

Rydberg-Stark States of Positronium



Adam Deller

AMOPP, Department of Physics and Astronomy University College London

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Positronium

The n = 1 ground-state is separable into a short lived (0.125 ns) s = 0singlet, and long lived (140 ns) s = 1triplet states,



¹R. Ley (2002) Appl. Surf. Sci., 194:301



Gravity Measurements with Positronium?

Rydberg Ps Drift Tube



Fig. 1. Experiment to measure the gravitational free fall of n = 25 Rydberg positronium.

²A.P. Mills Jr. and M. Leventhal (2002) *Nucl. Instr. and Meth. in Phys. Res. B*, **192**:102



Interrogating Ps with lasers





Interrogating Ps with lasers





Interrogating Ps with lasers

a brief history (pre 2008)

1982 Excitation of the Positronium1 ${}^3S \rightarrow 2{}^3S$ Two-Photon Transition

Steven Chu and Allen P. Mills, Jr. Phys. Rev. Lett. 48, 1333

1990 Optical saturation of the $1^3 S \rightarrow 2^3 P$ transition in positronium

K. P. Ziock, C. D. Dermer, R. H. Howell, F. Magnotta, K.M. Jones, J. Phys. B 23, 329 First observation of resonant excitation of high—n states in positronium

K. P. Ziock, R. H. Howell, F. Magnotta, R. A. Failor, and K. M. Jones, Phys. Rev. Lett. 64, 2366

1993 Measurement of the positronium $1^3S \rightarrow 2^3S$ interval by continuous-wave two-photon excitation

M. S. Fee, S. Chu, A. P. Mills, Jr., R. J. Chichester, D. M. Zuckerman, E. D. Shaw, and K. Danzmann, Phys. Rev. A48, 192



- Low-energy positron beam
- Buffer-gas positron trap
- Radial compression and time focusing



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- Radial compression and time focusing
- Positronium Spectroscopy
 - Cool positronium production
 - REMPI spectroscopy via 1s2p (UV + Green)
 - Rydberg Ps (UV + IR)
 - Positronium Stark-states
 - Manipulation of Rydberg-Stark states of Ps



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- Colder positronium production



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- (Anti)gravity measurements with positronium



UCL Positronium Lab

January 2014





UCL Positronium Lab

October 2014





UCL Positronium Lab





Solid Neon Moderator



³E.M. Gullikson and A.P. Mills, Jr. (1986) Phys. Rev. Lett., 57:376



Buffer-gas Positron Trap





⁴C.M. Surko et al. (1988) Hyperfine Interact., 44:185



Rotating quadrupole electric field radially compresses the cloud and improves accumulation by reducing diffusion to the walls, however, cooling gas (CF_4) is needed to counter heating.





⁵A. Deller et al. (2014) New J. Phys., **16**:073028



Pulses applied to the electrodes ejects the positron cloud from the trap.

A fast PbF_2 Cerkenov radiator + PMT is used to measure time-width of the annihilation gamma-ray pulse generated when the positrons impact a solid target near the focus.





Fit function: $V(t) = \frac{1}{2} A \exp \frac{\sigma^2 - 2l\tau + 2x_0 \tau}{2\tau^2} \operatorname{erfc}(\frac{-l\tau + x_0 \tau + \sigma^2}{\sqrt{2}\sigma\tau})$ $\sigma \approx 4 \text{ ns (positron pulse width)}$

Positronium Production

Compressed and bunched positrons transported to target chamber.

Positrons attracted to target by electric potential and embedded into porous silica.

The positrons form Ps in the bulk, then diffuse through the pores to vacuum.

Conversion efficiency of \sim 30 %.





⁶L. Liszkay, et al. (2008) Appl. Phys. Lett., **92**(6):063114.



Target Mount





Target mount and $(PbWO_4 + PMT)$ gamma ray detector.



Target Mount



Target mount and (PbWO₄ + PMT) gamma ray detector.



Broadband UV Laser





243 nm (UV): 0.1 - 4.0 mJ $\Delta t = 6 \text{ ns}$ $\Delta v \sim 85 \text{ GHz}$ 532 nm (Green): 10 - 40 mJ

 $\Delta t = 6 \text{ ns}$

 $\Delta\nu\sim \text{20 GHz}$



Positronium Lifetime Spectroscopy

Single-Shot Positron Annihilation Lifetime Spectroscopy



⁷D. Cassidy, et al. (2006) Appl. Phys. Lett., 88(19):194105





⁸D. Cassidy, et al. (2006) Appl. Phys. Lett., 88(19):194105



1s2p Doppler Linewidth





Energy of the emitted Ps decreases with increasing e^+ implantation energy - minimum around 40 meV.

⁹D. B. Cassidy, et al. Phys. Rev. A (2014) 81:012715



Aarhus transmission films - Doppler shift





¹⁰S. L. Andersen, et al. Eur. Phys. J. D (2014) 68: 124



Time-of-Flight



¹¹A. Deller, et al. New. J. Phys. (2015) 17: 043059



Time-of-Flight



¹²A. Deller, et al. New. J. Phys. (2015) 17: 043059



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¹³A. Deller, et al. New. J. Phys. (2015) 17: 043059



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Time-of-Flight + Doppler



UCL

Positronium Spectroscopy

Magnetic Quenching

Exciting the Lyman- α transition in a magnetic field mixes the singlet and triplet spin states, which generally reduces the overall lifetime; ergo, a significant difference in SSPALS spectra can be observed using only the UV laser.

The mixing depends strongly on the magnetic field strength, the states populated within the 2P manifold (UV polarization), and also the electric field.





Magnetic Quenching







Magnetic Quenching







Magnetic Quenching







Magnetic Quenching





¹⁵Stephen M. Curry (1973) *Phys. Rev. A*, **7**:447



Magnetic Quenching





¹⁵Stephen M. Curry (1973) Phys. Rev. A, 7:447



Rydberg Positronium

Positronium can be excited to Rydberg (high-*n*) states using lasers.

"Traditional" excitation scheme (UV + IR): $1S \rightarrow 2P \rightarrow nS/D$.

Self-annihilation in Rydberg states is practically negligible. This makes feasible a range of interesting applications, including gravity measurements.



¹⁶K.P. Ziock, *et al.* (1990) *Phys. Rev. Lett.*, **64**:2366
¹⁷D. Cassidy, *et al.* (2012) *Phys. Rev. Lett.*, **108**:043401
¹⁸A. C. L. Jones, *et al.* (2014) *Phys. Rev. A*, **90**:012503



UV + IR Laser





243 nm (UV): 0.1 - 4.0 mJ $\Delta t = 6 \text{ ns}$ $\Delta v \sim 85 \text{ GHz}$ 730 - 750 nm (IR): 1 - 15 mJ $\Delta t = 6 \text{ ns}$ $\Delta v \sim 5 \text{ GHz}$



Rydberg Ps (n = 11)



¹⁹T. Wall, et al. (2015) Phys. Rev. Lett., 114:173001



Rydberg Ps (n = 11)



²⁰T. Wall, et al. (2015) Phys. Rev. Lett., 114:173001



Rydberg Ps



²¹T. Wall, et al. (2015) Phys. Rev. Lett., 114:173001



Rydberg Ps



²¹T. Wall, et al. (2015) Phys. Rev. Lett., 114:173001



Rydberg Stark Filter (n = 18)



²²T. Wall, et al. (2015) Phys. Rev. Lett., 114:173001



Rydberg Stark states Ps (n = 11)



²³T. Wall, et al. (2015) Phys. Rev. Lett., 114:173001



Stark Deceleration of Ps*

Rydberg Ps beams can be decelerated, deflected or focused with electric field gradients.

Ps: $\mu_{\textit{max}} = \frac{3}{2}\textit{n}^2\textit{e}(2\textit{a}_0)$





²⁴S. Hogan, and D. Cassidy (2014) Int. J. Mod. Phys. Conf. Ser. 30, 1460259
²⁵S. Hogan, et al. (2012) Phys. Rev. Lett. 108, 063008



n = 3

Stark Deceleration of Ps*

Rydberg Ps beams can be decelerated, deflected or focused with electric field gradients.

Ps: $\mu_{max} = \frac{3}{2}n^2e(2a_0)$



(a)

^{1.0}Γ(vi)

^{1.0}[(v)

²⁶S. Hogan, and D. Cassidy (2014) Int. J. Mod. Phys. Conf. Ser. **30**, 1460259 ²⁷S. Hogan, et al. (2012) Phys. Rev. Lett. 108, 063008



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- Still to do:
 - Manipulation of Rydberg-Stark states of Ps (ask Stephen)
 - Colder positronium production
 - (Anti)gravity measurements with positronium



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UCL Positronium Group (II)





Thank you for your attention!

http://antimattergravity.com/



a.deller@ucl.ac.uk

