

Spacetime curvature and Higgs stability during inflation

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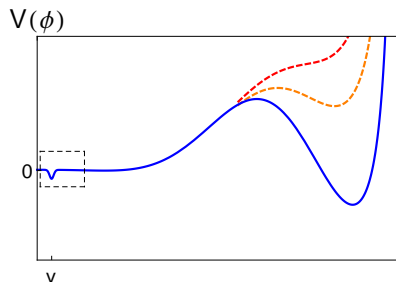
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Particle Physics and Cosmology
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- 1 Introduction
- 2 Higgs stability during inflation (QFT in Minkowski)
- 3 Higgs stability during inflation (QFT in curved space)
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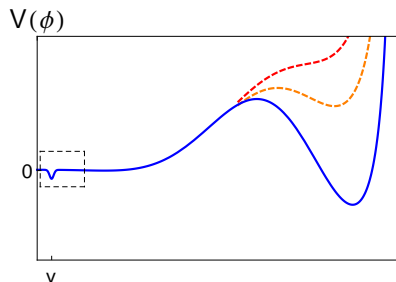
Standard Model Higgs potential

- $V(\phi)$ has a minimum at $\phi = v$
- Very sensitive to M_h and M_t
- A vacuum at $\phi \neq v$ incompatible with observations



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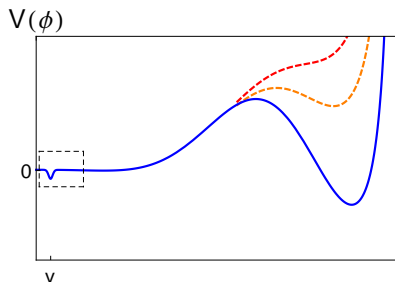
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 - *Meta* stable at 99% CL [1]
 - Lifetime much longer than 13.8×10^9 years
 - Is this also true for the early Universe (*inflation*)?



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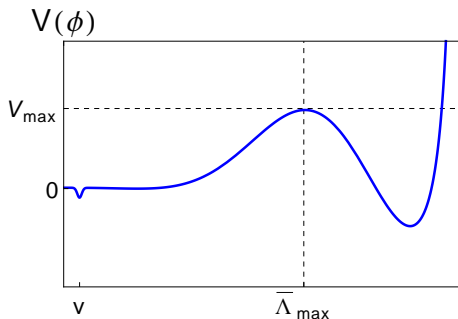


New physics needed to stabilize the vacuum?

[1] Buttazzo et al. (2013); Spencer-Smith (2014)

Inflation and the Standard Model

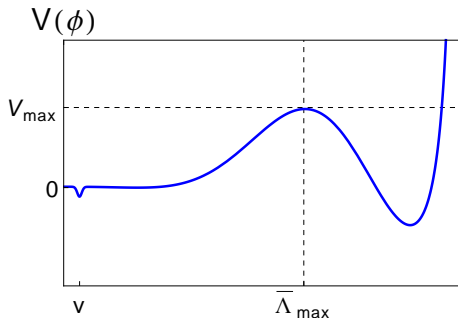
- In principle we can assume the SM to be valid
 - Energy-density is dominated by decoupled physics
- Inflation induces fluctuations to the Higgs field $\Delta\phi \sim H$
 - Important if $\bar{\Lambda}_{\max} \lesssim H$
 - State of the art calculations [2]: $\bar{\Lambda}_{\max} \sim 10^{11} \text{ GeV}$



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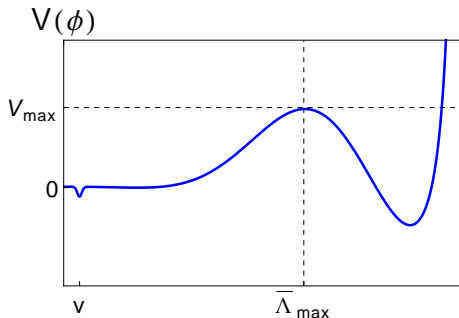
BICEP2/Keck/Planck

$$H \lesssim 10^{14}\text{GeV}$$

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BICEP2:

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Renormalization group improvement

- The perturbative $V_{\text{eff}}(\phi)$ suffers from large logarithms

Example: $V(\phi) = (1/2)m^2\phi^2 + (\lambda/4!)\phi^4$

$$M(\phi)^2 \equiv m^2 + \frac{\lambda}{2}\phi^2$$

$$V_{\text{eff}}(\phi) = \frac{1}{2}m^2\phi^2 + \frac{\lambda}{4!}\phi^4 + \frac{\overbrace{M(\phi)^4}^{\text{bracketed}}}{64\pi^2} \left[\log \left(\frac{M(\phi)^2}{\mu^2} \right) - \frac{3}{2} \right]$$

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- The physical result must not depend on μ :

$$\frac{d}{d\mu} V_{\text{eff}}(\phi) = 0 \quad \Leftrightarrow \quad \left\{ \mu \frac{\partial}{\partial \mu} + \beta_\lambda \frac{\partial}{\partial \lambda} + \gamma_\phi \phi \frac{\partial}{\partial \phi} \right\} V_{\text{eff}}(\phi) = 0$$

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⇒ can be used to improve the perturbative result [3]

optimal choice

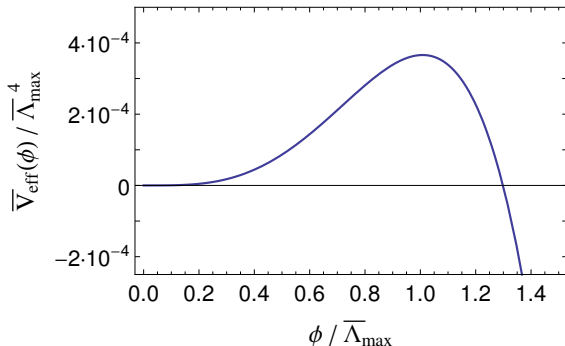
$$\mu \sim \phi$$

No large logarithms!

$$\Rightarrow V_{\text{eff}}(\phi) \approx \frac{\lambda(\phi)}{4}\phi^4$$

[3] Ford et. al. (1993)

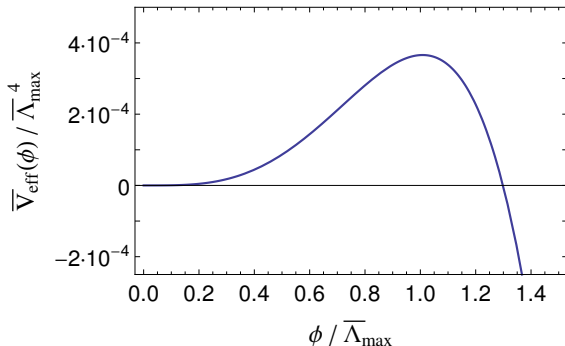
Stability results (Minkowski)



- For large H ($\sim 10^3 \bar{\Lambda}_{\text{max}}$), the SM is not stable [4]

[4] Kobakhidze & Spencer-Smith (2014); Hook et. al. (2014); Fairbairn & Hogan (2014); Enqvist, Meriniemi & Nurmi (2014); Zurek et. al. (2015); Espinosa et. al. (2015)

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Does including spacetime curvature in the quantum calculation change this?

[4] Kobakhidze & Spencer-Smith (2014); Hook et. al. (2014); Fairbairn & Hogan (2014); Enqvist, Meriniemi & Nurmi (2014); Zurek et. al. (2015); Espinosa et. al. (2015)

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Curved space effective action

- What is usually meant by "curved space QFT"?
 - No loops with $g^{\mu\nu}$
- A "curvature mass" is generated for SM

$$\Rightarrow V_{\text{eff}}(\phi) \approx \frac{\lambda(\mu)}{4} \phi^4 + \frac{\xi(\mu)}{2} R \phi^2$$

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- Solve modes in FRW space

$$\left[-\square + M(\phi)^2 + \xi R \right] \hat{\phi} = 0; \quad \hat{\phi} = \int \frac{d^3 k}{a(t)^{3/2}} \left[\hat{a}_{\mathbf{k}} u_{\mathbf{k}} + \hat{a}_{\mathbf{k}}^\dagger u_{\mathbf{k}}^* \right],$$
$$u_{\mathbf{k}} = \frac{1}{\sqrt{W}} e^{-i \int^t W dt'} e^{i \mathbf{k} \cdot \mathbf{x}}$$

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$$u_{\mathbf{k}} = \frac{1}{\sqrt{W}} e^{-i \int^t W dt'} e^{i\mathbf{k} \cdot \mathbf{x}}$$

$$W^2 = \frac{k^2}{a(t)^2} + M(\phi)^2 + \left(\xi - \frac{1}{6} \right) R + \mathcal{O}(k^{-2})$$

in curved space

$$\mu^2 \sim \phi^2 + R$$

1-loop Effective potential in curved space

$$V_{\text{eff}}(\phi, R) = -\frac{1}{2}m^2(\mu)\phi^2 + \frac{1}{2}\xi(\mu)R\phi^2 + \frac{1}{4}\lambda(\mu)\phi^4$$

$$+ \sum_{i=1}^9 \frac{n_i}{64\pi^2} M_i^4(\phi) \left[\log \frac{|M_i^2(\phi)|}{\mu^2} - c_i \right] \quad ; \quad \begin{aligned} M_i^2(\phi) &= \kappa_i \phi^2 - \kappa'_i + \theta_i R \\ \mu^2 &= \phi^2 + R \end{aligned}$$

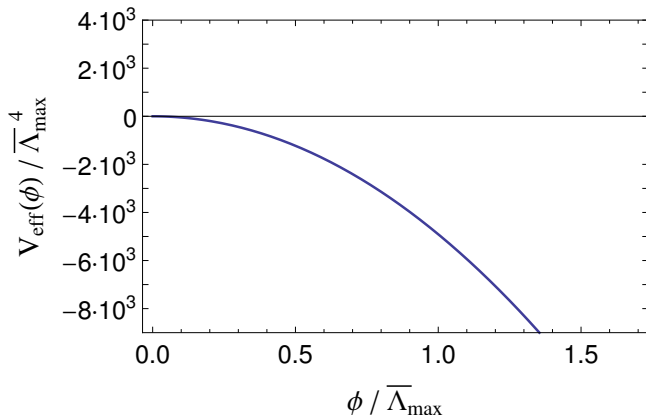
Φ	i	n_i	κ_i	κ'_i	θ_i	c_i
W^\pm	1	2	$g^2/4$	0	1/12	3/2
	2	6	$g^2/4$	0	-1/6	5/6
	3	-2	$g^2/4$	0	-1/6	3/2
Z^0	4	1	$(g^2 + g'^2)/4$	0	1/12	3/2
	5	3	$(g^2 + g'^2)/4$	0	-1/6	5/6
	6	-1	$(g^2 + g'^2)/4$	0	-1/6	3/2
t	7	-12	$y_t^2/2$	0	1/12	3/2
ϕ	8	1	3λ	m^2	$\xi - 1/6$	3/2
χ_i	9	3	λ	m^2	$\xi - 1/6$	3/2

Stability results (curved space) I

- First attempt, set $\xi_{EW} = 0$ and $H \sim 10^{10} \text{ GeV}$ ($\sim 10^3 \bar{\Lambda}_{\text{max}}$)

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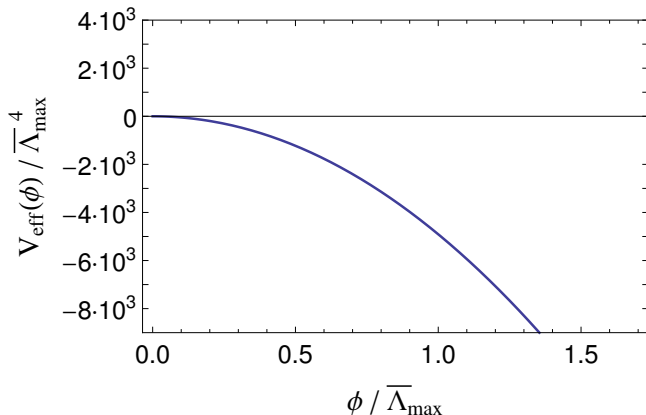
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- Potential negative *everywhere*

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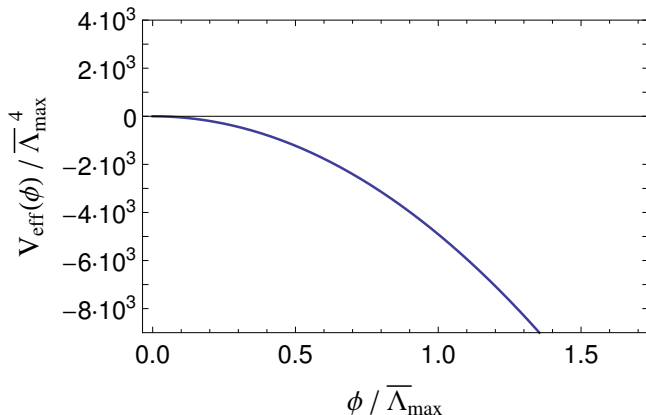
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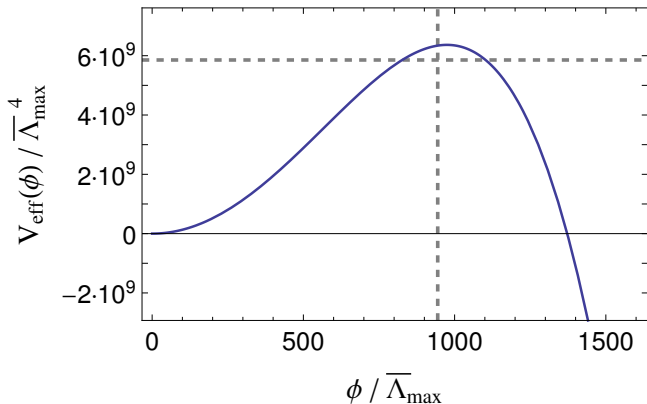
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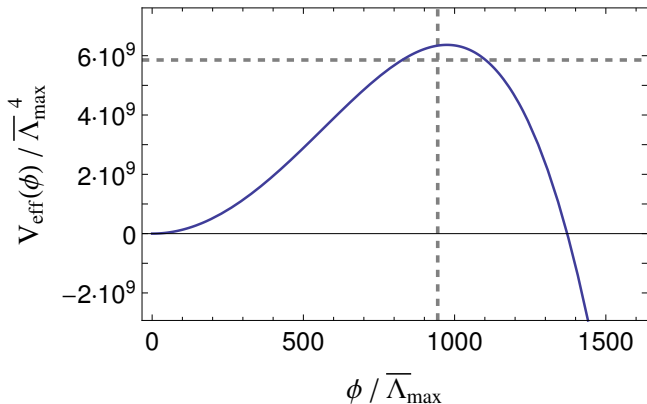
$$V_{\text{max}}^{1/4} \simeq H \frac{(6\xi)^{1/2}}{|\lambda|^{1/4}}$$

- $V_{\text{max}}(\text{curved}) \gg V_{\text{max}}(\text{flat})$ (and at a higher scale)

[5] Espinosa, Giudice & Riotto (2008)

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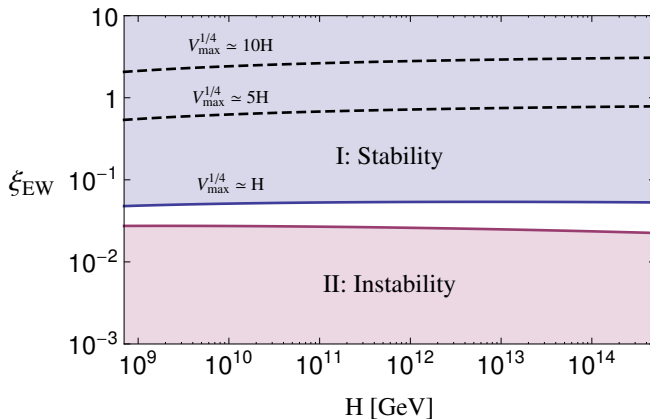
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- $V_{\text{max}}(\text{curved}) \gg V_{\text{max}}(\text{flat})$ (and at a higher scale)

$$P \sim \exp\left[-8\pi^2 (V_{\text{max}}/3H^4)\right] \Rightarrow \text{Stable!}$$

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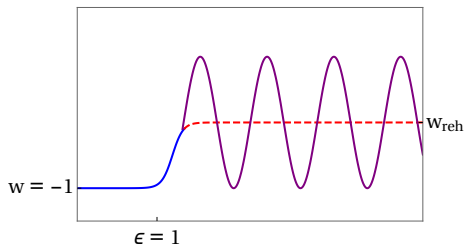
- The stability/instability of the potential is determined by ξ_{EW}



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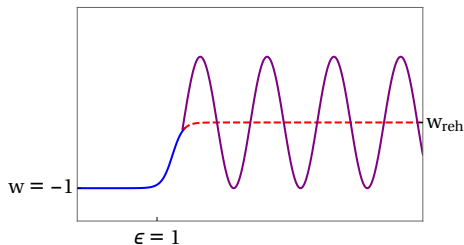
Reheating

- Energy of inflation is transferred to SM degrees of freedom, $T = 0 \rightarrow T_{\text{reh}}$
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- ϕ feels the dynamics of inflation via ξ !

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- More loops and smaller H
 - In flat space Λ_{\max} changes by orders of magnitude
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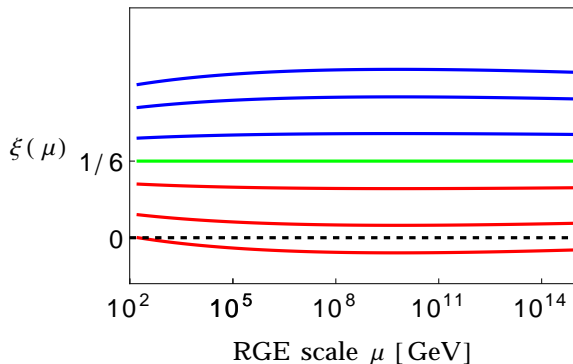
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Thank You!

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ξ_{EW}
0, 0.05, 0.12, 1/6,
0.22, 0.28, 0.33

Sensitivity to the choice of μ

- A loop calculation is never fully scale invariant
- How dependent is the result on the choice $\mu(t)^2 = \phi(t)^2 + R$?

$$\mu(t)^2 = \alpha\phi(t)^2 + \beta R \quad \alpha, \beta \in \{0.1 \dots 10\}$$

