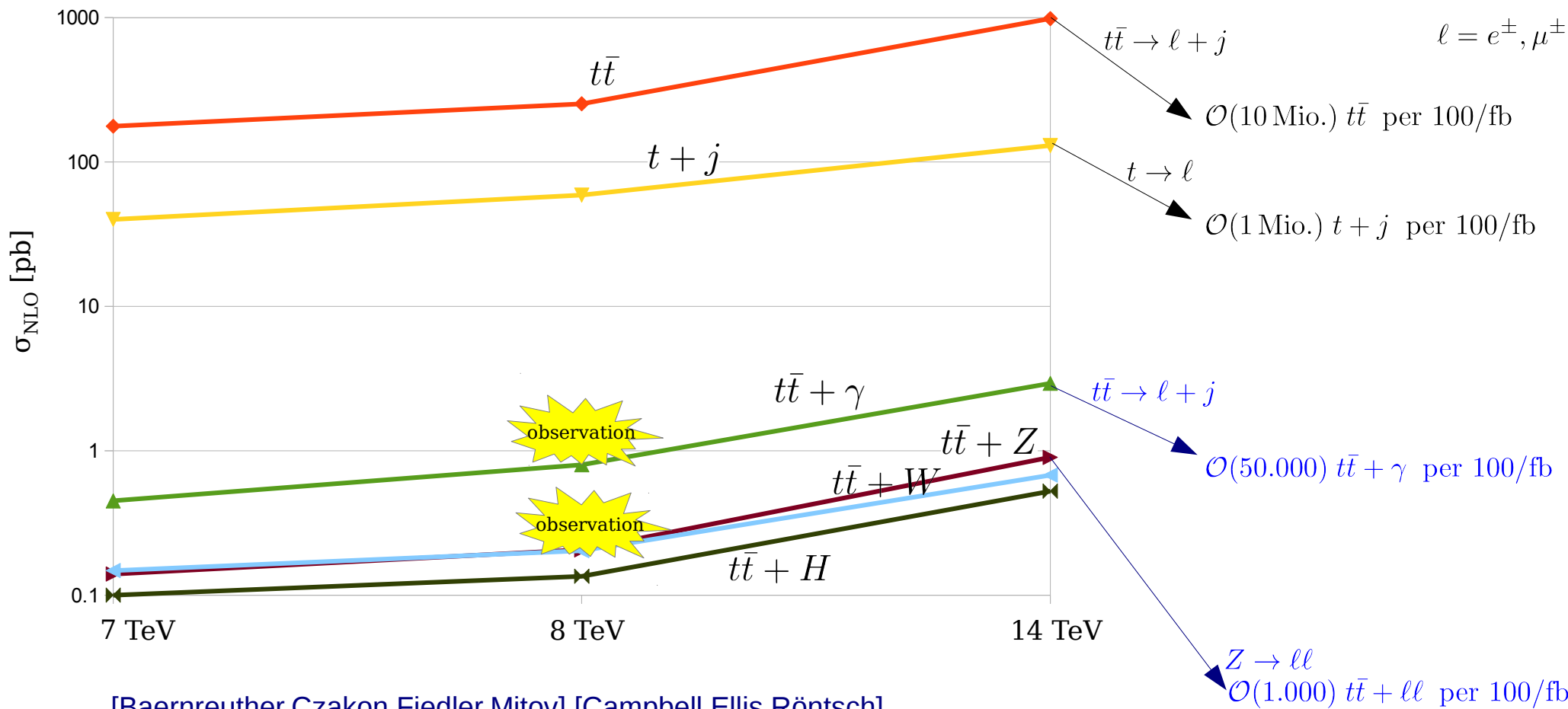
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Spectrum of cross sections



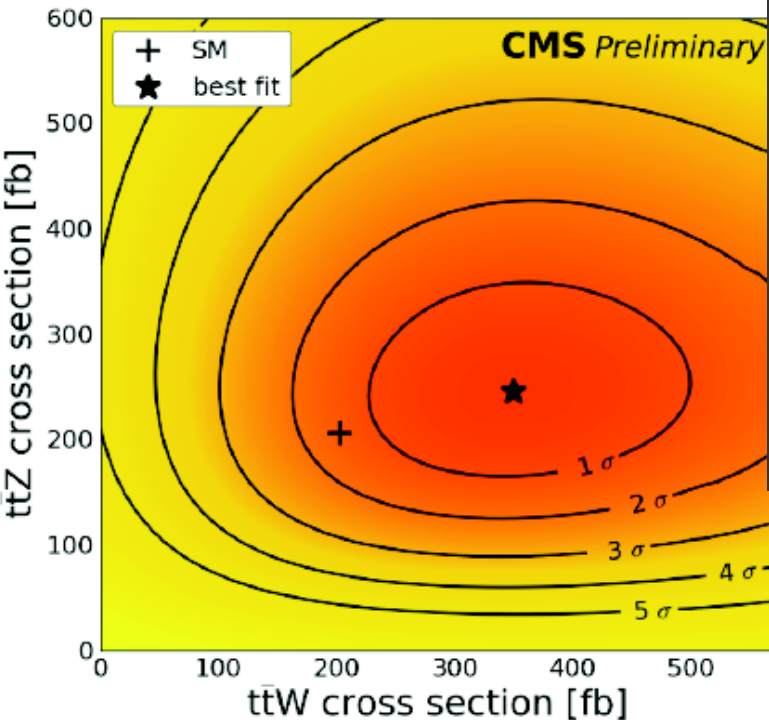
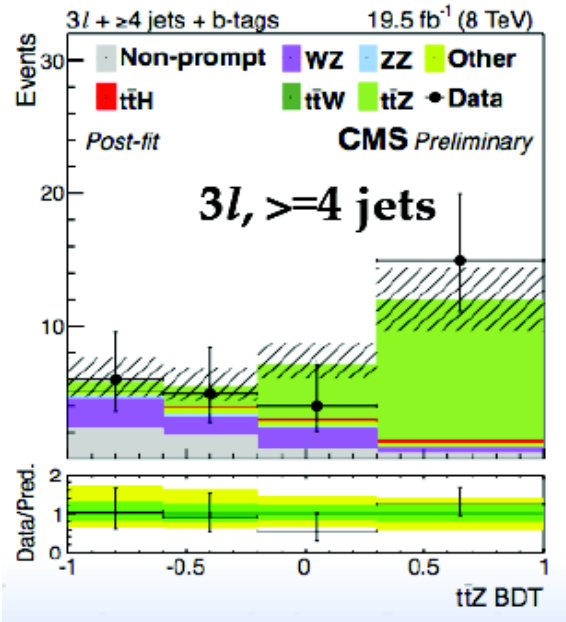
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 [Melnikov, Scharf, M.S.], [Garzelli, Kardos, Papadopoulos, Trocsanyi],
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Experimental results from Run-1



Summary

<i>ttW and ttZ measurements</i>		<i>ttW</i>				<i>ttZ</i>			
		<i>Cross section</i>		<i>Significance</i>		<i>Cross section</i>		<i>Significance</i>	
<i>Data</i>	<i>Analysis</i>	<i>Theory*</i>	<i>Obs.</i>	<i>Exp.</i>	<i>Obs.</i>	<i>Theory*</i>	<i>Obs.</i>	<i>Exp.</i>	<i>Obs.</i>
7 TeV (5 fb^{-1})	CMS ^[1]	147^{+14}_{-16}	-	-	-	137^{+12}_{-16}	280^{+150}_{-110}	?	3.3
	ATLAS ^[2]		-	-	-		< 710	-	-
8 TeV (20 fb^{-1})	CMS ^[3]	203^{+20}_{-22}	170^{+110}_{-100}	2.0	1.6	206^{+19}_{-24}	200^{+90}_{-80}	3.1	3.1
	ATLAS ^[4]		300^{+140}_{-110}	2.3	3.1		150^{+58}_{-54}	3.4	3.1
	CMS Preliminary		382^{+117}_{-102}	3.5	4.8		242^{+65}_{-55}	5.7	6.4

Preliminary documentation: twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsTOP14021

[1] CMS: *Phys. Rev. Lett.* 110 (2013) 172002

[2] ATLAS-CONF-12-126

[3] CMS: EPJ C74 (2014) 3060

[4] ATLAS-CONF-2014-038

* NLO cross sections with scale uncertainties from

Garzelli et. al., *JHEP* 11 (2012) 056



Anomalous couplings

Minimal set of top quark anomalous couplings:

[Aguilar-Saavedra]

$$\begin{aligned}\mathcal{L}_{Wtb} &= -\frac{g}{\sqrt{2}}\bar{b}\gamma^\mu(V_L P_L + V_R P_R)t W_\mu^- - \frac{g}{\sqrt{2}}\bar{b}\frac{i\sigma^{\mu\nu}q_\nu}{M_W}(g_L P_L + g_R P_R)t W_\mu^- + \text{H.c.}\dots \\ \mathcal{L}_{Ztt} &= -\frac{g}{2c_W}\bar{t}\gamma^\mu(X_{tt}^L P_L + X_{tt}^R P_R - 2s_W^2 Q_t)t Z_\mu - \frac{g}{2c_W}\bar{t}\frac{i\sigma^{\mu\nu}q_\nu}{M_Z}(d_V^Z + id_A^Z\gamma_5)t Z_\mu, \\ \mathcal{L}_{\gamma tt} &= -eQ_t\bar{t}\gamma^\mu t A_\mu - e\bar{t}\frac{i\sigma^{\mu\nu}q_\nu}{m_t}(d_V^\gamma + id_A^\gamma\gamma_5)t A_\mu.\end{aligned}$$

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$$\mathcal{L}_{Ztt} = -\frac{g}{2c_W} \bar{t} \gamma^\mu (X_{tt}^L P_L + X_{tt}^R P_R - 2s_W^2 Q_t) t Z_\mu - \frac{g}{2c_W} \bar{t} \frac{i\sigma^{\mu\nu} q_\nu}{M_Z} (d_V^Z + id_A^Z \gamma_5) t Z_\mu,$$

$$\mathcal{L}_{\gamma tt} = -e Q_t \bar{t} \gamma^\mu t A_\mu - e \bar{t} \frac{i\sigma^{\mu\nu} q_\nu}{m_t} (d_V^\gamma + id_A^\gamma \gamma_5) t A_\mu.$$

$$\mathcal{L}^{\text{eff}} = \sum \frac{C_x}{\Lambda^2} O_x + \dots,$$

$$\delta g_L = \sqrt{2} C_{dW}^{33*} \frac{v^2}{\Lambda^2}, \quad \delta V_L = C_{\phi q}^{(3,33)*} \frac{v^2}{\Lambda^2}, \quad \delta d_V^\gamma = \frac{\sqrt{2}}{e} \text{Re} [c_W C_{uB\phi}^{33} + s_W C_{uW}^{33}] \frac{vm_t}{\Lambda^2},$$

$$\delta g_R = \sqrt{2} C_{uW}^{33} \frac{v^2}{\Lambda^2}, \quad \delta V_R = \frac{1}{2} C_{\phi\phi}^{33} \frac{v^2}{\Lambda^2}, \quad \delta d_A^\gamma = \frac{\sqrt{2}}{e} \text{Im} [c_W C_{uB\phi}^{33} + s_W C_{uW}^{33}] \frac{vm_t}{\Lambda^2}.$$

$$\delta X_{tt}^L = \text{Re} [C_{\phi q}^{(3,33)} - C_{\phi q}^{(1,33)}] \frac{v^2}{\Lambda^2}, \quad \delta d_V^Z = \sqrt{2} \text{Re} [c_W C_{uW}^{33} - s_W C_{uB\phi}^{33}] \frac{v^2}{\Lambda^2},$$

$$\delta X_{tt}^R = -\text{Re} C_{\phi u}^{33} \frac{v^2}{\Lambda^2}, \quad \delta d_A^Z = \sqrt{2} \text{Im} [c_W C_{uW}^{33} - s_W C_{uB\phi}^{33}] \frac{v^2}{\Lambda^2}.$$

$$\delta Y_t^V = -\frac{3}{2} \text{Re} C_{u\phi}^{33} \frac{v^2}{\Lambda^2},$$

$$\delta Y_t^A = -\frac{3}{2} \text{Im} C_{u\phi}^{33} \frac{v^2}{\Lambda^2}.$$

$$O_{\phi q}^{(3,ij)} = i(\phi^\dagger \tau^I D_\mu \phi)(\bar{q}_{Li} \gamma^\mu \tau^I q_{Lj}),$$

$$O_{\phi q}^{(1,ij)} = i(\phi^\dagger D_\mu \phi)(\bar{q}_{Li} \gamma^\mu q_{Lj}),$$

$$O_{\phi\phi}^{ij} = i(\tilde{\phi}^\dagger D_\mu \phi)(\bar{u}_{Ri} \gamma^\mu d_{Rj}),$$

$$O_{\phi u}^{ij} = i(\phi^\dagger D_\mu \phi)(\bar{u}_{Ri} \gamma^\mu u_{Rj}),$$

$$O_{uW}^{ij} = (\bar{q}_{Li} \sigma^{\mu\nu} \tau^I u_{Rj}) \tilde{\phi} W_{\mu\nu}^I,$$

$$O_{dW}^{ij} = (\bar{q}_{Li} \sigma^{\mu\nu} \tau^I d_{Rj}) \phi W_{\mu\nu}^I,$$

$$O_{uB\phi}^{ij} = (\bar{q}_{Li} \sigma^{\mu\nu} u_{Rj}) \tilde{\phi} B_{\mu\nu},$$

$$O_{Du}^{ij} = (\bar{q}_{Li} D_\mu u_{Rj}) D^\mu \tilde{\phi},$$

$$O_{Du}^{ij} = (D_\mu \bar{q}_{Li} u_{Rj}) D^\mu \tilde{\phi},$$

$$O_{Dd}^{ij} = (\bar{q}_{Li} D_\mu d_{Rj}) D^\mu \phi,$$

$$O_{Dd}^{ij} = (D_\mu \bar{q}_{Li} d_{Rj}) D^\mu \phi,$$

$$O_{qW}^{ij} = \bar{q}_{Li} \gamma^\mu \tau^I D^\nu q_{Lj} W_{\mu\nu}^I,$$

$$O_{qB}^{ij} = \bar{q}_{Li} \gamma^\mu D^\nu q_{Lj} B_{\mu\nu},$$

$$O_{uB}^{ij} = \bar{u}_{Ri} \gamma^\mu D^\nu u_{Rj} B_{\mu\nu},$$

Anomalous couplings

Minimal set of top quark anomalous couplings:

[Aguilar-Saavedra]

$$\mathcal{L}_{Vf_i f_j}^{\text{OS}} = \bar{f}_j \gamma^\mu (\mathcal{A}_L P_L + \mathcal{A}_R P_R) f_i V_\mu + \bar{f}_j i\sigma^{\mu\nu} q_\nu (\mathcal{B}_L P_L + \mathcal{B}_R P_R) f_i V_\mu + \text{H.c.},$$

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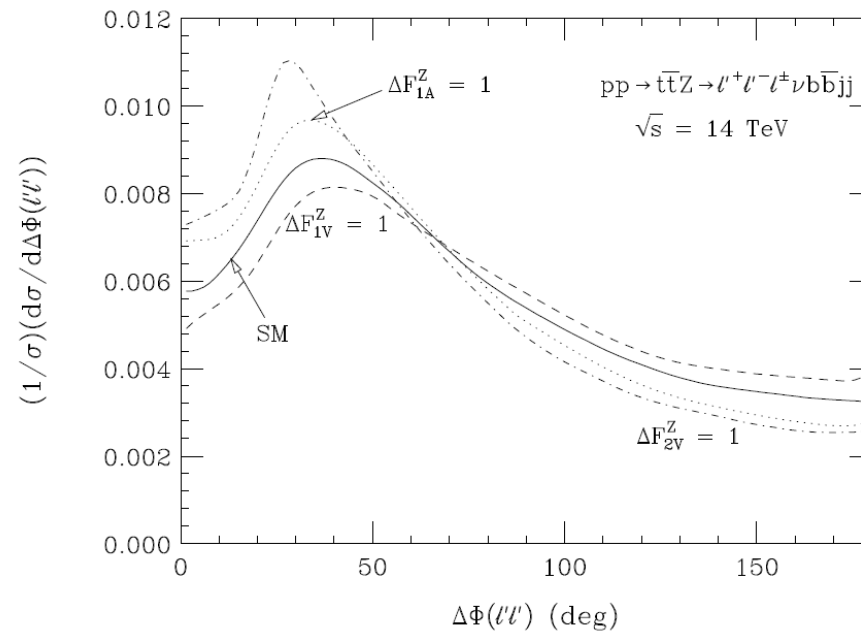
modifies strength
of SM couplings •

introduces anomalous •
weak *magnetic* and *electric*
dipole moments

$t\bar{t} + Z$

$$\Gamma_{\mu}^{ttV}(k^2, q, \bar{q}) = -ie \left\{ \gamma_{\mu} (F_{1V}^V(k^2) + \gamma_5 F_{1A}^V(k^2)) + \frac{\sigma_{\mu\nu}}{2m_t} (q + \bar{q})^{\nu} (iF_{2V}^V(k^2) + \gamma_5 F_{2A}^V(k^2)) \right\}$$

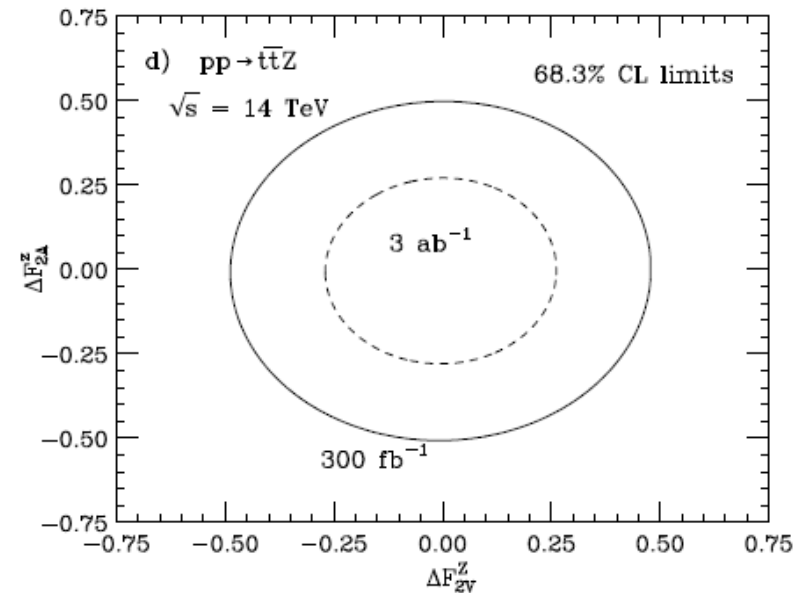
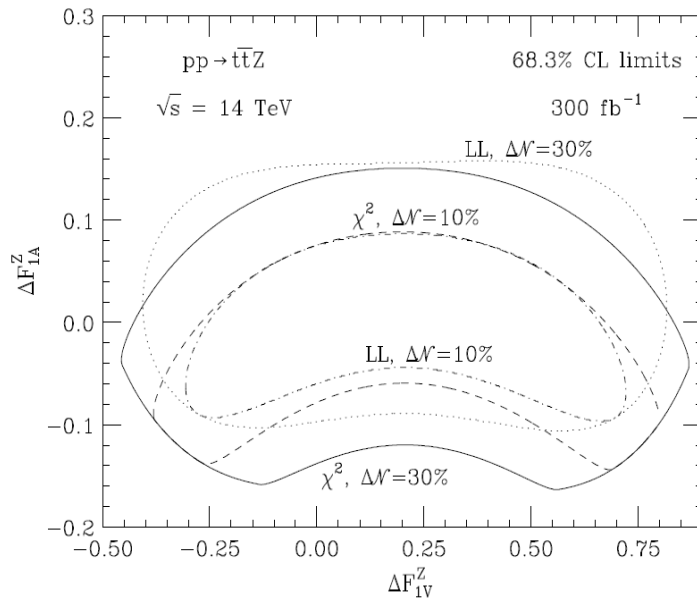
[Baur,Juste,Orr,Rainwater] (2004)



LO coupling constraints

$$\Gamma_{\mu}^{ttV}(k^2, q, \bar{q}) = -ie \left\{ \gamma_{\mu} (F_{1V}^V(k^2) + \gamma_5 F_{1A}^V(k^2)) + \frac{\sigma_{\mu\nu}}{2m_t} (q + \bar{q})^{\nu} (iF_{2V}^V(k^2) + \gamma_5 F_{2A}^V(k^2)) \right\}$$

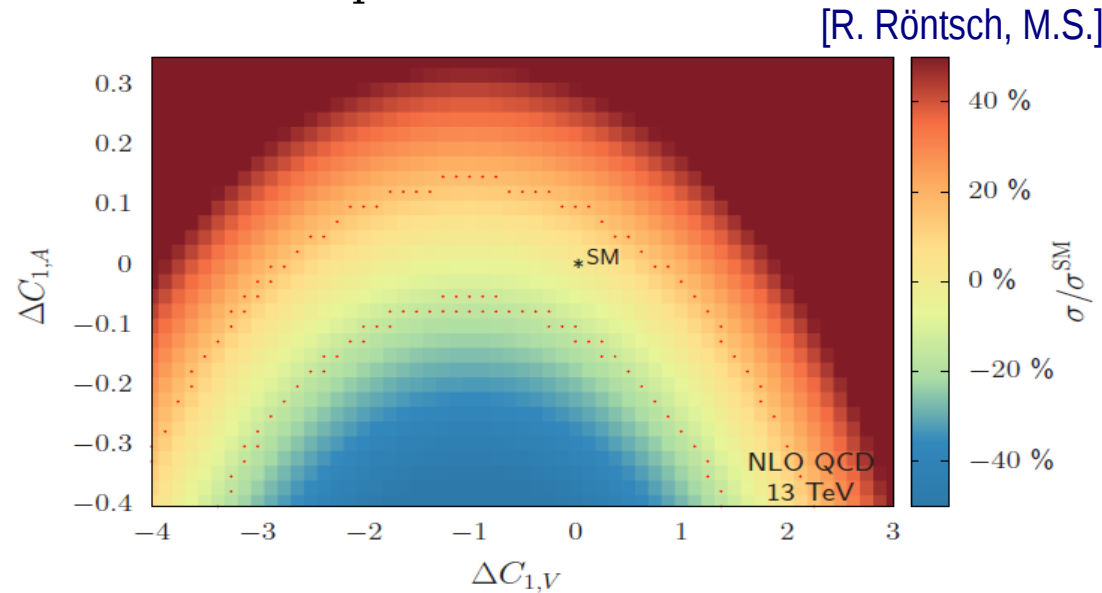
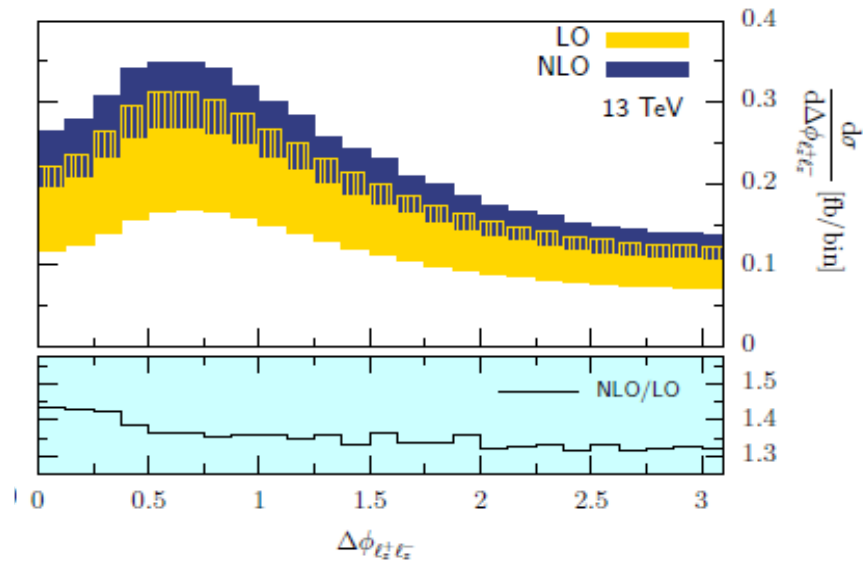
[Baur,Juste,Orr,Rainwater] (2004)



→ “Residual scale uncertainty is biggest limiting factor in these studies.”

NLO coupling constraints

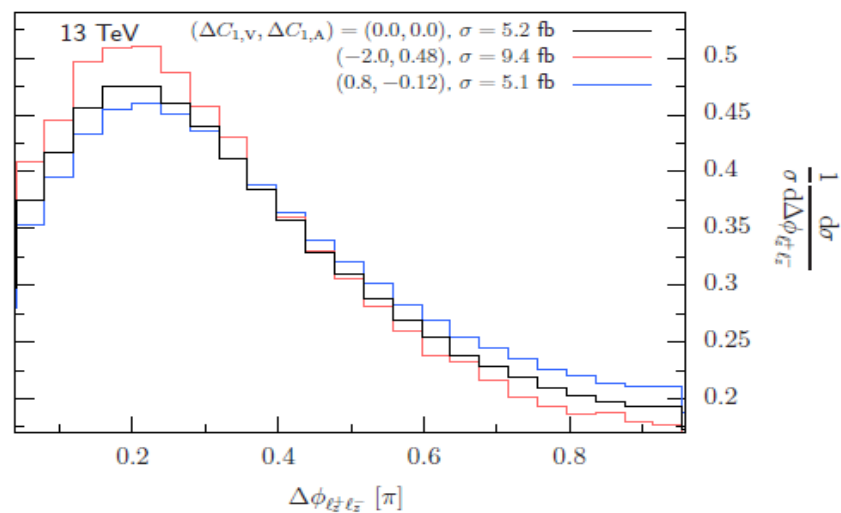
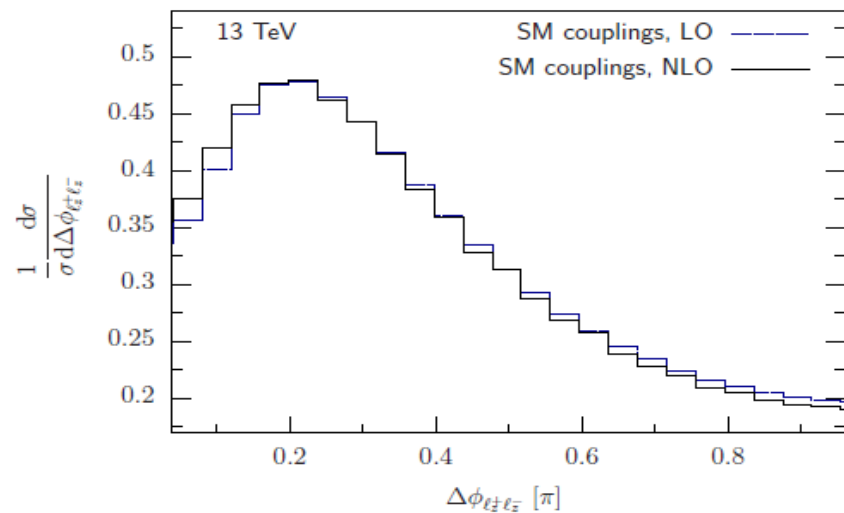
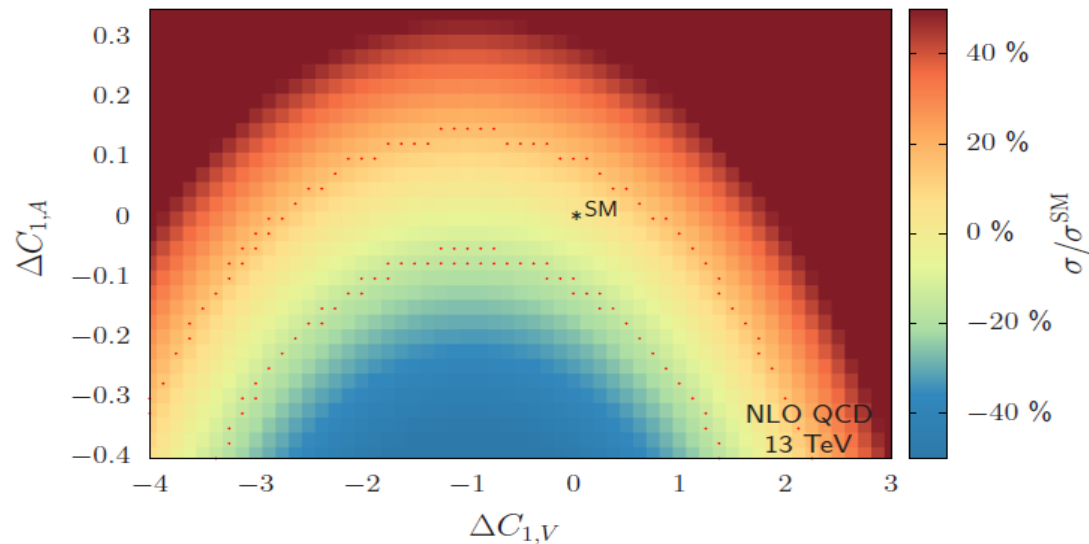
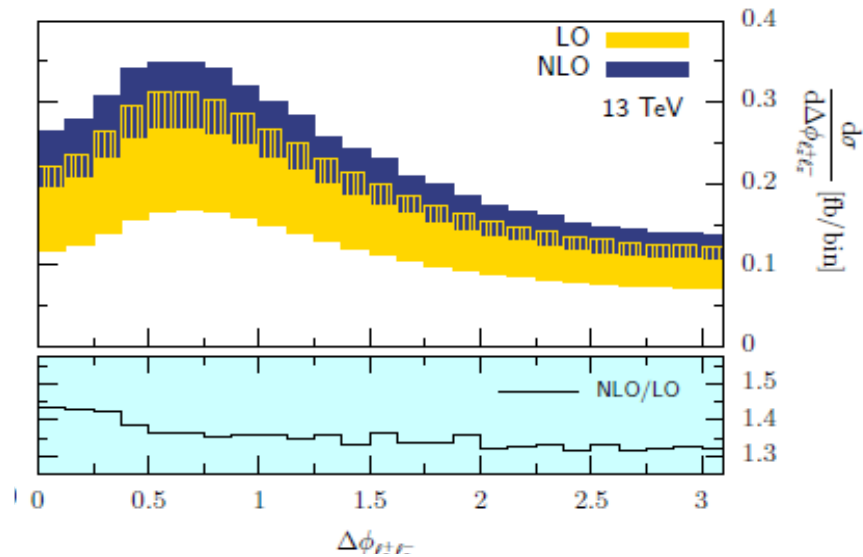
- Realistic final state with top and Z decays (here: 3 leptons + 4 jets)
- NLO QCD corrections in top quark production and decay with anomalous couplings
- Fully differential predictions with CMS and ATLAS acceptance cuts



NLO coupling constraints

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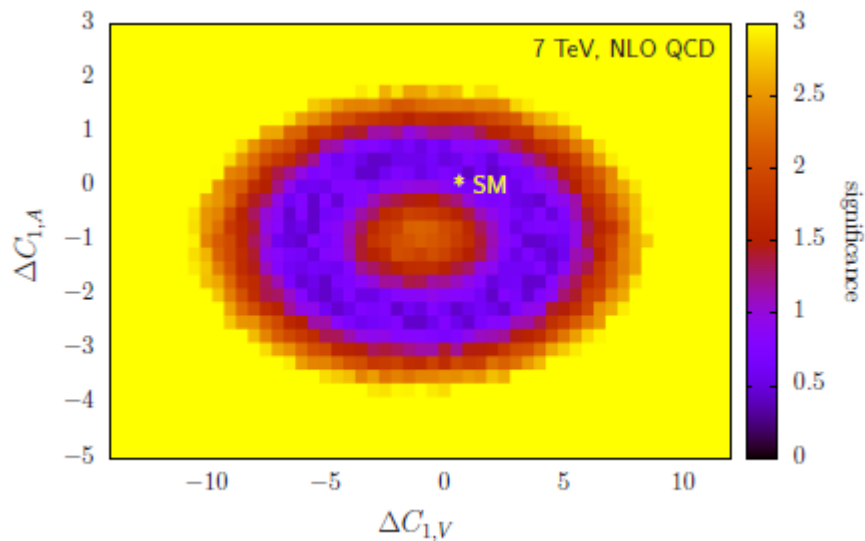
[R. Röntsch, M.S.]



NLO coupling constraints

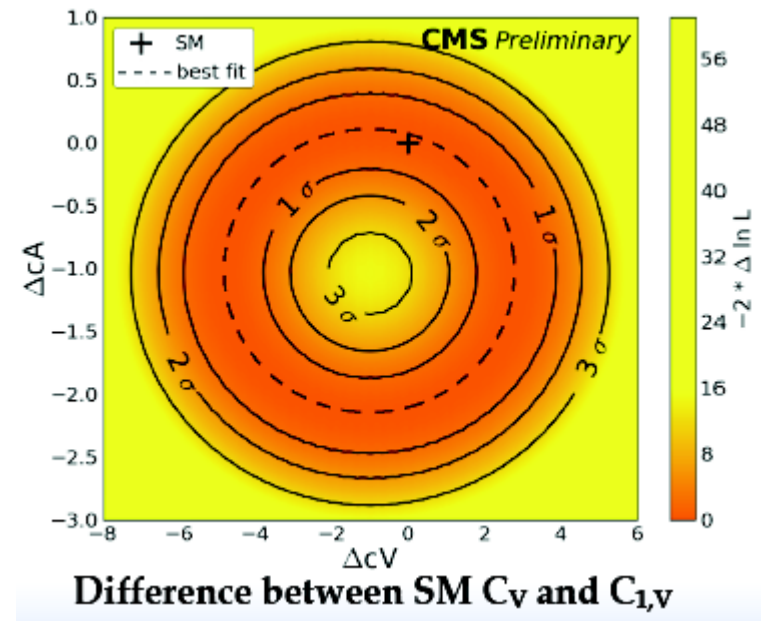
- Realistic final state with top and Z decays (here: 3 leptons + 4 jets)
- NLO QCD corrections in top quark production and decay with anomalous couplings
- Fully differential predictions with CMS and ATLAS acceptance cuts
- log-Likelihood ratio test to calculate exclusion of non-SM hypotheses (incl. uncertainties)

Our estimates for 7 TeV



[R. Röntsch, M.S.]

Application in CMS at 8 TeV

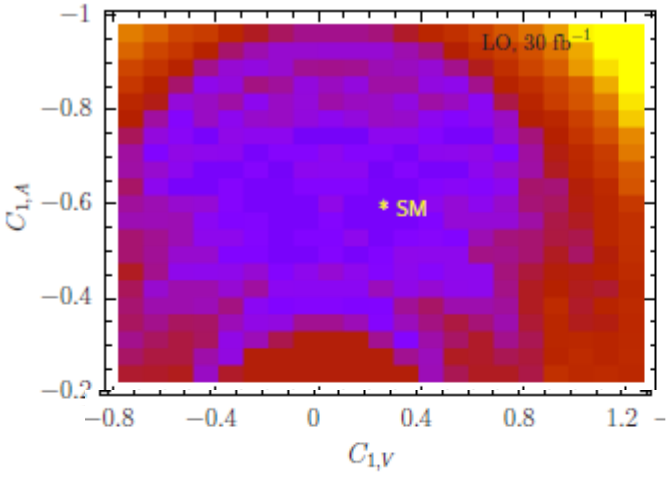


- Our estimates are compatible with experimental findings in [\[CMS \(no PAS\) TOP-14-021\]](#)

(no shape information used)

Constraints for LHC run-2

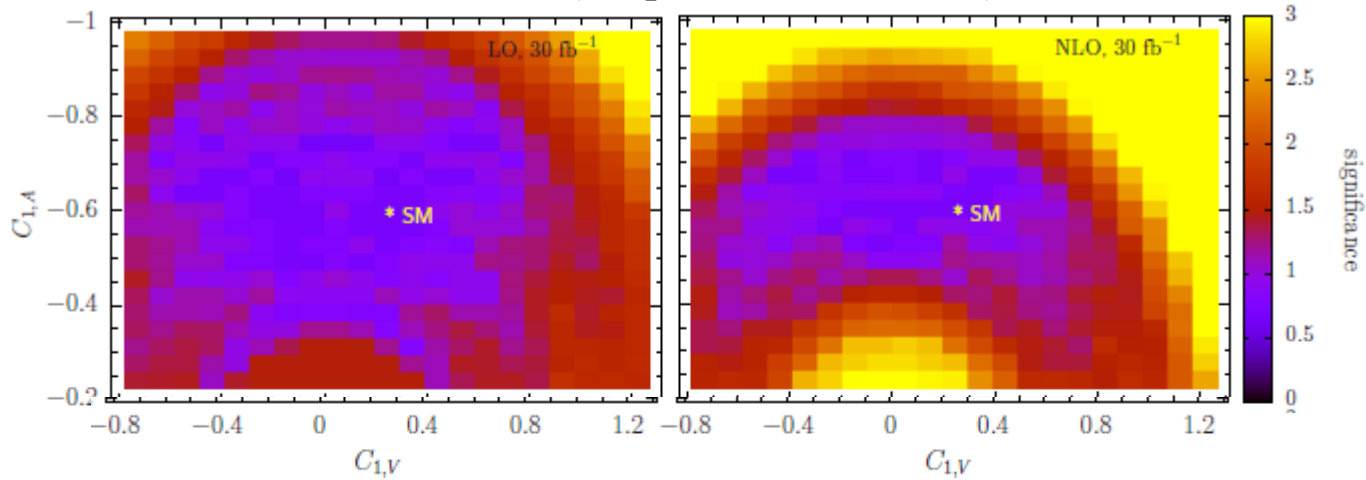
LHC 13 TeV (shape+normalization)



LO 30 fb⁻¹

Constraints for LHC run-2

LHC 13 TeV (shape+normalization)



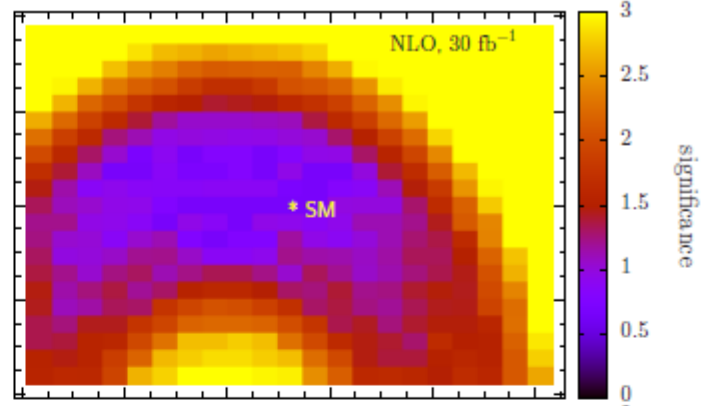
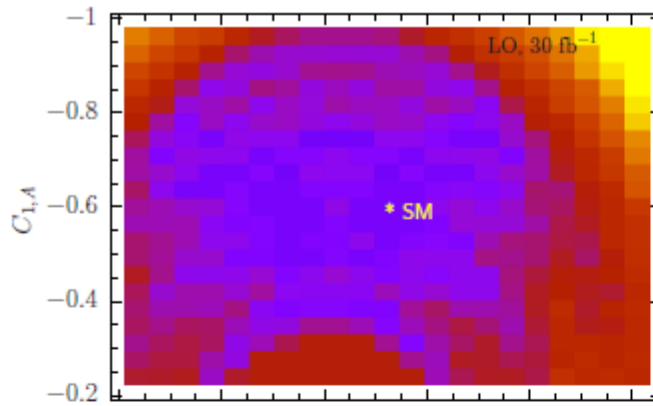
LO 30 fb⁻¹

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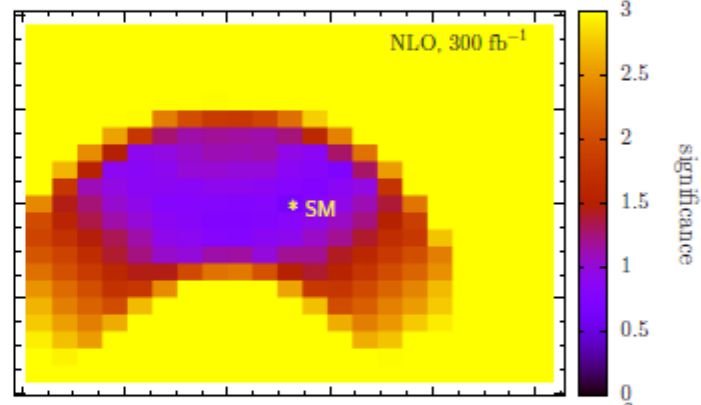
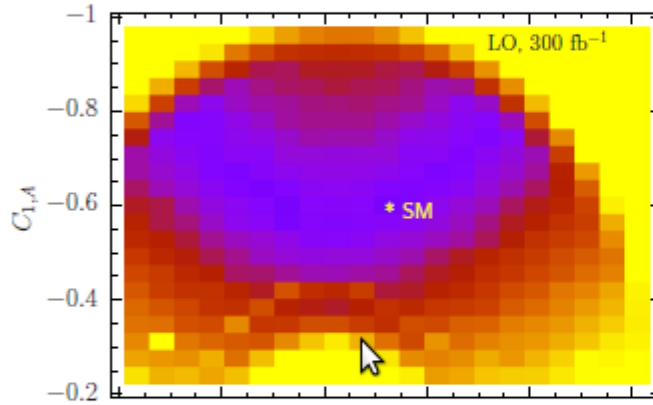
LHC 13 TeV (shape+normalization)

LO 30 fb⁻¹



NLO 30 fb⁻¹

LO 300 fb⁻¹

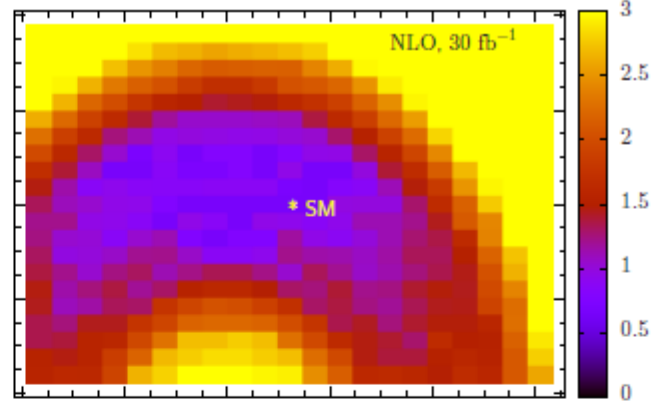
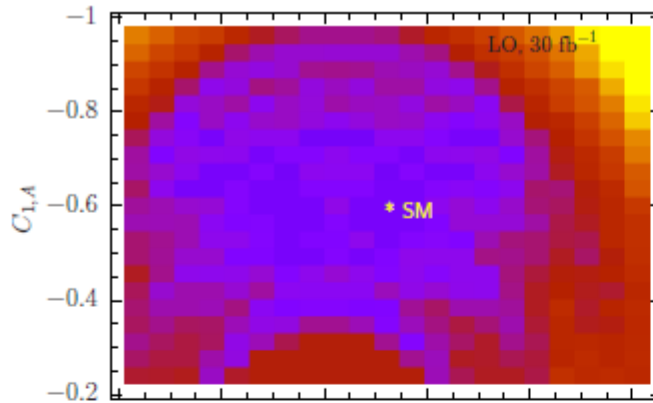


NLO 300 fb⁻¹

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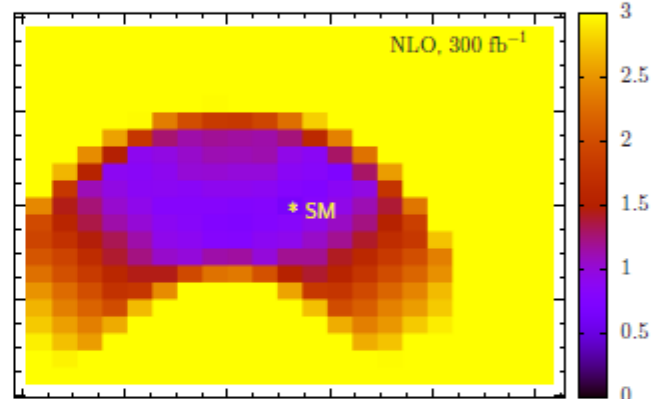
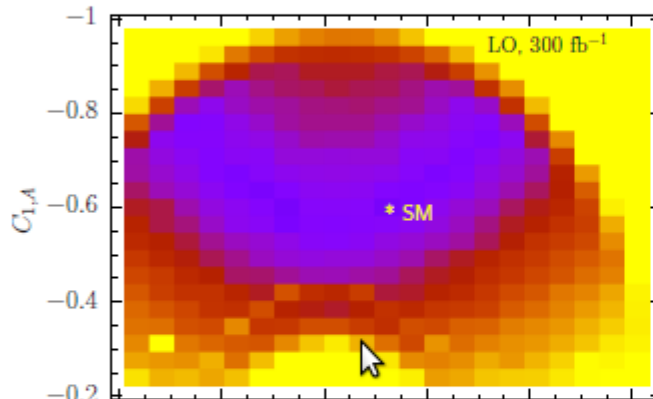
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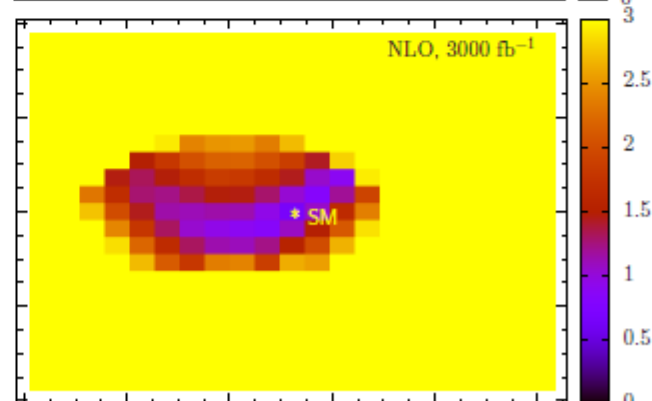
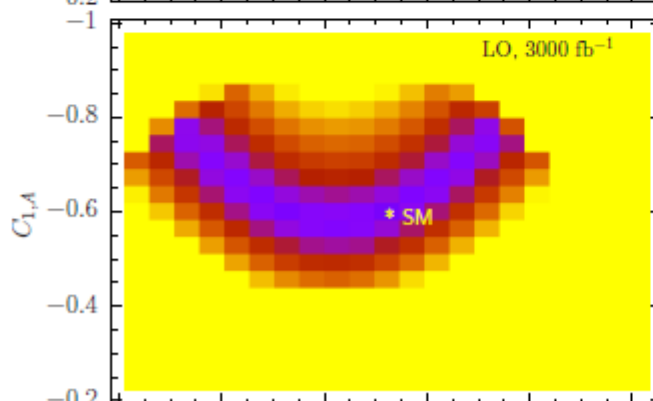
NLO 30 fb⁻¹

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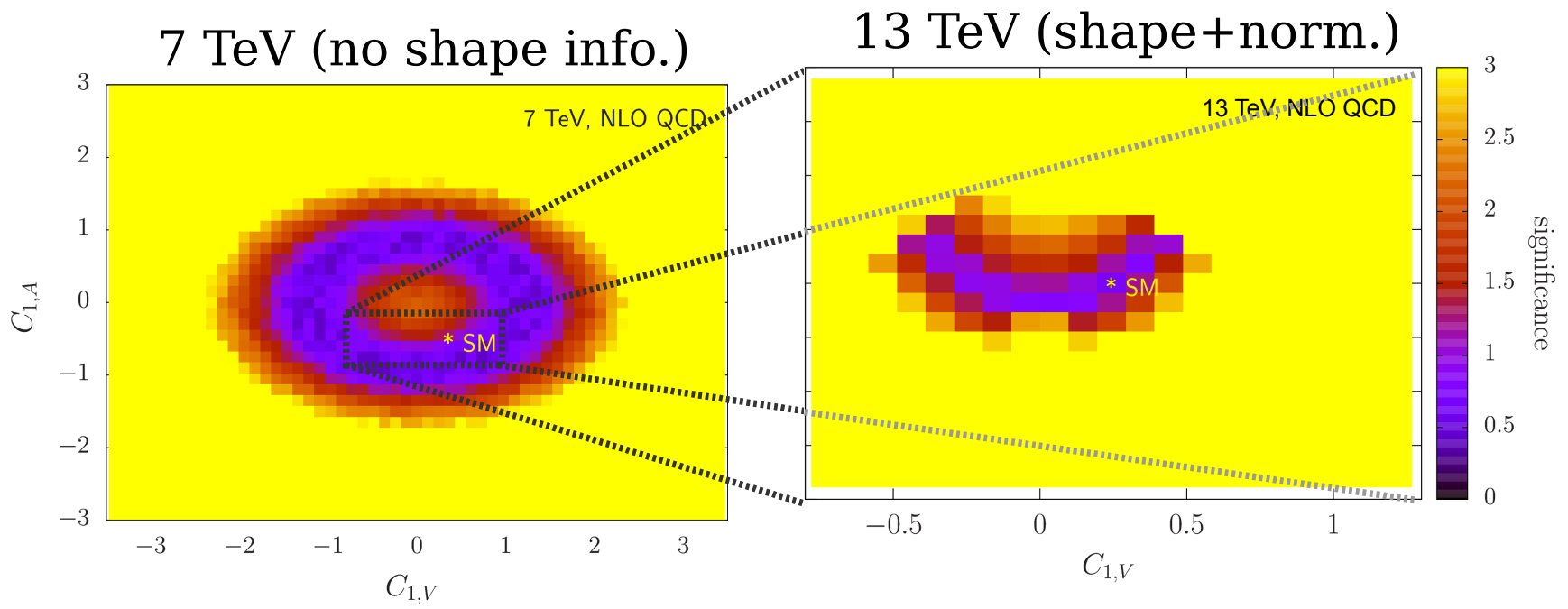


NLO 300 fb⁻¹

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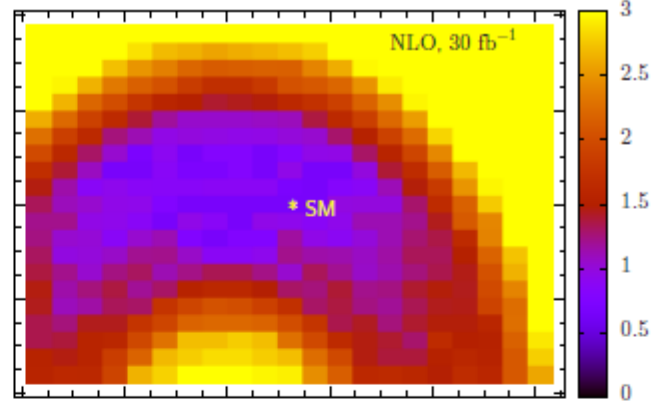
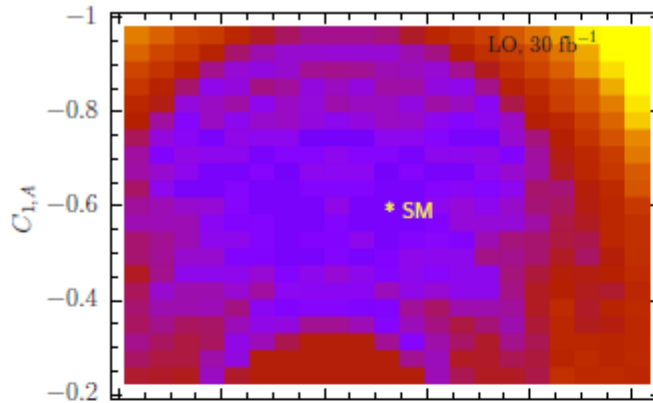
NLO 3000 fb⁻¹



Constraints for LHC run-2

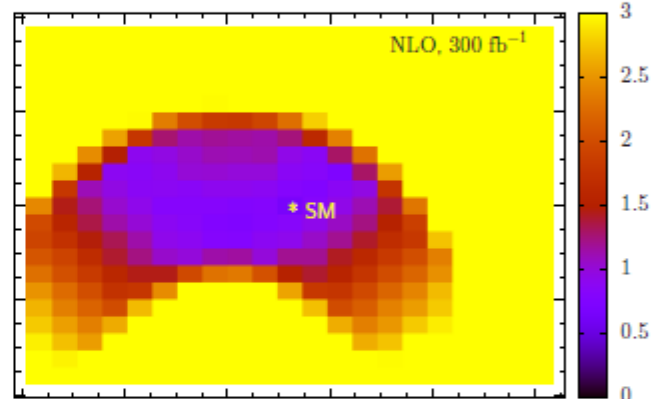
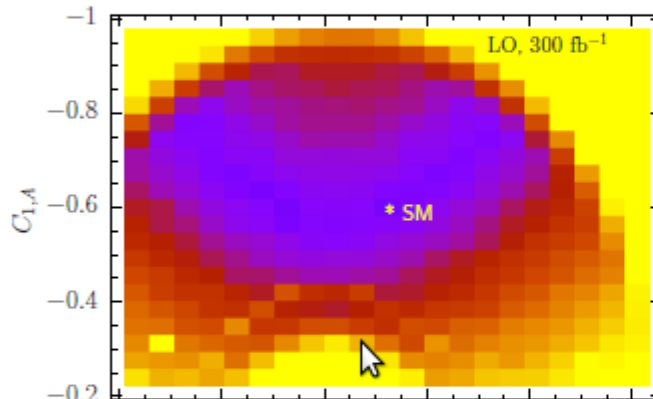
LHC 13 TeV (shape+normalization)

LO 30 fb⁻¹



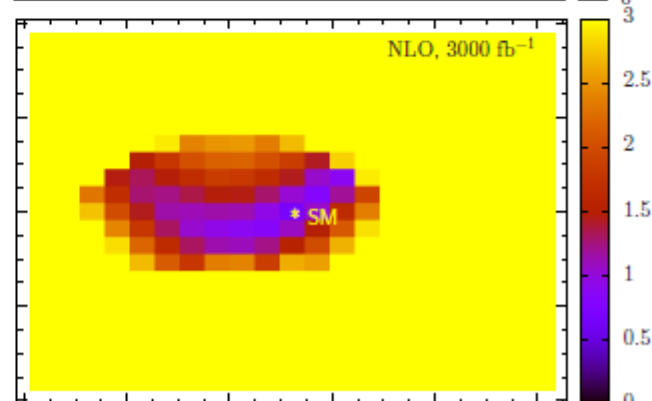
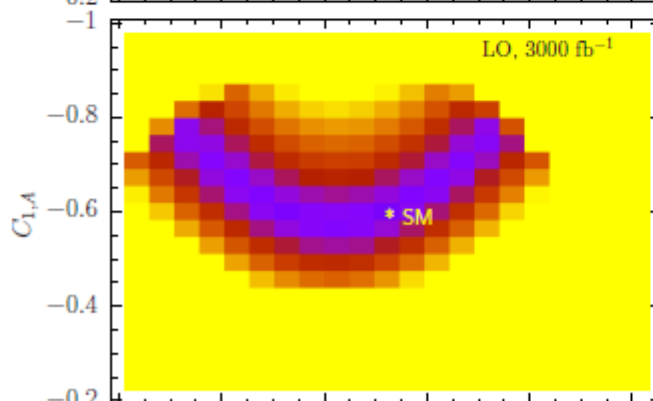
NLO 30 fb⁻¹

LO 300 fb⁻¹



NLO 300 fb⁻¹

LO 3000 fb⁻¹

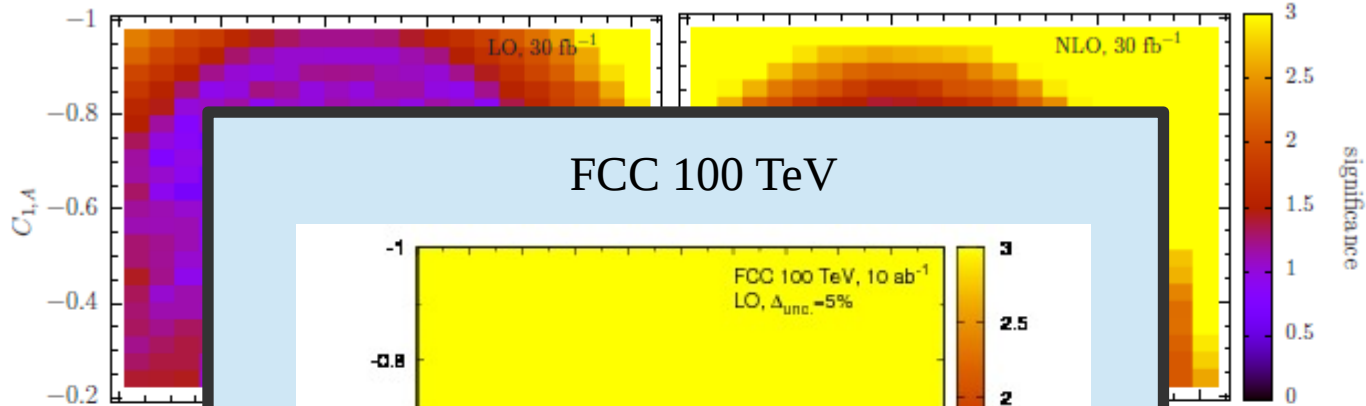


NLO 3000 fb⁻¹

Constraints for LHC run-2

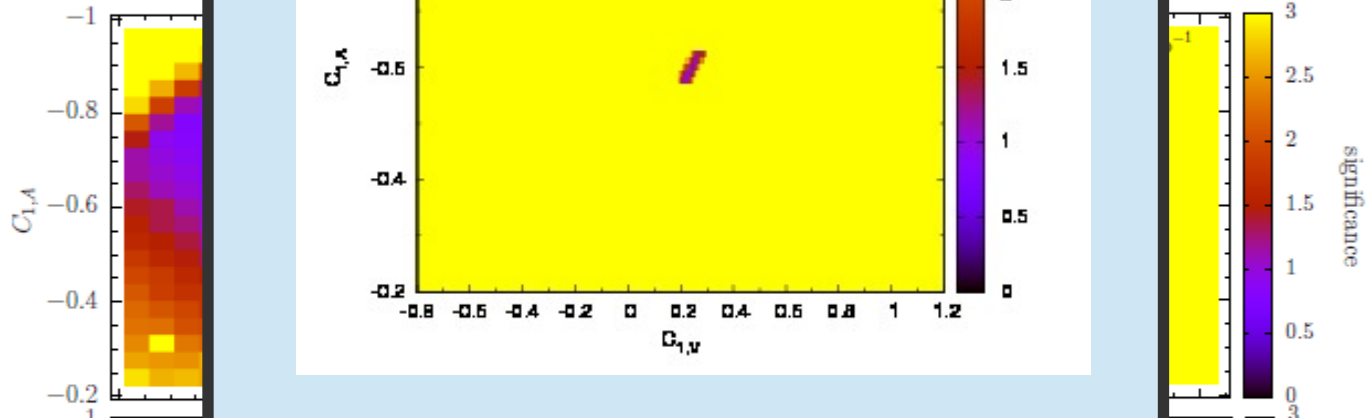
LHC 13 TeV (shape+normalization)

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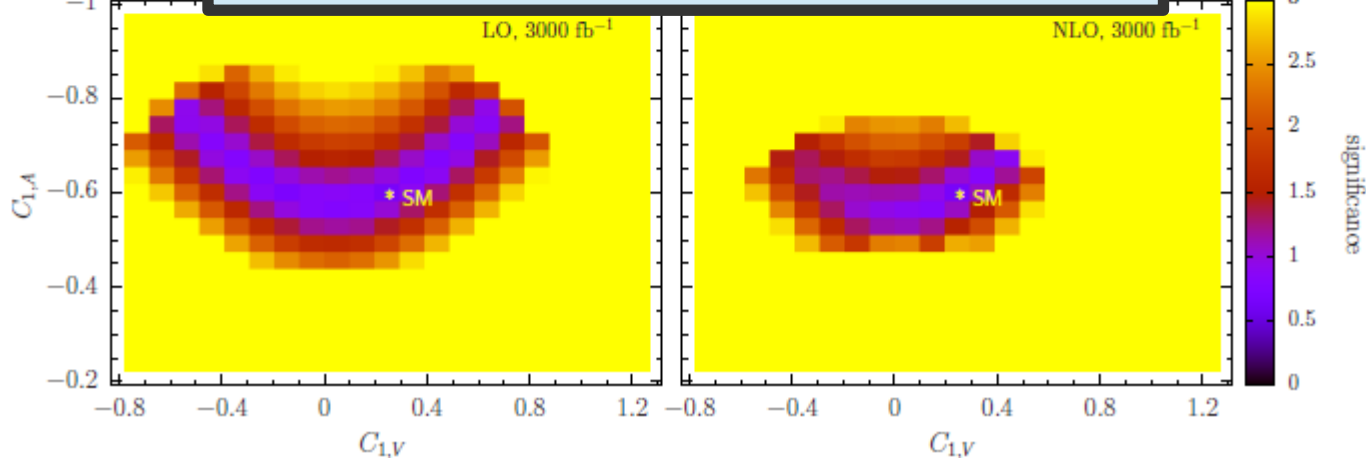
NLO 30 fb⁻¹

LO 300 fb⁻¹

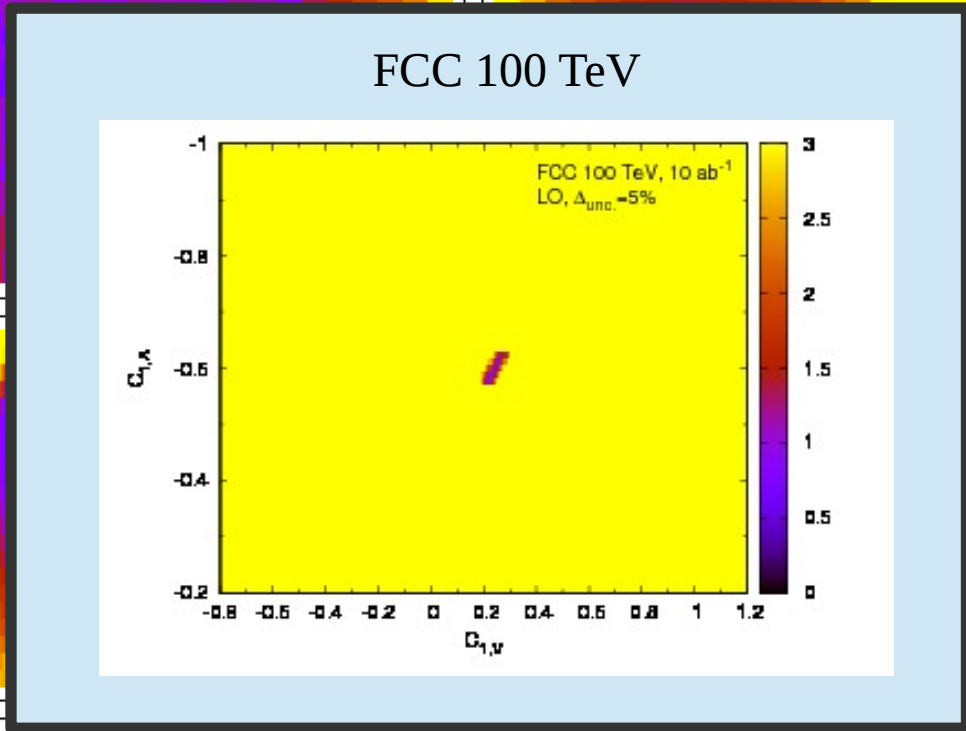


NLO 300 fb⁻¹

LO 3000 fb⁻¹



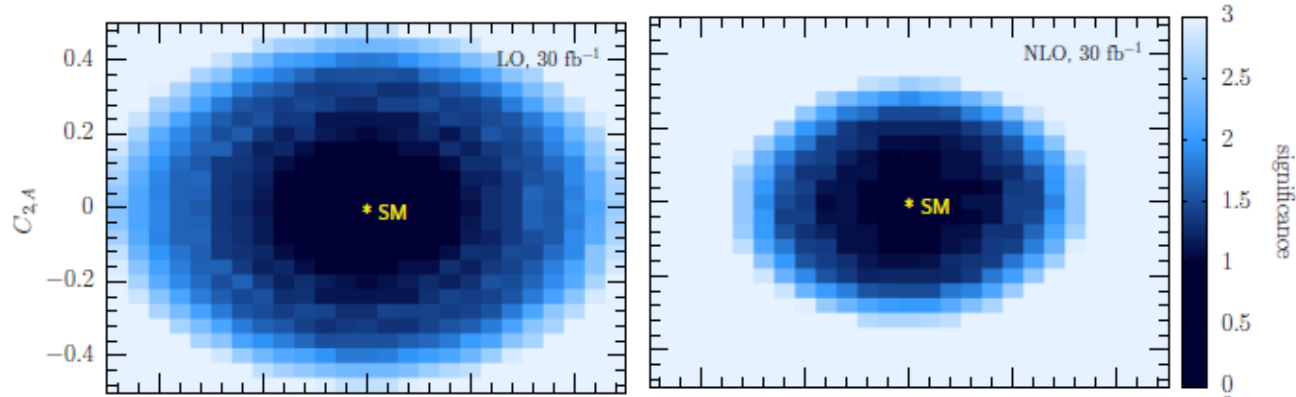
NLO 3000 fb⁻¹



Constraints for LHC run-2

Weak dipole moments (\sim zero in the SM)

LO 30 fb^{-1}

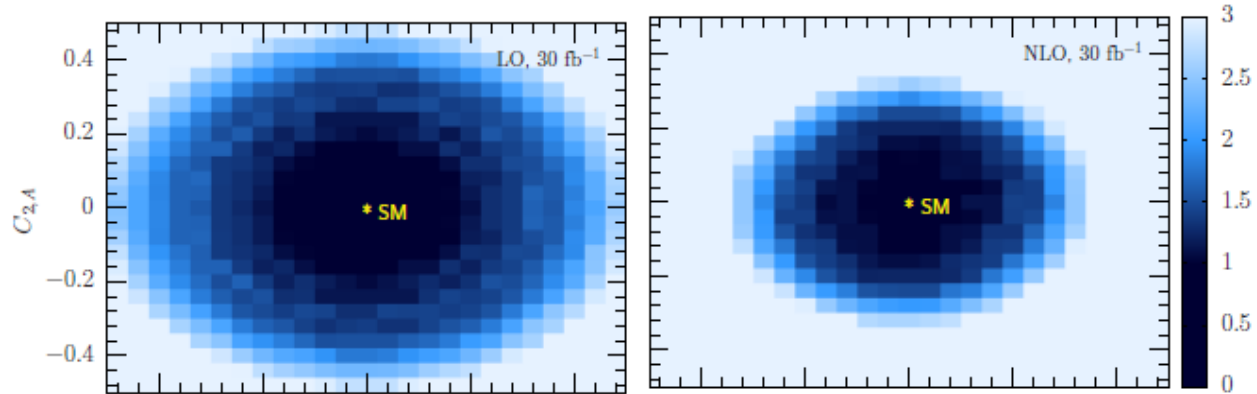


NLO 30 fb^{-1}

Constraints for LHC run-2

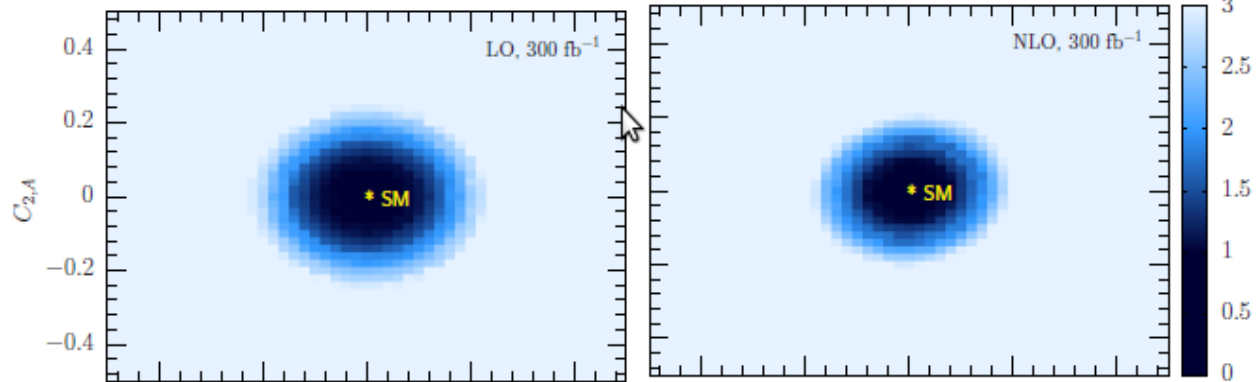
Weak dipole moments (\sim zero in the SM)

LO 30 fb⁻¹



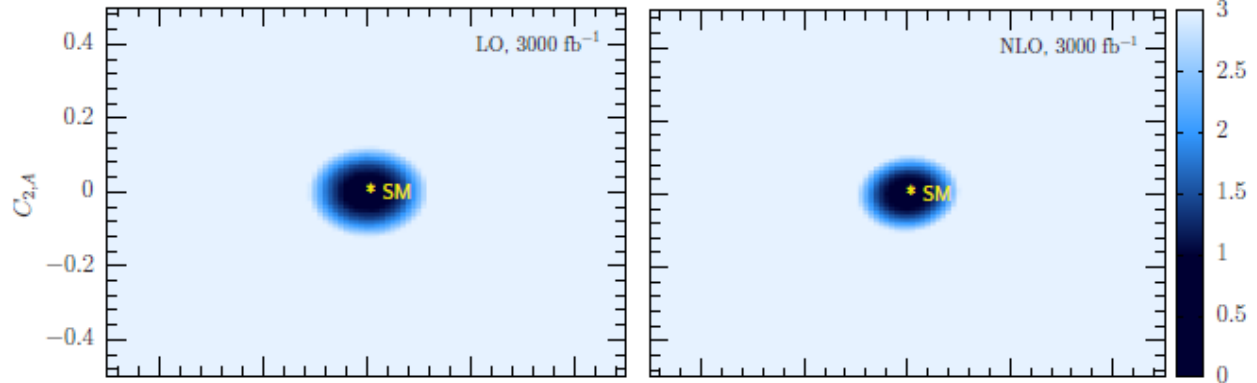
NLO 30 fb⁻¹

LO 300 fb⁻¹



NLO 300 fb⁻¹

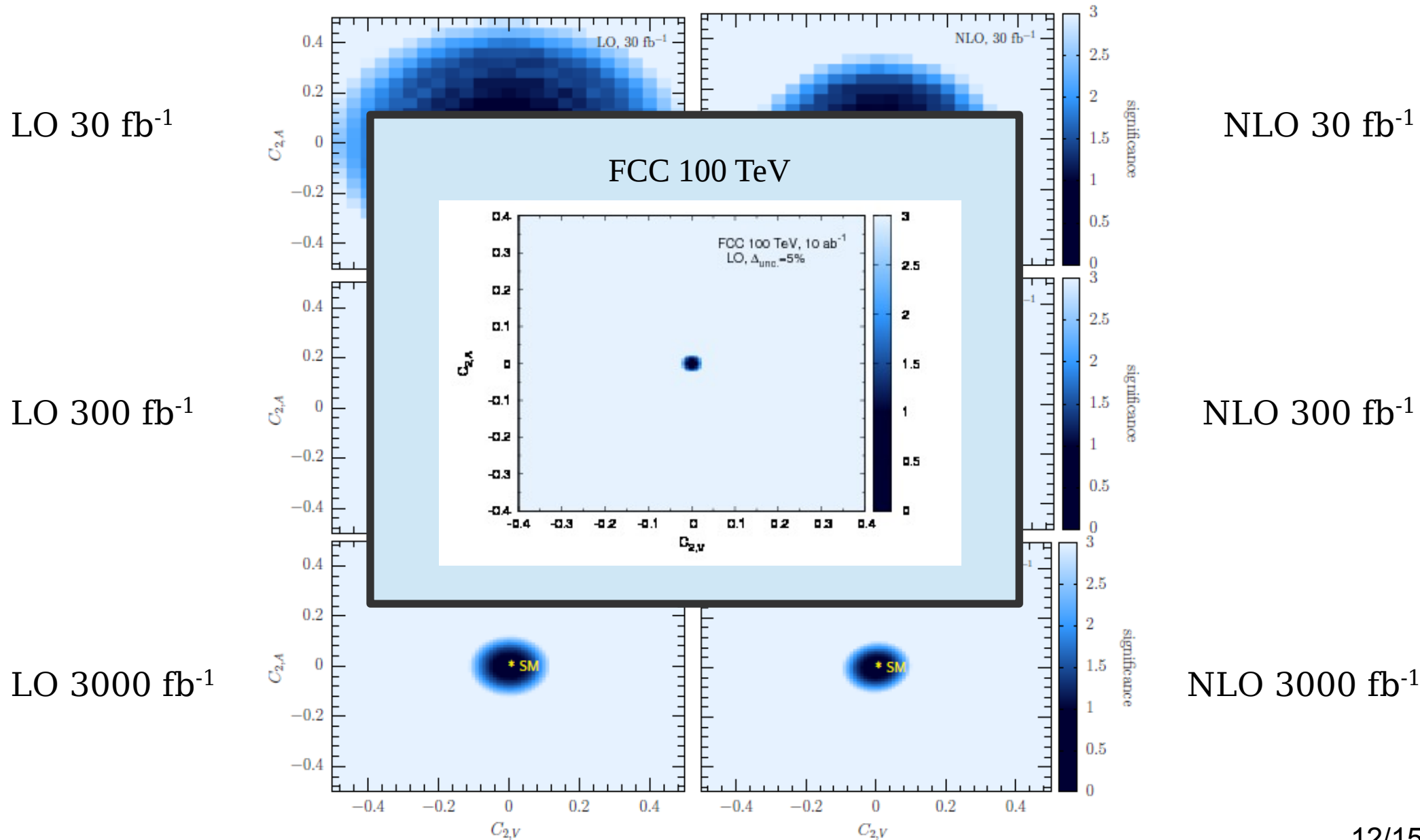
LO 3000 fb⁻¹



NLO 3000 fb⁻¹

Constraints for LHC run-2

Weak dipole moments (\sim zero in the SM)



Correlations with other processes

- Constraints can be further strengthened by combining results from $t\bar{t}b$, $t\bar{t}b+Z$, $t\bar{t}b+\gamma$, single-top.

$$\mathcal{L}_{Wtb} = -\frac{g}{\sqrt{2}}\bar{b}\gamma^\mu(V_L P_L + V_R P_R)t W_\mu^- - \frac{g}{\sqrt{2}}\bar{b}\frac{i\sigma^{\mu\nu}q_\nu}{M_W}(g_L P_L + g_R P_R)t W_\mu^- + \text{H.c.}$$

$$\mathcal{L}_{Ztt} = -\frac{g}{2c_W}\bar{t}\gamma^\mu(X_{tt}^L P_L + X_{tt}^R P_R - 2s_W^2 Q_t)t Z_\mu - \frac{g}{2c_W}\bar{t}\frac{i\sigma^{\mu\nu}q_\nu}{M_Z}(d_V^Z + id_A^Z\gamma_5)t Z_\mu,$$

$$\mathcal{L}_{\gamma tt} = -eQ_t\bar{t}\gamma^\mu t A_\mu - e\bar{t}\frac{i\sigma^{\mu\nu}q_\nu}{m_t}(d_V^\gamma + id_A^\gamma\gamma_5)t A_\mu.$$

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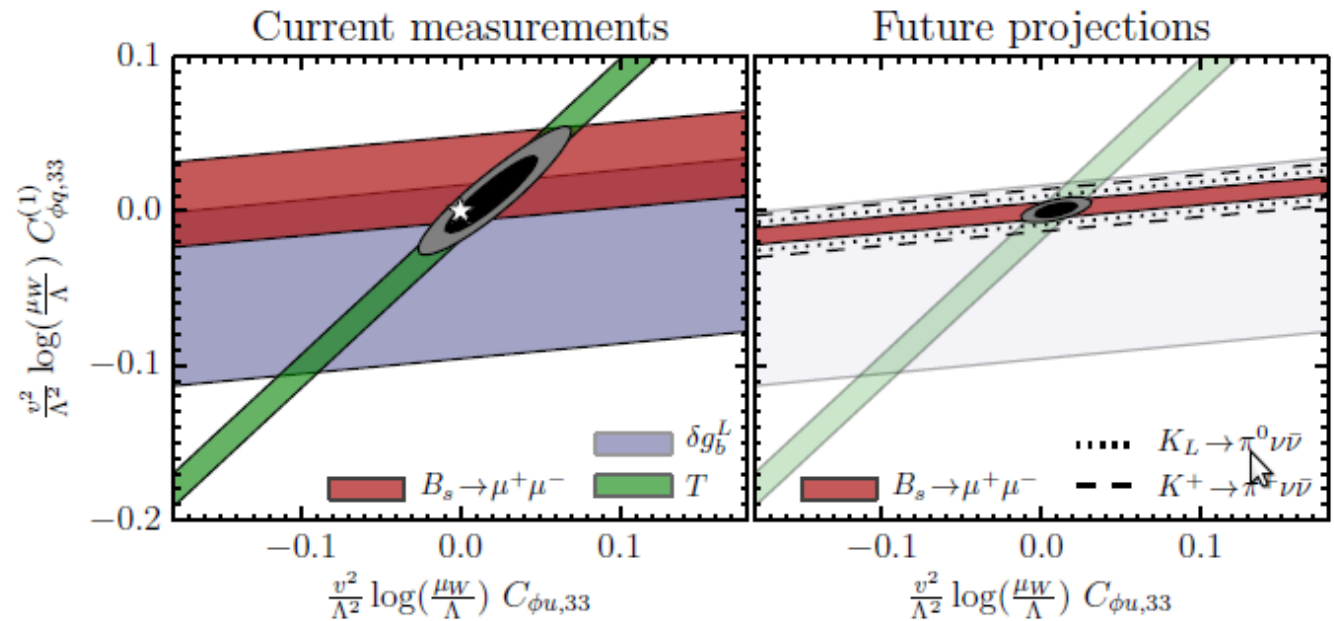
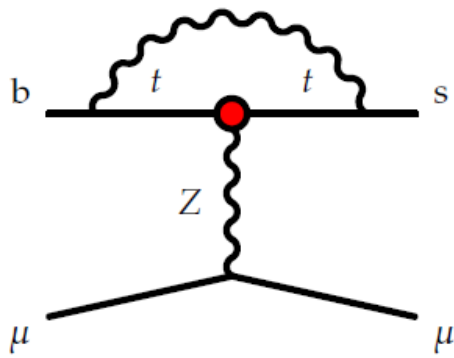
$$\mathcal{L}_{\gamma tt} = -eQ_t\bar{t}\gamma^\mu t A_\mu - e\bar{t} \frac{i\sigma^{\mu\nu} q_\nu}{m_t} (d_V^\gamma + id_A^\gamma \gamma_5) t A_\mu.$$

“A minimal set of top anomalous couplings” [J.A. Aguilar-Saavedra]

Indirect constraints

Probing anomalous top-Z interactions with rare meson decays

[Brod, Greljo, Stamou, Uttayarat]



Order of magnitude stronger bounds than direct reach with 3000/fb!

T	0.08 ± 0.07	[Ciuchini et al., arxiv:1306.4644]
δg_L^b	0.0016 ± 0.0015	[Ciuchini et al., arxiv:1306.4644]
$\text{Br}(B_s \rightarrow \mu^+ \mu^-)$ [CMS]	$(3.0_{-0.9}^{+1.0}) \times 10^{-9}$	[CMS, arxiv:1307.5025]
$\text{Br}(B_s \rightarrow \mu^+ \mu^-)$ [LHCb]	$(2.9_{-1.0}^{+1.1}) \times 10^{-9}$	[LHCb, arxiv:1307.5024]
$\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$	$(1.73_{-1.05}^{+1.15}) \times 10^{-10}$	[E949, arxiv:0808.2459]

Conclusions

- LHC run-2 allows detailed studies of $t\bar{t}b+Z$ final states for the 1st time
- Subsequent coupling constraints yield *direct* sensitivity to physics beyond the SM and complement indirect determinations from low energy experiments
- State-of-the-art NLO QCD predictions for anomalous electroweak top couplings are available and significantly improve coupling constraints
- Combination with other processes ($t\bar{t}b$, $t\bar{t}b+\gamma$, *single-top Z*) will further strengthen the limits

Extras

$$\mathcal{L}_{t\bar{t}Z} = ie\bar{u}(p_t) \left[\gamma^\mu (C_{1,V} + \gamma_5 C_{1,A}) + \frac{i\sigma_{\mu\nu}q_\nu}{M_Z} (C_{2,V} + i\gamma_5 C_{2,A}) \right] v(p_{\bar{t}}) Z_\mu,$$

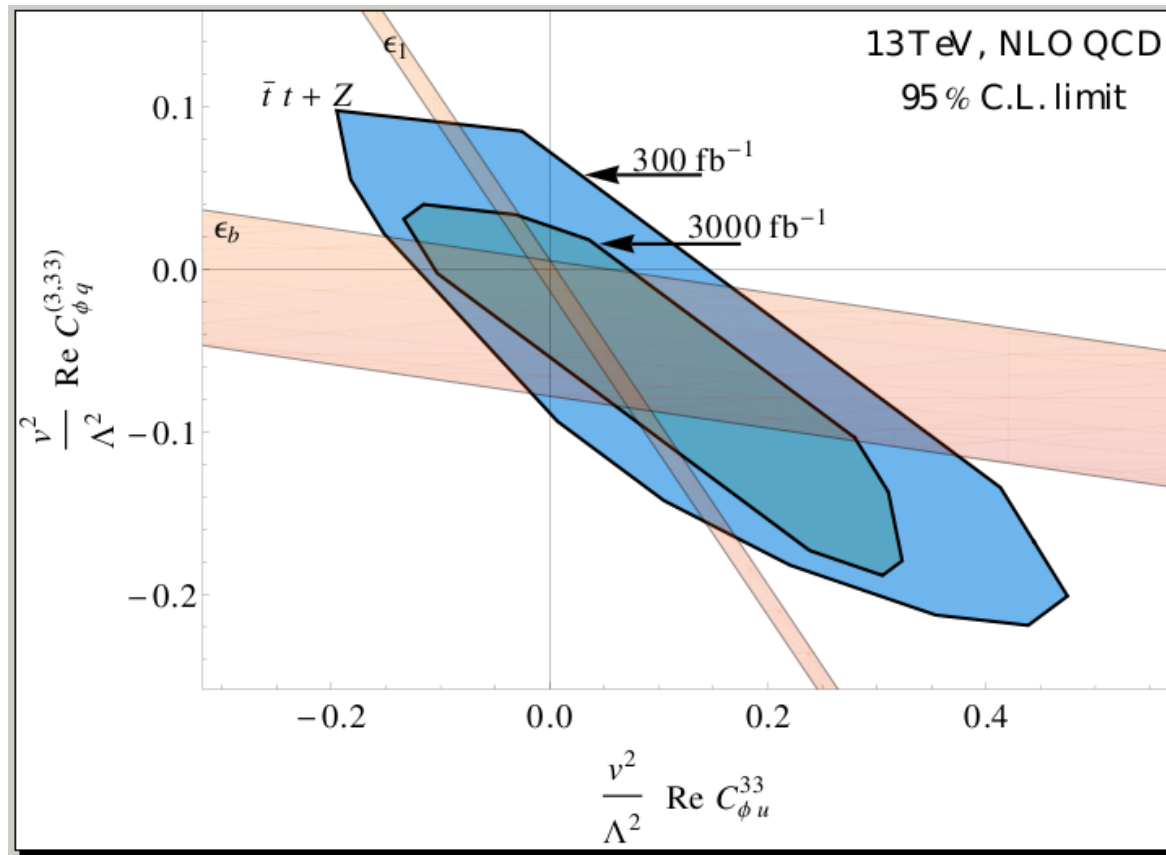
$$C_{1,V} = C_{1,V}^{\text{SM}} + \left(\frac{v^2}{\Lambda^2} \right) \text{Re} \left[C_{\phi q}^{(3,33)} - C_{\phi q}^{(1,33)} - C_{\phi u}^{33} \right],$$

$$C_{1,A} = C_{1,A}^{\text{SM}} + \left(\frac{v^2}{\Lambda^2} \right) \text{Re} \left[C_{\phi q}^{(3,33)} - C_{\phi q}^{(1,33)} + C_{\phi u}^{33} \right],$$

$$C_{\phi q}^{(3,33)} = i(\phi^\dagger \tau^a D_\mu \phi) (\bar{t}_L \gamma^\mu \tau_a t_L),$$

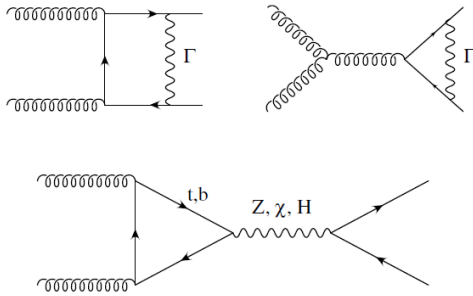
$$C_{\phi q}^{(1,33)} = i(\phi^\dagger D_\mu \phi) (\bar{t}_L \gamma^\mu t_L),$$

$$C_{\phi u}^{33} = i(\phi^\dagger D_\mu \phi) (\bar{t}_R \gamma^\mu t_R).$$



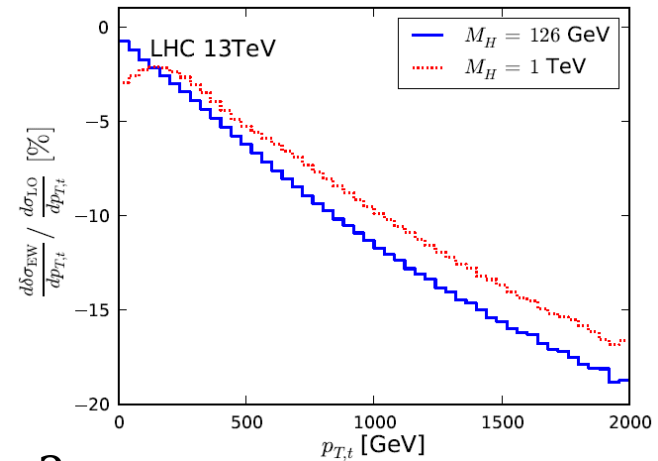
Electroweak corrections

- El.weak corrections to $t\bar{t}$ [Kühn,Scharf,Uwer], [Bernreuther,Fücker,Si]



inclusive (13 TeV): $\sim -3+0.2(\text{real})\% = -2.5\%$ [Baur]

$$\frac{\delta\sigma_{\text{EW}}^{13\text{TeV}}}{\sigma_{\text{LO}}} = (-2.63 + 0.0029g_Y + 0.63g_Y^2)\%$$



Sensitivity to anomalous $top\text{-}Z/W$ couplings?

- El.weak corrections to $t\bar{t}+Z,W,H$ [Frixione,Hirschi,Pagani,Shao,Zaro]

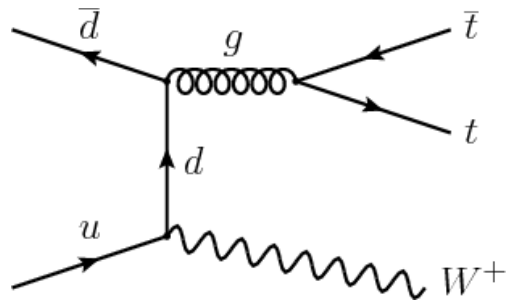
inclusive $t\bar{t}+Z$ (13 TeV): $\sim -3.8+0.9(\text{real})\% = -2.9\%$

inclusive $t\bar{t}+W$ (13 TeV): $\sim -7.7+3.9(\text{real})\% = -3.8\%$

inclusive $t\bar{t}+H$ (13 TeV): $\sim 0.0+0.9(\text{real})\% = +0.9\%$

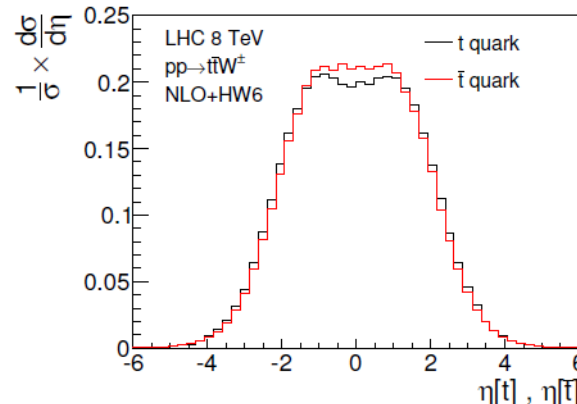
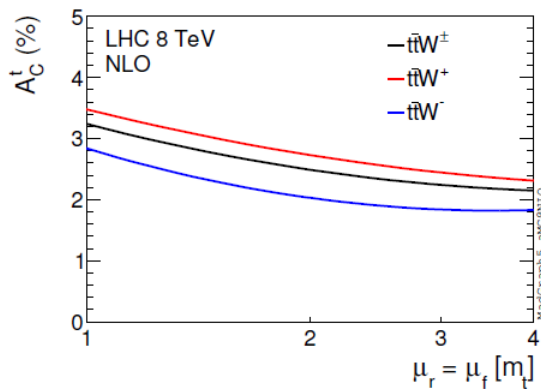
t - b - W Interactions

- $t\bar{t} + W^\pm$ at the 14TeV LHC has a total NLO cross section of 0.7 pb [Campbell, Ellis] [Garzelli, Kardos, Papadopoulos, Trocsanyi]



→ No sensitivity to t - b - W couplings at the LHC

- The absence of a gluon-gluon channels enhances sensitivity to (radiatively generated) asymmetry. [Maltoni, Mangano, Tsinikos, Zaro]
- Moreover, the W -boson polarizes the IS and enhances sensitivity

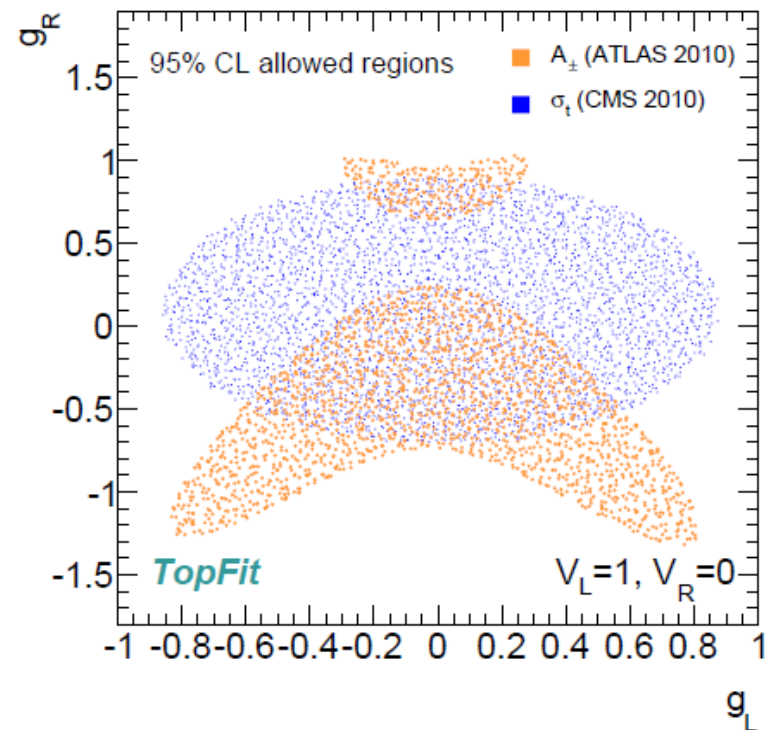
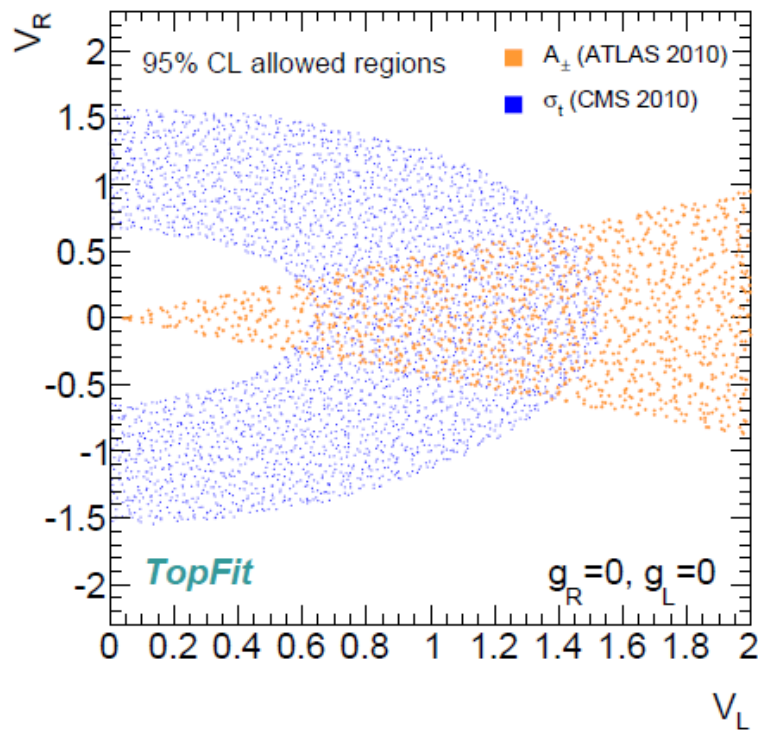


relative precision:

$$\left\{ \begin{array}{l} 45\% \text{ with } 14 \text{ TeV} + 300 \text{ fb}^{-1} \\ 14\% \text{ with } 14 \text{ TeV} + 3000 \text{ fb}^{-1} \\ 3\% \text{ with } 100 \text{ TeV} + 3000 \text{ fb}^{-1} \end{array} \right.$$

t - b - W Interactions

- t - b - W can also be probed in *single-top production* and *top decay in $t\bar{t}$* [> 30 papers]
- one example: Recent combination of both [Aguilar-Saavedra,Castro,Onofre]

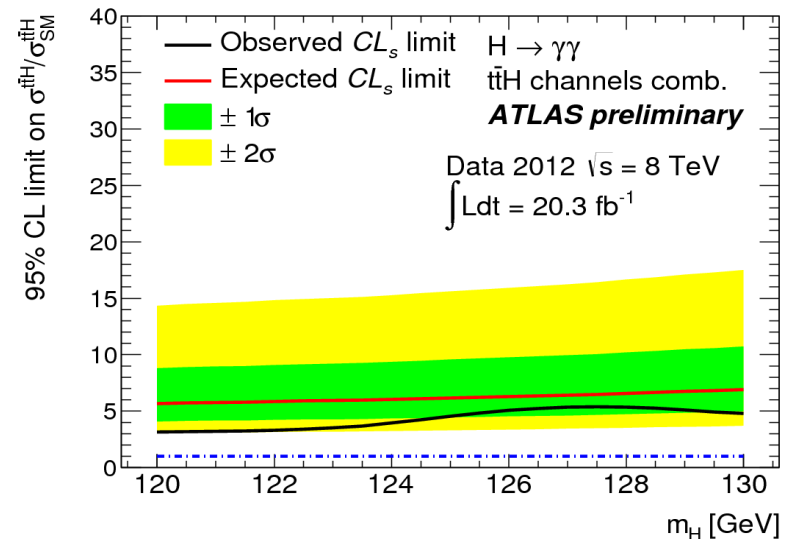
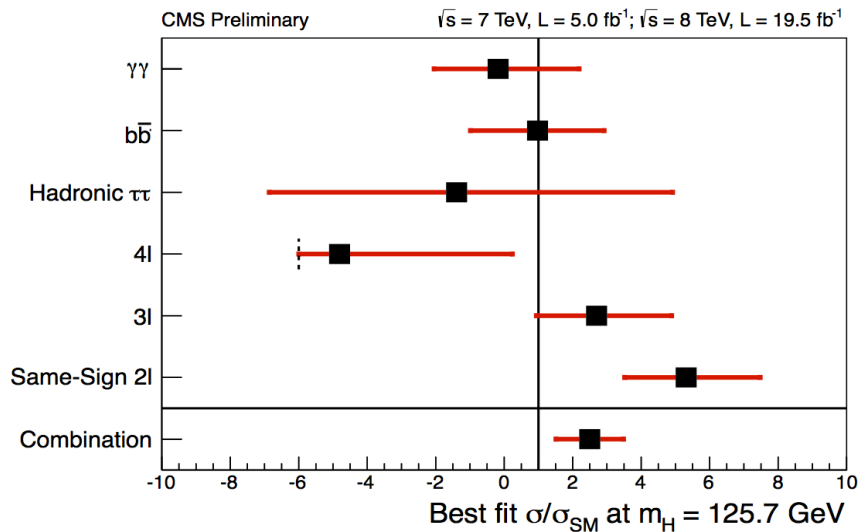


- $t\bar{t}$ process is also ideal for studying chromo-dipole moments

$t\text{-}t\text{-}H$ Interactions

- $t\bar{t} + H$ at the 14TeV LHC has a total NLO cross section of 0.7 pb
 [Beenakker,Dittmaier,Kramer,Plumper,Spira,Zerwas], [Dawson,Jackson,Orr,Reina,Wackerroth]
 [Frederix,Frixione,Hirschi,Maltoni,Pittau,Torielli], [Garzelli,Kardos,Papadopoulos,Trocsanyi]
- Large QCD backgrounds in $H \rightarrow b\bar{b}$. Lots of work on improving S/B.
 E.g. *boosted jets*: [Seymour] [Butterworth et.al], [Plehn,Salam,Spannowsky]
matrix element method: [Artoisenet,Aquino,Maltoni,Mattelaer]
top spin correlations: [Biswas,Frederix,Gabrielli,Mele]

→ No observation yet (7+8 TeV). Upper limit $\sim 3x$ SM cross section.



t - t - H Interactions

- Studies of coupling sensitivity: [Ellis,Hwang,Sakurai,Takeuchi]

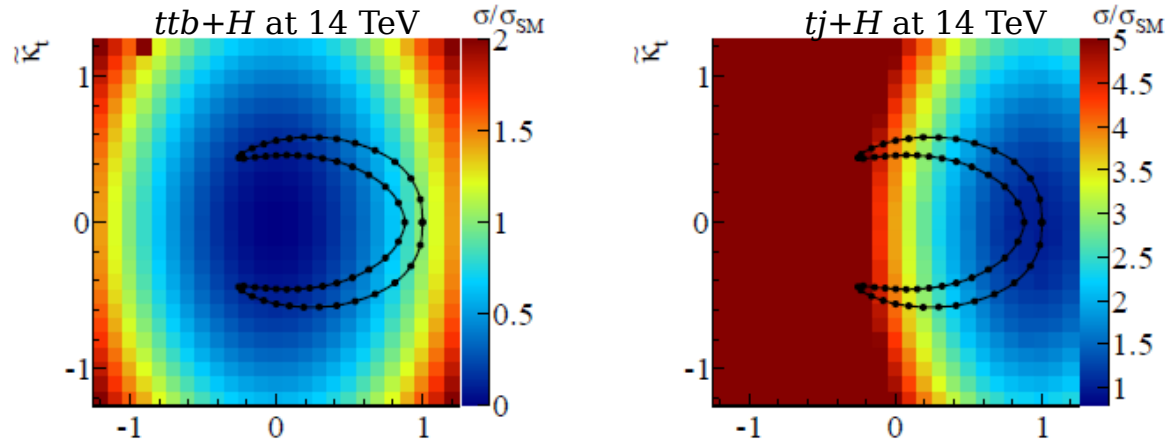
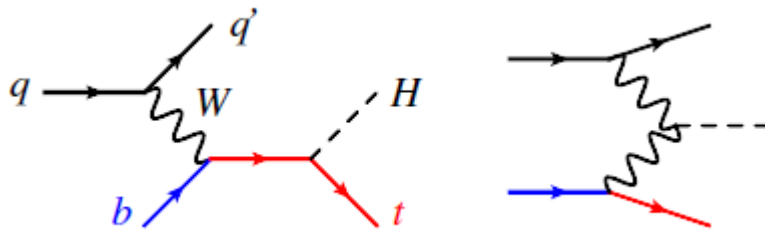
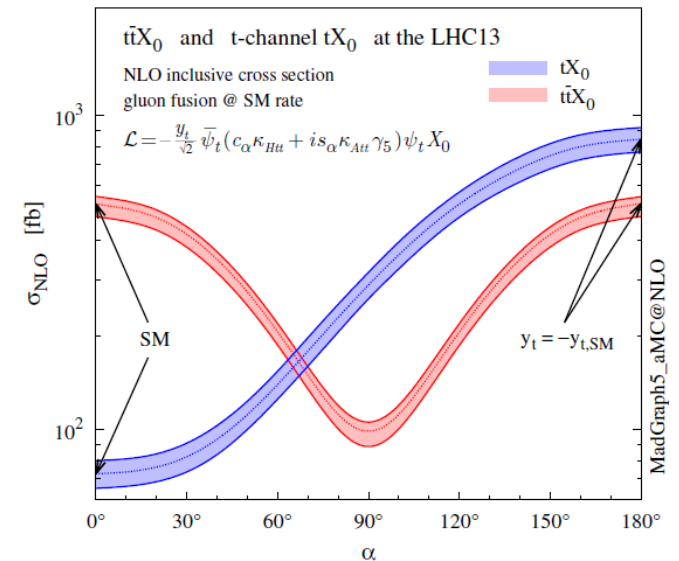


Figure 3: The ratios of $\sigma(\bar{t}tH)$ to the Standard Model value (left panel) and of $\sigma(tH)$ to the Standard Model value (right panel) are shown using the indicated colour codes. Also shown is the crescent-shaped region in Fig. 1 that is allowed by present data at the 68% CL.

- single top + Higgs: [Campbell, Ellis, Röntsch]
[Demartin, Maltoni, Mawatari, Page, Zaro]



destructive interference leads to very small SM cross section

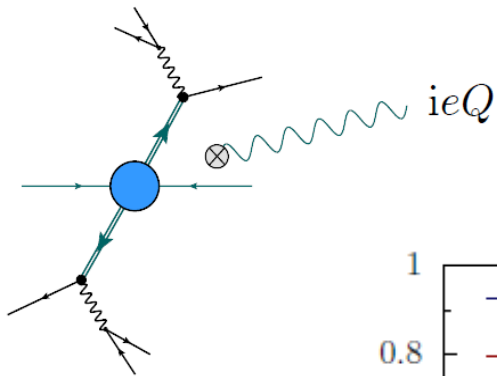


$t\bar{t} + \text{photon}$

- $t\bar{t} + \gamma$ at the 14 TeV LHC has a total NLO cross section of 2.9 pb

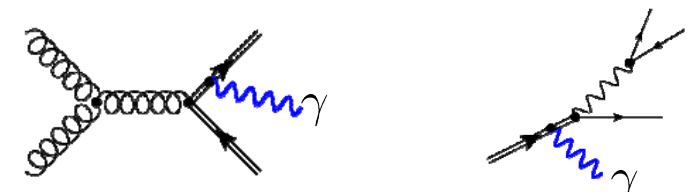
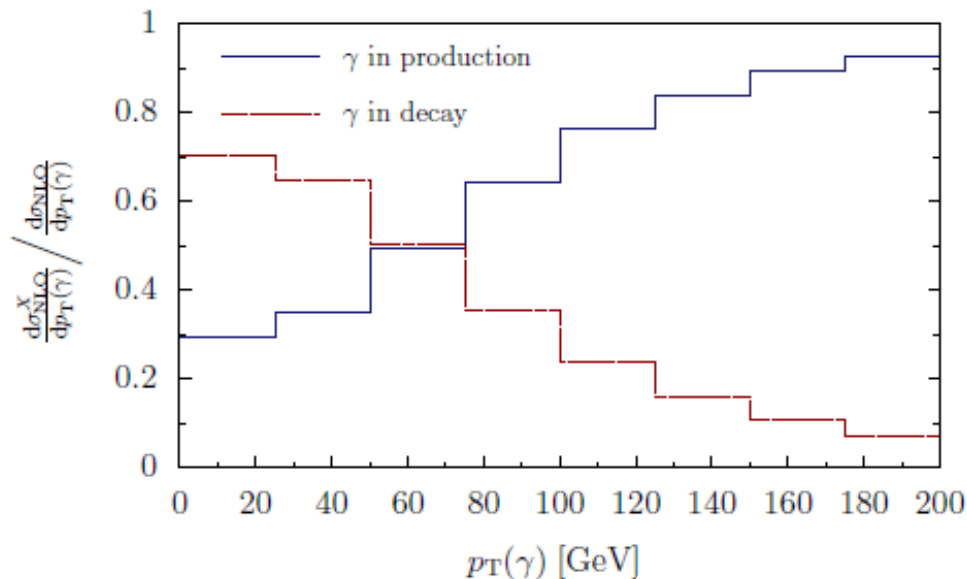
[Duan, Ma, Zhang, Han, Guo, Wang], [Melnikov, Scharf, M.S.], [Kardos, Trocsanyi]

- In contrast to $t\bar{t} + H, Z, W$, photon can be radiated off top decay products



Lepton+Jets Channel + acceptance cuts: $p_T^\gamma > 20 \text{ GeV}$

$$\sigma_{\ell+j}^{\text{NLO}} = 138 \text{ fb} = 61(\text{production}) + 77(\text{decay}) \text{ fb}$$



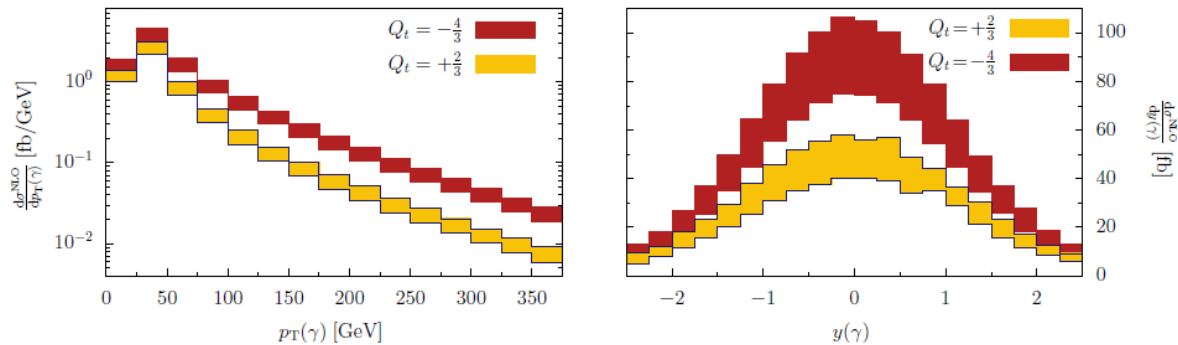
→ Photons with $p_{T\gamma} < 50 \text{ GeV}$ are dominantly emitted in the decay

[Melnikov, Scharf, M.S.]

$t\bar{t} + \text{photon}$

Effects of radiative top decays:

- Simple BSM study: top quark charge measurement in $t\bar{t} + \text{photon}$



$$\sigma_{t\bar{t}\gamma}^{\text{NLO}} = 138 \text{ fb} \xrightarrow{Q_t = \frac{2}{3} \rightarrow -\frac{4}{3}} \sigma_{t\bar{t}\gamma}^{\text{NLO}} = 243 \text{ fb}$$

→ Naive expectation of Q_t^2 scaling fails: $\mathcal{R}^{\text{NLO}} = \frac{\sigma_{\text{NLO}}^{Q_t=-4/3}}{\sigma_{\text{NLO}}^{Q_t=2/3}} = 1.76_{-0.02}^{+0.01}$.

- Implications for FB asymmetry:

$$A_{t,\text{FB}}^{\text{LO}} = -17\%$$

$$A_{t,\text{FB}}^{\text{NLO}} = -12\%$$

$$A_{\ell,\text{FB}}^{\text{NLO}} \approx \text{few}\%$$

stable tops

decayed tops
+ radiative decays

