# Dark Matter Simplified models

Alessandro Vichi



# Physics at the High-Luminosity LHC

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## EFT not an option..

- EFT are a powerful and general tool, but must used cum grano salis
- At LEP and Teveatron the use of EFT was perfectly legitimate
- At LHC8, and even more at LHC14, the regime of validity of EFT is shrinking
- Quantitatively well illustrated by the quantity R<sub>Λ</sub>: the ratio of the cross section obtained in the EFT with the requirement Q<sub>tr</sub> < Λ over the total cross section obtained in the EFT. It gives a measure of the fraction of events with momentum transfer lower than the EFT cutoff scale: R<sub>Λ</sub> < 1 signals the failure of EFT description.</p>
- EX:  $D5 = \bar{\chi} \gamma_{\mu} \chi \bar{q} \gamma^{\mu} q$



[G.Busoni, A.De Simone, J.Gramling, E.Morgante, A.Riotto]

### Dark Matter at colliders: what to look for?

- Dark matter: particle stable on collider time scales
- ► If dark matter is the lightest state odd under an (approximate) Z<sub>2</sub> symmetry it will pair produced at collider
- Leading signal:  $pp \rightarrow MET$



- Most common strategy is to dress the leading process with bosons and look for mono-X (X: jet, Z, W, γ, Higgs)
- This not necessarily the best strategy: for a given simplified model the search for the mediator is complementary and sometimes more constraining

EX: t-channel scalar mediator (squark like) + fermion DM Dijets+MET are more constraining than mono-jet or any other mono-X search ex: [M. Papucci, AV, K. Zurek]

### Vector s-channel mediator: Leptophobic Z'

$$\mathcal{L} \supset g_q Z'_\mu \sum_{i=1,2} \left( \bar{\mathcal{Q}}_L^i \gamma^\mu \mathcal{Q}_L^i + \bar{u}_R^i \gamma^\mu u_R^i + \bar{d}_R^i \gamma^\mu d_R^i \right) + g_{DM} Z'_\mu \bar{\chi} \gamma^\mu \chi + \mu_{Z'} h Z'_\mu Z'^\mu.$$

- $m_{DM} > m_{Z'}/2$ : poor sensitivity at 8 TeV
- $m_{DM} \leq m_{Z'}/2$ : Jets+MET ( $\equiv$  monojet) win



#### [S.P.Liew, M. Papucci, AV, K. Zurek]

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► What about HL-LHC?



[O.Buchmueller, M.J. Dolan, S.A. Malik, C.McCabe]

- ► So far only mono-jet studies have been performed: sensible increase in reach for  $m_{DM} \le m_{Z'}/2$ . Wish list:
  - ► ATLAS/CMS reinterpretation of dijet resonance search and angular distribution search as limits on g<sub>q</sub>

Susy-like simplified models: squarks+DM

$$\mathcal{L} \supset g_{DM} \sum_{i=1,2} \left( \widetilde{\mathcal{Q}}_L^i \overline{\mathcal{Q}}_L^i + \widetilde{u}_R^i \overline{u}_R^i + \widetilde{d}_R^i \overline{d}_R^i \right) \chi + \text{mass terms} + h.c.$$

 model well studied by susy searches: jets+MET are the most constraining search in the bulk of parameter space



- Projections show poor sensitivity in the compressed region
- mono-X will help in the compressed region

# Mono-Z

- When Z bosons are emitted from the initial state the process is not expected to have sizeable cross sections.
- ► Z boson radiated from final (or internal) states can have sizeable signal



- ▶ In a squarks+DM model, generically this analysis still gives negligible limits w.r.t. a jets+MET search
- When  $m_{SO} \simeq m_{DM}$ , cross section enhanced by the process:

 $gg \longrightarrow \widetilde{q}\widetilde{q}^{\dagger}Z \longrightarrow 2(\text{soft})j + Z + \text{MET}$  mono-Z final state

HL-LHC projections shows exclusion limits up to  $m_{SQ} \simeq 500 GeV$ ,  $(m_{SQ} - m_{DM} = 10 GeV)$ . [S.P.Liew, M. Papucci, AV, K. Zurek]

#### Wish list:

Study the impact of an hybrid analysis Z+MET+(not so soft)jets

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# Susy-like simplified models: sbottoms+DM

Similar to previous slides:

model well studied by susy searches: jets+MET are the most constraining search in the bulk of parameter space



- Projections show poor sensitivity in the compressed region
- mono-jet and mono-Z will help in the compressed region
- Solution can have sizeable coupling to the Higgs without introducing flavour issues  $(g_H \sim m_b)$
- HL-LHC starts being sensitive to mono-Higgs signals, but still not competitive using current analysis

# **Mono-Higgs**

(Not so) Simplified model where mono-Higgs can play a fundamental role: [A.Berlin, T.Lin, L.T.Wang]

- 2 Higgs doublet model
- Heavy Z' coupled to quarks and higgs sector only
- ▶ DM coupled to pseudo-scalar A<sub>0</sub>







# **Mono-Higgs**

- At present no dedicated mono-Higgs analysis exists.
- Current limits are obtained recasting the limits on the signal strength:

$$\mu = \frac{\sigma(pp \to hZ(Z \to \nu n\bar{u}))}{\sigma_{SM}}$$

Not binned on missing energy or Higgs p<sub>T</sub>

#### Wish list:

Optimise a mono-Higgs analysis for different simplified models

## Neutralinos

For the first time we will be sensitive to Simplified Models where the Dark Matter is the only light state.

- Can be though as a limit of MSSM
- Also, in the framework of minimal dark matter, one introduces one electroweak multiplet at the time and uses the neutral component as dark matter candidate
- Explaining the relic abundance fixes the DM mass
- ► Higgsino, Wino are most considered examples, but other can be studied



[M.Low, L.T.Wang]

# **Disappearing tracks**

- A standard and agnostic method to look for neutralinos is the use of mono-X (mostly mono-jet) searches.
- ► Starting from HL-LHC this is not the most effective technique
- In the minimal scenario where only one multiplet is light, charged and neutral components are split by loop effects:

 $m_{\chi^{\pm}} - m_{\chi^0} \simeq 166 \text{ MeV} (355 \text{MeV})$  for a Wino (Higgsino)

The small splitting results is a macroscopic lifetime of the chargino, which then can leave a trace in the detector before decaying: disappearing tracks are the smoking gun of this Simplified Model.



#### [M.Low, L.T.Wang]

- Soft leptons and VBF not as constraining.
- ► To saturate the relic abundance  $m_W \sim 3.1$ TeV,  $m_H \sim 1$ TeV: HL-LHC can't close the parameter space but it can start carving it out.

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### Conclusions

- ► HL-LHC will definitively improve the reach of DM searches
- Complementarity of different channels will be important, since they are sensitive to different combination of simplified models couplings
- Optimise the searches based on Simplified models will be a key aspect

In this talk I only focused on a small subset of simplified models. A more complete list of Simplified models for mono-X (other than mono-jet) could be:

| Model            | mono-Higgs                                     | mono-Z                                                     | Model            | mono-Higgs                                                               | mono-Z                           |
|------------------|------------------------------------------------|------------------------------------------------------------|------------------|--------------------------------------------------------------------------|----------------------------------|
|                  | q<br>$Z; Z'_{Z}; Z' \qquad \chi; \phi$         |                                                            |                  | q<br>2' x' x                                                             | $q \xrightarrow{\chi} \chi \chi$ |
| s-channel vector | $q = \chi; \phi$                               |                                                            | Inelastic DM     | q h                                                                      | q Z                              |
|                  |                                                | $\overline{q}$ $Z$ $K$ | Squark/shottom   | $q$ $q$ $q$ $h$ $h$ $\chi$                                               |                                  |
| s-channel scalar | q <sup>γ</sup> <sup>1</sup> χ                  | 9 X                                                        | Squark/sbouoin   |                                                                          |                                  |
| 2HDM             | $q$ $q$ $Z; Z_{h;S}$ $\chi; \phi$ $\chi; \phi$ |                                                            | Inelastic squark | $q = \frac{\chi}{\hat{q}} + \frac{\chi}{\hat{h}} + \frac{\chi}{\hat{h}}$ | $q$ $\chi$ $\chi$ $\chi$ $\chi$  |

<sup>[</sup>S.P.Liew, M. Papucci, AV, K. Zurek]