



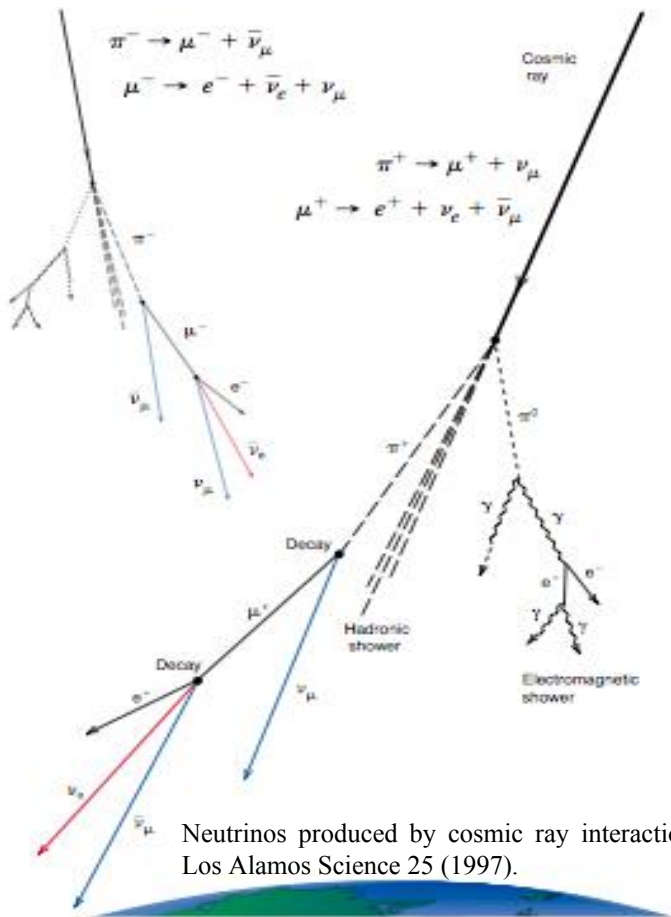
# A first look at NOvA's upward-going muon trigger data

Rob Mina - University of Virginia - DPF  
2015

# The NOvA Program

- NOvA (NuMI Off-axis  $\nu_e$  Appearance) is a two-detector accelerator neutrino oscillation experiment that uses Fermilab's NuMI, the most intense high-energy neutrino beam on the planet.
- It will say something about the  $\theta_{23}$  octant,  $\Delta m^2_{23}$ , and  $\delta_{cp}$ .
- But I'm going to talk about **something else entirely**.

# Atmospheric neutrinos



Neutrinos produced by cosmic ray interactions, from Los Alamos Science 25 (1997).

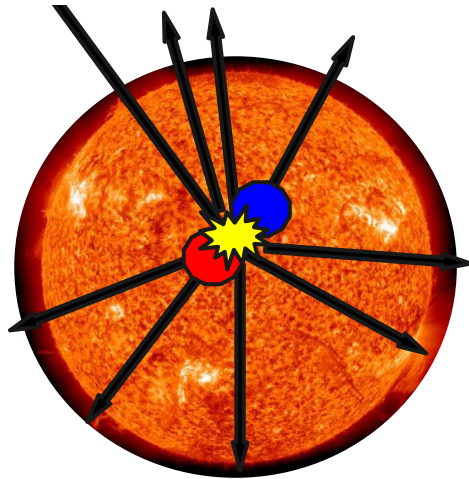
High-energy cosmic rays interact in Earth's atmosphere and produce particles that decay into neutrinos (among other things).

It is well established that atmospheric neutrinos oscillate<sup>[1]</sup>. We should see this oscillation in NOvA.

# WIMPs



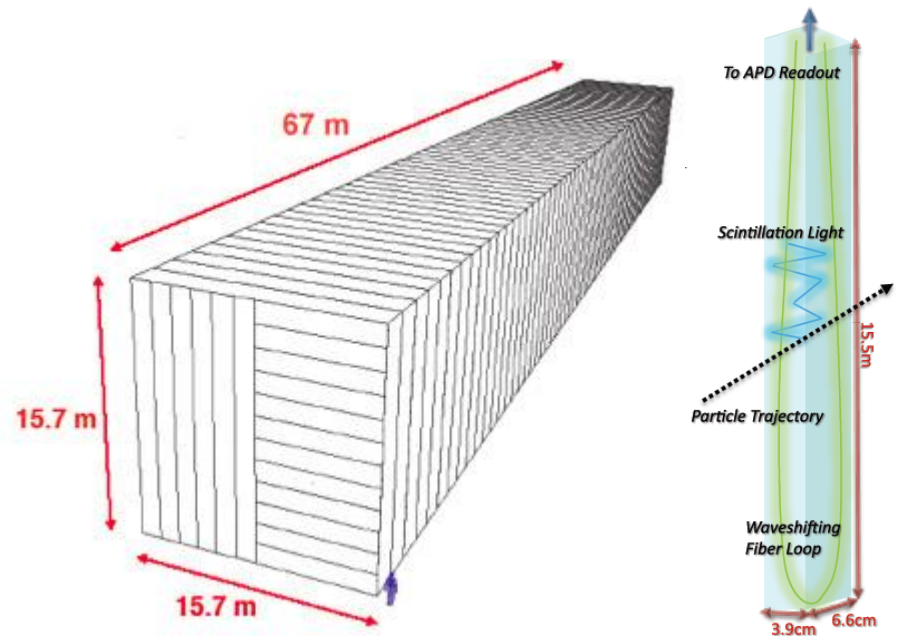
WIMPs annihilate in the Sun, producing neutrinos, some of which may be detectable on Earth.



- Dark matter contributes about 25% to the energy density of the observable universe, while “normal” (Standard Model) matter contributes only about 4% [2].
- One important class of theoretical candidate particles for Dark Matter is the Weakly Interacting Massive Particle (WIMP).
- WIMPs can become trapped in the gravitational well of the Sun, where they annihilate at a rate equal to the capture rate. These annihilations produce a neutrino signal to which NOvA may be sensitive.

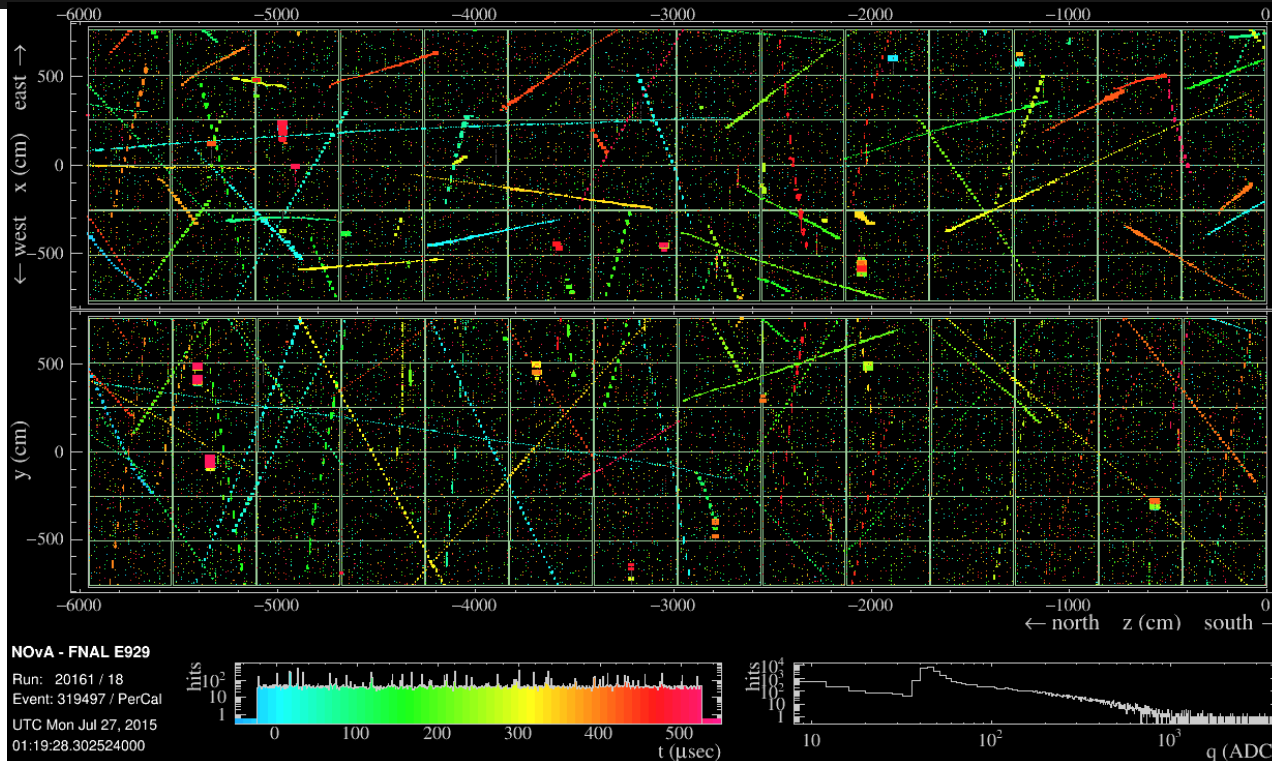
# The NOvA Far Detector

- Located 14 mrad off-axis of the beam, 810 km from the source, in Ash River, MN.
- Large - 14 kTon - highly active scintillating tracking calorimeter.
- Fine spatial granularity ( $\sim 3\text{cm}$ ) and time resolution ( $\sim 10\text{ns}$ )  $\rightarrow$  an **excellent muon tracker** (among other things).
- $> 300,000$  independently active channels, each read out at  $2 \times 10^6 \text{Hz}$   $\rightarrow$  **very high data throughput**.



NOvA's alternating planes allow particle tracking in 3D. A NOvA cell

# The NOvA UpMu Trigger



A 550  $\mu\text{sec}$  exposure of the far detector showing the high rate of cosmic ray muons.

- The rate of downward-going muons from cosmic ray interactions is very high ( $\sim 100\text{kHz}$ ) in the far detector because it is on the surface.
- **Upward-going muons are relatively rare.** The expected rate of upward-going muons from cosmic ray-induced neutrinos is  $\sim 1$  per day.

# The NOvA UpMu Trigger

- To say anything about these atmospheric neutrinos, NOvA must overcome this (cosmic) **background-to-signal ratio of  $\sim 10^{10}$** .
- Predictions of the WIMP neutrino flux are often quite small even compared to the atmospheric neutrino flux (which is a background in the indirect dark matter search).
- Data storage and processing is expensive, so we can't afford to record every muon track.
- We have written two upward-going muon triggers to produce signal-enriched samples:
  - *The “through-going” trigger is optimized to search for through-going muons from interactions in the rock surrounding the detector.*
  - The “contained” trigger searches for muon tracks from neutrino interactions within the detector.

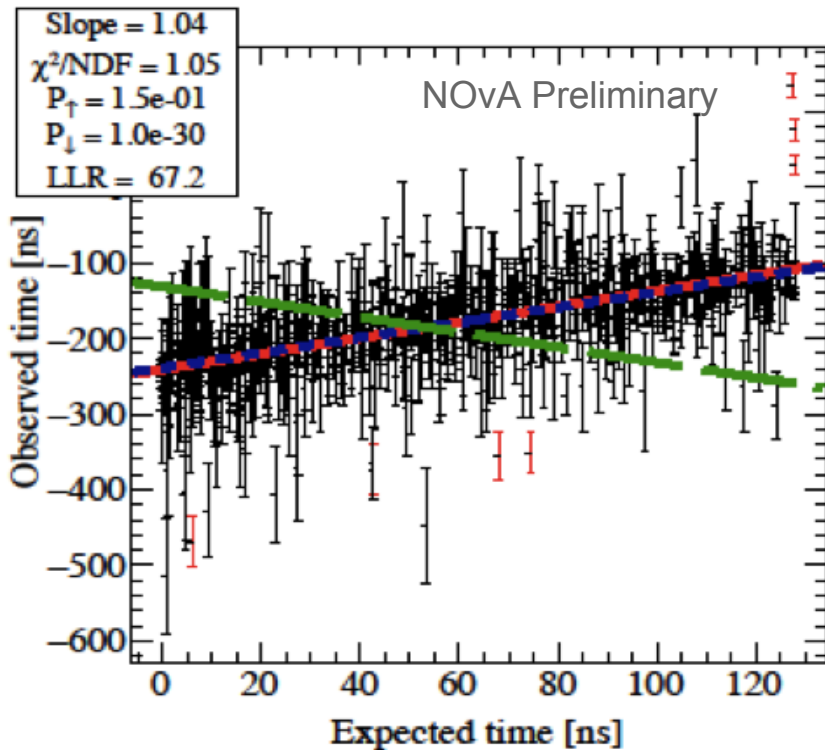
# General strategies for reducing cosmic background

Neutrino experiments employ several methods to select “interesting” upward-going muons:

- geometric veto - Cosmic ray muons originating outside the detector must travel through the top or sides. Vetoing events with activity near the edges of the detector will eliminate this class of background. The contained trigger searches for fully contained events, but **the through-going trigger can't use this.**
- topology - particular signal interactions often produce distinctive event topologies in detectors. This search can benefit especially from vertex-finding and Michel electron identification. **Work is ongoing to exploit these features.**
- **timing** - a finely segmented detector with sufficiently small hit time resolution (eg. NOvA) can use hit timing along a track to infer the direction of the particle. **This is the primary strategy used in this search.**



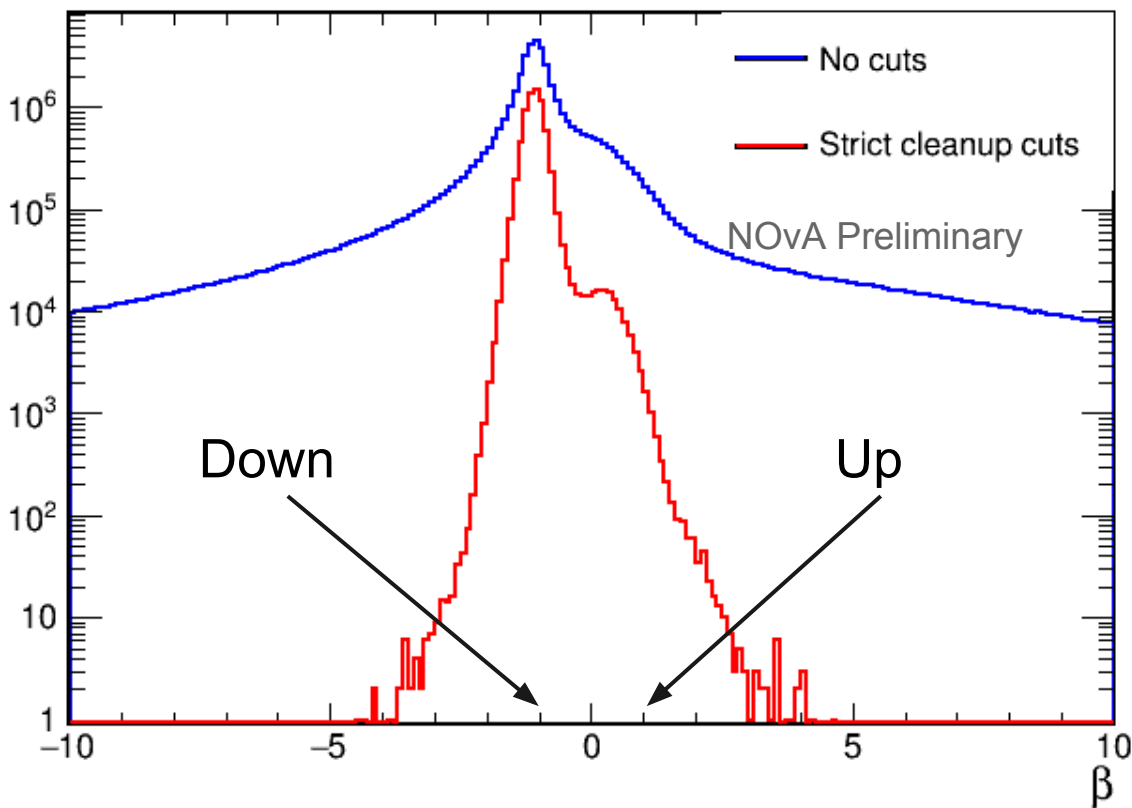
# $\beta$



The measured times vs expected times for cell hits along a promising upward-going muon candidate. Uncertainties are taken from a parameterization of single-hit timing resolution as a function of energy. The best-fit slope ( $\beta$ , red line) of 1.04 indicates a likely upward-going particle.

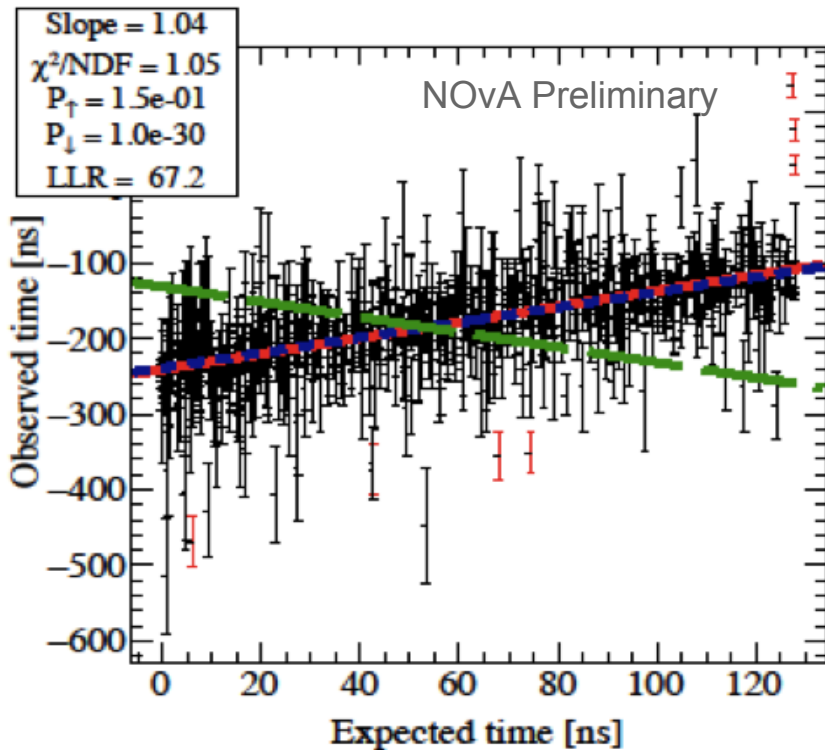
- To determine directionality, we calculate a speed,  $\beta$ , for each muon-like (linear) track, normalized by  $c$ . 1.0 means upward-going and -1.0 means downward-going.
- This is simply the slope from a linear fit to the measured times of the track cell hits vs the expected times assuming an upward-going relativistic particle.

# $\beta$



Even with strict cleanup cuts designed to maximize the discriminating power of  $\beta$ , we cannot achieve separation between downward- and upward-going muons in the data. A more powerful discriminator is needed.

# Log Likelihood Ratio



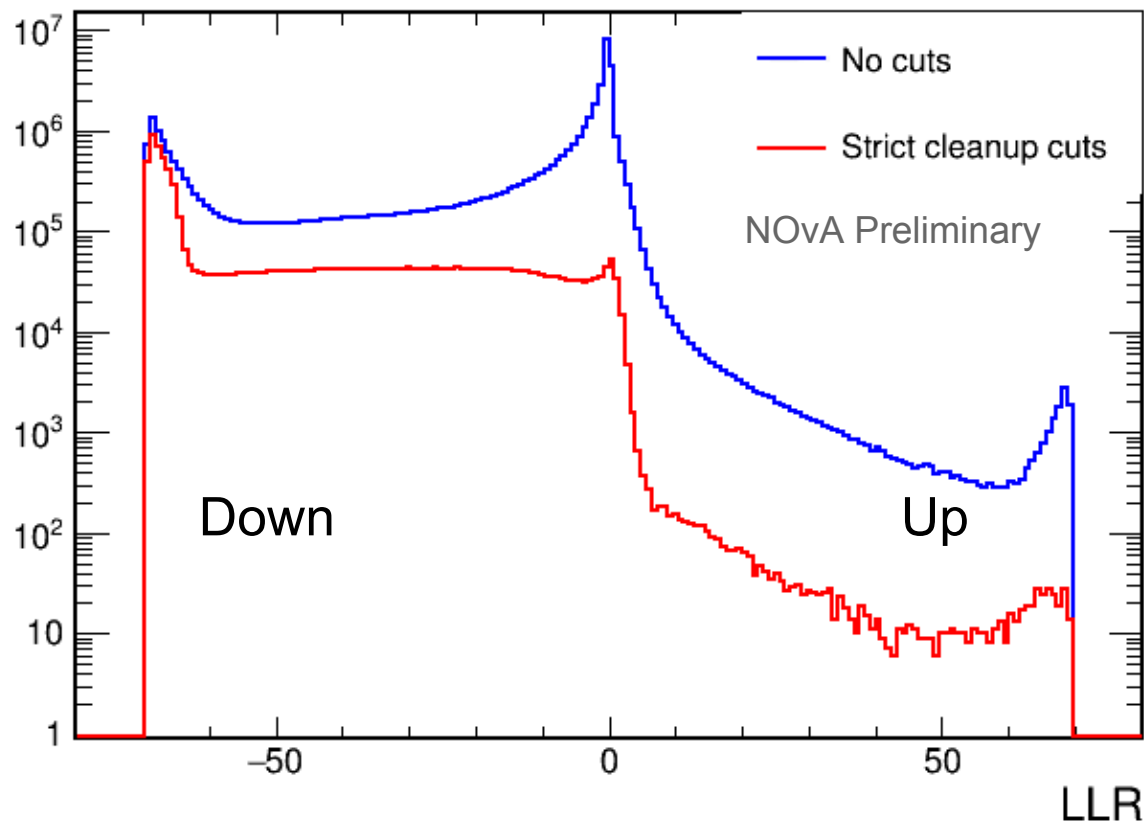
The measured times vs expected times for cell hits along a promising upward-going muon candidate. The fit with slope 1 is in blue and that with slope -1 is in green. The LLR of 67.2 indicates a likely upward-going particle.

If a track is not perfectly horizontal (we place a minimum on the absolute value of  $\Delta Y$ ), then the true value of  $\beta$  is either 1 or -1.

We leverage this simplifying constraint using the LLR:

- Perform a linear fit to the measured vs expected times with a fixed slope of 1
- Perform another fit with a fixed slope of -1
- Calculate the probability of each fit using the  $\chi^2$ . Set a nonzero minimum for the probabilities ( $10^{-30}$ ).
- $\text{LLR} = \ln(\text{Prob}_{\text{up}} / \text{Prob}_{\text{dn}})$
- LLR close to 0 means the track was neither conclusively upward- nor downward-going.

# LLR



Combined with cleanup cuts, the LLR provides a powerful discriminator for track directionality. The drop of 3 orders of magnitude between tracks with negative (downward-going) and positive (upward-going) LLR reflects this.

# UpMu sample

Today we'll take a first look at some of the data taken by the “through-going” trigger, which fires at  $\sim 1$  Hz, thus reducing the background in the triggered sample by  $\sim 5$  orders of magnitude.

The sample consists of 695 runs spanning a period of 164 days with total livetime 83.98 days.

# Selecting candidates

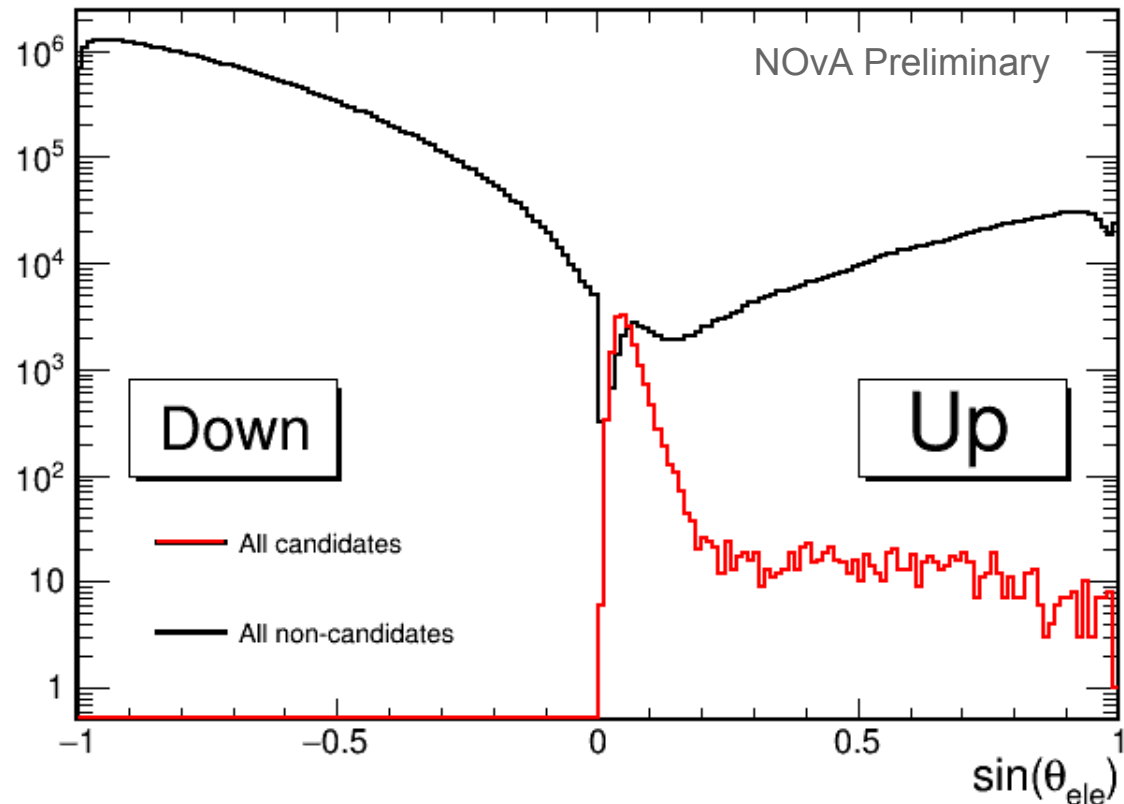
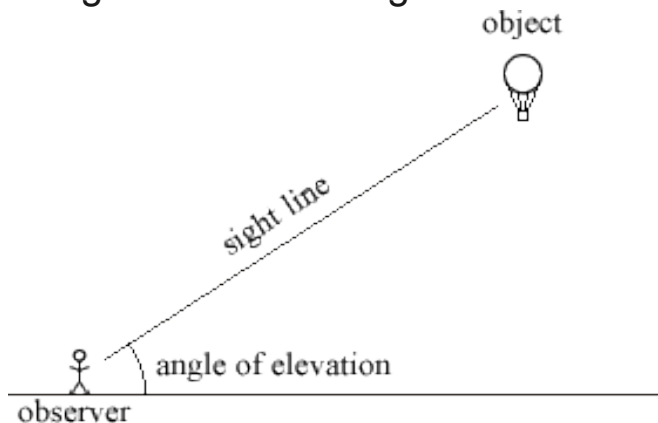
To select upward-going candidates from this sample (which contains ~43 million tracks), I began with cuts used in the trigger, and tightened the requirement on LLR:

- Length > 5m
- $\Delta X > 5$  cell widths
- $\Delta Y > 10$  cell widths
- $\Delta Z > 5$  plane widths
- LLR > 10
- LLR per view > 5
- $0 < \text{best-fit } \beta < 2$
- $\chi^2$  of linear fit mT vs eT < 1.5
- $R^2$  of linear fit to X, Y vs Z per view > 0.99
- Number of hits used in timing fits > 60
- Number of hits per view used in timing fits > 15

# Elevation angle

## Most of the candidate tracks are horizontal.

Cosmic ray muons travelling slightly upward but mostly horizontally pass the timing cuts. This is the primary component of the candidate pool with elevation angles below  $10^\circ$ . The “interesting” signal should be easier to extract in the region with higher elevation angle.

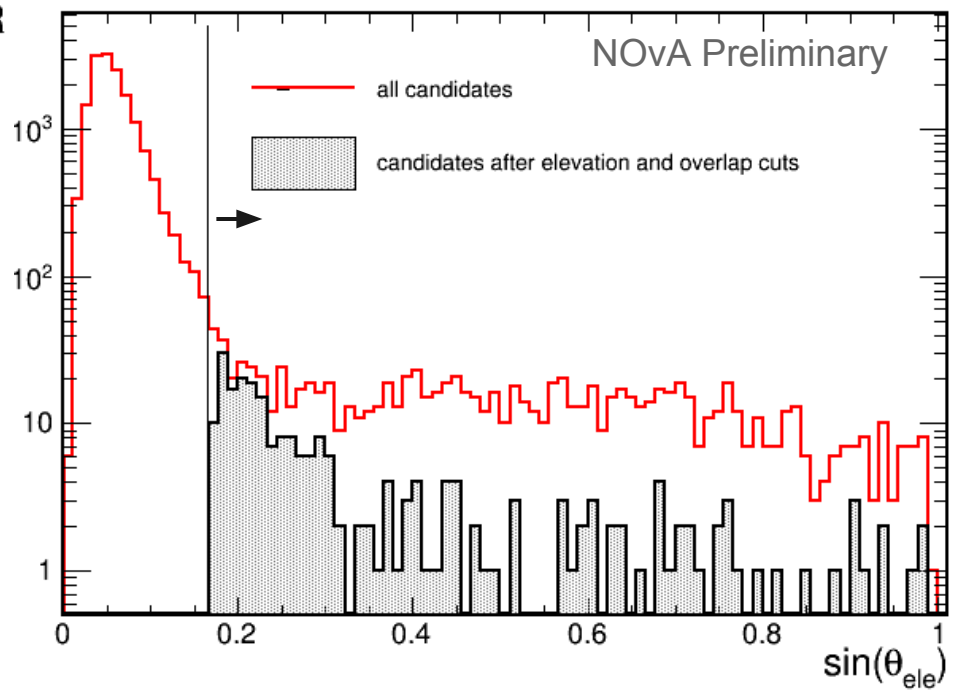
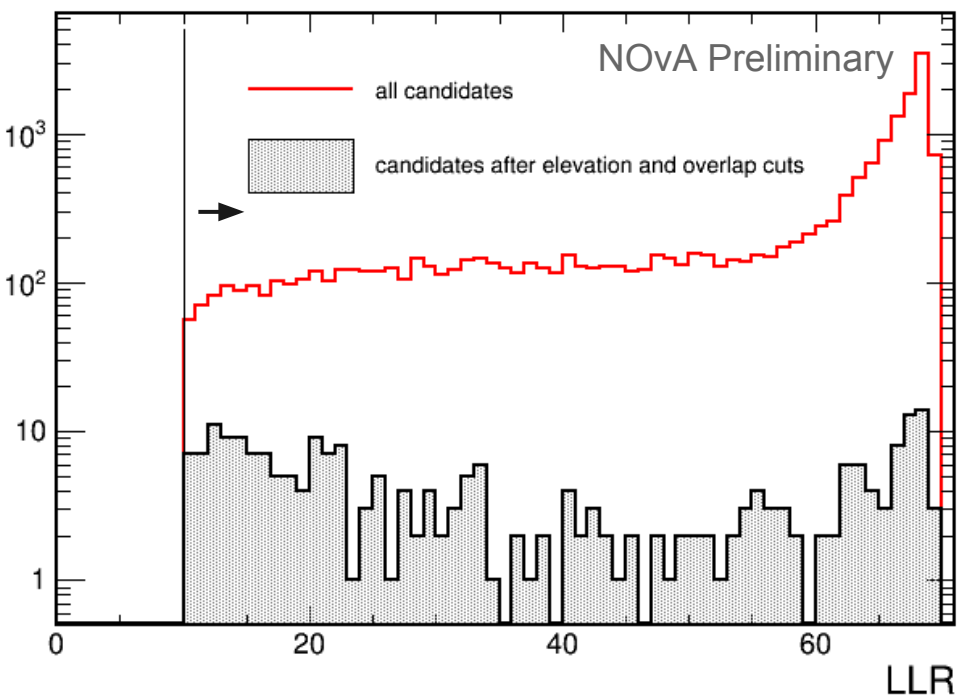


# Timing + Geometry

Combining the timing information from the LLR with a cut on elevation angle, I further reduced the pool to 1,051 tracks. Of those that remain, ~75% are not conclusively upward-going due to a misreconstruction in which two unrelated overlapping muon tracks make reconstruction tricky.

This class of misreconstruction was mostly eliminated using a simple “overlap” cut, leaving 255 tracks. This represents a reduction of ~3 billion to 1 in the number of muons, with the trigger and these cuts combined.





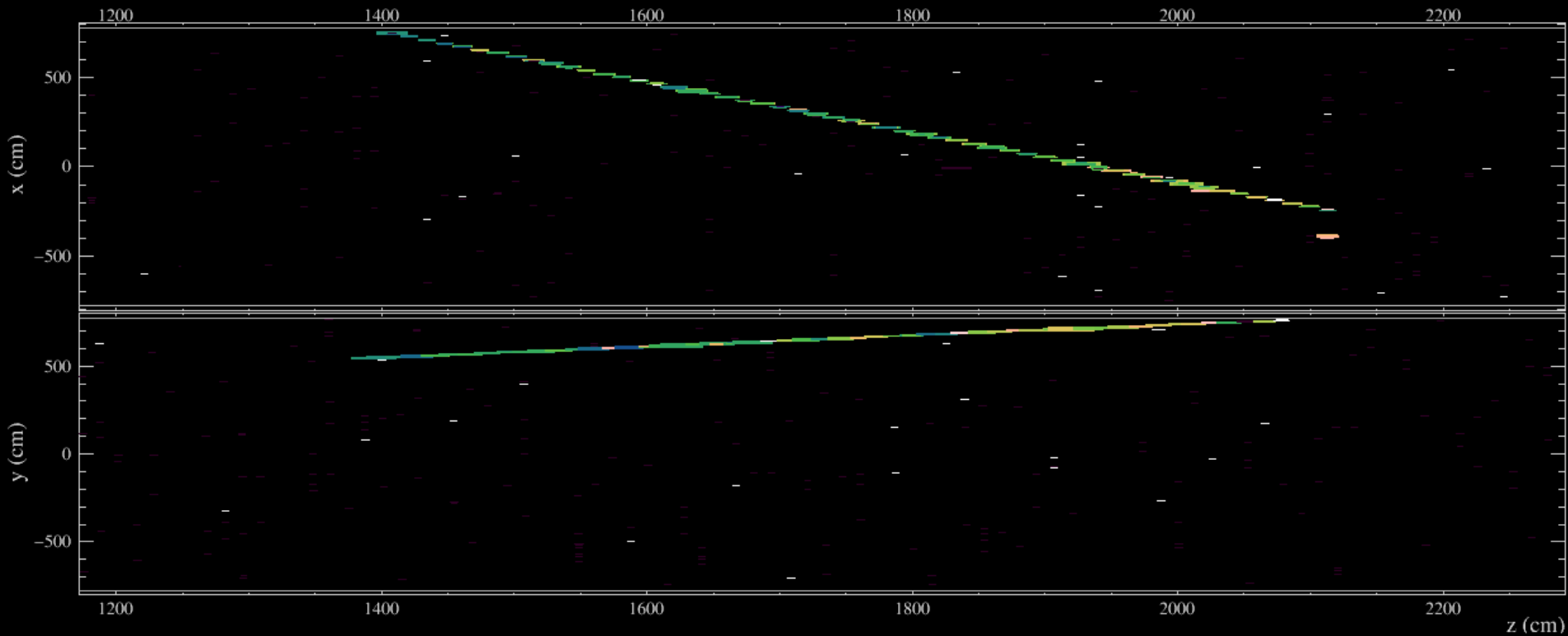
# Visual search

I visually examined each of these candidate events, categorizing them:

- Through-going: 105
- Stopping (seemingly upward-going): 75
- Overlapping muons: 34
- Up-scattered (within the detector) cosmics: 23
- Likely downward-going (backup): 1
- In-produced: 1
- Likely caused by a temporary timing miscalibration: 14\*
- Uncategorized (backup): 2

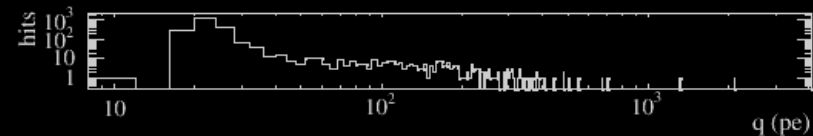
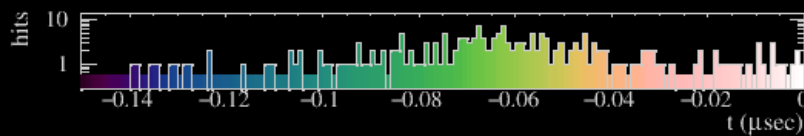
\* do not appear in the following plot

# Through-going

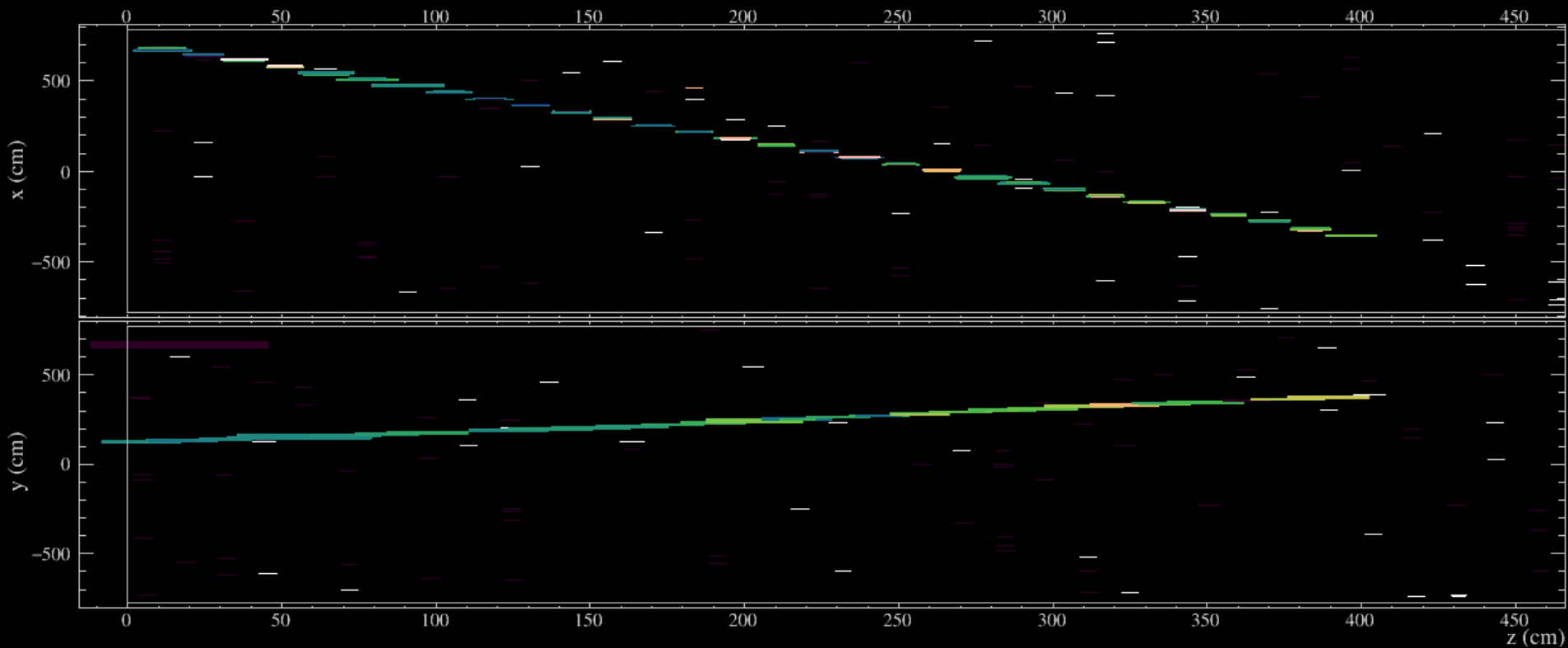


NOvA - FNAL E929

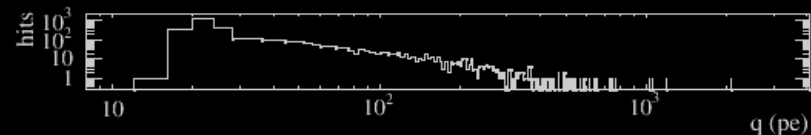
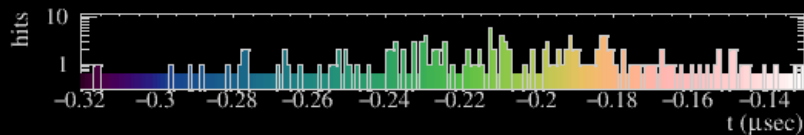
Run: 18631 / 5  
Event: 66953 / DDupmu  
UTC Sat Jan 10, 2015  
13:28:59.154943840



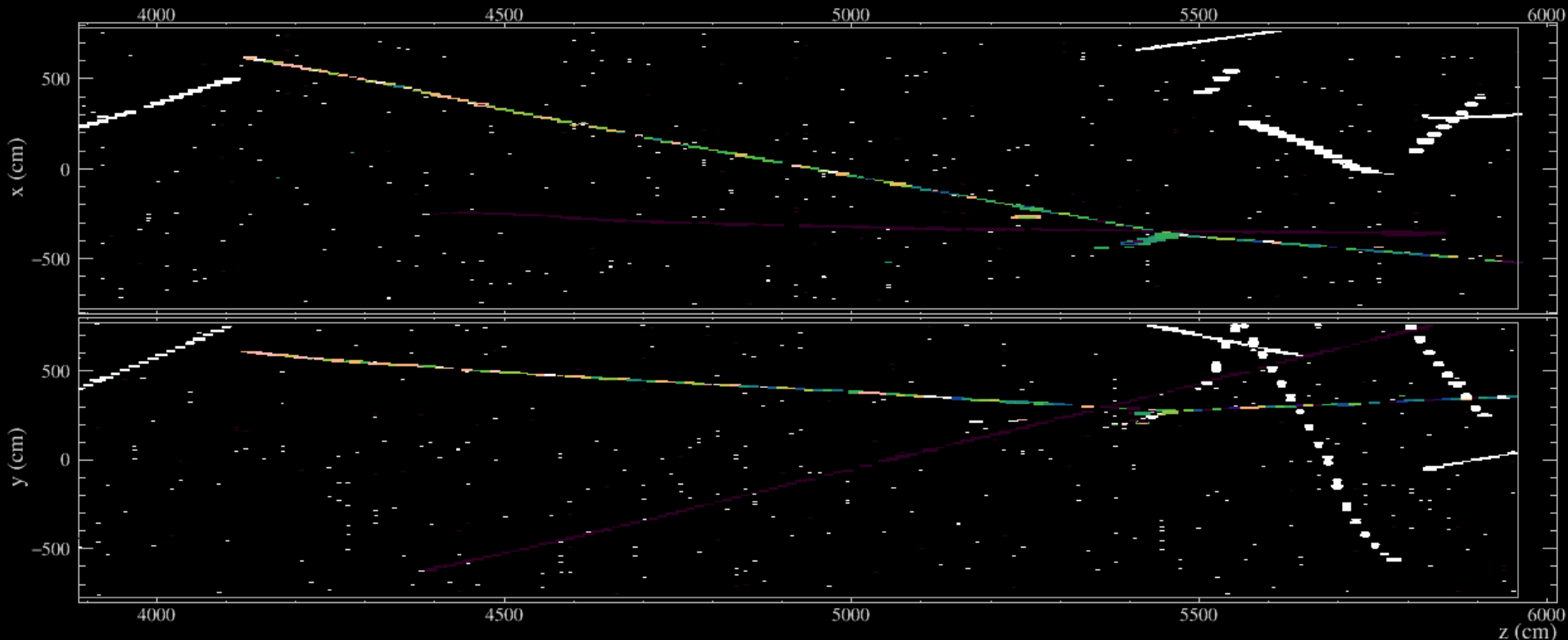
# Stopping



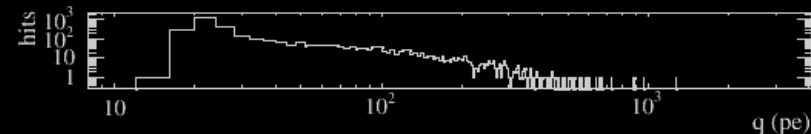
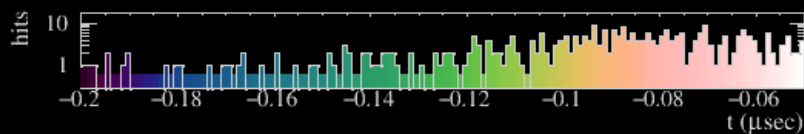
NOvA - FNAL E929  
Run: 19318 / 51  
Event: 742096 / DDupmu  
UTC Sat Apr 11, 2015  
04:43:58.417080288



# Up-scattered cosmics

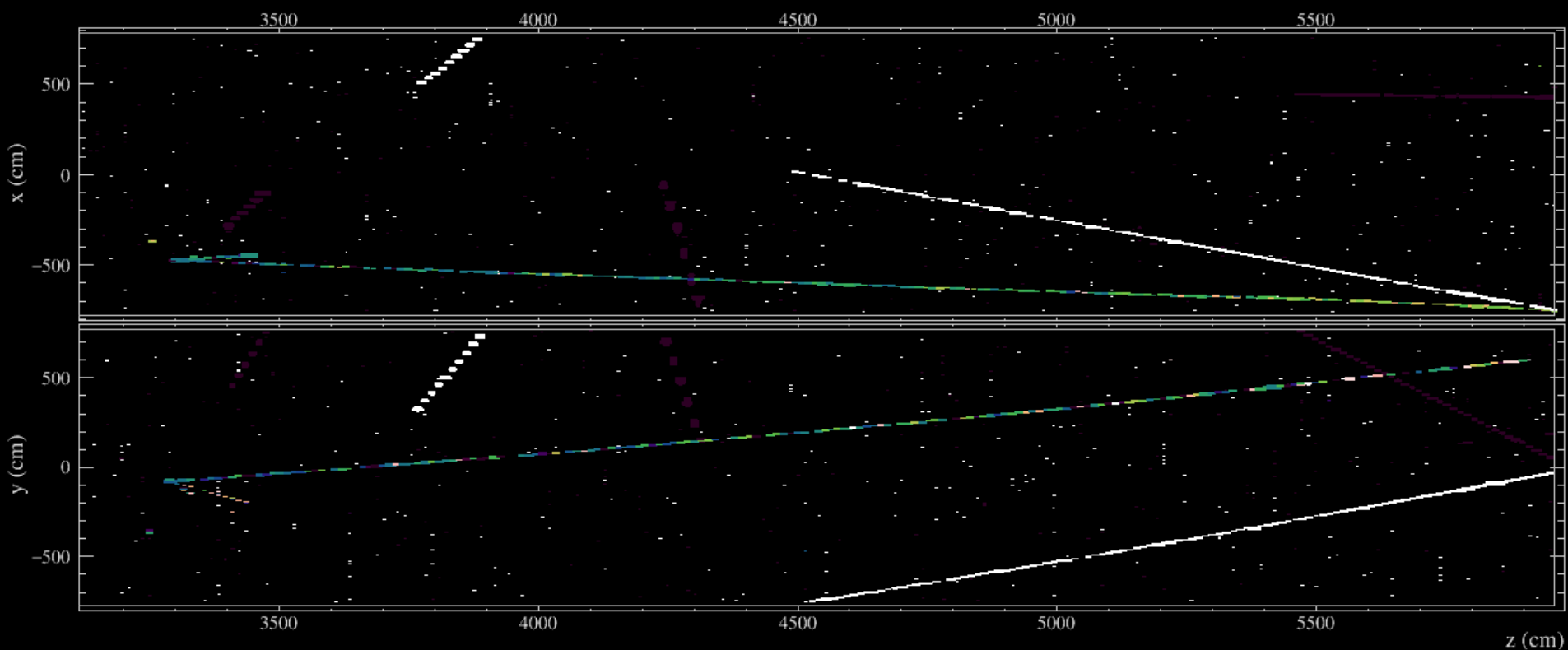


NOvA - FNAL E929  
Run: 18568 / 13  
Event: 174556 / DDupmu  
UTC Fri Jan 2, 2015  
04:04:17.055810780



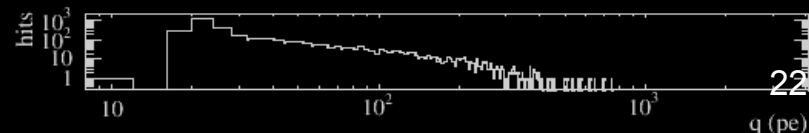
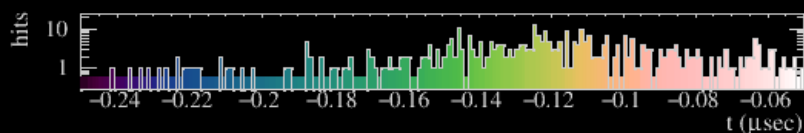
# In-produced

## Beam?

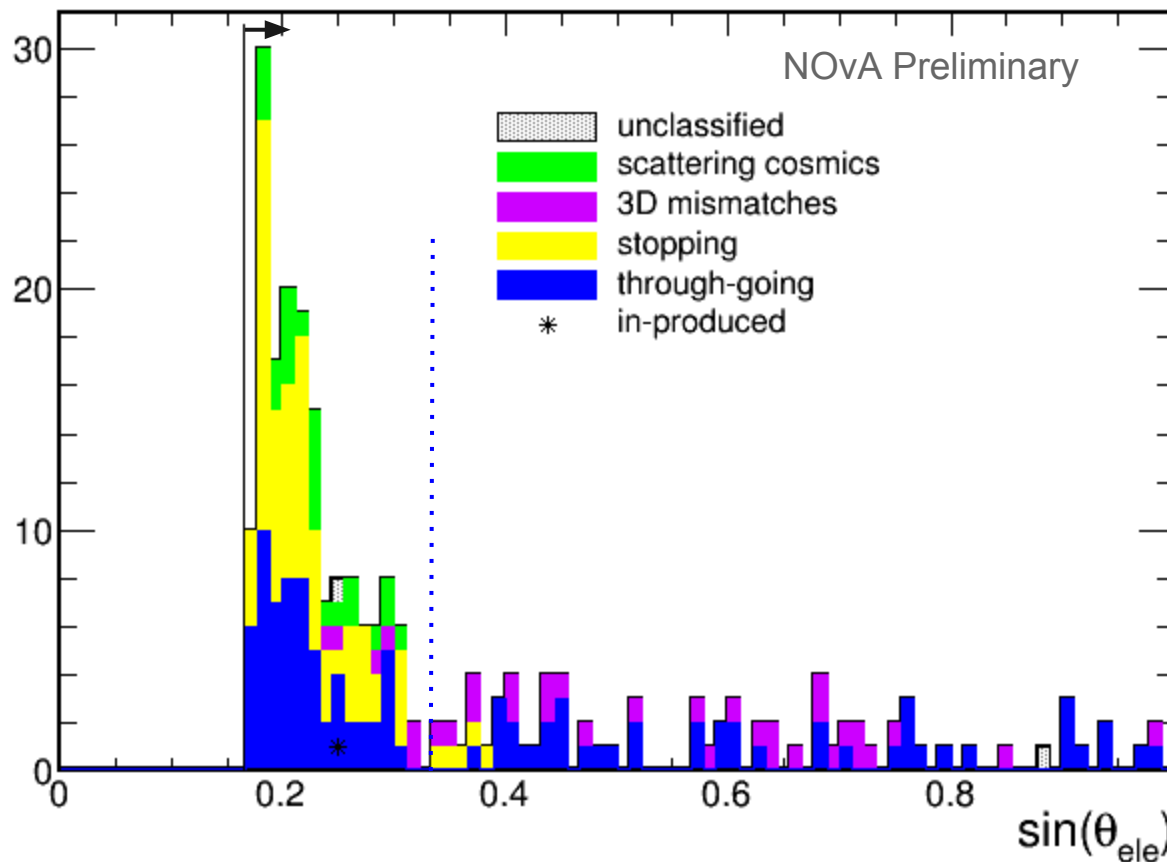


NOvA - FNAL E929

Run: 18585 / 19  
Event: 257078 / DDupmu  
UTC Sun Jan 4, 2015  
14:30:29.048571016



# Elevation angle revisited



Cutting at  $20^\circ$  ( $\sin(\theta) \sim 0.35$ ) should eliminate almost all the background from cosmes scattering in or around the detector.

\* 14 tracks that seem to be caused by a temporary timing miscalibration are not included

# Summary

- NOvA has a working upward-going muon trigger.
- We can reduce the cosmic ray muon background by nine orders of magnitude, using timing and geometry, to produce a sample sufficiently small to allow visual scanning.
- These techniques will allow an atmospheric neutrino oscillation study and potentially an indirect Dark Matter search.
- Stay tuned for an update on the contained-event sample.



# References

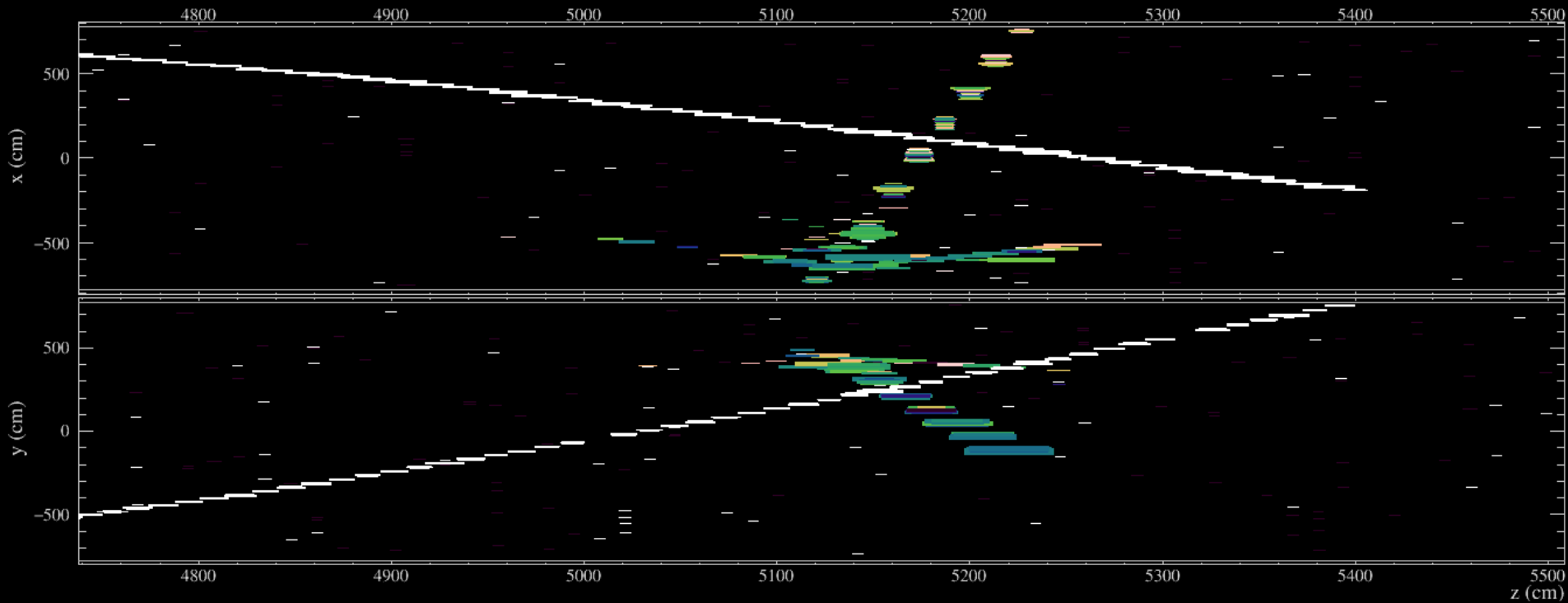
[1] The Super-Kamiokande Collaboration, “Measurement of the flux and zenith-angle distribution of upward through-going muons by Super-Kamiokande.” [arXiv:hep-ex/9812014](#)

The Super-Kamiokande Collaboration, “Neutrino-induced upward-going stopping muons in Super-Kamiokande.” [arXiv:hep-ex/9908049](#)

[2] C. L. Bennett et al., “Nine Year WMAP Observations: Final Maps and Results.” [arXiv:1212.5225](#) [[astro-ph.CO](#)]

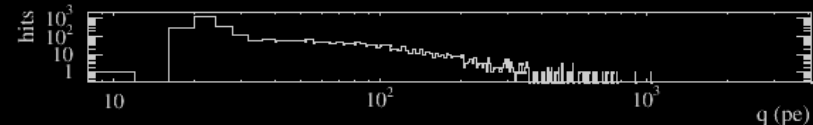
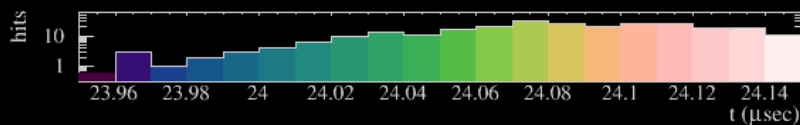
# Backup slides

# Likely downward-going

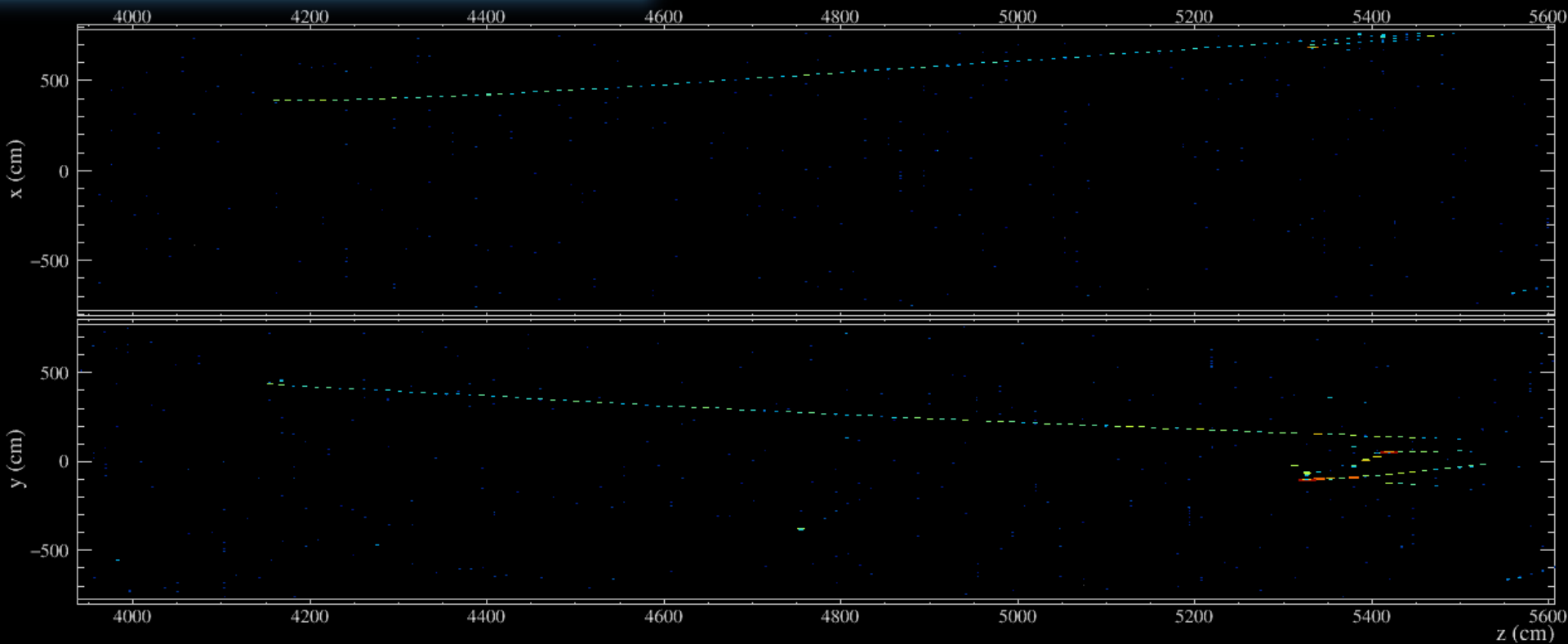


NOvA - FNAL E929

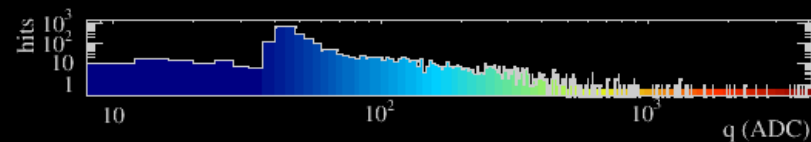
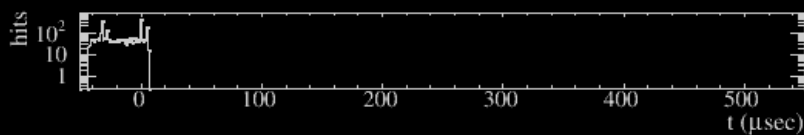
Run: 18593 / 1  
Event: 16699 / DDupmu  
UTC Mon Jan 5, 2015  
16:10:42.199410224



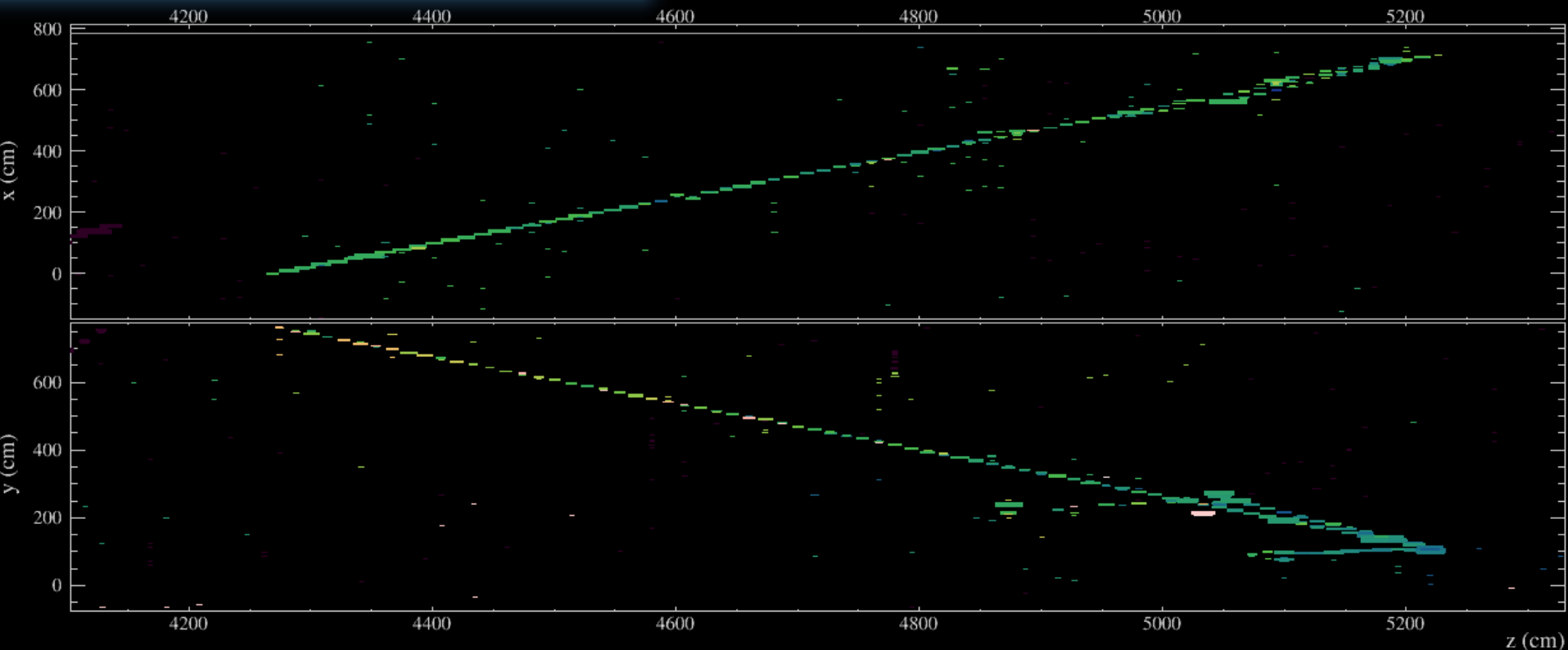
# Uncategorized candidate 1



**NOvA - FNAL E929**  
Run: 18640 / 12  
Event: 167520 / DDupmu  
UTC Sun Jan 11, 2015  
19:24:54.311743456



# Uncategorized candidate 2



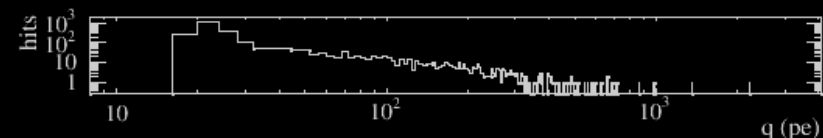
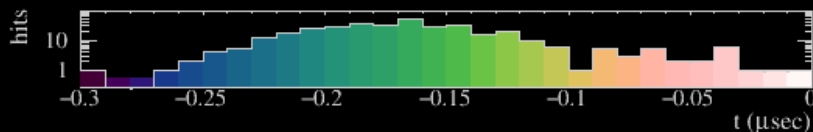
NOvA - FNAL E929

Run: 19316 / 38

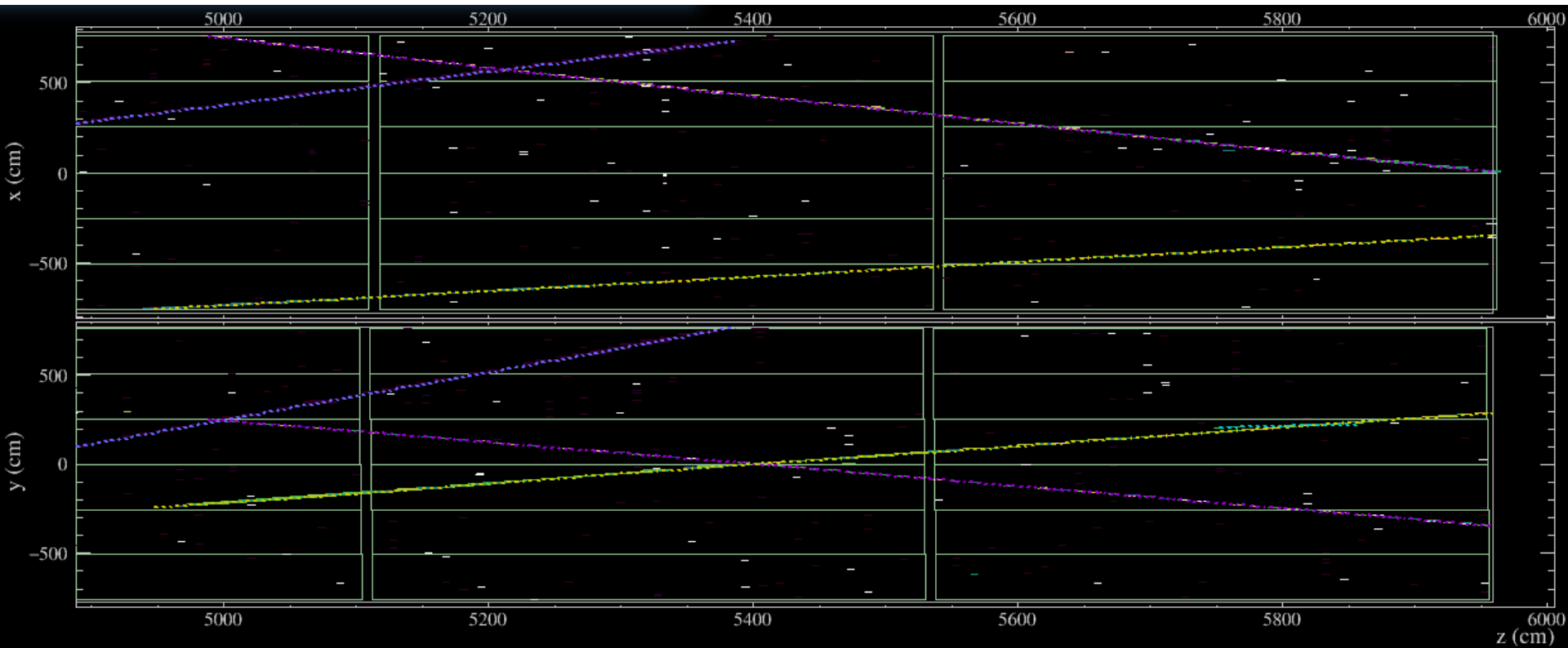
Event: 548344 / DDupmu

UTC Fri Apr 10, 2015

20:40:24.042806436

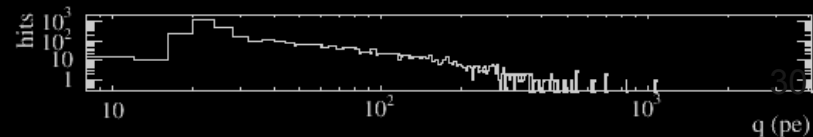
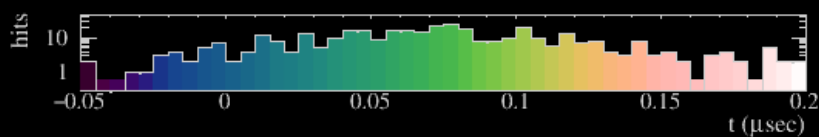


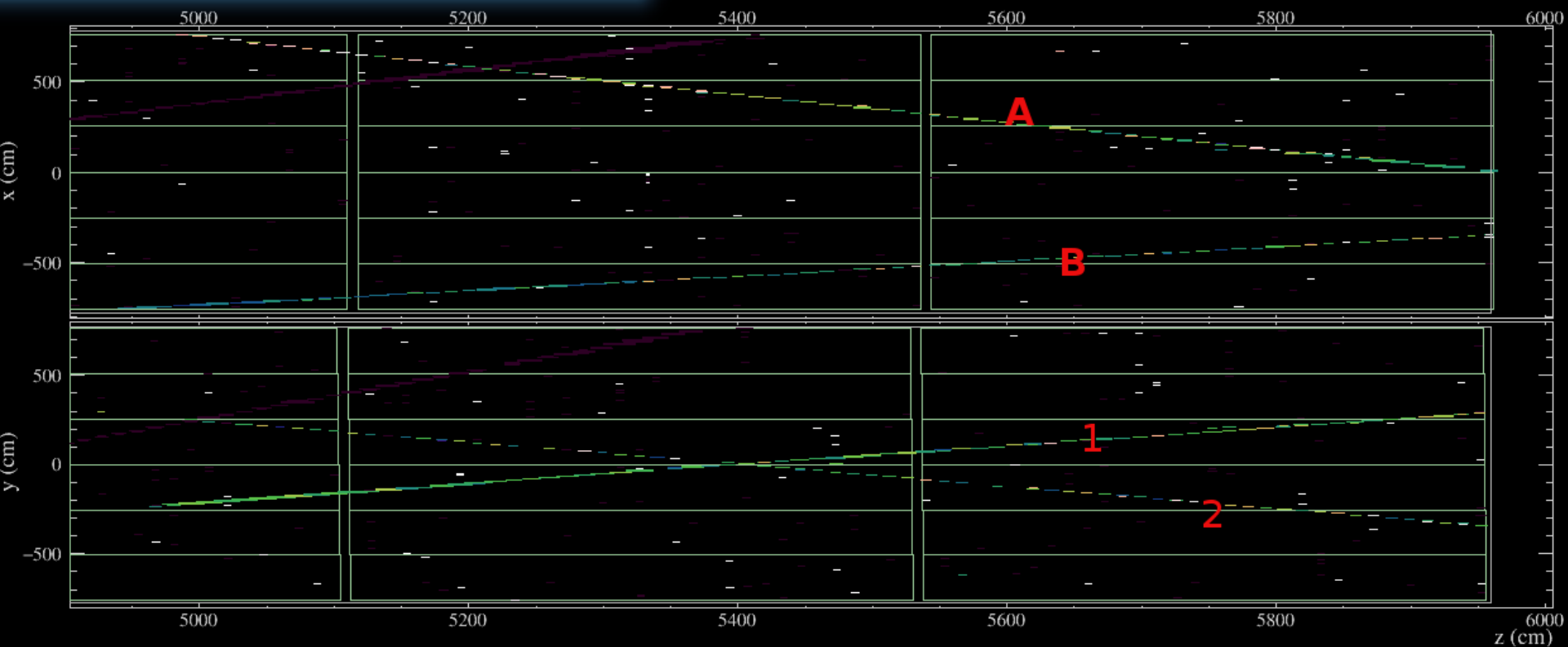
# 3D track mismatch... What is it?



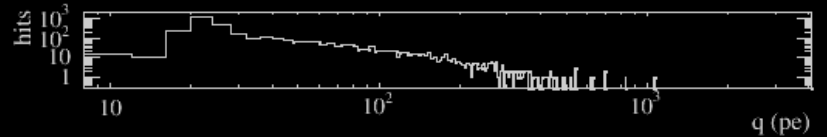
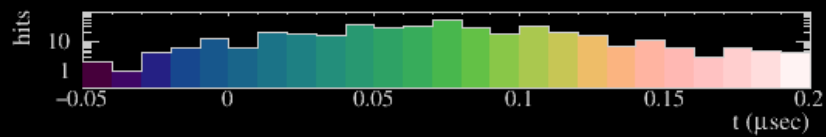
NOvA - FNAL E929

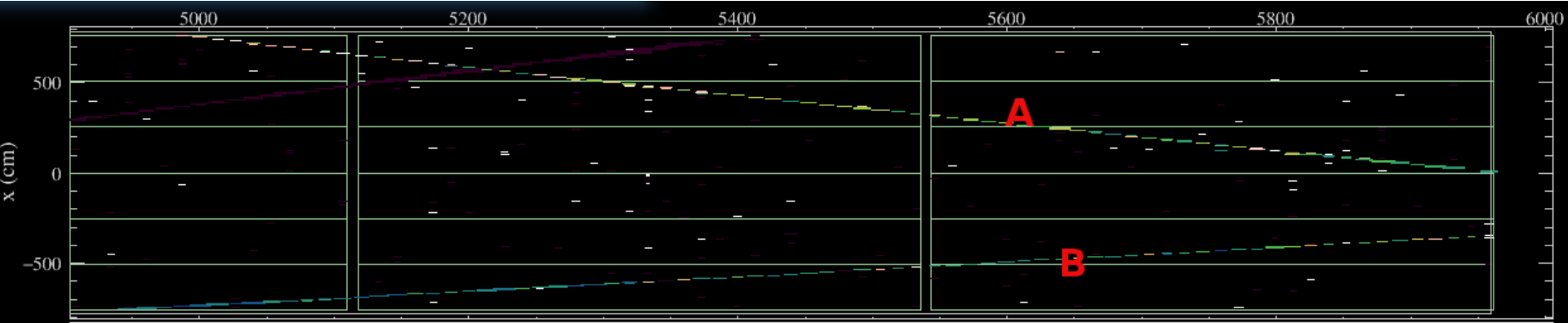
Run: 19625 / 40  
Event: 714542 / DDupmu  
UTC Thu May 21, 2015  
22:42:39.602736128





NOvA - FNAL E929  
 Run: 19625 / 40  
 Event: 714542 / DDupmu  
 UTC Thu May 21, 2015  
 22:42:39.602736128





Two meaningful ways to merge the four 2D tracks to get two 3D tracks:  $\{ A1, B2 \}$  or  $\{ A2, B1 \}$ .

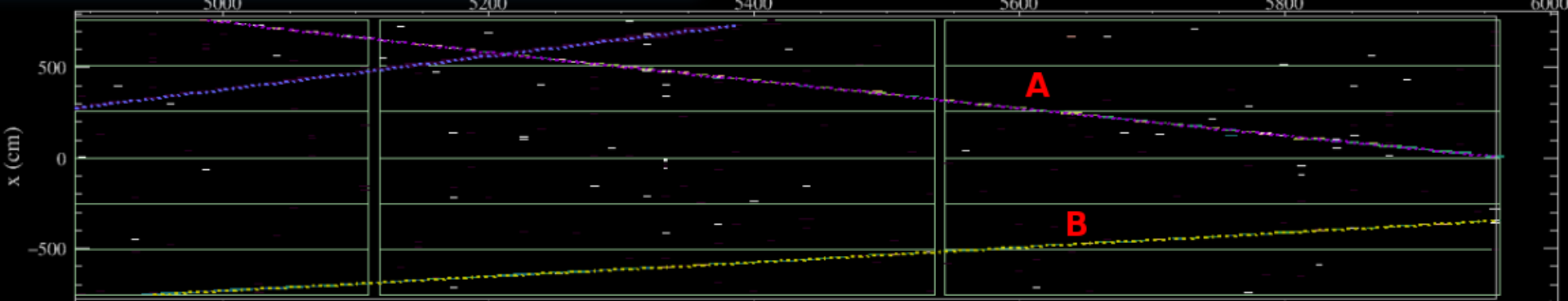
Four nonsensical ways:  $\{ A1, A2 \}$ ,  $\{ B1, B2 \}$ ,  $\{ A1, B1 \}$ ,  
 $\{ A2, B2 \}$

^^ I've never seen it produce nonsense. Impossible?

NOvA - FN  
 Run: 1962  
 Event: 7145  
 UTC Thu M  
 22:42:39.60

6000  
 z (cm)  
 q (pe)





KalmanTrackMerge has chosen { A2, B1 }

But is it correct?

NOvA - F  
Run: 1962  
Event: 714  
UTC Thu M  
22:42:39.6

6000  
z (cm)  
q (pc)

# The trickiness with timing

- The last step in the calibration of hit time for a single cell hit is to estimate the time of propagation of the scintillation light down the fiber to the readout.
- That is:  $\text{real\_time} = \text{apparent\_time} - (\text{distance\_from\_readout} / \text{speed\_of\_propagation})$
- See docdb 12570, page 19, item 1

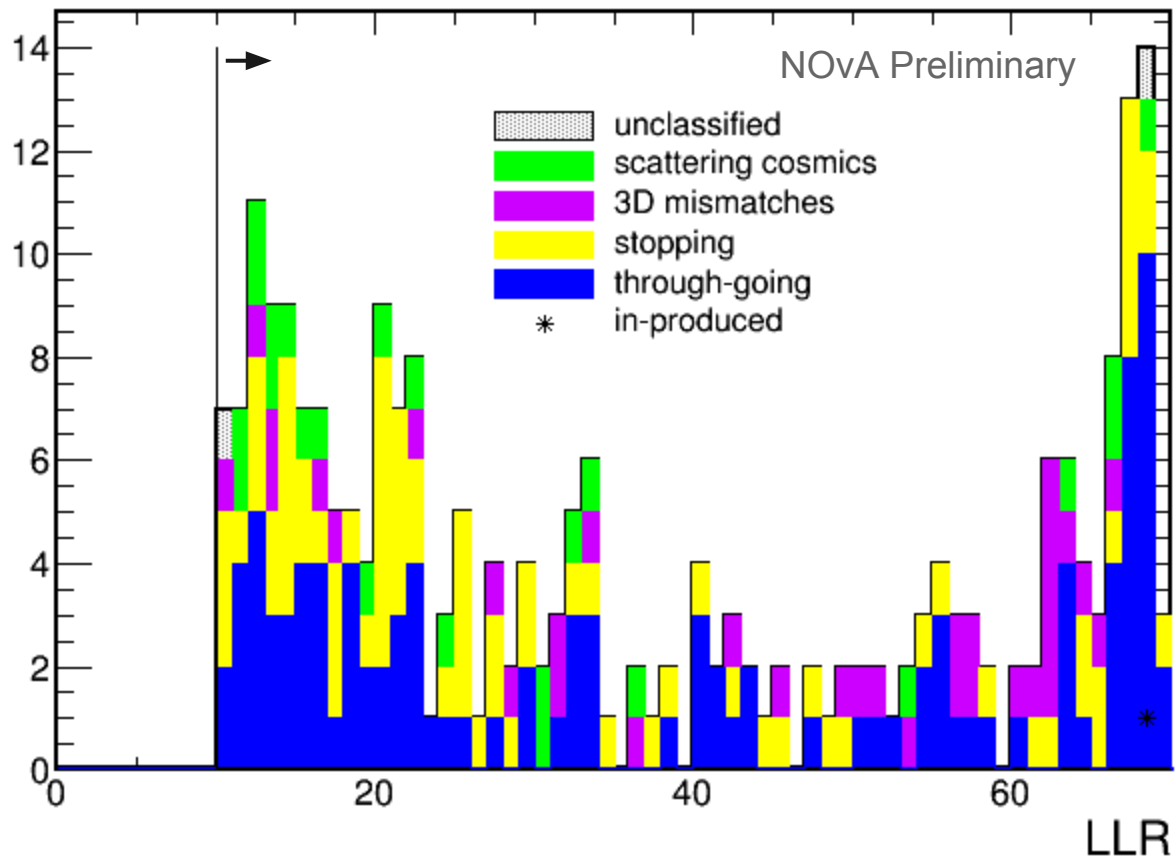
# The trickiness with timing

- The distance to the readout for a hit can only come from 3D track information, so the reconstructed times of the hits in these tracks depend on which pairing the track merger makes.
- The maximum size of this effect is 16 meters (detector width/height) / 15.3 cm/ns =  $\sim 100$  ns. Most cosmic muons travel through the detector in less than this much time.
- This is why these tracks are the dominant misreconstruction in the UpMu sample after cuts have eliminated almost all cosmics.

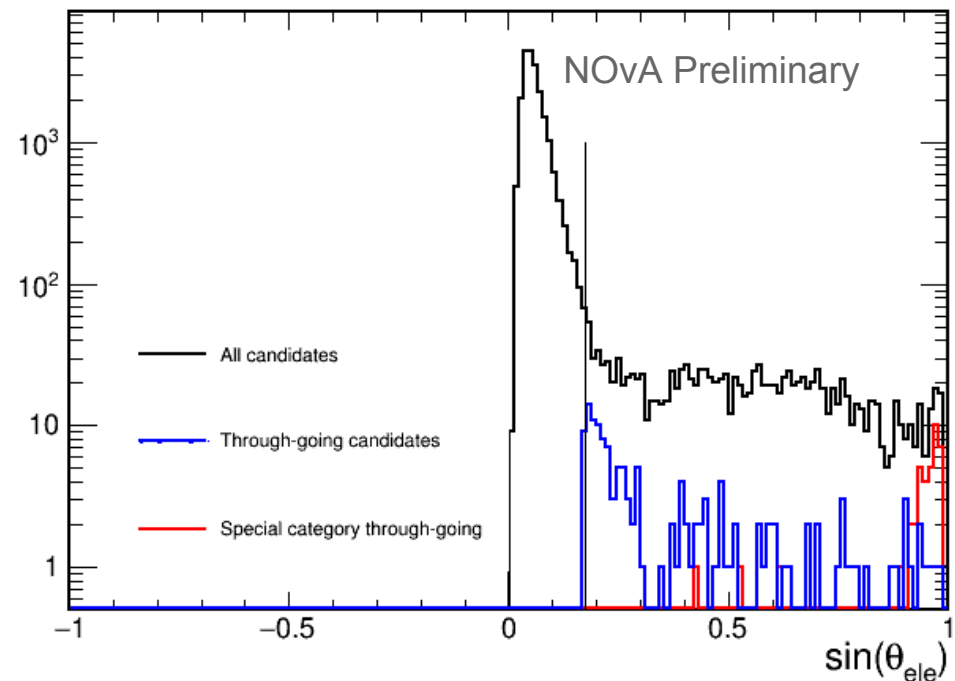
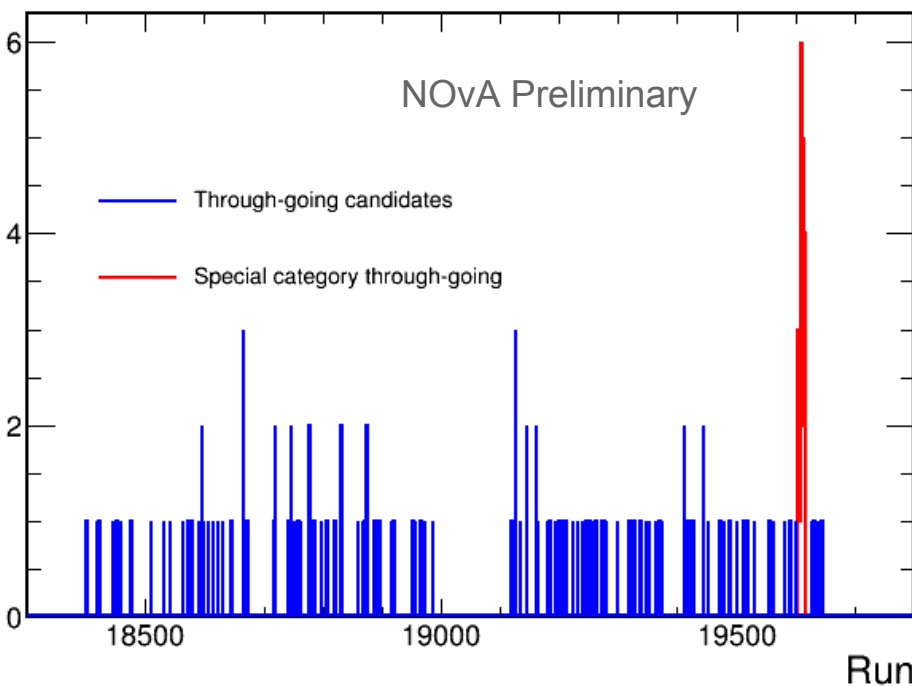
# What are we doing about it?

- For the ddupmu sample, I designed a slice-level cut that eliminates most of this contamination without losing much signal efficiency:
  - Pass all slices with 1 or  $> 3$  tracks (in order not to cut contained-vertex events)
  - For each track in a slice with 2 or 3 tracks, compute the differences in  $Z$  between the candidate track's start and end and that of the other track. If *both* the start and end differ by less than 5 plane widths, call this a match. Cut slices with more than 1 match.
  - Count the number of tracks in the slice with length  $> 8\text{m}$ . Cut slices with more than 1 such track.
- For the ddcontained sample, the track containment requirement should eliminate almost all of these events.

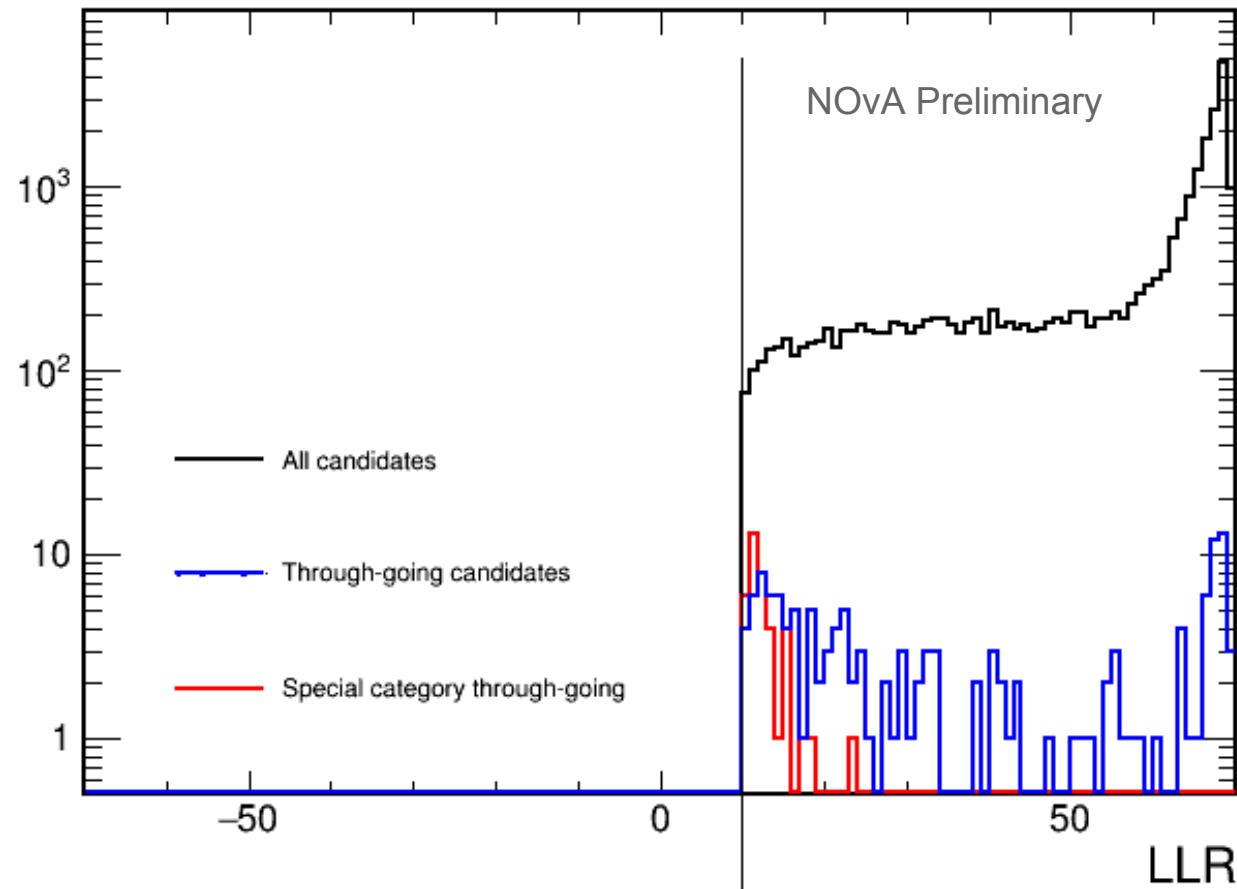
# LLR revisited



# Separate “special” through-going sample



# Separate “special” through-going sample



- High rate per run
- Highly vertical
- Low LLR compared to other candidates
- (not pictured) contained in Diblock 14