Contribution of Dark Matter annihilations to the low-redshift metagalactic ionization rate

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DPF2015, Ann Arbor, MI

06 August 2015

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1 Introduction

2 Connection between Lyman- α forest and DM annihilation

8 Non-thermal p-wave annihilation

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- \checkmark Direct detection

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- ✓ Production at colliders
- ✓ Direct detection
- ✓ Indirect detection

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 \checkmark



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This talk will be about an indirect² detection method.

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²No, it is not a footnote. I really mean squared.

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- Energy injection due to DM annihilation/decay can potentially modify the thermal history of the Universe.
- ✓ At high redshift ($z \sim 1000$) CMB anisotropies put strong constraints on the WIMP mass and annihilation cross section. But it is cosmic variance limited.
- ✓ Looking forward to 21-cm cosmology ($50 \le z \le 200$) and CMB spectral distortions ($10^3 \le z \le 10^5$)!
- At the low-redshift end, Lyman-α forest has dirty astrophysics. But the background characterization is improving (Haardt & Madau, 2012).

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Inject UV/X-ray (10-100 eV) photons directly into the IGM? Inject low energy e^{-}/e^{+} and let them collisionally ionize the IGM? Inject relativistic e^{-}/e^{+} and let them upscatter the CMB to the UV/X-ray band! 1000.000 otal Bremsstrahlung (whole galaxy) 100.000 (inner aalaxy) ion Time Scale 10.000 τ [Gyrs] 1.000 0.100 0.010 0.001 10^{2} 10⁰ 10¹ 10^{3} 10^{4} 10⁵ Energy [MeV]

✓ Assume that thermally frozen-out WIMPs, $M_{\chi} = O(100)$ GeV, annihilate to pairs of $b\bar{b}$, which then produce showers of $\sim O(1)$ GeV e^{-}/e^{+} .

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- Resulting UV and X-ray photons can potentially contribute to the ionization state of the IGM.

Injection spectrum

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Injection spectrum

 e^{-}/e^{+} multiplicity (green) peaks at $\sim 0.01 M_{\chi}$ in hadronic annihilations channels and $\sim 0.1 M_{\chi}$ in leptonic channels (Cirelli et al. 2012).



DM annihilations

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 \checkmark DM annihilation emissivity is given by,

$$j(\vec{r}, z, E_e) \equiv E_e^2 \frac{\mathrm{d}N}{\mathrm{d}V \,\mathrm{d}t \,\mathrm{d}E_e} (\vec{r}, z, E_e) = \times E_e^2 \frac{\mathrm{d}N_{ann}}{\mathrm{d}E_e} \times \langle \sigma v \rangle_{\vec{v}} \times n^2(\vec{r}, z)$$
$$= \frac{\langle \sigma v \rangle_{\vec{v}}}{M_\chi^2} \times E_e^2 \frac{\mathrm{d}N_{ann}}{\mathrm{d}E_e} \times \rho^2(\vec{r}, z).$$

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✓ After DM particles start clustering the spatially averaged signal is averaged over a collection of halos,

$$\langle j(z, E_e) \rangle_{\vec{r}} \approx \langle j \rangle_{\text{halos}}(z, E_e) = \int \mathrm{d}M \, \frac{\mathrm{d}N}{\mathrm{d}M} \int_{V_{vir}} \mathrm{d}^3 \vec{r} \, j(r, z, E_e).$$

Using the CMB as a calorimeter for annihilation products

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Using the CMB as a calorimeter for annihilation products

The UV/X-ray photon emissivity is collected as the integral ICS power radiated by the e^-/e^+ .

$$j_{\gamma}(E_{\gamma},z) = \int_{m_e}^{M_{DM}} dE_e \frac{dP_{ic}}{dE_e}(E_{\gamma},E_e,z) \Psi_e^{igm}(E_e,z)$$

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$$10^{0} - 0.10 \text{ GeV}, z = 0.1$$

$$10^{-5} - 1.6 \text{ GeV}, z = 3.1$$

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The UV/X-ray emissivity can be projected on the sky to get the intensity

$$I(E_{\gamma},z) = \frac{c}{4\pi} \int_{z}^{\infty} \mathrm{d}z' \, \frac{\mathrm{d}t}{\mathrm{d}z'} \frac{(1+z)^{3}}{(1+z')^{3}} j_{\gamma}(E_{\gamma}',z') e^{-\tau_{\mathrm{eff}}(E_{\gamma}',z,z')}.$$

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Summary of requirements

- ✓ The mass of the WIMP should be low ($M_{\chi} \approx 10 100$ GeV).
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For the DM enthusiast: what can we do to make it work?

- Mess up with the mean concentration-mass relation or level of substructure around the scale radius?
- ✓ Change the production history of the WIMP?

✓ The relic abundance of a thermally produced WIMP is set by the annihilation cross section.

$$\Omega_{CDM} \approx 0.25 \left(\frac{3 \times 10^{-26} \text{ cm}^3/\text{s}}{\langle \sigma v \rangle} \right)$$

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✓ Similarly gamma-rays from stacked dwarf galaxies constrain thermal production to $M_{\chi} \gtrsim 100$ GeV (Fermi-LAT and DES collaborations, 2015).

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- ✓ The reheat of the Universe follows the freeze-out ($T_r < T_f \sim 100$ GeV), but should precede the BBN ($T_{BBN} \sim 1 \text{ MeV} < T_r$).
- However boosting the s-wave annihilation cross section violates the constraints!

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✓ In constrast p-wave annihilation

$$\langle \sigma v
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✓ Given the DM velocity variance during the freeze-out, $\langle v_f^2 \rangle \approx 0.1$, and the observed velocity variance in galaxy clusters, $\langle v_c^2 \rangle \approx \times 10^{-5}$, a p-wave velocity suppression of 10^{-4} is more than compensated by a non-thermal production boost of $T_f/T_r \sim 10^5$.

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- ✓ s-wave annihilations of thermal WIMPs cannot contribute to the photoionization of the IGM at low-redshift.
- p-wave annihilations of non-thermally produced DM can do better, avoiding CMB anisotropy constraints.
- This scenario was initially motivated by the gamma-ray excess in the inner galaxy (Goodenough and Hooper, 2009).
- ✓ However this is no longer the case if DM is non-thermally produced and is predominantly annihilating in the p-wave. (See Stephen's talk on Tuesday for the other reason.)