

# *EXTRA DIMENSION* VERSUS *SUPERSYMMETRY* AT THE LHC

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[Work done in collaboration with Kirtiman Ghosh and Durmus Karabacak **JHEP**  
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# Overview

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# minimal Universal Extra Dimension

In minimal UED  $\Rightarrow$  mUED, 5D space-time is  $M^4 \times S_1/Z_2$

- There are **KK-excitations** of every SM particles

KK-parity conserved



- Can only be pair produced
- Gives Dark Matter

$$m_n = \frac{n}{R}$$



Radiative corrections  
breaks mass  
degeneracy.



Spectra depends on the  
cut-off scale  $\Lambda$  and  $R^{-1}$

- $\Lambda$  can not be too large, otherwise gauge couplings blow off.
- Spectra of KK-particles is pretty degenerate at each level

**With only level-1 KK excitations, similar to compressed SUSY**

# Extra Dimension vs SUSY at the LHC

However, most SUSY breaking scenarios do not give such a compressed spectra

- Pair production of SUSY particles and their decay gives high  $p_T$  jets plus missing transverse energy.
  - Classic signals of SUSY
  - Will distinguish SUSY from UED

However, LHC Higgs data do not agree with mUED at the  $1\sigma$  level

Need to go beyond mUED  $\implies$  non-minimal UED  $\implies$  nmUED

# Model: nmUED

$$S = S_{bulk} + S_{BLKT}$$

$$S_{bulk} = \int d^4x \int_{-L}^L dy \left[ \sum_{\mathcal{A}}^{G,W,B} -\frac{1}{4} \mathcal{A}_{MN} \mathcal{A}^{MN} + \sum_{\Psi}^{Q,U,D,L,E} i \bar{\Psi} \overleftrightarrow{D}_M \Gamma^M \Psi - M_{\Psi} \bar{\Psi} \Psi \right]$$

$$S_{bdry} = \int d^4x \int_{-L}^L dy \left( \sum_{\mathcal{A}}^{G,W,B} -\frac{r_{\mathcal{A}}}{4} \mathcal{A}_{\mu\nu} \mathcal{A}^{\mu\nu} + \sum_{\Psi=Q,L} i r_{\Psi} \bar{\Psi}_L D_{\mu} \gamma^{\mu} \Psi_L \right. \\ \left. + \sum_{\Psi=U,D,E} i r_{\Psi} \bar{\Psi}_R D_{\mu} \gamma^{\mu} \Psi_R \right) \times [\delta(y-L) + \delta(y+L)].$$

## Parameters:

- Coefficients of the BLKT terms
  - $r_A \implies r_g, r_W, r_B$
  - $r_{\psi} \implies r_U, r_D, r_E$
- Also involve  $M_{\psi} = \mu \theta(y) \implies$  5D fermion bulk mass.

Masses of the KK-gauge bosons and fermions are determined by solving transcendental equations

# Constraints on the parameters

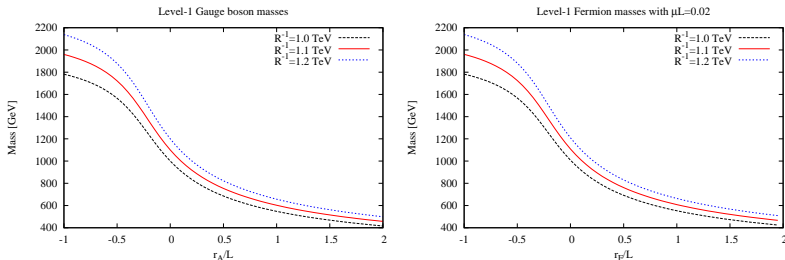
- $\frac{r_\Psi}{L} > \frac{\exp^{-2\mu L} - 1}{2\mu L}$  to avoid ghosts and/or tachyons in the fermion sector.
- $r_A/L > -1$  to avoid ghosts and/or tachyons in the gauge sector.
- The bounds on the parameters are also obtained from the low-energy observables. [T. Flacke, K. Kong and S. C. Park, JHEP **1305**, 111 (2013)]
  - KK-parity conserving interactions,  $\mathcal{L}_{002n}$
  - $Z^{(2)}$  contribute to 4-fermi interactions
    - $r_A > 0.5L$  for  $\mu L = -0.1$  and fixed  $R^{-1}$
  - However, for  $0 > \mu L > -0.03$  and  $R^{-1} \approx 850$  GeV,  $g_{200}$  will be small and  $M_{Z^{2n}}$  will be heavy enough to escape this bound
- EW precision test and the collider searches are insensitive to small values of  $\mu$ , for example  $\mu L = -0.02$

# Phenomenology

We discuss the implications of nmUED in the context of the Higgs data and multijets plus  $E_T$  searches at the 8 TeV LHC

- Parameters of the model:
  - $r_\psi, \Psi = Q, U, D, L, E$
  - $r_A, A = G, W, B$
  - $\mu \implies$  Bulk fermion mass term
- Take universal boundary parameters for all quarks and leptons  
 $\implies r_F$
- For gauge sector, we choose  $r_g \neq r_W = r_B$
- Also choose  $\mu L = -0.02$

# Mass spectra for level-1 KK-excitations



**Figure :** Level-1 gauge boson (left panel) and fermion (right panel) as a function of  $r/L$  for three different values of  $R^{-1}$ . For the level-1 fermion masses in the right panel, we consider  $\mu = -0.02L$ .

## NOTE:

- Masses of KK-excitations are very sensitive to the BLKT parameters  $\implies$  Large splitting even at the tree level
- Both level-1 KK-fermions and gauge bosons masses increase if we decrease  $r/L$ .



# Higgs Phenomenology

- $m_H = 125 \text{ GeV} \implies \lambda = 0.129$  in the SM
- $\lambda \rightarrow 0$  at  $Q = 10^{11} \text{ GeV} \implies$  Vacuum instability
- For mUED,  $\lambda$  evolves much faster because of KK-excitations
  - $\lambda \rightarrow 0$  at  $4\text{-}6 R^{-1}$
  - For subsequent analysis, we take  $\Lambda = 5R^{-1}$
- Higgs production measured in different channels at the LHC

$$\mu_i = \frac{(\sigma \times \text{BR})_i}{(\sigma \times \text{BR})_i^{\text{SM}}}$$

- $H \rightarrow gg$ : Only KK-tower of **top** contributes
- $H \rightarrow \gamma\gamma$ : KK-towers of **top and W** both contribute

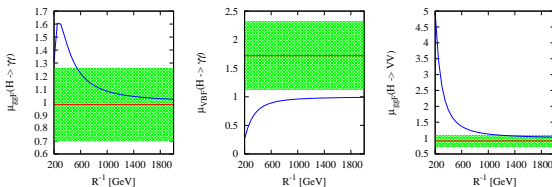
# Combined best-fit values for $\mu_i$

[G. Belanger, B. Dumont, U. Ellwanger, J. F. Gunion and S. Kraml, PRD **88**, 075008 (2013)]

$\gamma\gamma$ decay channel		$VV$ decay channel	
$\hat{\mu}^{ggF}$	$\hat{\mu}^{VBF}$	$\hat{\mu}^{ggF}$	$\hat{\mu}^{VBF}$
$0.98 \pm 0.28$	$1.72 \pm 0.59$	$0.91 \pm 0.16$	$1.01 \pm 0.49$

**Table :** Combined best-fit Higgs signal strengths for different Higgs production and decay modes.

Results for mUED: has only 2 parameters  $\Rightarrow \Lambda$  and  $R^{-1}$



**Figure :** Higgs signal strengths relative to the SM expectations:  $\mu_{ggF}^{H \rightarrow \gamma\gamma}$  (left panel),  $\mu_{VBF}^{H \rightarrow \gamma\gamma}$  (middle panel) and  $\mu_{ggF}^{H \rightarrow VV}$  (right panel), in the context of mUED scenario as a function of  $R^{-1}$ . The combined best-fit values (from Table 1) of the abovementioned Higgs signal strengths are also presented.

Can not fit data at  $1\sigma$  level (see middle panel)

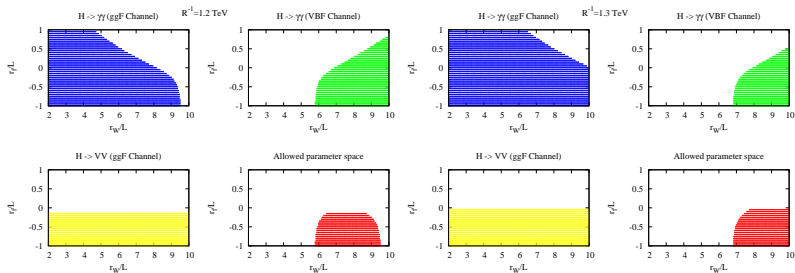
# Results for nmUED

In addition to  $R^{-1}$  and  $\Lambda$ , we have several additional parameters coming from the BLK terms

- $r_\psi, \Psi = Q, U, D, L, E$
- $r_A, A = G, W, B$
- $\mu \implies$  Bulk fermion mass term
- We choose cut-off scale  $\Lambda = 5R^{-1}$
- Universal  $r_F$
- But, for gauge sector, we choose  $r_g \neq r_W = r_B$
- Also choose  $\mu L = -0.02$

We have scanned the parameter space  $(r_W, r_F)$  to fit the Higgs data in different channels for 2 values of  $R^{-1}$  (1.2 and 1.3 TeV)

## Continued.



**Figure :** Scattered points in  $r_W/L$ - $r_F/L$  plane which are consistent with the combined best-fit results of  $\mu_{ggF}^{H \rightarrow \gamma\gamma}$ ,  $\mu_{VBF}^{H \rightarrow \gamma\gamma}$ ,  $\mu_{ggF}^{H \rightarrow VV}$  and all three together for  $R^{-1} = 1.2$  TeV (left panel) and  $R^{-1} = 1.3$  TeV (right panel).

All combined best-fit Higgs data is consistent with nmUED for large  $r_W$  and negative  $r_F$

# Collider Phenomenology at the LHC

- Multijets at high  $p_T$  plus large  $\cancel{E}_T$  signal
- Use the LHC data for SUSY search and if we can reproduce the exact limits for nmUED using the allowed parameter space
- We use 8 TeV LHC data with  $L = 20.3 \text{ fb}^{-1}$
- SUSY limits: For  $m_{\tilde{q}} = m_{\tilde{g}} < 1.7 \text{ TeV}$  is excluded from jets +  $\cancel{E}_T$  channel.
- Two questions:
  - Can we reproduce any SUSY signals from the level-1 KK particles using the allowed parameter space of nmUED ?
  - What limits we can put on  $q^{(1)}$  and  $g^{(1)}$  masses in nmUED ?

# nmUED Benchmark point for multijets analysis

Benchmark Point (BP)					
$R^{-1}$	$\Lambda R$	$\mu L$	$r_g/L$	$r_F/L$	$r_W/L$
1.2 TeV	5	-0.02	-0.05	-0.42	7.4
Masses in GeV					
$m_{Q(1)}$	$m_{L(1)}$	$m_{G(1)}$	$m_{W(1)\pm}$	$m_{Z(1)}$	$m_{\gamma(1)}$
1800	1800	1265	275	275	260

- Produce  $q^1 q^1$ ,  $g^1 g^1$ ,  $q^1 g^1$
- Decay these using the spectra for the benchmark point
- Apply the same cuts as ATLAS multijets SUSY searches

## Results: ATLAS cuts

Cuts	A (2-jets)		B (3-jets)		C (4-jets)		D	E (6-jets)		
	L	M	M	T	M	T	(5-jets)	L	M	T
$\cancel{E}_T > [\text{GeV}]$	160									
$p_T^1 > [\text{GeV}]$	130									
$p_T^2 > [\text{GeV}]$	60									
$p_T^3 > [\text{GeV}]$	-		60		60		60		60	
$p_T^4 > [\text{GeV}]$	-		-		60		60		60	
$p_T^5 > [\text{GeV}]$	-		-		-		60		60	
$p_T^6 > [\text{GeV}]$	-		-		-		-		60	
$\Delta\phi(j_i, \vec{\cancel{E}}_T)_{min} >$	0.4 {i=1,2,3 if $p_T^3 > 40 \text{ GeV}$ }				0.4 {i=1,2,3}, 0.2 $p_T^{j_i} > 40 \text{ GeV}$					
$\cancel{E}_T / M_{eff}(N_j) >$	0.2	-	0.3	0.4	0.25	0.25	0.2	0.15	0.2	0.25
$m_{eff}(incl.) [\text{TeV}]$	1.0	1.6	1.8	2.2	1.2	2.2	1.6	1.0	1.2	1.5
$\sigma_{BSM} [\text{fb}]$	66.07	2.52	0.73	0.33	4.00	0.12	0.77	4.55	1.41	0.41

**Table :** Cuts used by the ATLAS collaboration to define the signal regions.  $\Delta\phi(jet, \vec{\cancel{E}}_T)$  is the azimuthal separations between  $\vec{\cancel{E}}_T$  and the reconstructed jets.  $m_{eff}(N_j)$  is defined to be the scalar sum of the transverse momenta of the leading  $N$  jets together with  $\cancel{E}_T$ . However, for  $m_{eff}^{incl.}$ , the sum goes over all jets with  $p_T > 40$  GeV. Last column corresponds to the 95% C.L. observed upper limits on the non-SM contributions  $\sigma_{BSM}$ .

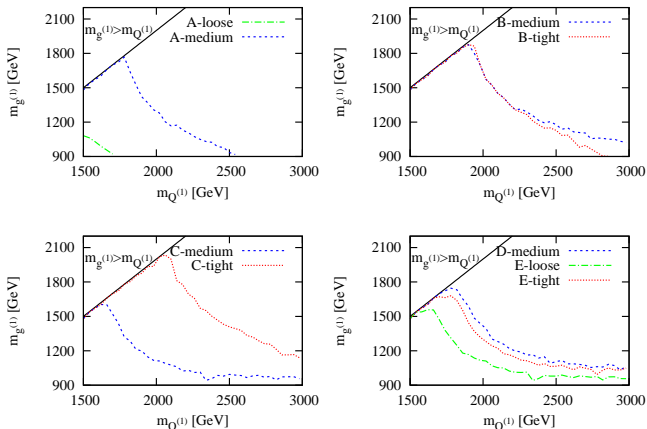
## Results: Cut-flow table.

Process	Supersymmetry		nmUED
	$\tilde{g}\tilde{g}$ one-step		$g^{(1)}g^{(1)}$ one-step
Point	$m_{\tilde{g}} = 1265$ GeV		$m_{g^{(1)}} = 1265$ GeV
	$m_{\tilde{\chi}_1^\pm} = 865$ GeV		$m_{\tilde{W}(1)\pm} = 865$ GeV
	$m_{\tilde{\chi}_1^0} = 465$ GeV		$m_{\tilde{\gamma}(1)} = 465$ GeV
Cuts (E-tight)	Absolute efficiency in %		
	ATLAS Appendix-C of [?]	Our Simulation	Our Simulation
0-lepton	63.5	66.1	57.3
$E_T > 160$ GeV	55.6	57.6	54.7
$p_T^{j1} > 130$ GeV	55.6	57.5	54.7
$p_T^{j2} > 60$ GeV	55.6	57.5	54.6
$p_T^{j3} > 60$ GeV	55.4	57.3	51.8
$p_T^{j4} > 60$ GeV	53.4	55.2	41.3
$p_T^{j5} > 60$ GeV	46.3	47.1	27.4
$p_T^{j6} > 60$ GeV	31.7	31.1	15.0
$\Delta\phi(j_i, E_T), i = 1, 2, 3$	26.5	26.1	12.2
$\Delta\phi(j, E_T), p_T^j > 40$ GeV	21.3	21.6	9.7
$E_T/m_{\text{eff}}(N_j) > 0.25$	12.0	12.7	4.7
$m_{\text{eff}}(\text{incl.}) > 1.5$ TeV	7.9	8.3	4.5

- Our simulation agrees very well with the ATLAS simulations
- nmUED signals are pretty close to the SUSY for similar cuts



# Results



**Figure :** The exclusion limits on  $m_{Q^{(1)}} - m_{g^{(1)}}$  plane from 8 TeV 20.3 inverse femtobarn integrated luminosity ATLAS data for different ATLAS defined signal regions. We have assumed fixed mass for the level-1 electroweak KK gauge bosons ( $m_W^{(1)\pm} = m_{Z^{(1)}} = 275$  GeV and  $m_{W\gamma^{(1)}} = 260$  GeV).

For  $m_{q^1} = m_{g^1}$ , the limit is 2.1 TeV.

# Conclusion

- Though mUED signals are very different from the SUSY signals at the LHC, nmUED signals are not.
- With switable choice of BLK terms, we can reproduce any multijets +  $\cancel{E}_T$  signal given by SUSY in nmUED.
- If enhancement in  $H \rightarrow \gamma\gamma$  persists in LHC Run 2, it is allowed in nmUED, but not in mUED.
- nmUED with the assumption  $m_{q1} = m_{g1}$ , 8 TeV LHC limit is  $\sim 2.1$  TeV.
- Since nmUED can reproduce any SUSY signals, the production of level-2 KK excitations will be the key to distinguish between extra-dimension and supersymmetry if any signal is seen.