

# ONE LOOP ELECTROWEAK CONSTRAINTS IN COMPOSITE HIGGS MODELS

R. Contino, M. Salvarezza, arXiv:1504.02750 [hep-ph]

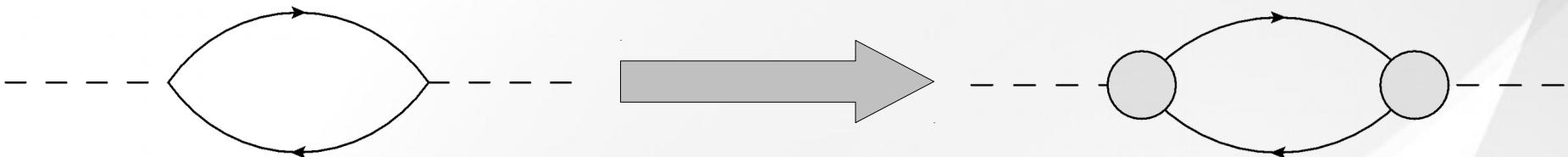
Matteo Salvarezza

Università di Roma “La Sapienza”, INFN Rome

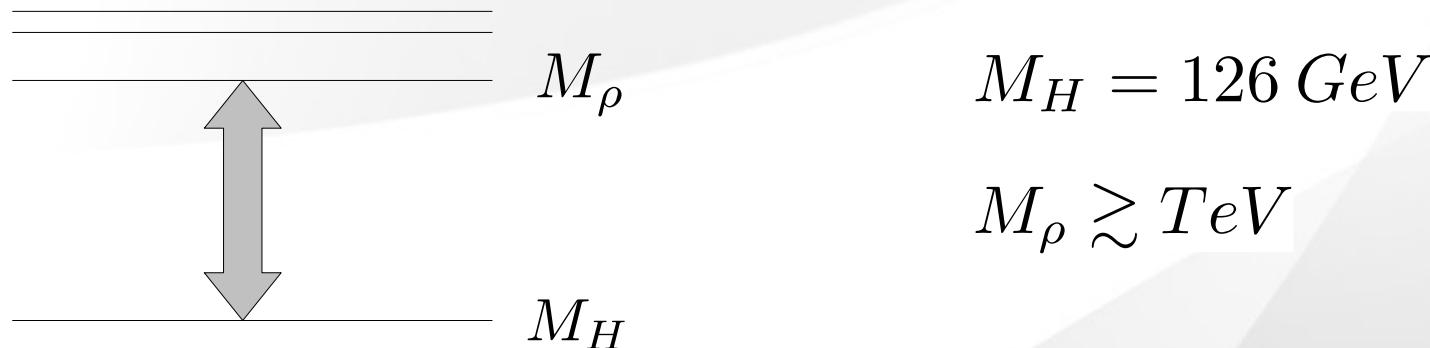
27th Rencontres de Blois  
Particle Physics and Cosmology  
03/06/2015

# STRONG EWSB: COMPOSITE HIGGS MODELS

- A heavy ( $\sim$ TeV) strong sector triggers the EWSB: the Higgs Boson is **composite**

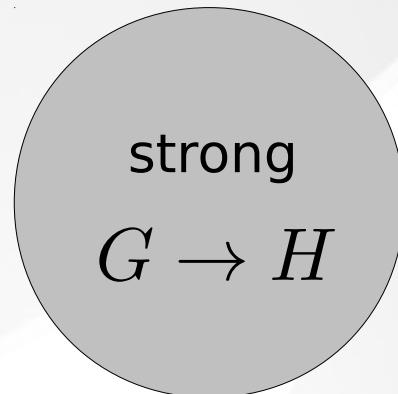


- **Problem**: big mass gap between the Higgs and the other resonances!



# STRONG EWSB: COMPOSITE HIGGS MODELS

- Learn from QCD: the Higgs as a **pseudo Nambu-Goldstone Boson** accounts naturally for such a picture [Georgi, Kaplan, 1984]
- General recipe:
  - EFT for NGB Higgs: spontaneous symmetry breaking  $G \rightarrow H$  at a scale  $f > v$



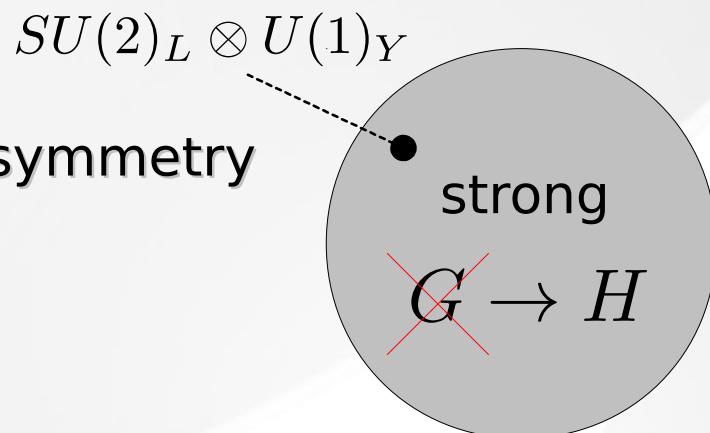
# STRONG EWSB: COMPOSITE HIGGS MODELS

- Learn from QCD: the Higgs as a pseudo Nambu-Goldstone Boson accounts naturally for such a picture

[Georgi, Kaplan, 1984]

- General recipe:

- EFT for NGB Higgs: spontaneous symmetry breaking  $G \rightarrow H$  at a scale  $f > v$
- Couple the SM to the strong sector  $\Rightarrow$  explicit breaking of  $G$



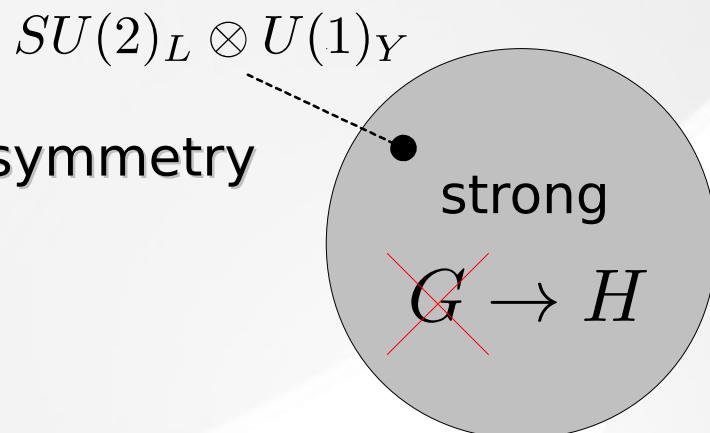
# STRONG EWSB: COMPOSITE HIGGS MODELS

- Learn from QCD: the Higgs as a pseudo Nambu-Goldstone Boson accounts naturally for such a picture

[Georgi, Kaplan, 1984]

- General recipe:

- EFT for NGB Higgs: spontaneous symmetry breaking  $G \rightarrow H$  at a scale  $f > v$
- Couple the SM to the strong sector  $\Rightarrow$  explicit breaking of  $G$
- A Higgs Potential is generated at one loop  $\Rightarrow$  Higgs mass, EWSB



# STRONG EWSB: COMPOSITE HIGGS MODELS

- Learn from QCD: the Higgs as a **pseudo Nambu-Goldstone Boson** accounts naturally for such a picture

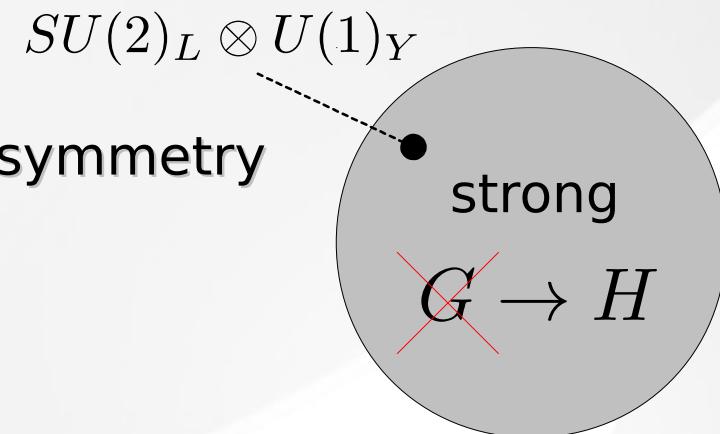
[Georgi, Kaplan, 1984]

- General recipe:

- EFT for NGB Higgs: spontaneous symmetry breaking  $G \rightarrow H$  at a scale  $f > v$
- Couple the SM to the strong sector  $\Rightarrow$  explicit breaking of  $G$
- A Higgs Potential is generated at one loop  $\Rightarrow$  Higgs mass, EWSB

- Separation of scales:

$$\xi = (v/f)^2$$



Decoupling:  $\xi \rightarrow 0$

# COMPOSITE HIGGS MODELS AND EWPT

- Oblique EWPO relevant for our purpose:

$$\epsilon_1 = \frac{1}{M_W^2} (A_{33}(0) - A_{W^+W^-}(0)) - M_Z^2 F'(M_Z^2)$$

$$\epsilon_3 = \frac{c}{s} F'_{3B}(M_Z^2) + c^2 (F_{\gamma\gamma}(0) - F_{\gamma\gamma}(M_Z^2)) - c^2 M_Z^2 F'_{ZZ}(M_Z^2)$$

$$\Pi_{ij}(q^2) = A_{ij}(0) + q^2 F_{ij}(q^2)$$

[Altarelli, Barbieri, Caravaglios, 1993]

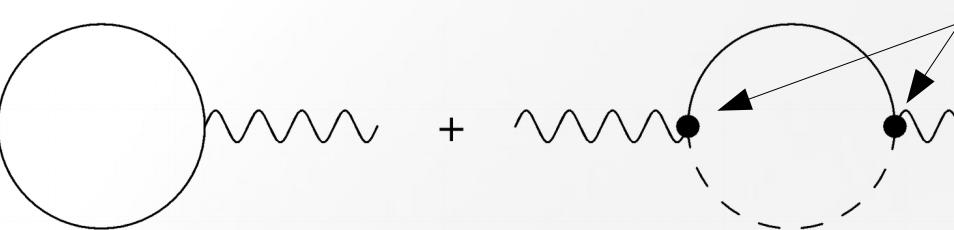
- Two sources of deviations  $\Delta\epsilon_i = \epsilon_i - \epsilon_i^{SM}$

- IR effect from modified Higgs dynamics

- UV effect from composite resonances

# COMPOSITE HIGGS MODELS AND EWPT

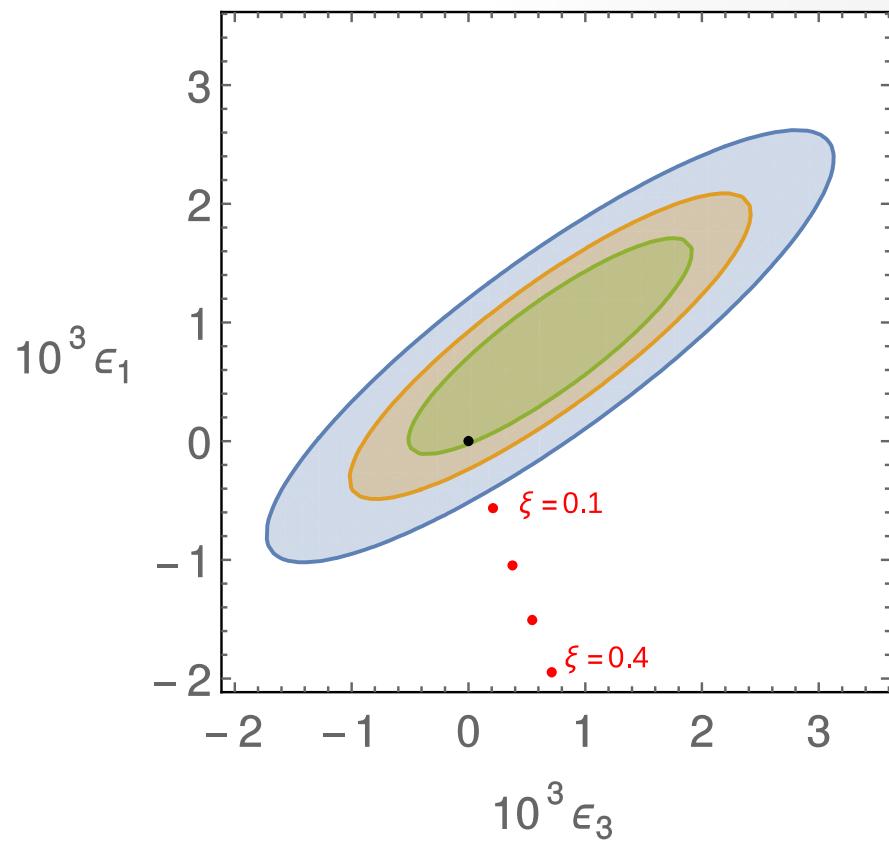
- IR effect on EWPO:

$$\epsilon_3 \supset \text{wavy loop} + \text{wavy loop with internal loop} \propto \xi \log\left(\frac{\Lambda}{M_H}\right)$$


# COMPOSITE HIGGS MODELS AND EWPT

- IR effect on EWPO:

$$\epsilon_3 \supset \text{wavy loop} + \text{wavy loop with internal loop} \propto \sqrt{1 - \xi} \log\left(\frac{\Lambda}{M_H}\right)$$



- Same effect for  $\epsilon_1$
- ~5-10% fine tuning on  $\xi$  required!
- Investigate if this picture can be improved (or not) with contributions from resonances

# VECTOR RESONANCES LAGRANGIAN

- Introduce a single spin-1 resonance lighter than the others

---

$\Lambda$

---

$M_\rho$

---

$M_H$

# VECTOR RESONANCES LAGRANGIAN

- Introduce a single spin-1 resonance lighter than the others

---

 $\Lambda$ 

---

 $M_\rho$ 

---

 $M_H$ 

- Goal:

Calculate the 1-loop contributions of such resonance to EWPO.

# VECTOR RESONANCES LAGRANGIAN

- Introduce a single spin-1 resonance lighter than the others

---

$\Lambda$

- Goal:

---

$M_\rho$

Calculate the 1-loop contributions of such resonance to EWPO.

---

$M_H$

➤ Estimates:  $\Delta\epsilon_i^{(\rho)} \sim \frac{1}{16\pi^2} \xi \log\left(\frac{\Lambda}{M_\rho}\right)$        $\Delta\epsilon_i^{(\Lambda)} \sim \frac{1}{16\pi^2} \xi$

- Not really a sizeable enhancement
- Not an actual precise calculation
- Rather a refined qualitative estimate

# VECTOR RESONANCES

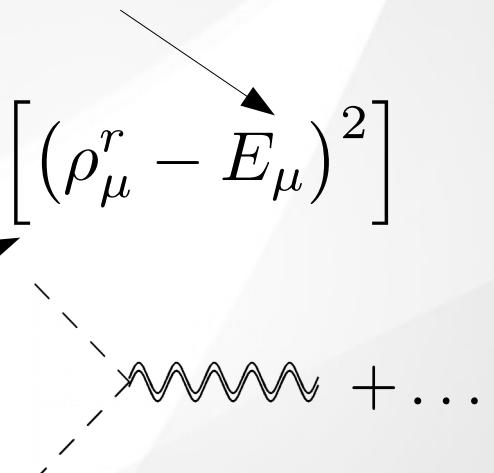
## LAGRANGIAN

- Coset G/H: focus on the minimal  $\text{SO}(5)/\text{SO}(4)$  Composite Higgs model
- Spin-1 resonance **in a  $(1,3)\oplus(3,1)$  of  $\text{SU}(2)_L \otimes \text{SU}(2)_R \sim \text{SO}(4)$**
- Two-derivatives lagrangian:

$$E^\mu = \pi D^\mu \pi + \dots$$

$$\mathcal{L}_\rho = \sum_{r=L,R} -\frac{1}{4g_{\rho r}^2} \text{Tr} [\rho_{\mu\nu}^r \rho^{r\mu\nu}] + \frac{1}{2} \frac{M_{\rho r}^2}{g_{\rho r}^2} \text{Tr} [(\rho_\mu^r - E_\mu)^2]$$

Mass + interactions:



- Interactions regulated by  $a_\rho \equiv \frac{M_\rho}{g_\rho f} \sim 1$

# VECTOR RESONANCES LAGRANGIAN

- Vector resonances mix linearly with EW gauge boson (partial compositeness)

$$\mathcal{L}_\rho \supset \frac{1}{2} \frac{M_\rho^2}{g_\rho^2} Tr \left[ (\rho^\mu - E^\mu)^2 \right] \supset M_\rho^2 \frac{g}{g_\rho} \rho_\mu^a W^{\mu a}$$

- Mixing is small and can be treated as a vertex
- Relevant four-derivative operators (additional tree-level contribution to  $\Delta\epsilon_3$ ):

$$\alpha_2 r Tr \left[ \rho_{\mu\nu}^r f^{\mu\nu} \right] \quad \left( f^{\mu\nu} = e^{-i\pi^a T^{\hat{a}}} F^{\mu\nu} e^{i\pi^a T^{\hat{a}}} \right)$$

# CALCULATION OF EW OBSERVABLES

➤ Some diagrams:

$$\epsilon_1 = \text{wavy line loop} + \text{wavy line loop with dashed outer boundary}$$

$B^\mu$

$$+ \text{wavy line loop} + \dots$$

---

$$\epsilon_3 = \text{wavy line loop with solid circle} + \text{wavy line loop with dashed circle}$$
$$+ \text{wavy line loop} + \text{wavy line loop with dashed outer boundary} + \dots$$

# CALCULATION OF EW OBSERVABLES

- Final expressions for the case of a single  $\rho$ :

$$\Delta\epsilon_1 = -\frac{3g'^2}{32\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_1\left(\frac{M_H^2}{M_Z^2}\right) \right] \quad \left( a_\rho = \frac{M_\rho}{g_\rho f} \right)$$


  
 IR Composite Higgs boson  
 contribution

$$+ \frac{3g'^2}{32\pi^2}\xi \frac{3}{4}a_\rho^2 \left[ \log\left(\frac{\Lambda}{M_\rho}\right) + \frac{3}{4} \right]$$

$$\Delta\epsilon_3 = \frac{g^2}{96\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_3\left(\frac{M_H^2}{M_Z^2}\right) \right]$$



$$+ g^2\xi \left( \frac{1}{4g_\rho^2} - \alpha_2 \right) - \frac{g^2}{96\pi^2}\xi \left[ \frac{3}{4} (a_\rho^2 + 28) \log\left(\frac{\Lambda}{M_{\rho L}}\right) + \frac{41}{16}a_\rho^2 + 1 \right]$$

# CALCULATION OF EW OBSERVABLES

- Final expressions for the case of a single  $\rho$ :

$$\Delta\epsilon_1 = -\frac{3g'^2}{32\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_1\left(\frac{M_H^2}{M_Z^2}\right) \right] \quad \left( a_\rho = \frac{M_\rho}{g_\rho f} \right)$$

$$+ \frac{3g'^2}{32\pi^2}\xi \frac{3}{4}a_\rho^2 \left[ \log\left(\frac{\Lambda}{M_\rho}\right) + \frac{3}{4} \right]$$

$$\Delta\epsilon_3 = \frac{g^2}{96\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_3\left(\frac{M_H^2}{M_Z^2}\right) \right]$$

UV Resonance contribution

$$+ g^2\xi \left( \frac{1}{4g_\rho^2} - \alpha_2 \right) - \frac{g^2}{96\pi^2}\xi \left[ \frac{3}{4} (a_\rho^2 + 28) \log\left(\frac{\Lambda}{M_{\rho L}}\right) + \frac{41}{16}a_\rho^2 + 1 \right]$$

# CALCULATION OF EW OBSERVABLES

- Final expressions for the case of a single  $\rho$ :

$$\Delta\epsilon_1 = -\frac{3g'^2}{32\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_1\left(\frac{M_H^2}{M_Z^2}\right) \right] \quad \left( a_\rho = \frac{M_\rho}{g_\rho f} \right)$$

↑

$<0$

$$+ \frac{3g'^2}{32\pi^2}\xi \frac{3}{4}a_\rho^2 \left[ \log\left(\frac{\Lambda}{M_\rho}\right) + \frac{3}{4} \right]$$

↑

$>0$

$$\Delta\epsilon_3 = \frac{g^2}{96\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_3\left(\frac{M_H^2}{M_Z^2}\right) \right]$$

$$+ g^2\xi \left( \frac{1}{4g_\rho^2} - \alpha_2 \right) - \frac{g^2}{96\pi^2}\xi \left[ \frac{3}{4} (a_\rho^2 + 28) \log\left(\frac{\Lambda}{M_{\rho L}}\right) + \frac{41}{16}a_\rho^2 + 1 \right]$$

# CALCULATION OF EW OBSERVABLES

- Final expressions for the case of a single  $\rho$ :

$$\Delta\epsilon_1 = -\frac{3g'^2}{32\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_1\left(\frac{M_H^2}{M_Z^2}\right) \right] \quad \left( a_\rho = \frac{M_\rho}{g_\rho f} \right)$$

$$+ \frac{3g'^2}{32\pi^2}\xi \frac{3}{4}a_\rho^2 \left[ \log\left(\frac{\Lambda}{M_\rho}\right) + \frac{3}{4} \right]$$

$>0$



$$\Delta\epsilon_3 = \frac{g^2}{96\pi^2}\xi \left[ \log\left(\frac{\Lambda}{M_Z}\right) + f_3\left(\frac{M_H^2}{M_Z^2}\right) \right]$$

$<0$



Any sign

$$\rightarrow +g^2\xi \left( \frac{1}{4g_\rho^2} - \alpha_2 \right) - \frac{g^2}{96\pi^2}\xi \left[ \frac{3}{4} (a_\rho^2 + 28) \log\left(\frac{\Lambda}{M_{\rho L}}\right) + \frac{41}{16}a_\rho^2 + 1 \right]$$

# PARAMETER SPACE CONSTRAINTS

➤ Fix parameters:

$$\alpha_2 = 0$$

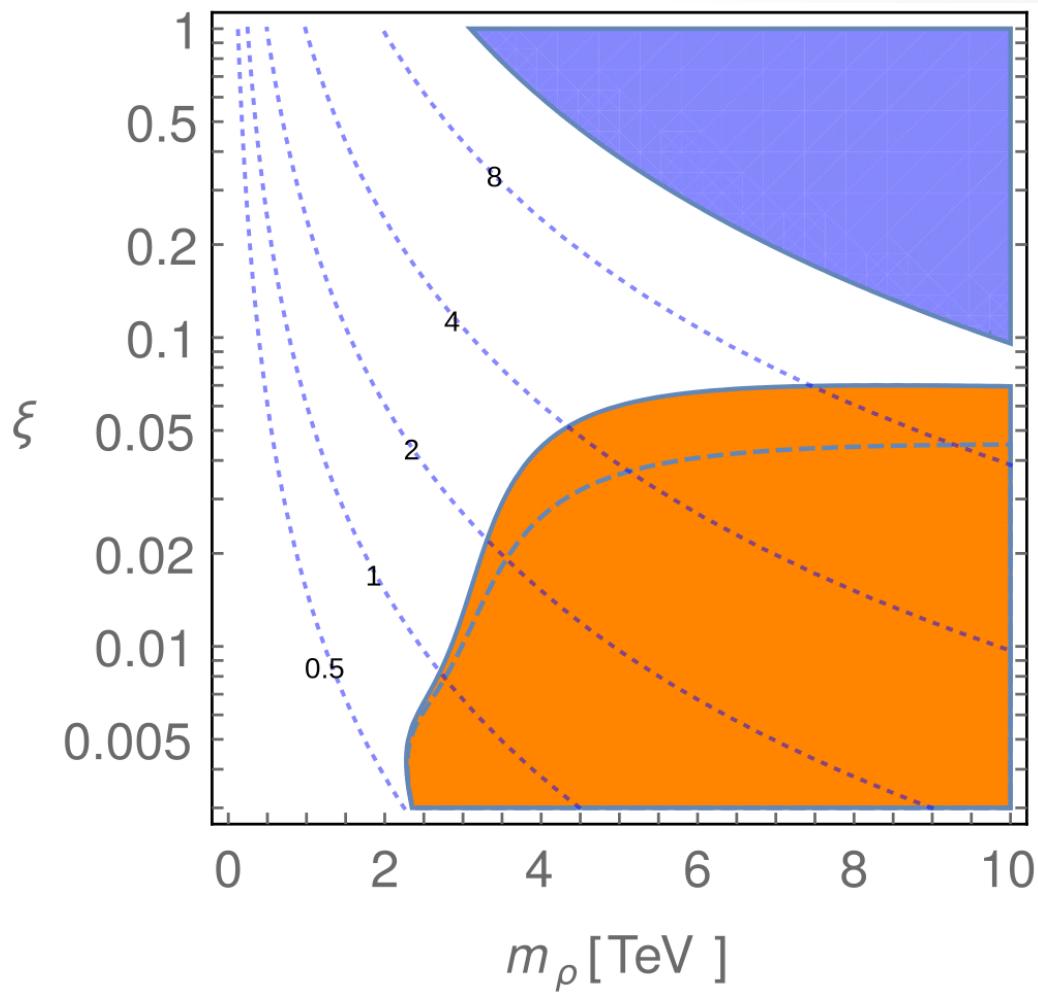
$$a_\rho = 1$$

$$\Lambda = 3M_\rho$$

➤ Constrain  $(\xi, M_\rho)$  @ 95% CL

---

# PARAMETER SPACE CONSTRAINTS



➤ Fix parameters:

$$\alpha_2 = 0$$

$$a_\rho = 1$$

$$\Lambda = 3M_\rho$$

➤ Constrain  $(\xi, M_\rho)$  @ 95% CL

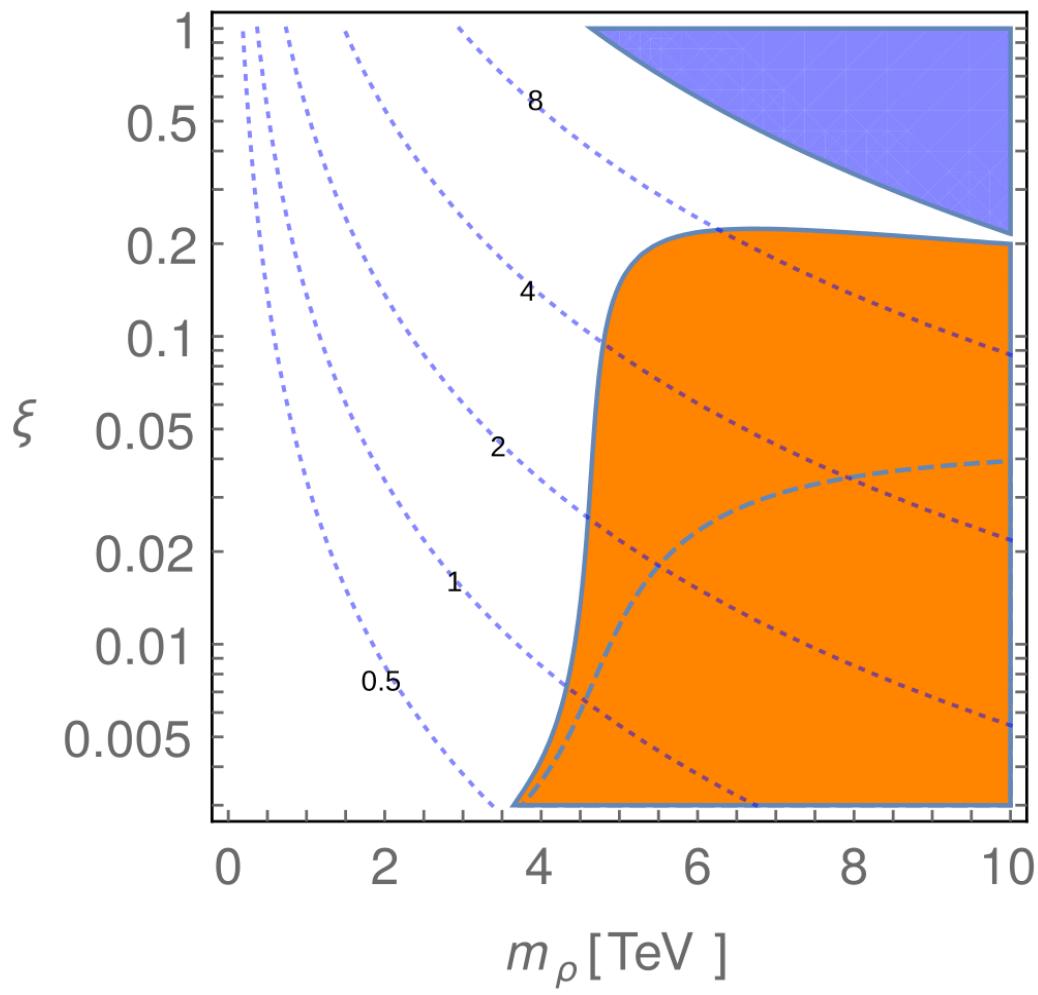
---

◻ IR + tree level  $\rho$

■ IR + one loop  $\rho$

■  $g_\rho > 4\pi$

# PARAMETER SPACE CONSTRAINTS



➤ Fix parameters:

$$\alpha_2 = 0$$

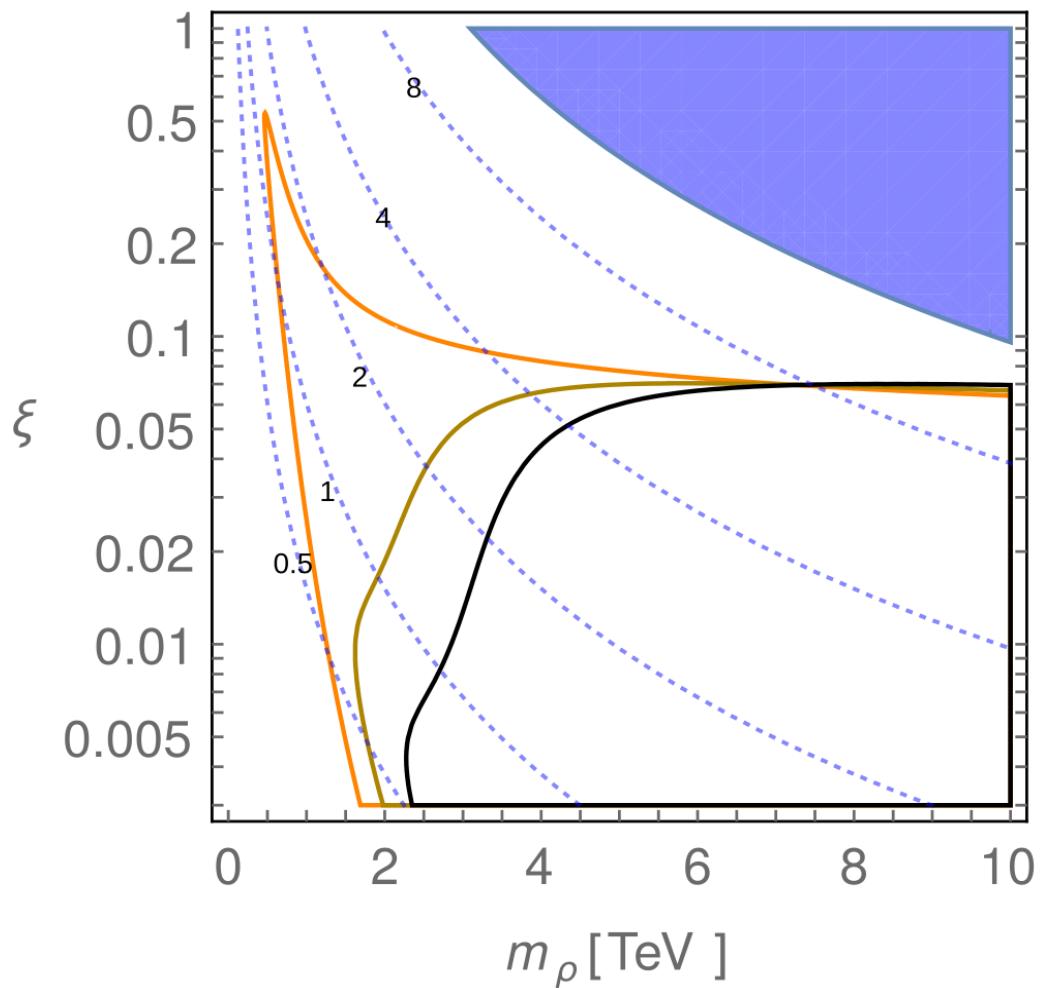
$$a_\rho = 1.5$$

$$\Lambda = 3M_\rho$$

➤ Constrain  $(\xi, M_\rho)$  @ 95% CL

- 
- IR + tree level  $\rho$
  - IR + one loop  $\rho$
  - $g_\rho > 4\pi$

# PARAMETER SPACE CONSTRAINTS



➤ Fix parameters:

$$\alpha_2 = 0, 1/8g_\rho^2, 1/4g_\rho^2$$

$$a_\rho = 1$$

$$\Lambda = 3M_\rho$$

➤ Constrain  $(\xi, M_\rho)$  @ 95% CL

---

■  $\alpha_2 = 0$

■  $\alpha_2 = 1/8g_\rho^2$

■  $\alpha_2 = 1/4g_\rho^2$

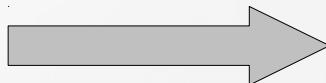
# TO DO (ALMOST DONE)

- Aim at a first “complete” description of EWPT in composite Higgs
  - Vector resonances
  - Fermionic resonances
- Include many constraints
  - Oblique EWPO
  - $Z\bar{b}b$  coupling
  - Higgs potential
  - Top mass
  - $Z$  mass

# CCWZ LAGRANGIAN FOR $\text{SO}(5) \rightarrow \text{SO}(4)$

- CCWZ: natural formalism for spontaneously broken effective theories, chiral expansion is built-in.

• **Global G invariance**



• **Local H invariance**

[Callan, Coleman, Wess, Zumino, 1977]

- Building blocks:

$$d_\mu = d_\mu^a T^{\hat{a}} = T^{\hat{a}} \left[ \frac{\sqrt{2}}{f} \partial_\mu \pi^a + \frac{1}{f\sqrt{2}} \epsilon^{abc} \pi^b (W_\mu^c + \delta^{c3} B_\mu) + \dots \right]$$

$$E_\mu = E_\mu^a T^a = T^a \left[ \frac{1}{2f^2} (\epsilon^{abc} \pi^b \partial_\mu \pi^c + \pi^a \partial_\mu \pi^4 - \pi^4 \partial_\mu \pi^a) + W_\mu^a + \dots \right]$$

$$d_\mu \xrightarrow[g \in SO(5)]{} h(x) d_\mu h^\dagger(x) \quad h \in SO(4)$$

$$E_\mu \xrightarrow[g \in SO(5)]{} h(x) d_\mu h^\dagger(x) - i h \partial_\mu h^\dagger$$

$$\mathcal{L} = \mathcal{L} (d_\mu, E_\mu)$$

$d_\mu, E_\mu \longleftrightarrow 1 \text{ derivative}$