

# Neutrino Nucleon Deep Inelastic Scattering

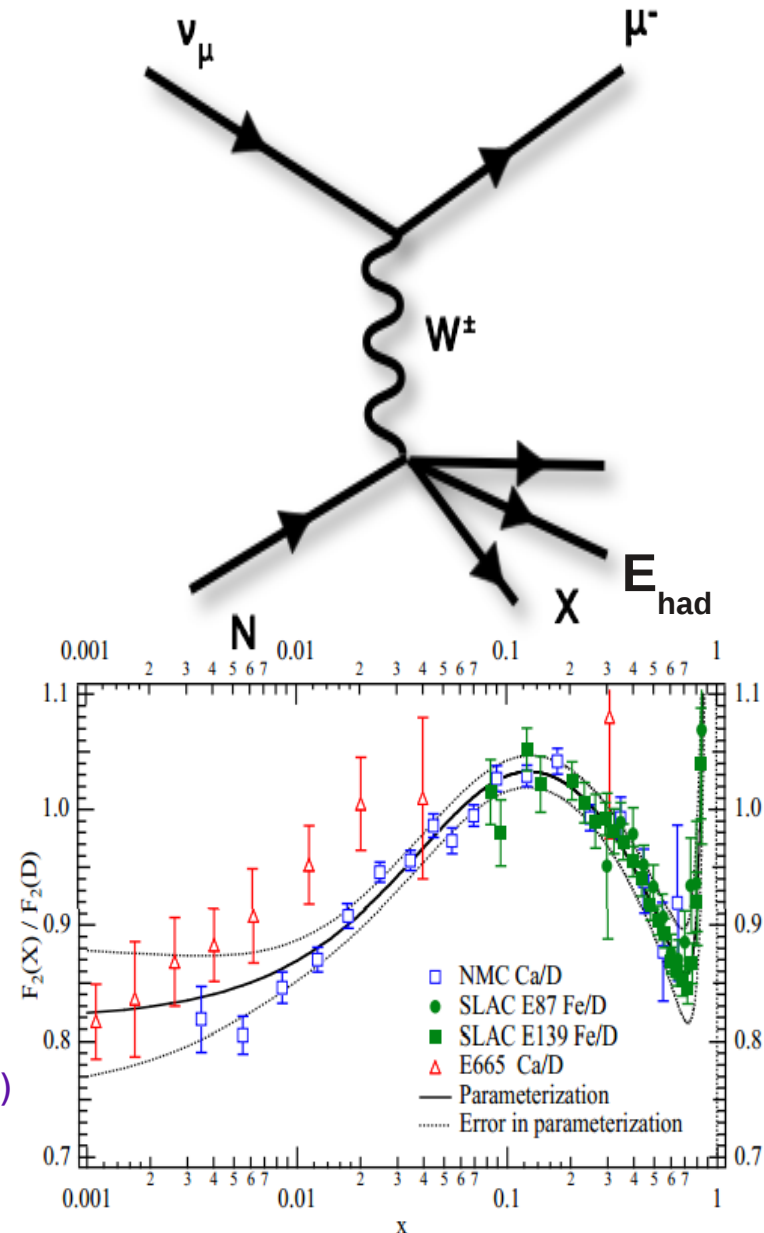
Joel Mousseau  
On Behalf of the MINERvA Collaboration  
University of Florida  
August 6<sup>th</sup> 2015  
DPF 2015, Ann Arbor, MI

# Neutrinos in Nuclear Media

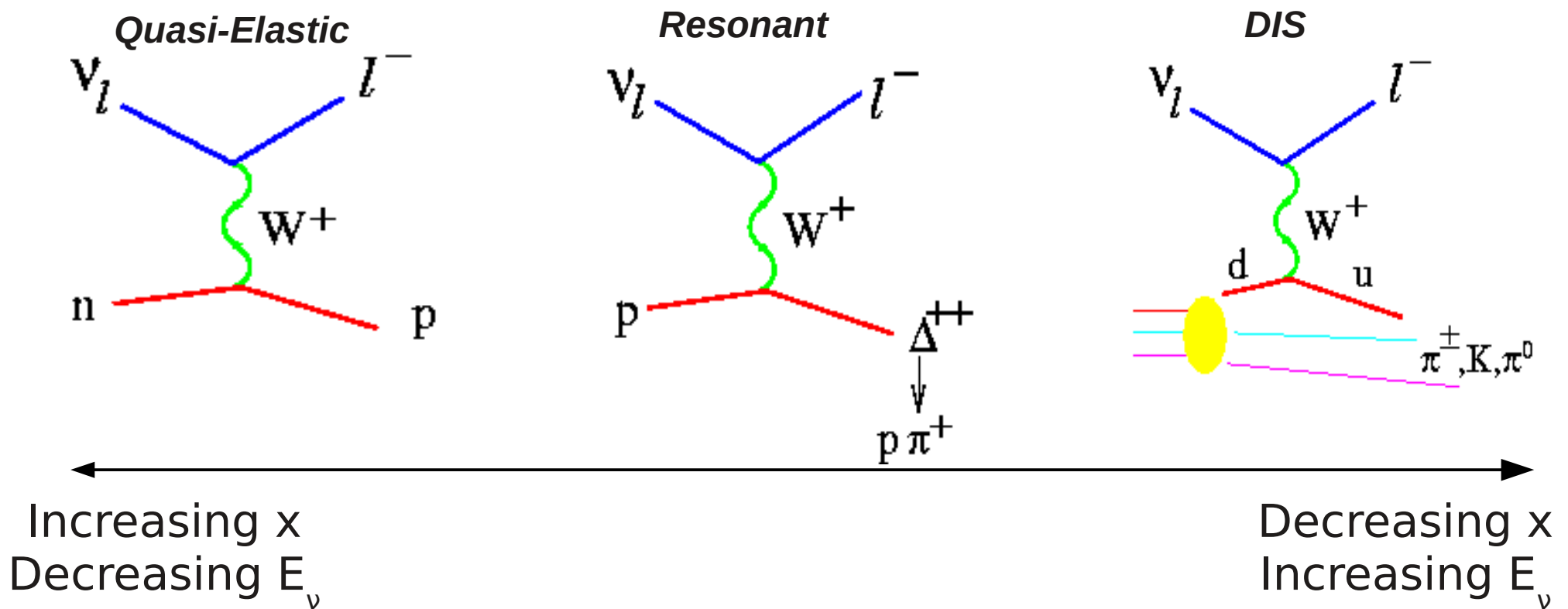
- One common theme of contemporary neutrino experiments: they rely on large A materials to supply adequate event rates (Fe, Ar, C, H<sub>2</sub>O etc.)
- Problem: nuclear effects caused by nucleons bound in a nucleus distort the energy reconstruction of the neutrinos.
- These effects manifest in *ratios* of  $d\sigma/dx$ .
- Effects not well understood in neutrino physics. General strategy has been to adapt electron scattering effects into neutrino scattering theory.

$$x = \frac{Q^2}{2ME_{had}}$$

A. Bodek, I. Park, and U.-K. Yang,  
Nucl. Phys. Proc. Suppl. 139, 113 (2005)



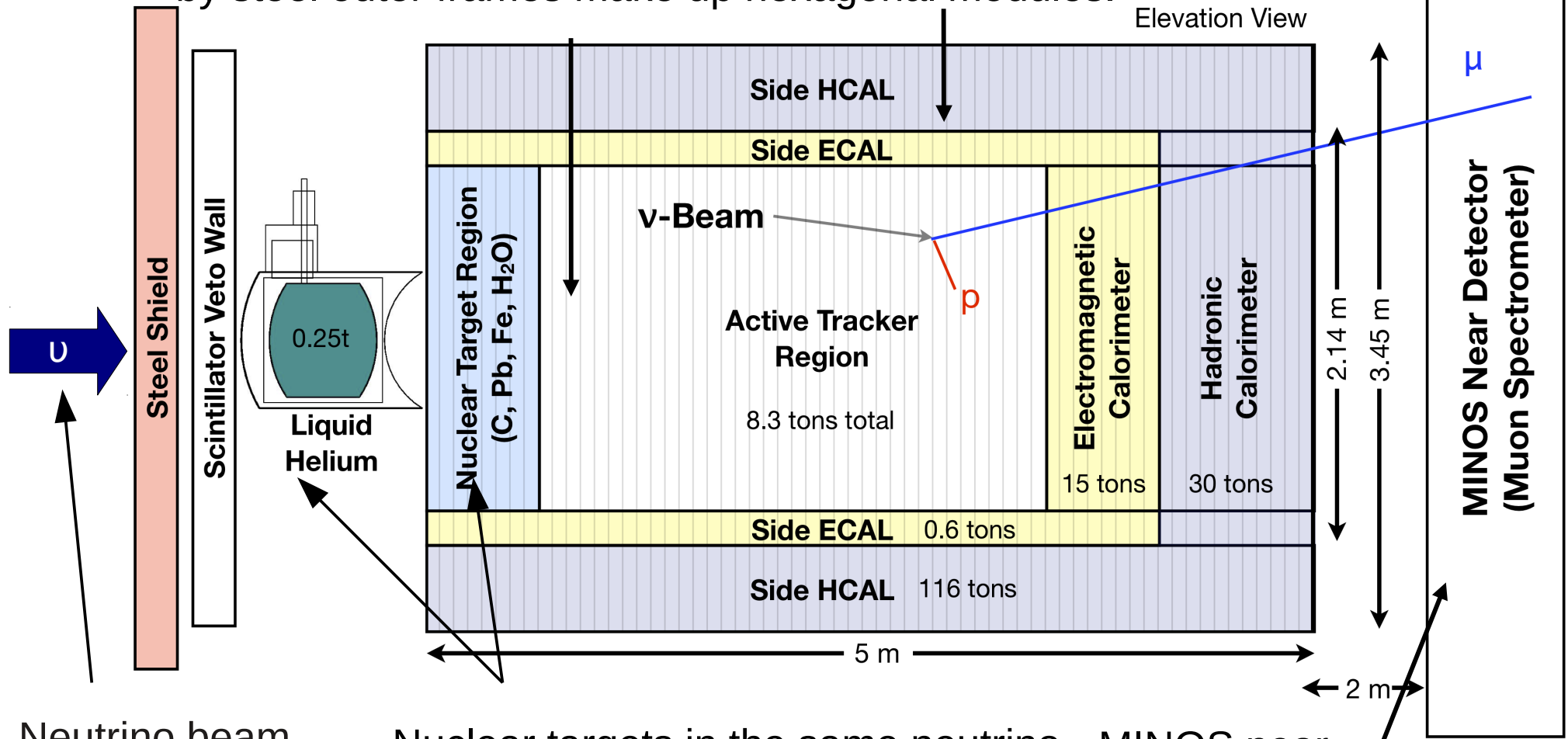
# Introduction to Neutrino Scattering



- Charged current neutrino + nucleon events are broadly categorized into quasi-elastic (single nucleon final state), resonant (multiple pion, single nucleon final state) and deeply inelastic (multiple hadron final state)
- Lower Bjorken- $x$  implies more inelastic events.
- Total neutrino cross section dominated by quasi-elastic up to  $E_\nu \sim 2.0$  GeV

# Enter MINERvA

Planes of scintillator strips, surrounded by steel outer frames make up hexagonal modules.



Neutrino beam created by Fermilab NuMI beamline

Nuclear targets in the same neutrino beam allow MINERvA to make A-dependent physics measurements.

MINOS near detector used for escaping muon ID and reconstruction.

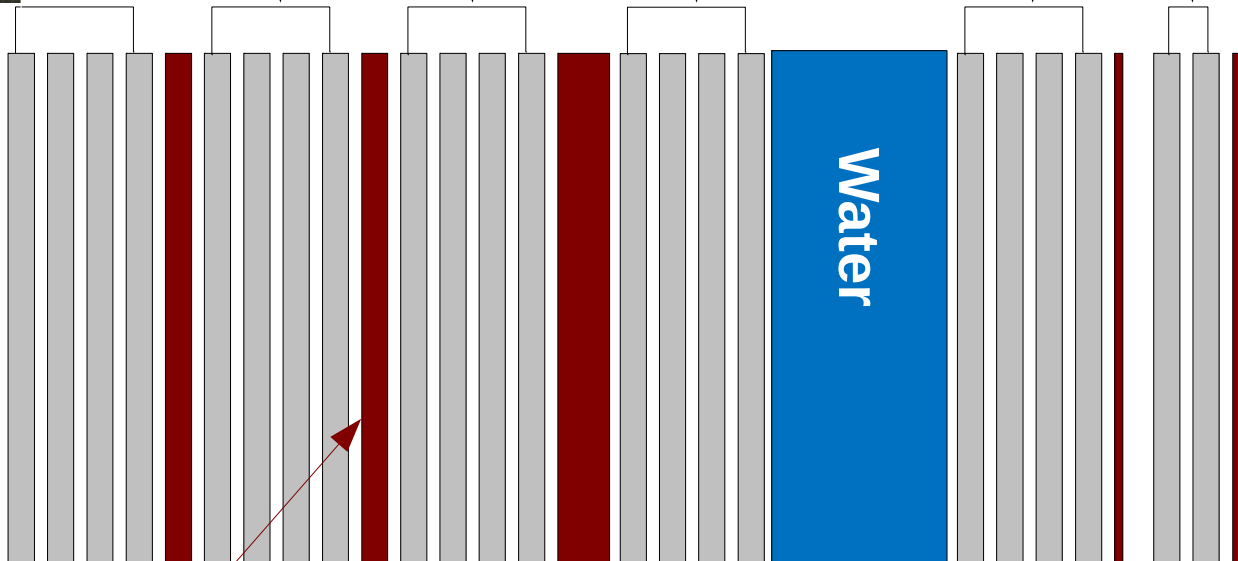
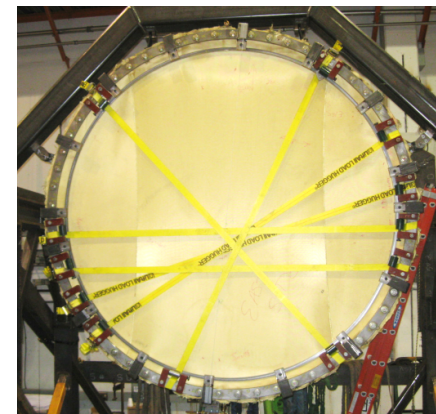




250 kg  
Liquid He:  
not used

## Active Scintillator Modules

500kg  
Water:  
Not used in today's analysis

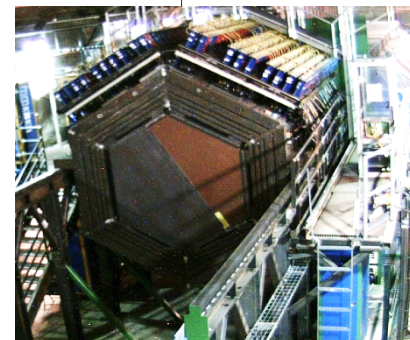
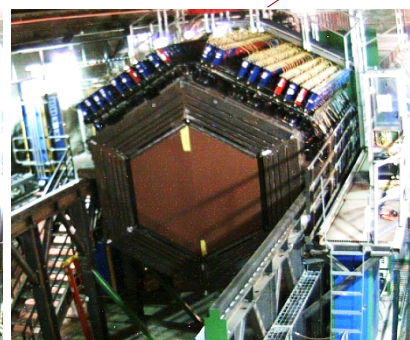
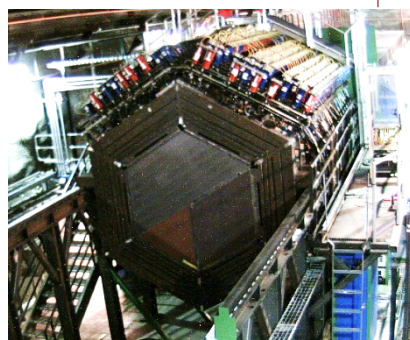
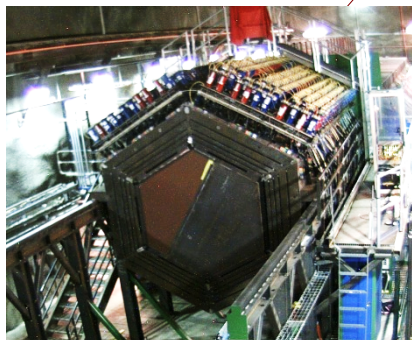


1" Pb / 1" Fe  
266kg / 323kg

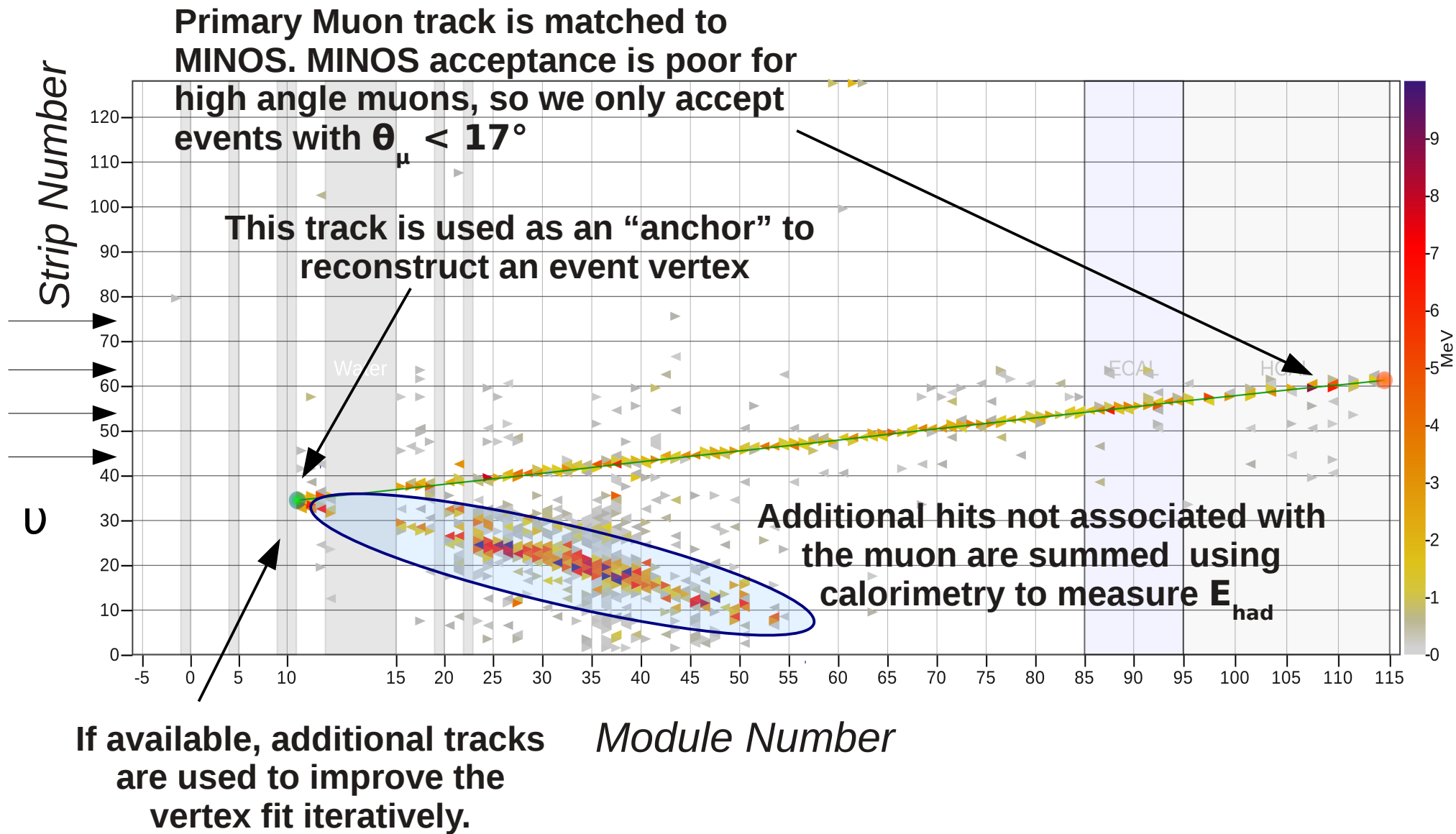
3" C / 1" Fe / 1" Pb  
166kg / 169kg / 121kg

0.3" Pb  
228kg

.5" Fe / .5" Pb  
161kg / 135kg

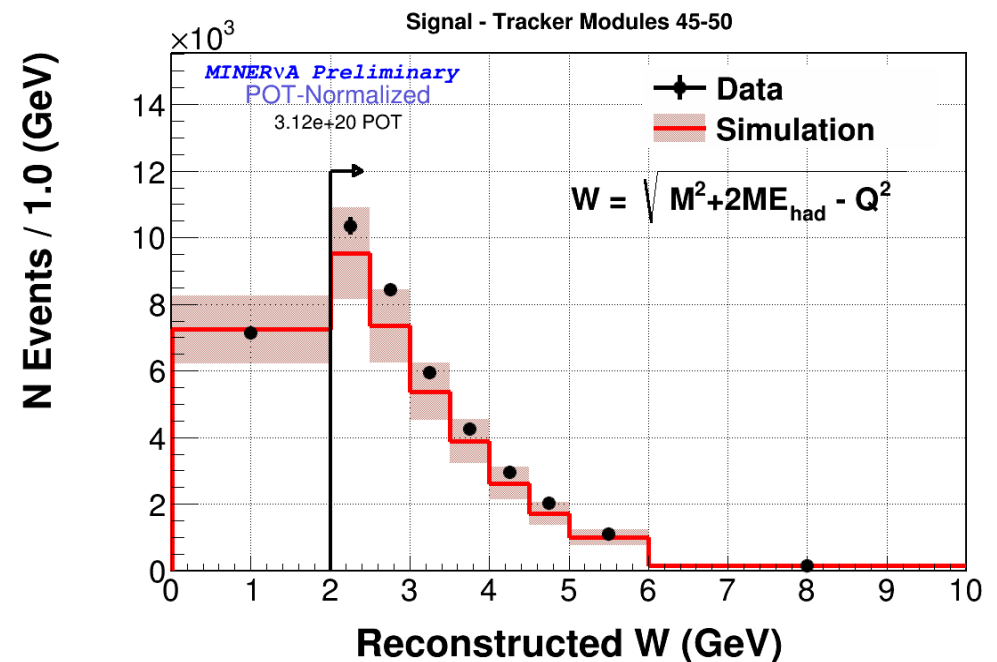
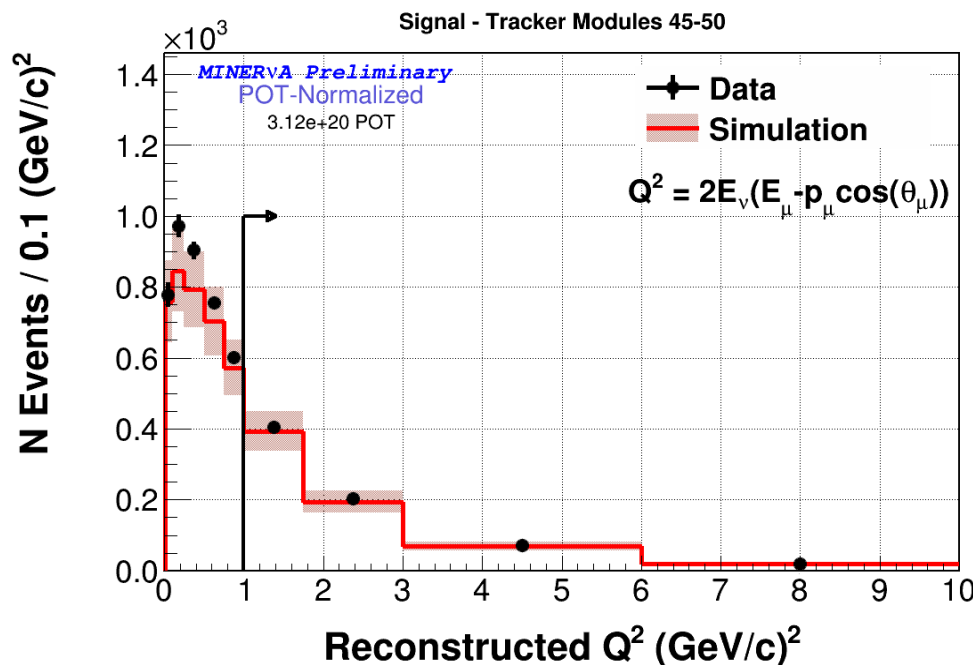


# Event Selection and Reconstruction

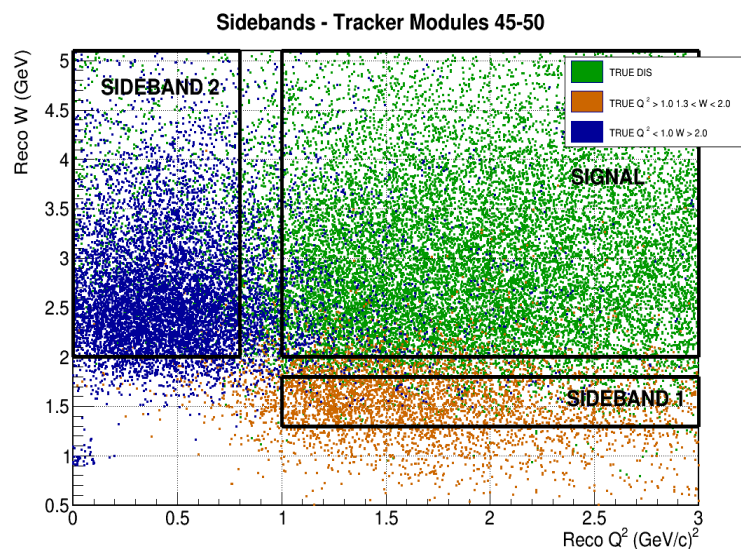
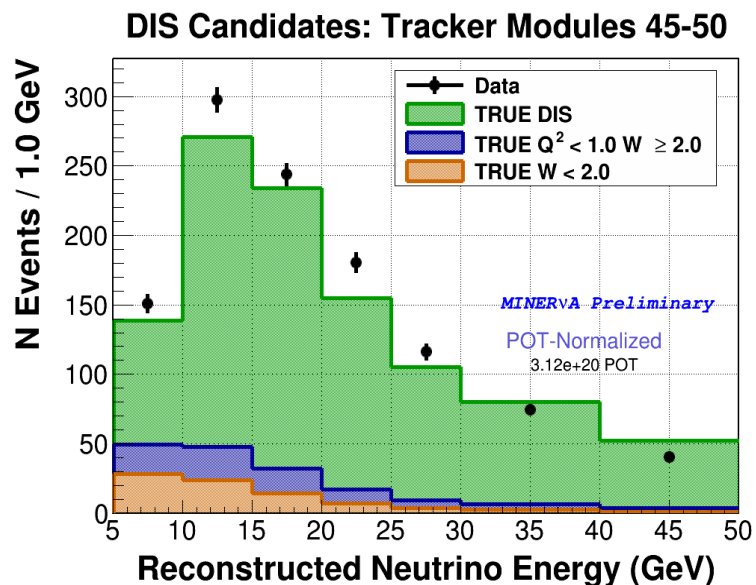


# From Inclusive to DIS

- We isolate a deeply inelastic sample by making cuts on the four momentum transfer ( $Q^2$ ) and final state invariant mass ( $W$ )
- Require  $Q^2 > 1.0 \text{ (GeV / c)}^2$  and  $W > 2.0 \text{ GeV / c}^2$ . These cuts remove the quasielastic and resonant events from the inclusive sample, and allow us to interpret our data on the quark level.
- Cuts are illustrated for CH events between 5 and 50  $\text{GeV } E_\nu$  and  $\theta_\mu < 17^\circ$ .



# Backgrounds (Kinematic):



- After making kinematic cuts on  $Q^2$  and  $W$ , we are left with a background of events with *true*  $Q^2 < 1.0$  (GeV/c)<sup>2</sup> and  $W < 2.0$  (GeV/c<sup>2</sup>) that smear into the sample.

- Estimate this background in the nuclear targets and scintillator using MC (left plots).

- MC is tuned to data using events adjacent to  $W = 2.0$  (GeV/c<sup>2</sup>) and  $Q^2 = 1.0$  (GeV/c)<sup>2</sup>

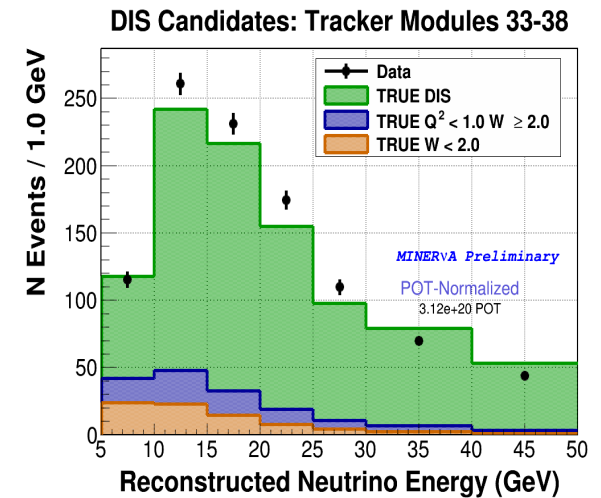


# Fitting Sidebands

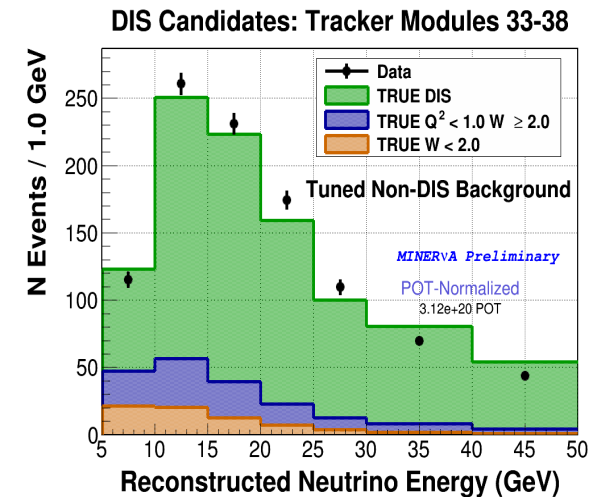
## Scale Factors Applied to Simulation (stat. Error only)

A	$W < 2.0$	$Q^2 < 1.0 \text{ } W > 2.0$
C	$0.87 \pm 0.07$	$1.42 \pm 0.10$
CH	$0.90 \pm 0.01$	$1.45 \pm 0.01$
Fe	$0.93 \pm 0.04$	$1.36 \pm 0.05$
Pb	$0.85 \pm 0.04$	$1.19 \pm 0.04$

- The MC of both sidebands are fit simultaneously over the region  $5 < E_\nu < 50$  GeV using a  $\chi^2$  minimization.
- The data and MC of each target is summed by material prior to fitting, so we end up with a scale factor for C, CH, Fe and Pb.
- Primarily, the data prefer *more* backgrounds.

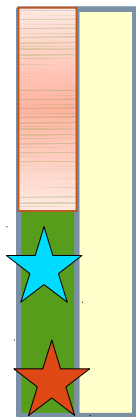


*Before Fitting*

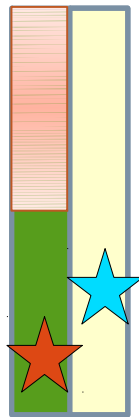


*After Fitting*

# Background Events (Wrong Nuclei)

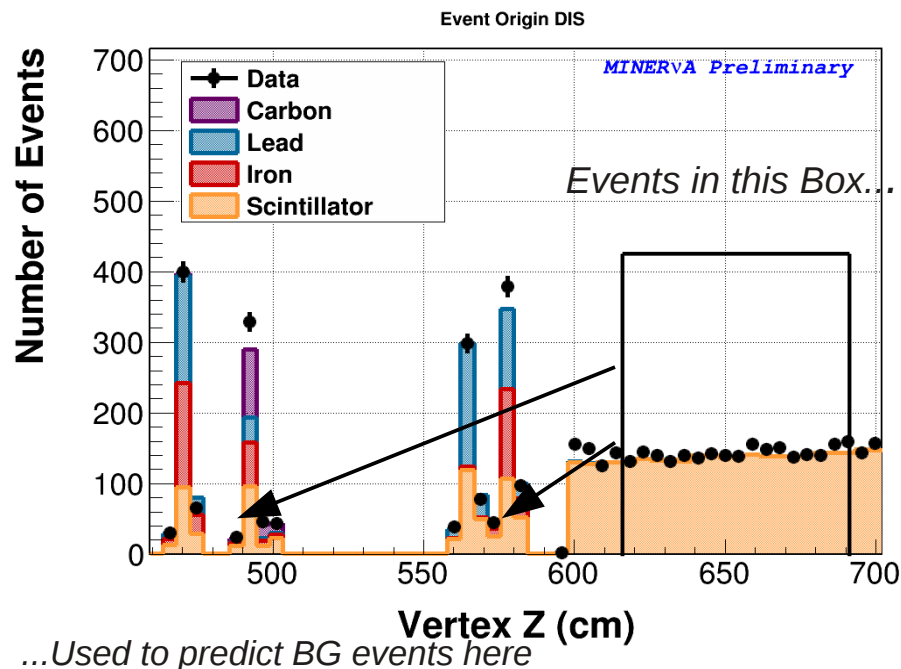


True vertex (blue star) is in the same material as the reconstructed vertex (orange star).



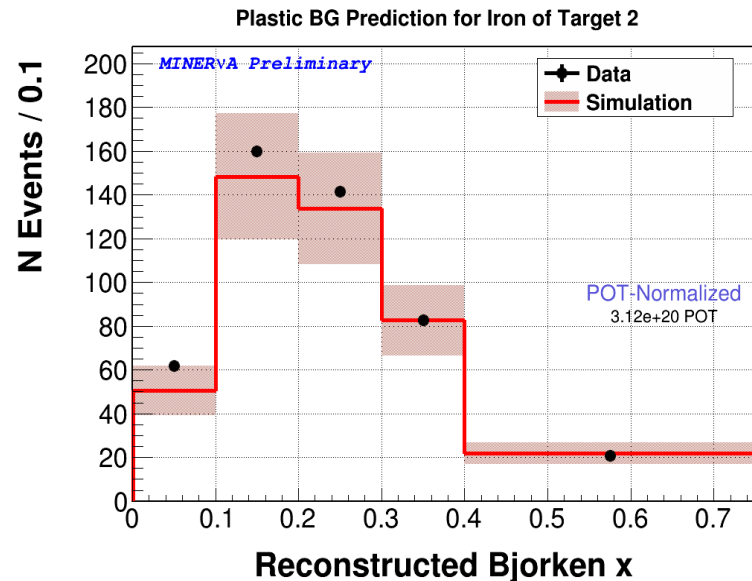
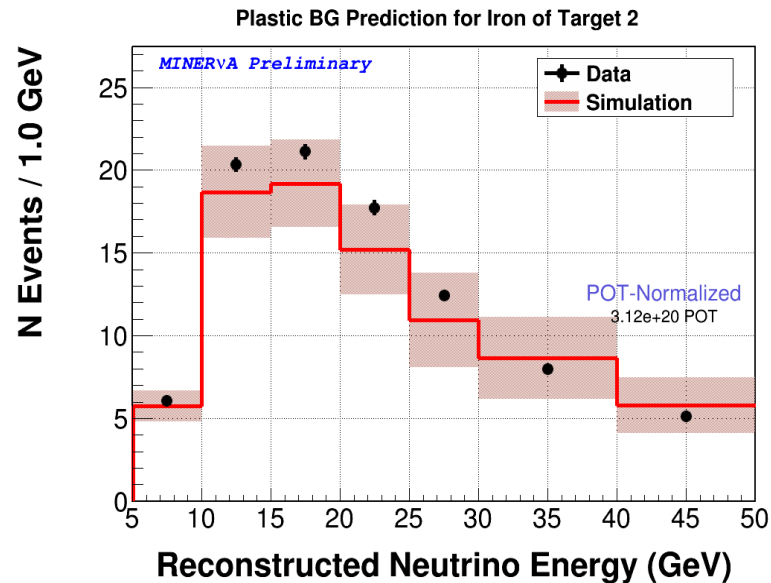
Vertex is reconstructed in the Fe (green). However, the true vertex of the event is in the scintillator (yellow).

- Events occasionally truly occur in the scintillator surrounding the nuclear target, but are reconstructed to the passive target. This makes up a second background.



- We subtract this background by measuring the event rates in the downstream tracker, and extrapolating these events upstream to the nuclear target region.
- Downstream events are weighted for MINOS acceptance based on  $E_\mu$ ,  $\theta_\mu$ .

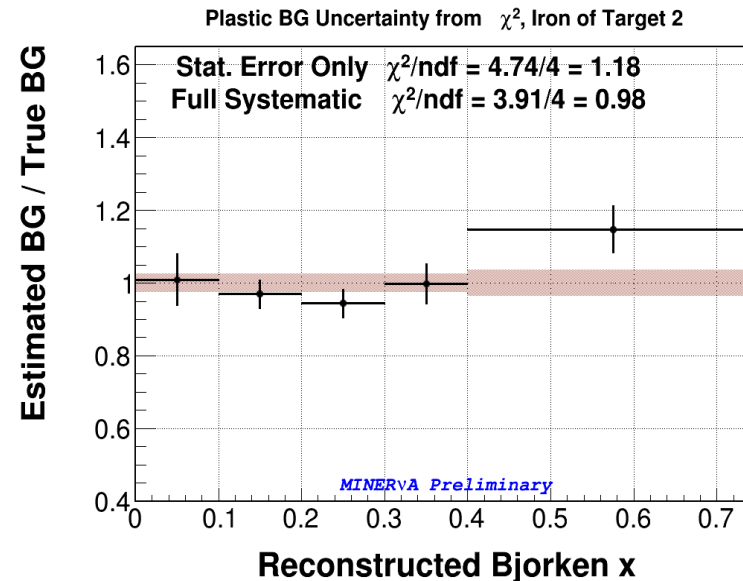
# Wrong Nuclei BG (Data / MC)



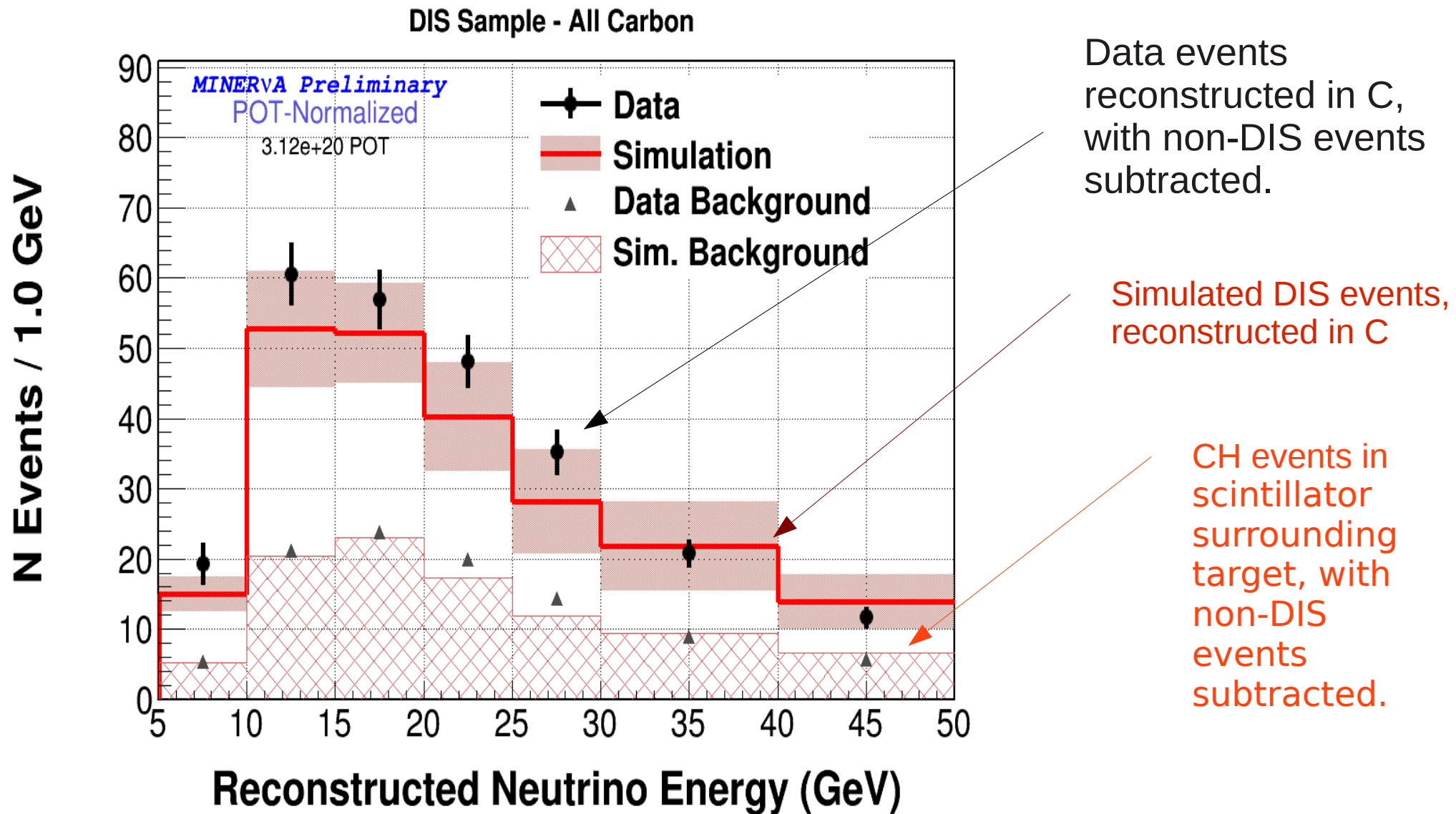
- Wrong nuclei backgrounds are extracted separately for data and MC, in each variable ( $E_\nu$ ,  $x$ , etc.)

- In each case, the non-DIS events have been subtracted using the procedure previously described.

- Prediction accuracy is measured from MC. Additional systematic uncertainty is calculated from the disagreement.



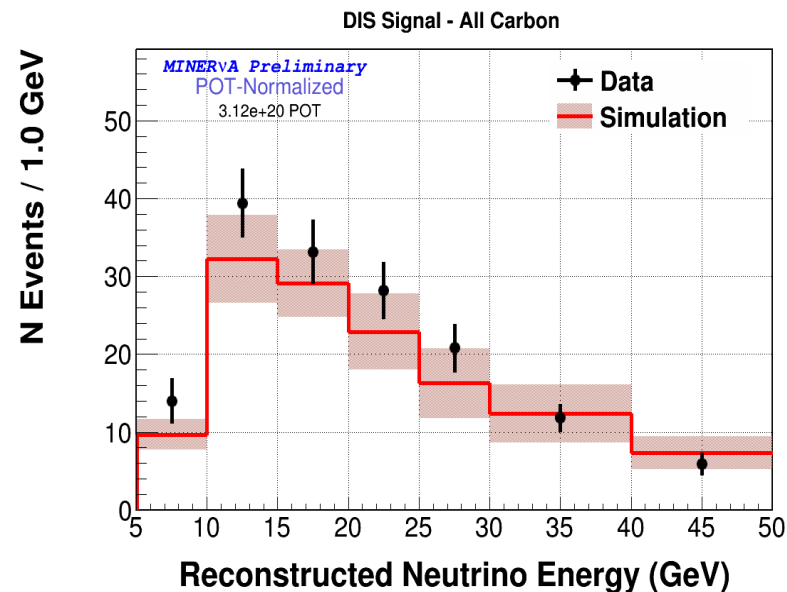
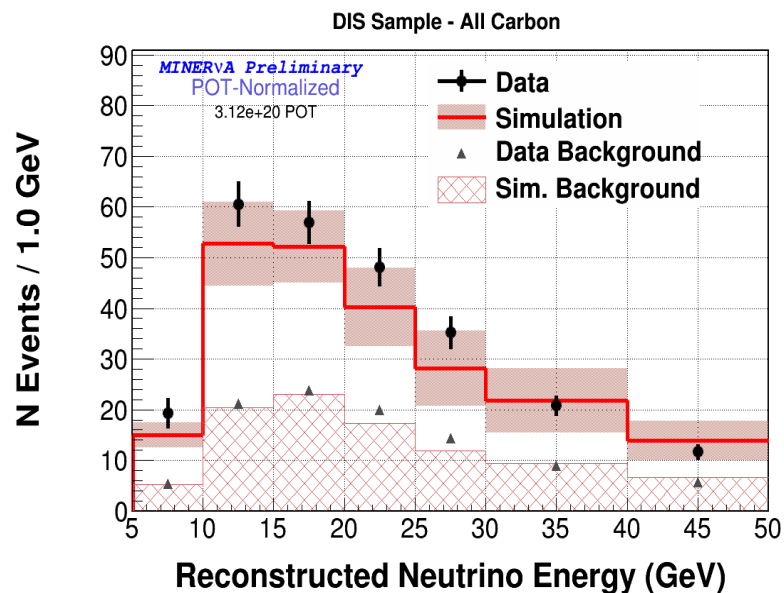
# Putting it Together





# Putting it Together

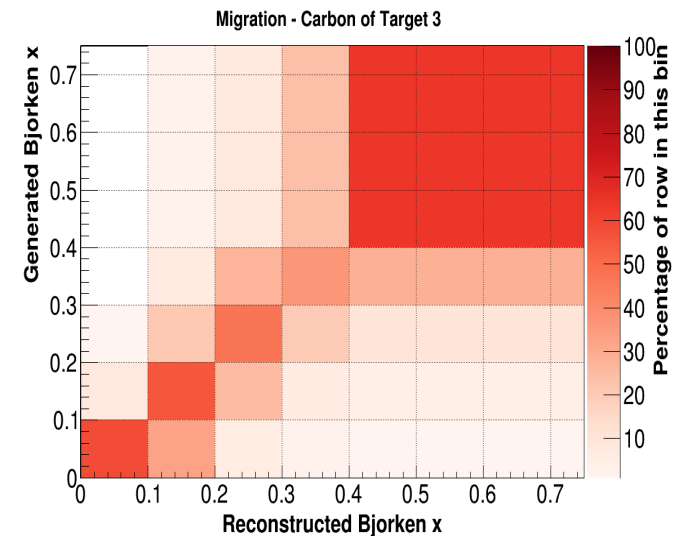
Take our sample of reconstructed DIS events in carbon with CH events...



...And subtract those events to obtain a sample of DIS on carbon in data and MC.  
Large uncertainties on neutrino flux,  
measure ratios of C, Fe and Pb /CH where  
flux will cancel.

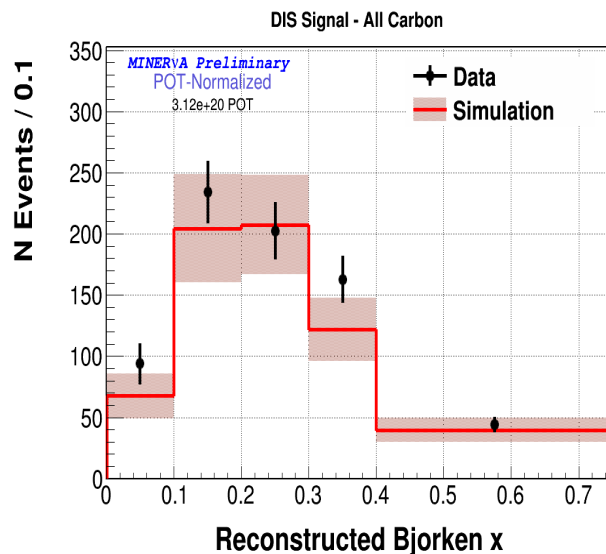
# Migration and Unfolding

- Detector resolution smears the reconstructed values of  $x$  and  $E_\nu$  from their generated quantities (right plot).
- Correct for this smearing using *unfolding*.
- Unfolded distributions are then efficiency corrected to account for detector effects prior to taking ratios.

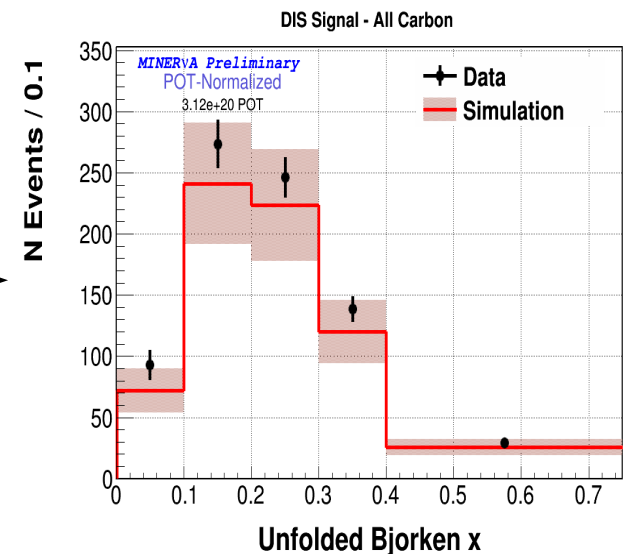


Migration  
Matrices used  
as input

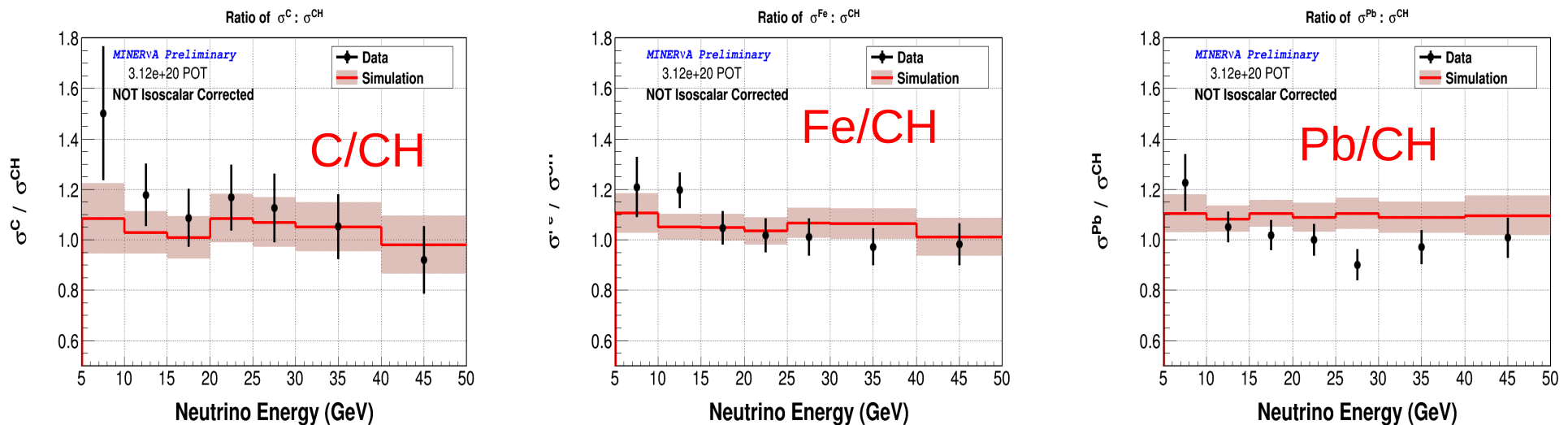
Warning! Unfolding introduces  
correlations between bins



Unfold

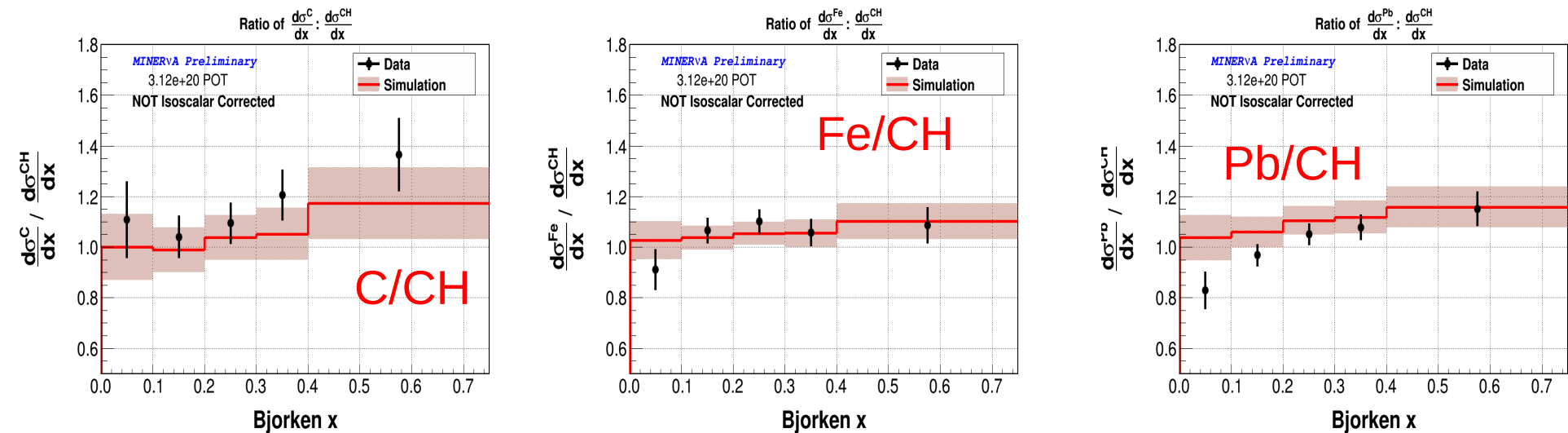


# DIS Ratios: $\sigma(E_\nu)$



- We measure *ratios* of cross sections to reduce systematic errors from the neutrino flux calculation.
- Ratios of the heavy nuclei (Fe, Pb) to lighter CH are evidence of nuclear effects.
- There is a general trend of the data being below the MC at high energy.
- This trend is larger in the lead than in the iron.

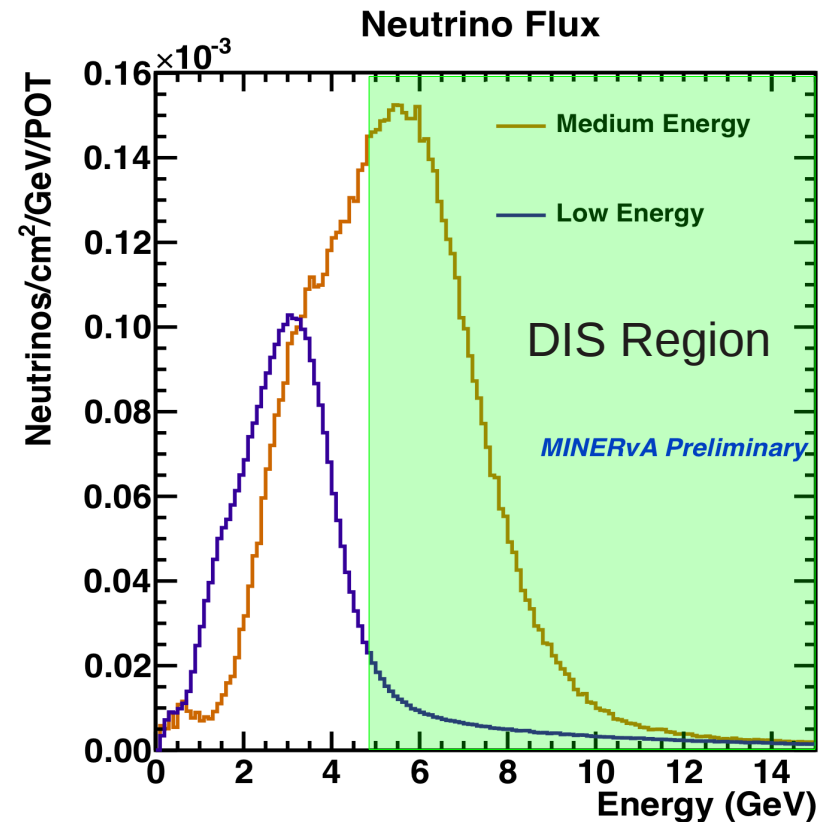
# DIS Ratios: $d\sigma/dx$



- $x$  dependent ratios directly translate to  $x$  dependent nuclear effects.
- Currently, our simulation assumes the *same*  $x$ -dependent nuclear effects for C, Fe and Pb tuned to  $e^-$  scattering.
- The shape of the data at low  $x$ , especially with lead is consistent with additional nuclear shadowing.
- The intermediate  $x$  range of ( $0.3 < x < 0.75$ ) shows good agreement between data and simulation.

# Future Directions

- Future studies of nuclear effects will benefit greatly from MINERvA's increased energy and intensity run, taking data as we speak.
- Expect much better sensitivity at high and low  $x$  with increased beam energy.
- Currently have a quasi-elastic analysis in nuclear targets in progress for Low Energy beam
- A bright future is expected!*



# Conclusions

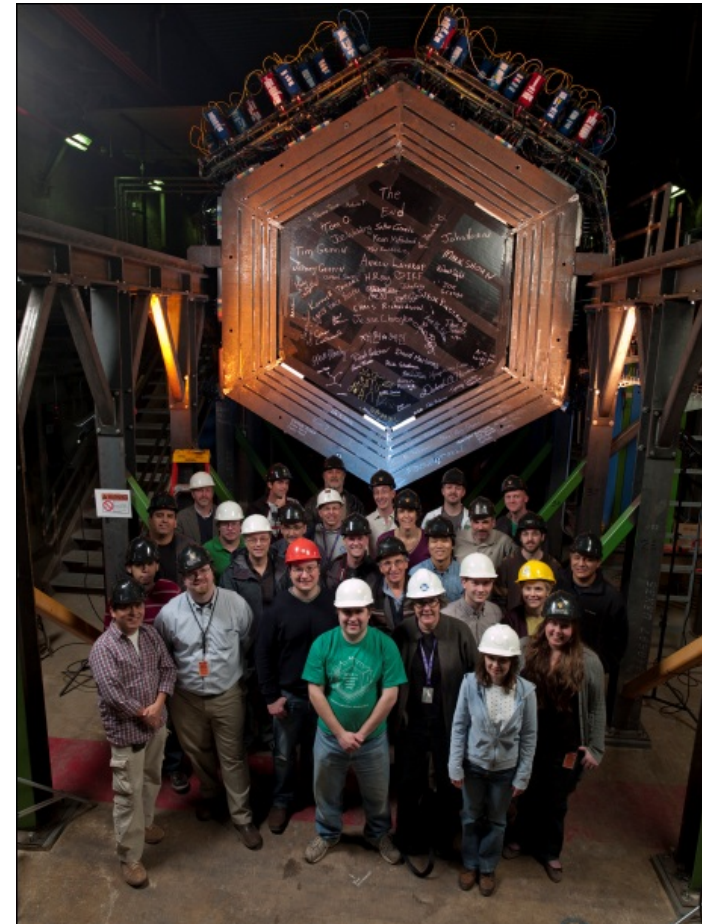
- MINERvA has made a measurement of neutrino DIS events on multiple nuclei in an identical neutrino beam.
- These measurements may be interpreted directly as DIS  $x$ -dependent nuclear effects.
- We currently observe a deficit in our lead data suggestive of additional nuclear effects.
- Our data in the intermediate  $x$  region shows no deviation from theory.
- Future higher energy measurements will be higher statistics as well as the ability to resolve larger  $x$  values ( $x > 0.75$ ).

Thank you for Listening!

# The MINERvA collaboration consists of approximately 65 Nuclear and Particle Physicists

University of California at Irvine  
Centro Brasileiro de Pesquisas Físicas  
University of Chicago  
Fermilab  
University of Florida  
Université de Genève  
Universidad de Guanajuato  
Hampton University  
Inst. Nucl. Reas. Moscow  
Massachusetts College of Liberal Arts  
University of Minnesota at Duluth

Universidad Nacional de Ingeniería  
Northwestern University  
Otterbein University  
Pontificia Universidad Católica del Perú  
University of Pittsburgh  
University of Rochester  
Rutgers, The State University of New Jersey  
Universidad Técnica Federico Santa María  
Tufts University  
William and Mary



**At the conclusion of  
detector construction**



# Backup Slides

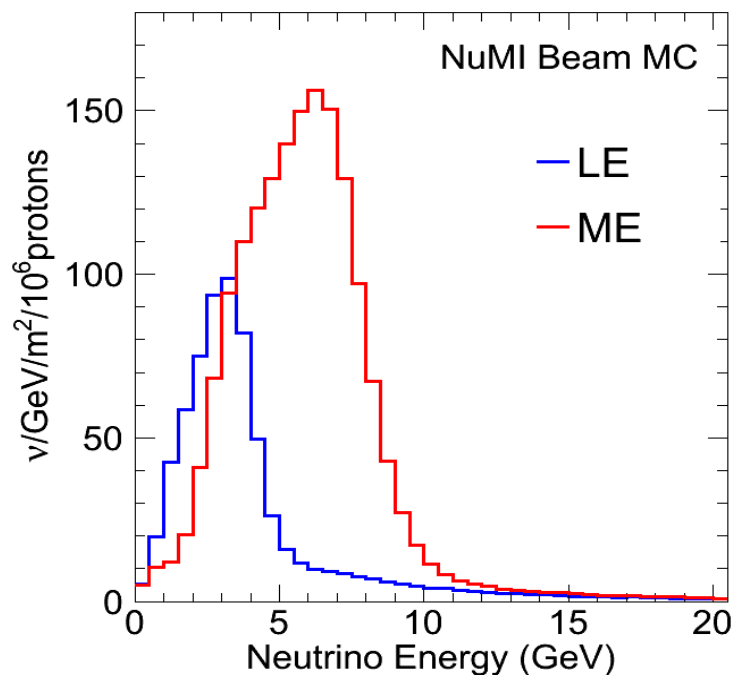
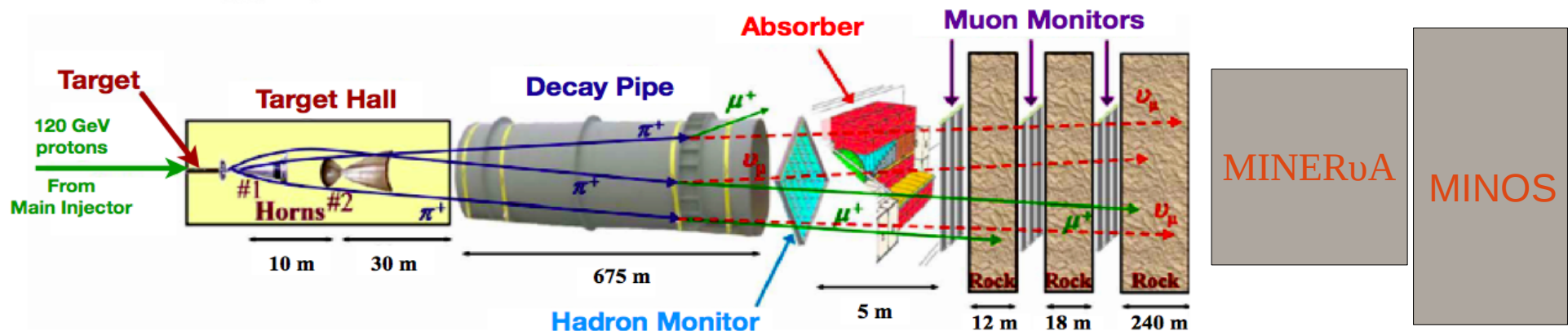
# Isoscalarity

- Heavier nuclei (Fe, Pb) are composed of an unequal number of protons and neutrons (e.g. Pb: 82 protons, 125 neutrons).
- The  $\nu_\mu + N$  cross section is *different* for protons and neutrons;  $\nu_\mu$  want to couple to  $d$  quarks, and the neutron contains more  $d$  than  $u$  quarks.
- This effect is  $x$  dependent (higher  $x \rightarrow$  more valence quarks  $\rightarrow$  more  $d$  quarks).
- Currently, the MINERvA data does not correct for this difference; this requires some theory input.

$$f_{iso} = (A/B) \times (Z_B/Z_A) \frac{1 + (N_B/Z_B) \frac{\sigma(\nu n_f)}{\sigma(\nu p_f)}}{1 + (N_A/Z_A) \frac{\sigma(\nu n_f)}{\sigma(\nu p_f)}}$$

*Isoscalar correction of two nuclei A and B with Z protons and N neutrons.*

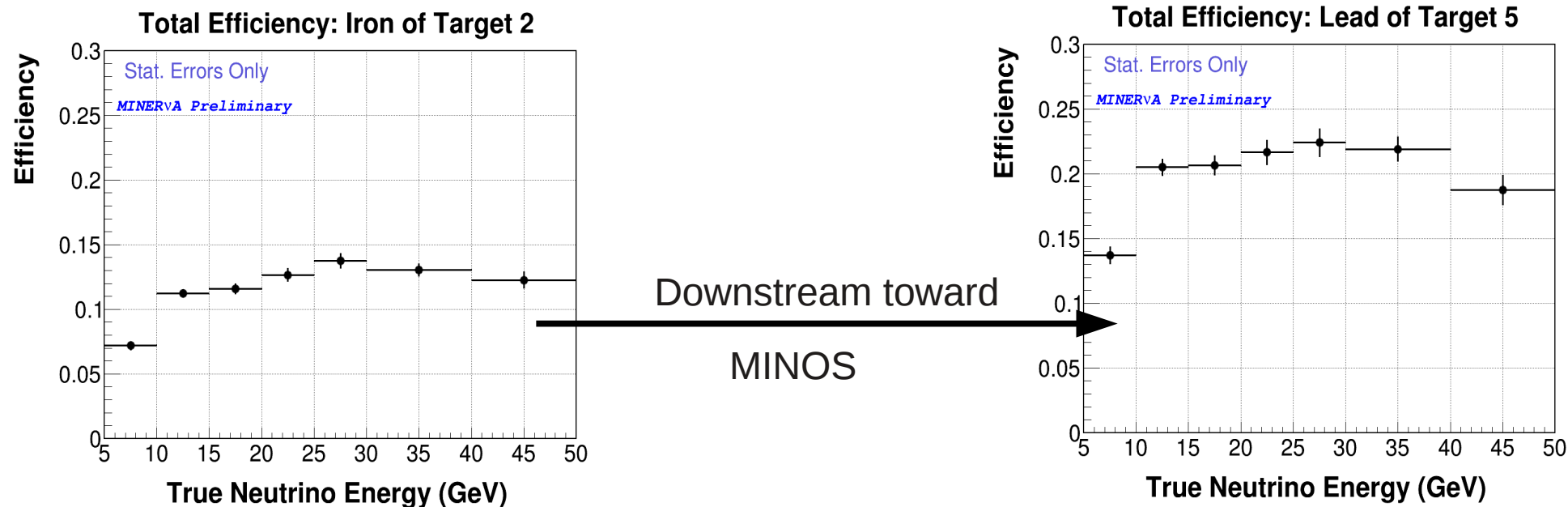
# The NuMI Beamline



- Low energy beam (blue plot) has a peak energy of approximately 3 GeV.
- Medium energy is 7 GeV; higher beam energy means more DIS events, and access to higher  $x$ .
- Flux measurement is constrained with external hadron production data (NA49).

# Efficiency Correction

- We correct for detector acceptance, using an efficiency correction derived from the MC. However, we only correct up to  $E_\nu = 50$  GeV and  $\theta_\mu < 17^\circ$ .
- Efficiency is corrected target by target, since it is a function of the distance from the target to MINOS.
- Largest source of inefficiency is MINOS matching requirement. This acceptance improves as we move downstream in the detector.

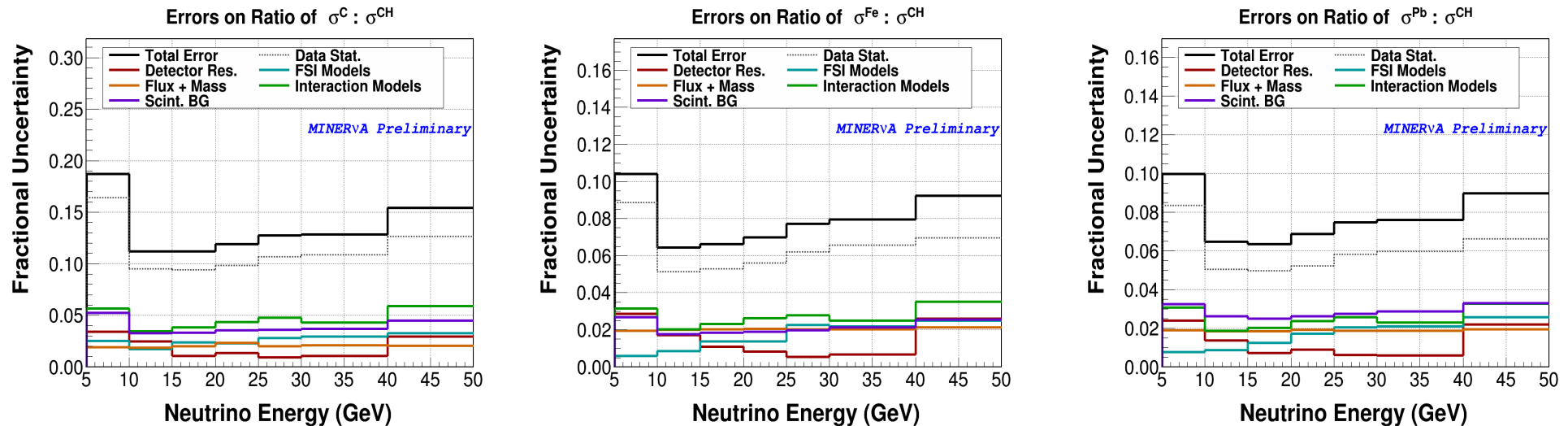


# Event Table

$X_{bj}$ (unfolded)	C	Fe	Pb
0 – 0.1	90	220	230
0.1 - 0.2	270	840	930
0.2 – 0.3	250	800	940
0.3 – 0.4	140	390	520
0.4 – 0.75	100	250	350
0.75+	1	1	1
TOTAL	850	2500	2970

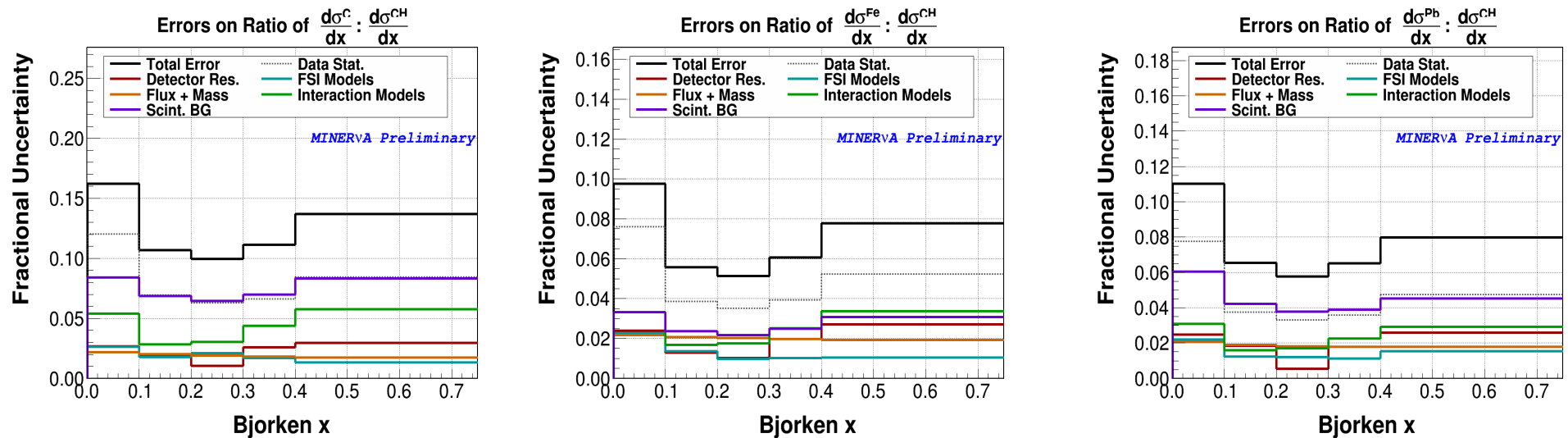
- Most of our events are in the anti-shadowing and shadowing region; a marginal number in the EMC region.

# Ratio Uncertainties



- Most of the uncertainty stems from data statistics.
- Higher intensity, higher energy beam will improve this substantially.
- Correlations in data introduced from unfolding are *NOT* accounted for in Data stat. uncertainty.

# Ratio Uncertainties



- Ratios are dominated by data statistics for the most part.
- Scintillator background is a larger uncertainty in  $x$ .
- Correlations in data introduced from unfolding are *NOT* accounted for in Data stat. uncertainty.

# Recoil Reconstruction

- Recoil energy = all non muon energy in a -25, 30 ns window of the vertex time.

$$E_{had} = \alpha \times \sum_i^{hits} c_i E_i$$

- Calibrated energy deposits ( $E_i$ ) in the detector weighed by the energy lost in passive material ( $c_i$ ; see table).

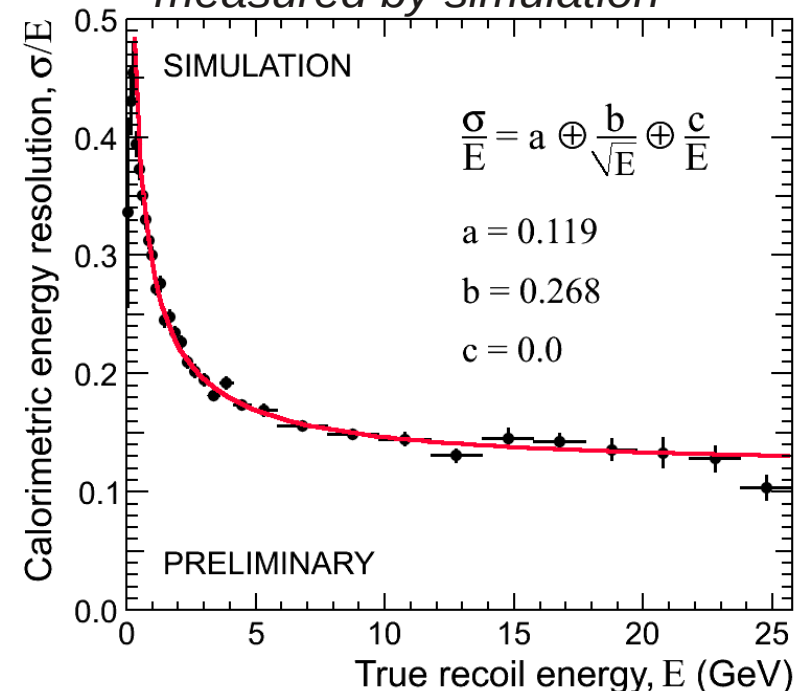
*Energy lost by a mip  
in each material*

Material	CH	C	Fe	Pb
dE/dx (MeV/g/cm <sup>2</sup> )	1.96	1.74	1.45	1.12

*Overall scale factor ( $\alpha$ )  
computed from simulation*

vertex	Tgt 2	Tgt 3	Tgt 4	Tgt 5	Trk
$\alpha$	1.81	1.71	1.60	1.59	1.62

*Tracker recoil energy resolution as  
measured by simulation*

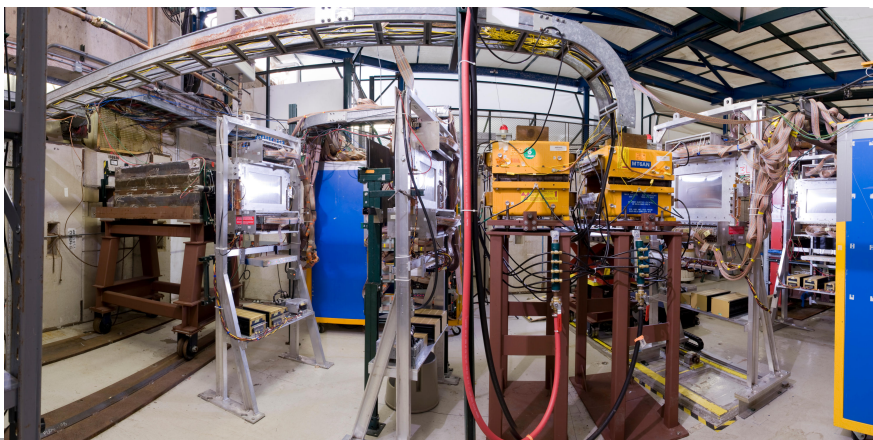




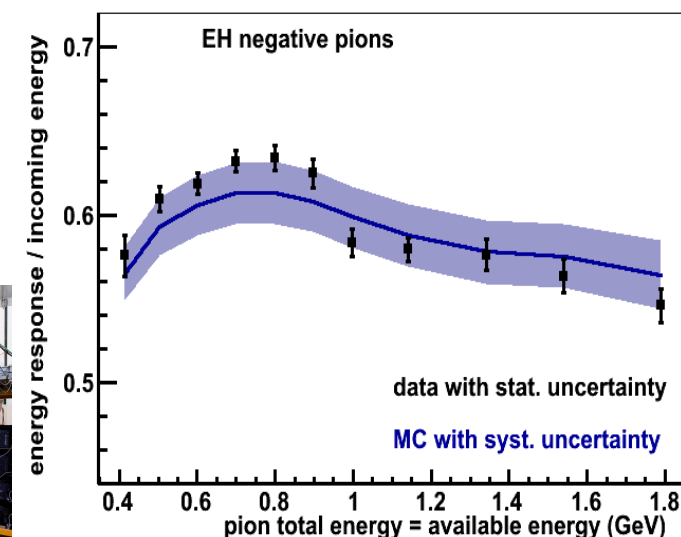
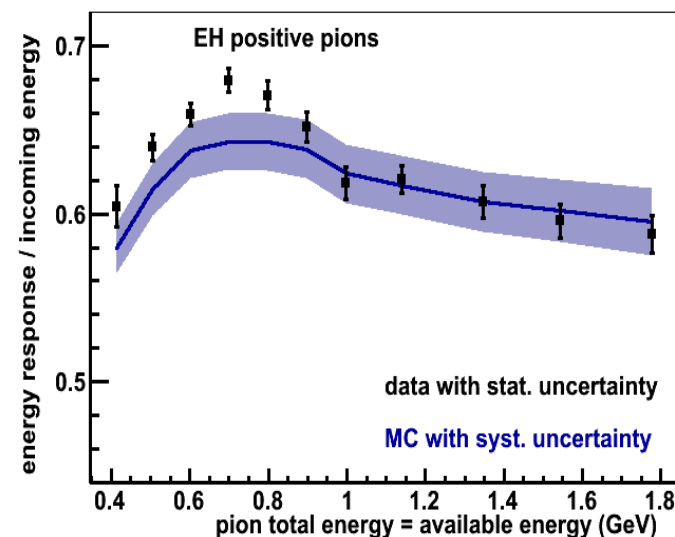
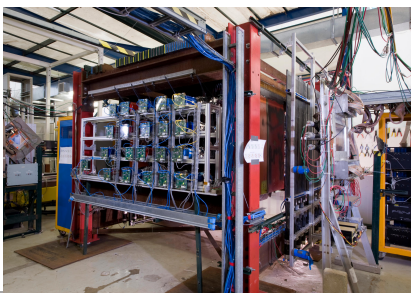
# Test Beam

- The MINERvA detector's hadronic energy response is measured using a dedicated test beam experiment at the Fermilab Test Beam Facility (FTFB)
- Custom built beamline took data during the summer of 2010.
- In addition to a Birk's Law calculation, hadronic energy reconstruction uncertainty is estimated from difference between test beam data and GEANT MC.

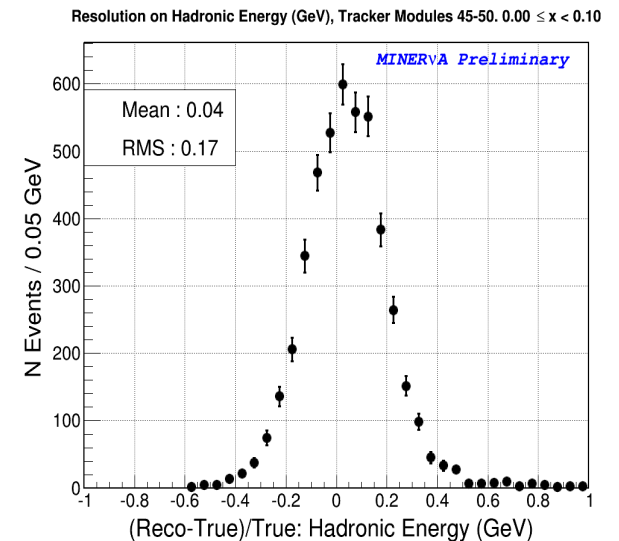
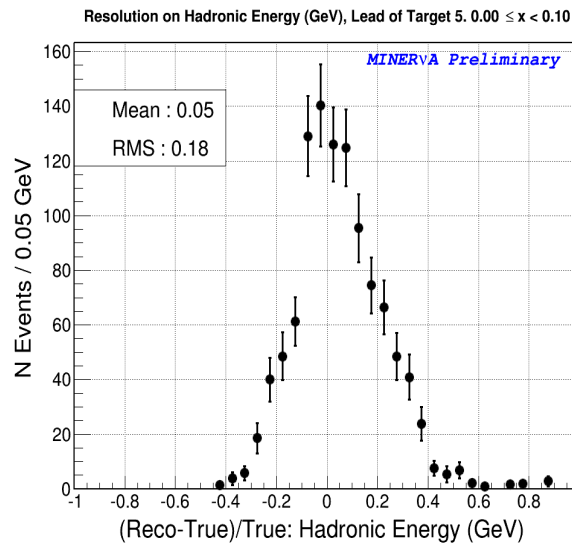
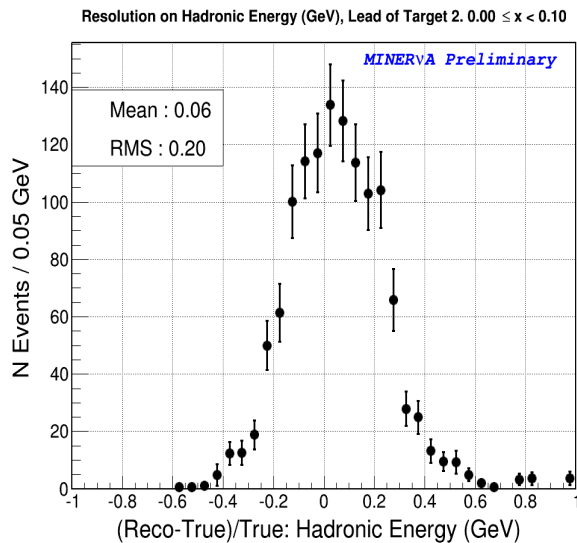
*Custom built beamline*



*Plus miniature detector*



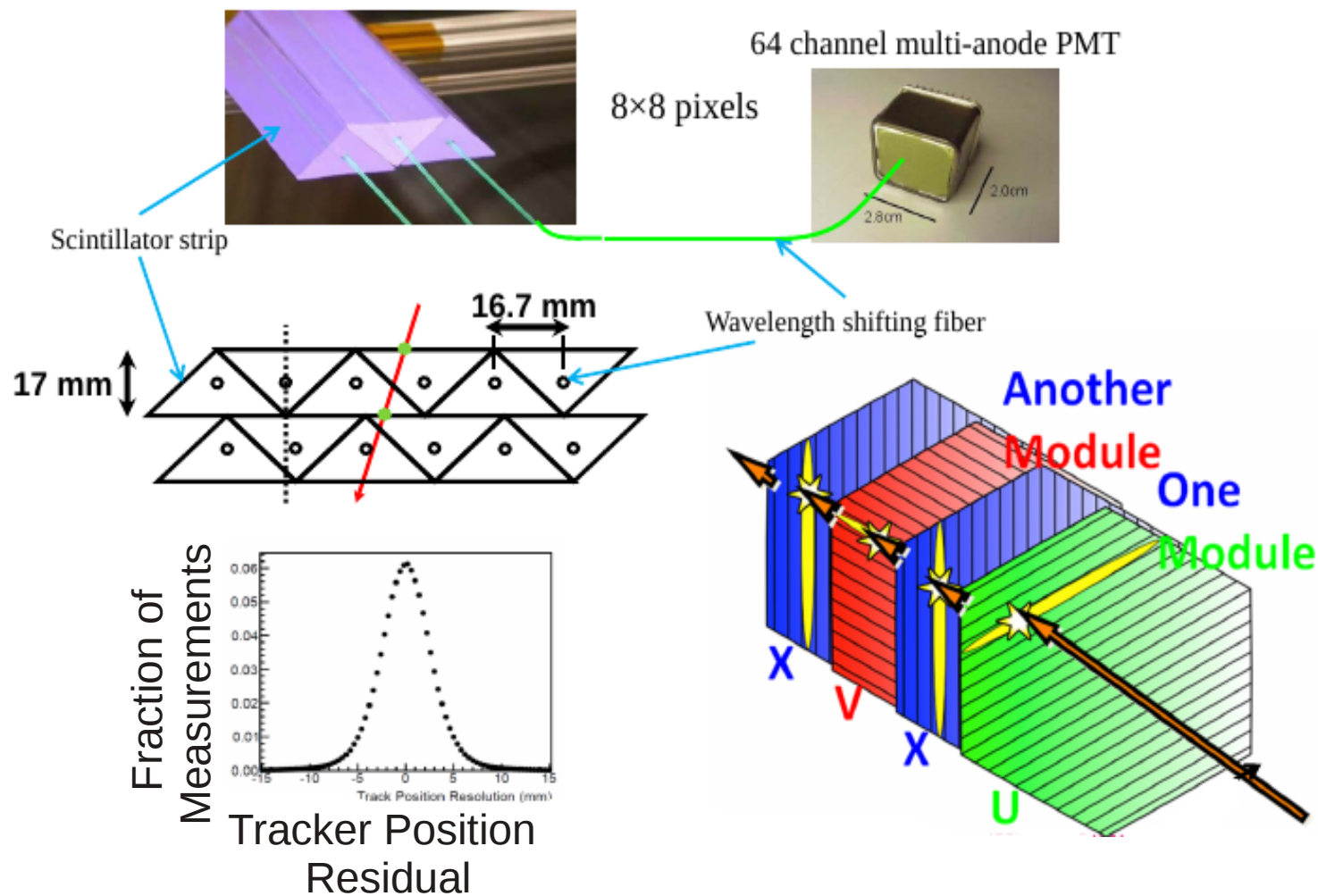
# Hardonic Energy Resolution



Downstream

- Accuracy of high-energy, low  $x$  hadronic showers is very similar between nuclear targets and tracker modules. Our simulation adequately accounts for the different geometry encountered by hadronic showers, regardless of where they originate.

# Detector Technology

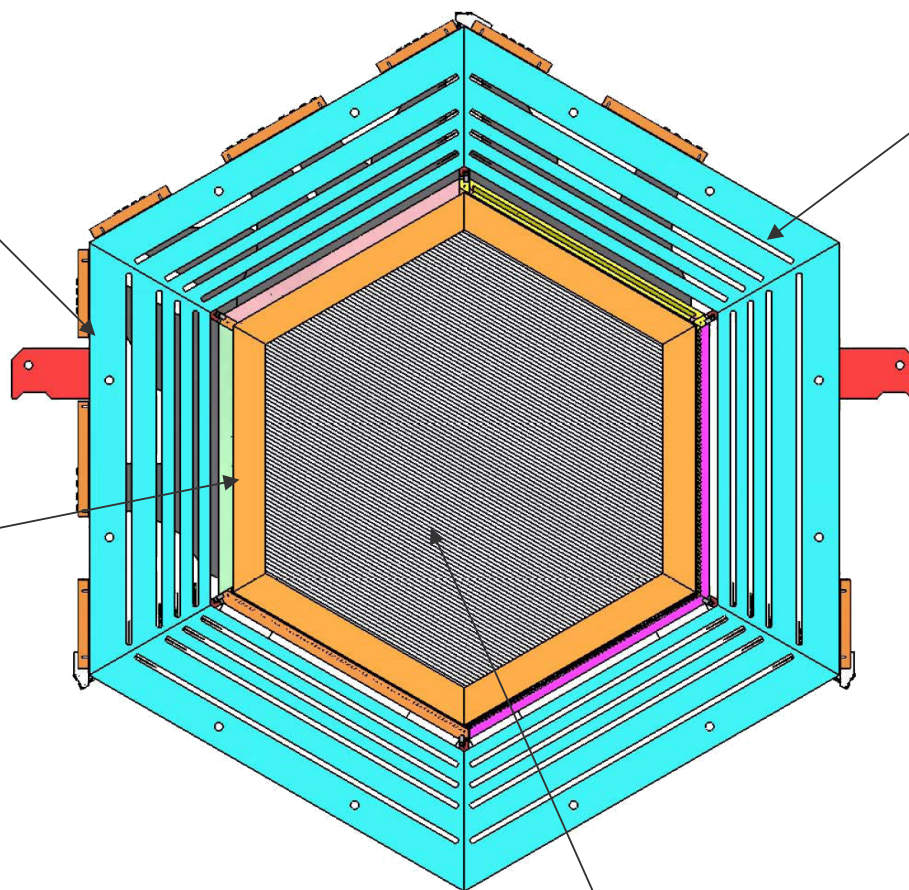


# A MINERvA Module

Outer Detector  
Frame

Scintillator  
Bars

Lead collar



Steel  
supports used  
for hanging  
modules on  
rails.

Inner Detector:  
Plastic scintillator strips