

# Central exclusive production at CMS: recent results and future prospects with the CT-PPS



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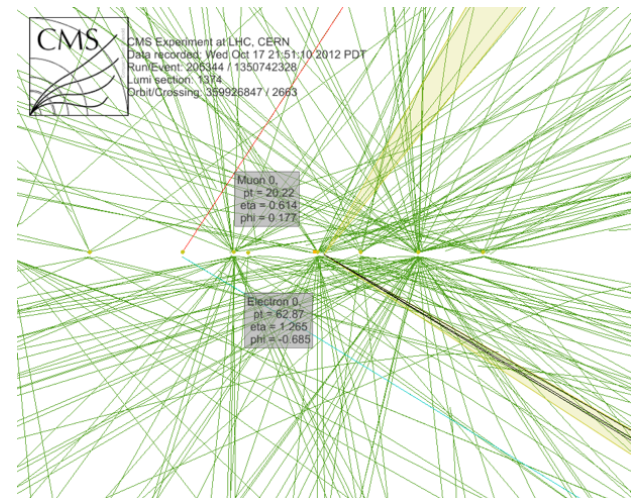
Università di Torino, INFN Torino

*on behalf of the CMS and TOTEM collaborations*



## EDS Blois 2015 : the 16th conference on Elastic and Diffractive scattering

*Borgo, Corsica, 2 July 2015*



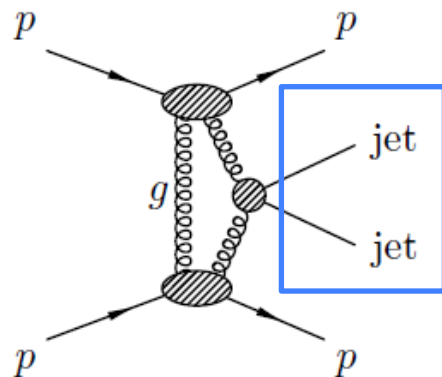
# Outline

- Central Exclusive Production
- Exclusive  $\gamma\gamma \rightarrow W^+W^-$  production and anomalous quartic gauge couplings:  
new CMS results at 8 TeV
- CEP at CMS with proton taggers (CT-PPS):
  - Exclusive WW production
  - Exclusive dijet production

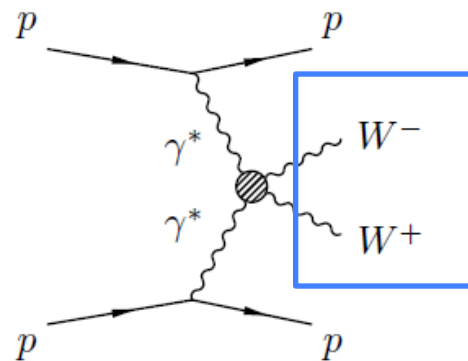
# Central Exclusive Production

**CEP:** process of type  $pp \rightarrow p+X+p$

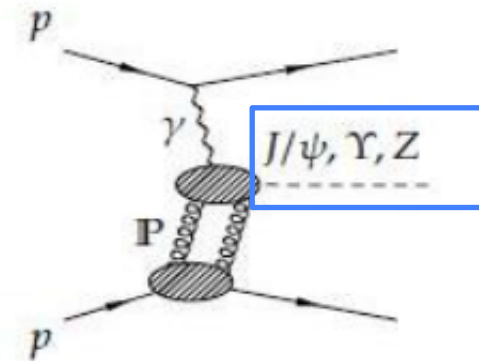
- **Scattered protons survive the collision intact**
- No additional activity between the outgoing protons and the **central system X** (exclusive)
- **Particularly clean experimental conditions** thanks to the absence of proton remnants



**PP exchanges**



**$\gamma\gamma$  interactions**



**P $\gamma$  fusions**

- ✓ Double pomeron exchange: pomeron structure → [Marta Ruspa's talk](#)
- ✓  $W+W^-$ : study of exclusive processes at high mass and constraint anomalous couplings
- ✓  $l+l^-$ : compare to precision QED predictions, calibrate integrated luminosity normalisation and study proton dissociation
- ✓  $J/\psi$  and  $Y$  photo-production: probe gluon distribution at low  $x$

# CMS RESULTS

- Search for exclusive or semi-exclusive  $\gamma\gamma$  production and observation of exclusive or semi-exclusive  $e^+ e^-$  production in pp collisions at  $\sqrt{s} = 7$  TeV [ [JHEP 11 \(2012\) 080](#) ]
- Exclusive  $\gamma\gamma \rightarrow \mu^+ \mu^-$  production in proton-proton collisions at  $\sqrt{s} = 7$  TeV [ [JHEP 01 \(2012\) 052](#) ]

Based on  
 $\sim 40 \text{ pb}^{-1}$   
collected  
in 2010

- Study of exclusive two-photon production of  $W^+W^-$  in pp collisions at  $\sqrt{s} = 7$  TeV [ [JHEP 07 \(2013\) 116](#) ]

*New*

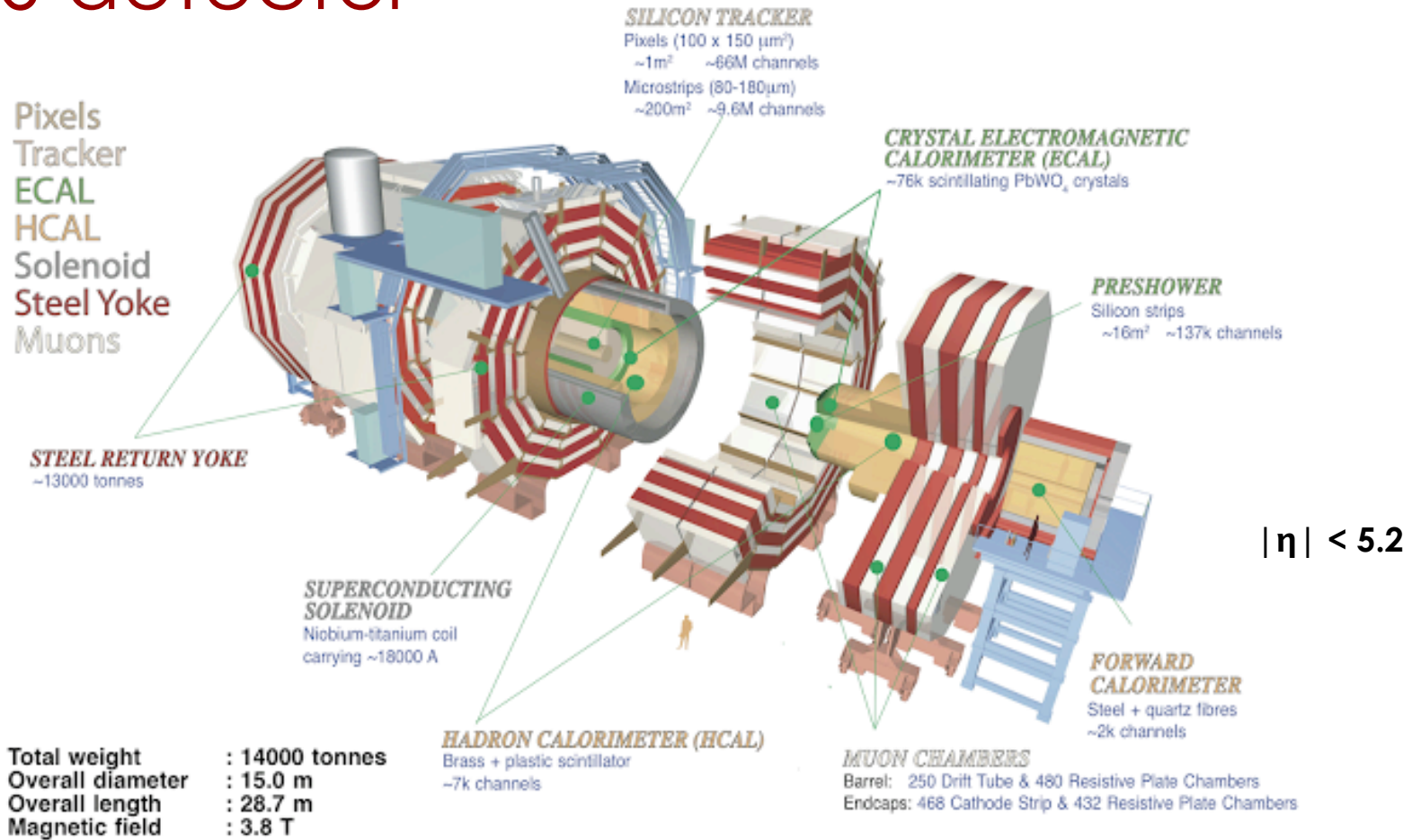
- Study of exclusive two-photon production of  $W^+W^-$  in pp collisions at  $\sqrt{s} = 8$  TeV [ [CMS PAS FSQ-13-008](#) ]

$\sqrt{s}$ (TeV)	Luminosity ( $\text{fb}^{-1}$ )	Pileup (PU)
7	5	7
8	19.7	21

**Results from 8 TeV data analysis shown here**



# CMS detector



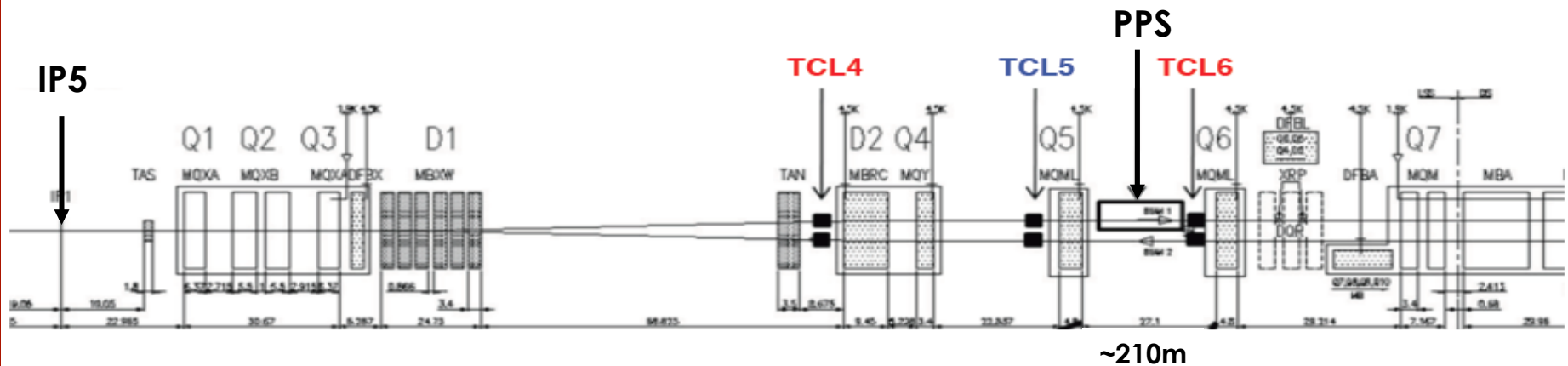
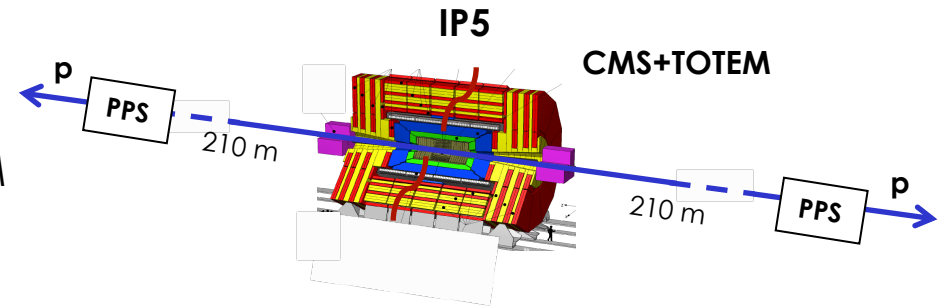
- **BSC, Beam Scintillator Counters** :  $3.2 < |\eta| < 4.7$  (in front of the Forward Calorimeter)
- **CASTOR calorimeter**:  $-6.6 < |\eta| < -5.2$  (14.4 m from IP, one side only )
- **Forward Shower Counters FSC**:  $6 < |\eta| < 8$  (59-114 m from IP )
- **Zero Degree Calorimeter**:  $|\eta| > 8.1$  (140 m from IP)

# Forward proton tagging: CT-PPS

The **CMS-TOTEM Precision Proton Spectrometer (CT-PPS)** [1] will allow precision proton measurements in the very forward regions on both sides of CMS during standard LCH running:

- Two stations for **tracking detectors** and two stations for **timing detectors** installed at ~210 m from the common CMS-TOTEM interaction point (IP5) on both sides of the central apparatus
- LHC magnets between IP5 and the detector stations used to bend out of the beam envelope protons that have lost a small fraction of their initial momentum in the interaction

→ **proton fractional longitudinal momentum loss ( $\xi$ ) between 2% and 10%**



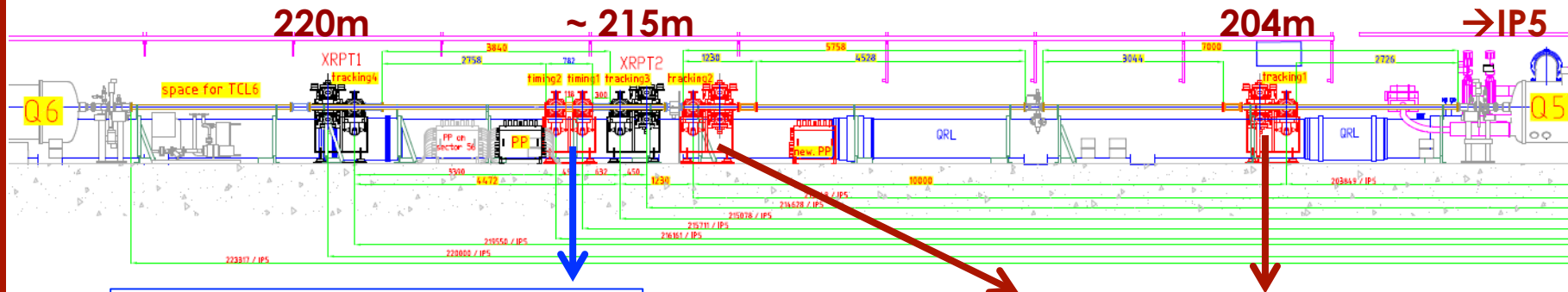
LHC lattice between IP5 and CT-PPS detector stations

[1] CT\_PPS TDR, CERN-LHCC-2014-021

# CT-PPS design performance and CEP

The measurement of the two scattered protons fully determines the kinematics of the central system X, irrespective of its decay mode.

$$M_X = \sqrt{s \cdot \xi_1 \cdot \xi_2}$$

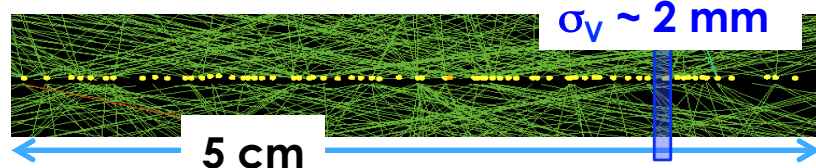


2 new horizontal cylindrical pots equipped with timing detectors

2 horizontal rectangular pots equipped with tracking detectors

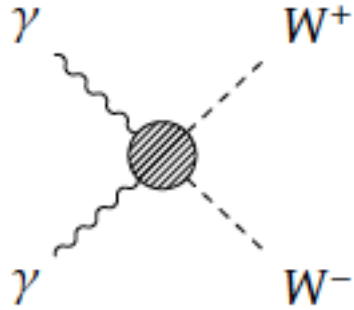
Proton timing measurement from both sides of CMS allows to **determine the primary vertex**, correlate it with that of the central detector and **reject pile-up**

- Time resolution  $\sim 10$  ps  
 → **Vertex z-by-timing:  $\sim 2$  mm**



Proton position and angle measurements, combined with the beam magnets, allow to **determine the momentum of the scattered protons**

- Position resolution of  $\sim 10 \mu\text{m}$
- Angular resolution of  $\sim 1-2 \mu\text{rad}$   
 →  $\Delta p/p \sim 2 \cdot 10^{-4}$   
**Mass resolution:  $\sim 5 \text{ GeV}/c^2$**



## CEP production of $W$ pairs

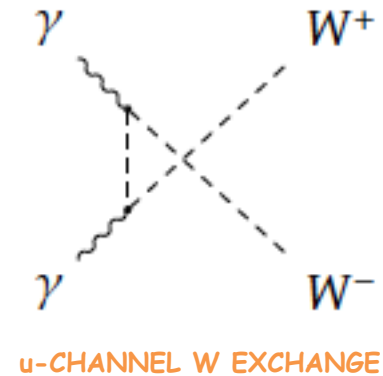
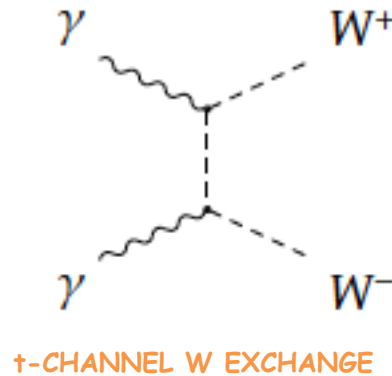
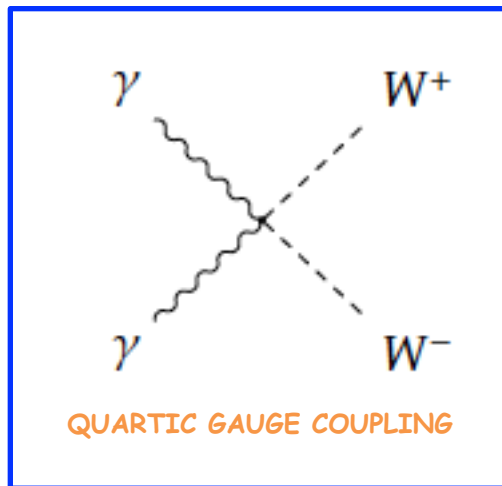
# Theoretical framework

## STANDARD MODEL PRODUCTION

SM Lagrangian density contains triple and quartic couplings between  $\gamma$  and W bosons

$$\begin{aligned}\mathcal{L}^{WW\gamma} &= -ie F_{\mu\nu} W^{+\mu} W^{-\nu} \\ \mathcal{L}^{WW\gamma\gamma} &= -e^2 (W_{\mu}^{+} W^{-\mu} A_{\nu} A^{\nu} - W_{\mu}^{+} W_{\nu}^{-} A^{\mu} A^{\nu})\end{aligned}$$

SM contribution to  $\gamma\gamma \rightarrow W^{+}W^{-}$  at leading order:



Measurements of the quartic  $WW\gamma\gamma$  coupling can be used to look for any deviation from the SM predictions, which would reveal a sign of new physics.

# Theoretical framework

## GENUINE ANOMALOUS QUARTIC COUPLINGS (AQGCs)

Potential deviation from SM can be quantified by introducing genuine anomalous quartic couplings (AQGC) of dimension-6 not related to SM triple or quartic couplings [1]

AQGCs introduced via an effective Lagrangian with 2 additional dimension-6 terms:

$$\begin{aligned}\mathcal{L}_6^0 &= \frac{e^2 a_0^W}{8 \Lambda^2} F_{\mu\nu} F^{\mu\nu} W^{+\alpha} W_{\alpha}^{-} - \frac{e^2}{16 \cos^2 \Theta_W} \frac{a_0^Z}{\Lambda^2} F_{\mu\nu} F^{\mu\nu} Z^{\alpha} Z_{\alpha} \\ \mathcal{L}_6^C &= \frac{-e^2 a_C^W}{16 \Lambda^2} F_{\mu\alpha} F^{\mu\beta} (W^{+\alpha} W_{\beta}^{-} + W^{-\alpha} W_{\beta}^{+}) - \frac{e^2}{16 \cos^2 \Theta_W} \frac{a_C^Z}{\Lambda^2} F_{\mu\alpha} F^{\mu\beta} Z^{\alpha} Z_{\beta}\end{aligned}$$

containing:

- parameters  $a_0^W$  and  $a_C^W$
- $\Lambda$  scale for new physics

A nonlinear representation of the spontaneously broken  $SU(2) \otimes U(1)$  symmetry assumed

[1] G. Belanger and F. Boudjema, "Probing quartic couplings of weak bosons through three vectors production at a 500-GeV NLC", Phys.Lett. B288 (1992) 201–209,

# Theoretical framework

## DIMENSION-8 FORMALISM

Assuming a linear representation of the spontaneously broken  $SU(2) \otimes U(1)$  symmetry, the lowest-order operators, where new physics may cause deviations in the purely quartic gauge boson couplings, are of dimension 8 [2].

Including the constraint that the  $WWZ\gamma$  vertex should vanish, a direct relationship between the dimension-8  $f_{M,0,1,2,3}/\Lambda^4$  couplings and the dimension-6  $a_{0,C}^W/\Lambda^2$  couplings is recovered

$$\frac{a_0^W}{\Lambda^2} = -\frac{4M_W^2 f_{M,0}}{g^2 \Lambda^4} - \frac{8M_W^2 f_{M,2}}{g'^2 \Lambda^4}$$

$$\frac{a_C^W}{\Lambda^2} = \frac{4M_W^2 f_{M,1}}{g^2 \Lambda^4} + \frac{8M_W^2 f_{M,3}}{g'^2 \Lambda^4}$$

with  $g = \frac{e}{\sin(\theta_W)}$  and  $g' = \frac{e}{\cos(\theta_W)}$

[2] M. Baak et al., "Working Group Report: Precision Study of Electroweak Interactions", arXiv:1310.6708.



# Theoretical framework

## DIPOLE FORM FACTOR

In both the dimension-6 and dimension-8 scenarios, the  $\gamma\gamma \rightarrow WW$  cross section in the presence of anomalous couplings increases rapidly with the  $\gamma\gamma$  center of mass energy ( $W_{\gamma\gamma}$ ).

→ dipole form factors introduced to preserve unitarity

$$a_{0,C}^W \rightarrow a_{0,C}^W(W_{\gamma\gamma}^2) = a_{0,C}^W \left( 1 + \frac{W_{\gamma\gamma}^2}{\Lambda_{\text{cutoff}}^2} \right)^{-2}$$

$\Lambda_{\text{cutoff}}$  : energy cutoff scale

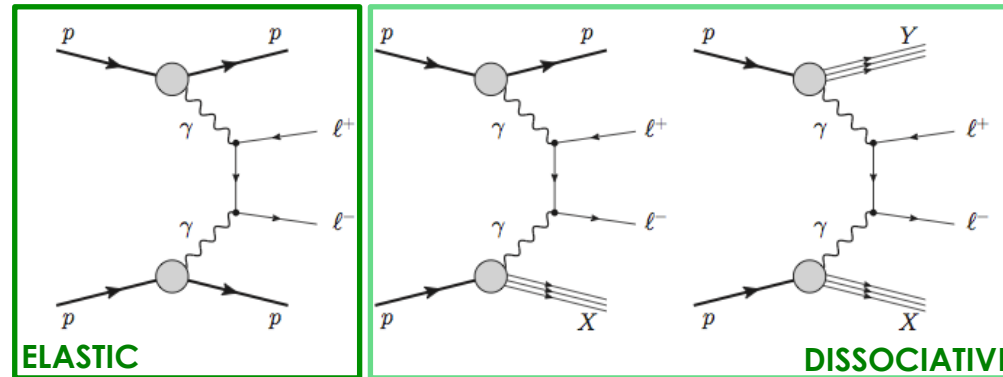
**Two scenarios considered:**

- $\Lambda_{\text{cutoff}} = 500 \text{ GeV}$
- $\Lambda_{\text{cutoff}} \rightarrow \infty$  (no form factor)

# Signal and background contributions

## Two contributions to signal:

- elastic production:  $pp \rightarrow pW^+W^-p$
- proton-dissociative production:  $pp \rightarrow p^*W^+W^-p^*$ ,  
(one or both protons dissociate into a low-mass system that escapes detection)



Can not be separated without tagging the forward protons

## Background sources:

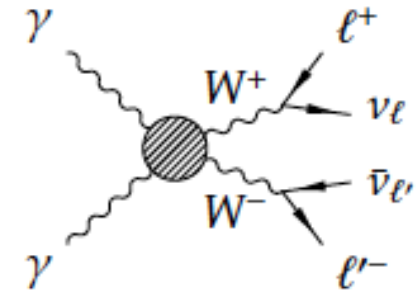
- Inclusive  $WW$ ,  $t\bar{t}$ ,  $W$ +jets processes
- $\tau\tau$  pairs produced via the Drell-Yan process
- Exclusive two photon processes:  $\gamma\gamma \rightarrow l\bar{l}$
- $WW$  production from single diffractive interactions
- $WW \rightarrow WW$  scattering (vector boson fusion)

All backgrounds estimated from Monte Carlo after comparison with data in control regions

# Event selection

**Unlike-flavor dilepton decay channel:**  $\gamma\gamma \rightarrow W^+W^- \rightarrow \mu^\pm e^\mp \nu\bar{\nu}$

(backgrounds due Drell-Yan and two-photon  $l^+l^-$  production more than an order of magnitude lower than in the same-flavor final states)



## Selection:

- Opposite sign leptons, with  $p_T > 20$  GeV and  $|\eta| < 2.4$ , matched to common primary vertex
- No extra tracks associated to the dilepton vertex
- Dilepton invariant mass greater than 20 GeV



## SM signal region:

$$p_T(\mu e) > 30 \text{ GeV}$$



## AQGCs search:

- 2 bins:  $p_T(\mu e) = 30-130$  GeV  
 $p_T(\mu e) > 130$  GeV

# Signal correction factors

**Two effects present in data difficult to describe in simulation:**

- mis-association of low  $p_T$ , forward tracks from pileup events to dilepton vertex
  - Efficiency of the “zero extra-tracks” cut overestimated in MC
  - Important effect in 8 TeV analysis (PU=21)
- proton dissociation contribution to  $\gamma\gamma$  interactions

**⇒  $\gamma\gamma \rightarrow ll$  control samples used to derive correction factors from data**

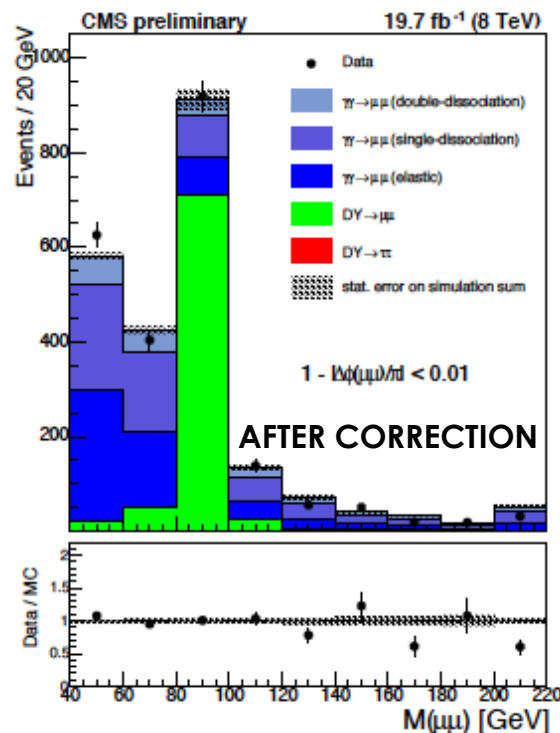
# Zero extra-tracks efficiency correction

Control region dominated by elastic  $\gamma\gamma \rightarrow ll$  sample used to estimate the correction to the efficiency of the zero extra-tracks cut

Requiring:

- dilepton vertex with 0 extra tracks
- Z-mass veto
- leptons back-to-back ( $1 - |\Delta\phi(ll)|/\pi < 0.01$ )

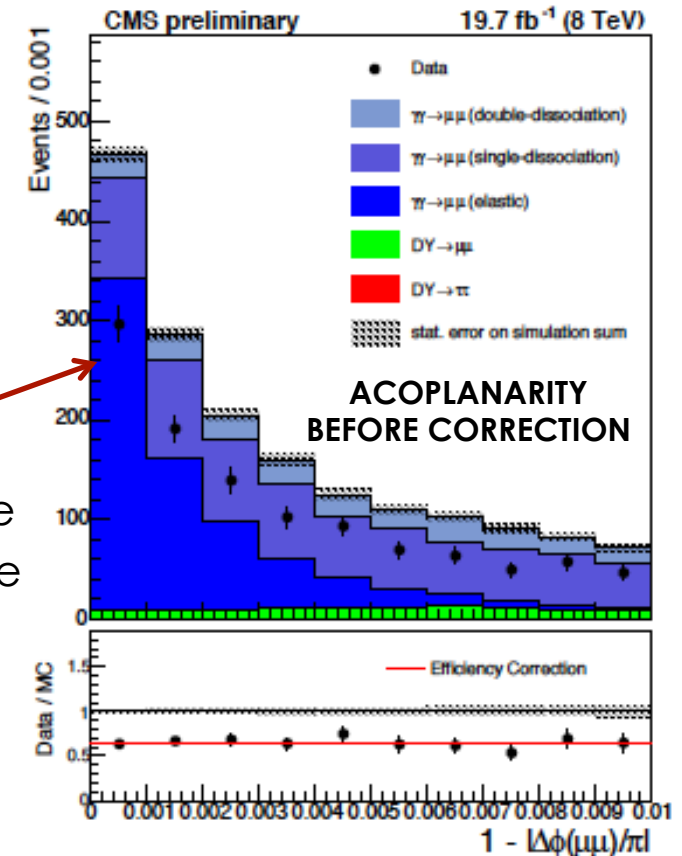
a deficit in the number of events found in data with respect to simulation



Data/MC in first bin:  
 $0.63 \pm 0.04$  in  $\gamma\gamma \rightarrow \mu\mu$  sample  
 $0.63 \pm 0.07$  in  $\gamma\gamma \rightarrow ee$  sample

This ratio, averaged over the  $\mu\mu$  and  $ee$  samples, applied to correct the efficiency derived on MC.

**Good data/MC agreement after correction**  
 (checked on several kinematic distributions)



# Proton dissociation factor

- Separation of elastic and proton-dissociative events in data not possible without proton tagging
- Simulation of single or double dissociation production difficult (involves soft interactions, only phenomenological models are available)
  - Signal MC samples are purely elastic
- **Strategy:** single/double proton dissociation contribution to  $\gamma\gamma \rightarrow WW$  estimated from the data is used to correct MC

Control samples:  $\gamma\gamma \rightarrow \mu\mu$  and  $\gamma\gamma \rightarrow ee$

- $m(l\bar{l}) > 160 \text{ GeV}$
- no extra tracks associated to dilepton vertex

→ Proton dissociation factor 
$$F = \left[ \frac{N_{ll \text{ data}} - N_{DY}}{N_{elastic}} \right]_{m(l^+l^-) > 160 \text{ GeV}}$$

$$F = 4.10 \pm 0.43$$

Stable as a function of the mass cut in both the dimuon and dielectron channels

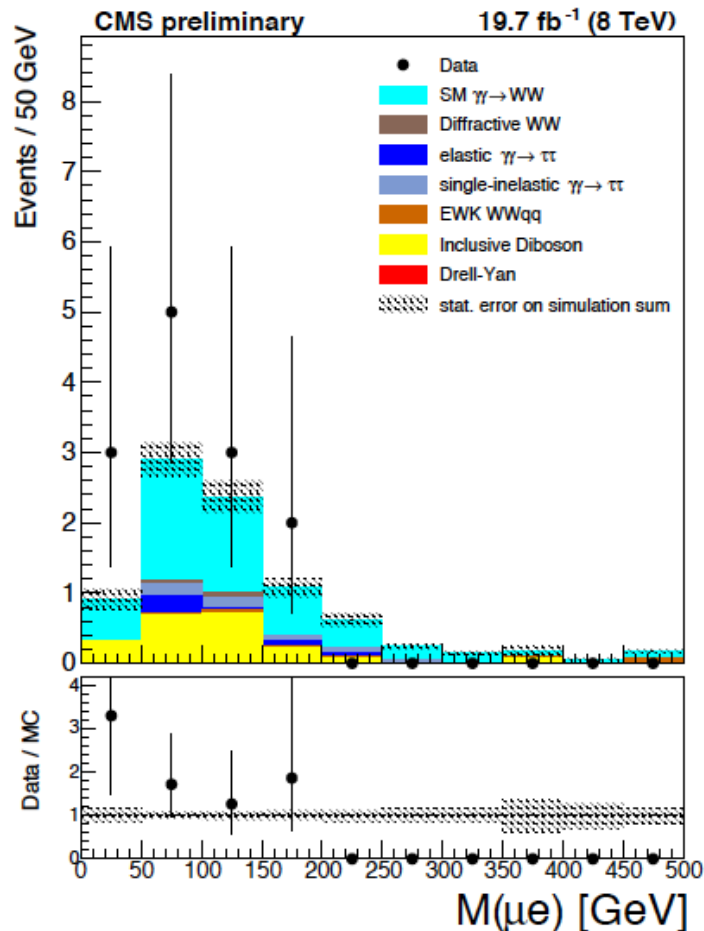
**F used to correct the elastic  $pp \rightarrow pW^+W^-p$  prediction to the total  $pp \rightarrow p^{(*)}W^+W^-p^{(*)}$  prediction, including proton dissociation**

# Results – SM signal region

## Number of expected signal and background events

Signal	All backgrounds	Inclusive WW	$\gamma\gamma \rightarrow \tau\tau$	DY $\rightarrow \tau\tau$	Diff. WW	Others
<b>5.3<math>\pm</math>0.1</b>	<b>3.5<math>\pm</math>0.5</b>	<b>2.0<math>\pm</math>0.4</b>	<b>0.9<math>\pm</math>0.2</b>	<b>0</b>	<b>0.1<math>\pm</math>0.1</b>	<b>0.3<math>\pm</math>0.2</b>

Mean expected signal significance: **2.4  $\pm$  0.5  $\sigma$**



**13 events are observed in data that pass all selection criteria**

**Properties of selected events ( $\mu e$  invariant mass, acoplanarity, missing transverse energy) consistent with the SM signal plus background prediction**

**Cross section measurement:**

$$\sigma(pp \rightarrow p^{(*)}W^+W^-p^{(*)} \rightarrow p^{(*)}\mu^\pm e^\mp p^{(*)}) = 12.3^{+5.5}_{-4.4} \text{fb}$$

**SM prediction: 6.9  $\pm$  0.6 fb**

**Observed significance above the background-only hypothesis: 3.6  $\sigma$**  (systematic uncertainties included)



# Results – anomalous couplings (I)

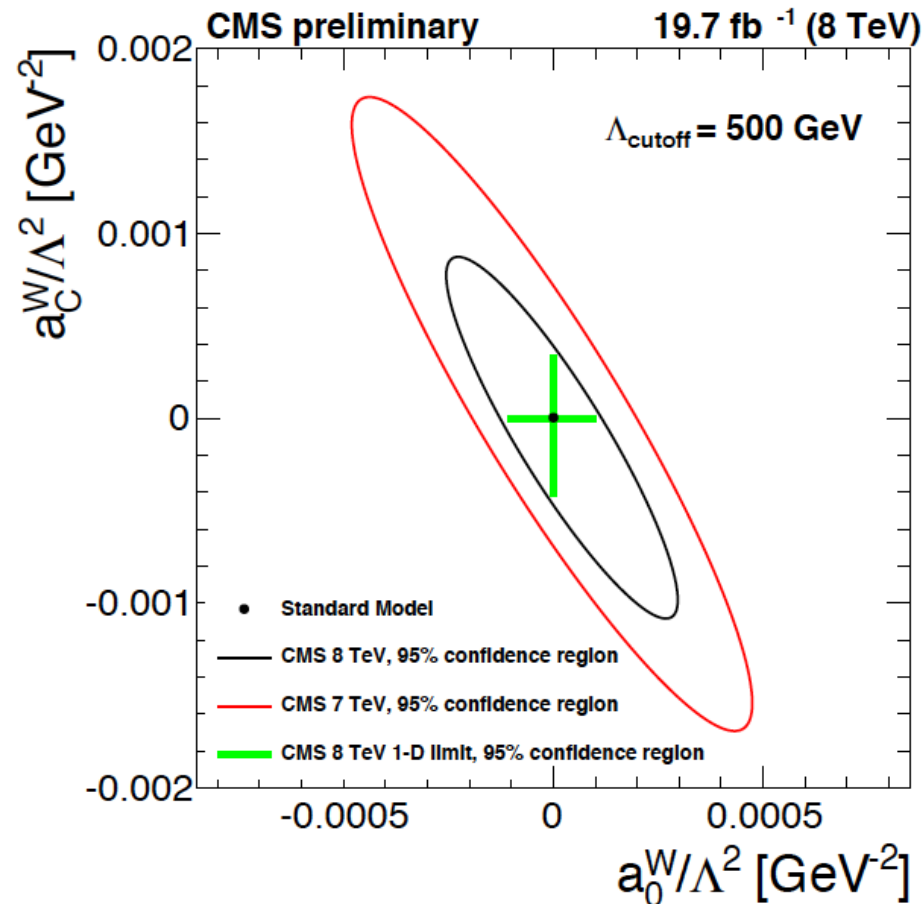
- **No significant deviation from SM observed** in the  $p_T(\mu e)$  distribution
- **95% CL limits on  $a_0^W/\Lambda^2$  and  $a_c^W/\Lambda^2$**  derived using Feldman-Cousins prescription with systematic uncertainties treated as log-normal nuisance parameters, for 2 scenarios:

		OPAL (2004)	DØ (2013)	CMS (2013)	CMS(2015)
$a_0^W / \Lambda^2 [GeV^{-2}]$	no form factor	$\pm 2 \times 10^{-2}$	$\pm 4.3 \times 10^{-4}$	$\pm 4.0 \times 10^{-6}$	$\pm 1.2 \times 10^{-6}$
	$\Lambda_{\text{cutoff}} = 500 \text{ GeV}$		$\pm 2.5 \times 10^{-3}$	$\pm 1.5 \times 10^{-4}$	$(-1.1 - 1.0) \times 10^{-4}$
$a_c^W / \Lambda^2 [GeV^{-2}]$	no form factor	$^{+3.7}_{-5.2} \times 10^{-2}$	$\pm 1.5 \times 10^{-3}$	$\pm 1.5 \times 10^{-5}$	$\pm 4.4 \times 10^{-6}$
	$\Lambda_{\text{cutoff}} = 500 \text{ GeV}$		$\pm 9.2 \times 10^{-3}$	$\pm 5.0 \times 10^{-4}$	$(-4.2 - 3.4) \times 10^{-4}$
				$\sqrt{s} = 7 \text{ TeV}$	$\sqrt{s} = 8 \text{ TeV}$

- ✓ Up to **two orders of magnitude improvement** to limits set by previous experiments
- ✓  $\Lambda_{\text{cutoff}} = 500 \text{ GeV}$  → **8 TeV limit ~25% better than 7 TeV limit**
- ✓ **No form factor** → **8 TeV limit ~3 times better than 7 TeV limit**

# Results – anomalous couplings (II)

Similar statistical procedure used to derive **two dimensional limits** in the  $a_0^W / \Lambda^2$ ,  $a_C^W / \Lambda^2$  parameter space with  $\Lambda_{\text{cutoff}} = 500 \text{ GeV}$ .



The area outside the solid contour excluded by this measurement at 95% CL

# Results – anomalous couplings (III)

The transformed dimension-8 limits derived from  $a^{W_0}/\Lambda^2$  and  $a^{W_C}/\Lambda^2$ :

## With Form Factor, cutoff = 500 GeV

$$-4.2 \times 10^{-10} < f_{M,0} / \Lambda^4 < 3.8 \times 10^{-10} \text{ GeV}^{-4}$$

$$-16 \times 10^{-10} < f_{M,1} / \Lambda^4 < 13 \times 10^{-10} \text{ GeV}^{-4}$$

$$-2.1 \times 10^{-10} < f_{M,2} / \Lambda^4 < 1.9 \times 10^{-10} \text{ GeV}^{-4}$$

$$-8.0 \times 10^{-10} < f_{M,3} / \Lambda^4 < 6.6 \times 10^{-10} \text{ GeV}^{-4}$$

## No Form Factor

$$-4.6 \times 10^{-12} < f_{M,0} / \Lambda^4 < 4.6 \times 10^{-12} \text{ GeV}^{-4}$$

$$-17 \times 10^{-12} < f_{M,1} / \Lambda^4 < 17 \times 10^{-12} \text{ GeV}^{-4}$$

$$-2.3 \times 10^{-12} < f_{M,2} / \Lambda^4 < 2.3 \times 10^{-12} \text{ GeV}^{-4}$$

$$-8.5 \times 10^{-12} < f_{M,3} / \Lambda^4 < 8.5 \times 10^{-12} \text{ GeV}^{-4}$$

**7-16 times more stringent than previous CMS limits on 8-dimensional parameters:**

“A Search for  $WW\gamma$  and  $WZ\gamma$  production in pp Collisions at  $\sqrt{s} = 8 \text{ TeV}$ “

[CMS-PAS-SMP-13-009]

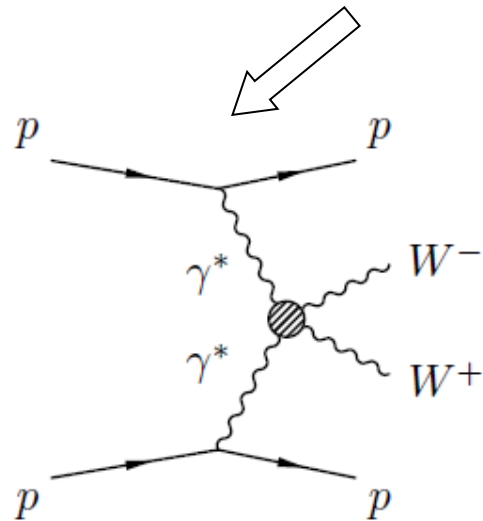
Observed Limits
$-77 \text{ (TeV}^{-4}) < f_{M,0} / \Lambda^4 < 81 \text{ (TeV}^{-4})$
$-131 \text{ (TeV}^{-4}) < f_{M,1} / \Lambda^4 < 123 \text{ (TeV}^{-4})$
$-39 \text{ (TeV}^{-4}) < f_{M,2} / \Lambda^4 < 40 \text{ (TeV}^{-4})$
$-66 \text{ (TeV}^{-4}) < f_{M,3} / \Lambda^4 < 62 \text{ (TeV}^{-4})$

# CEP in Run2 with CT-PPS

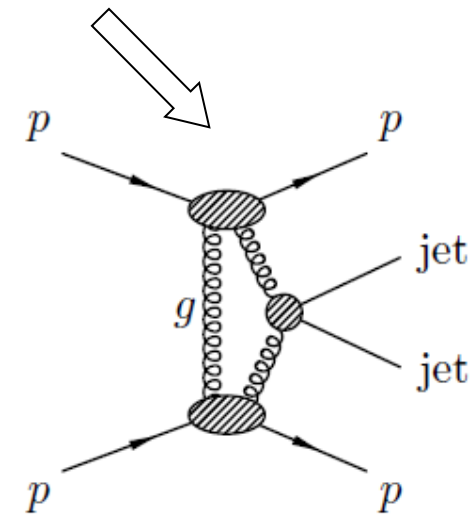
With the integrated luminosity expected in Run2 and the CT-PPS detector, the experimental reach of CMS can be extended. It will be possible to:

- improve current limits (e.g. on AQGC) by several orders of magnitude
- study new processes
- search for new physics

For the CT-PPS TDR two processes studied in detail



**Exclusive WW production**

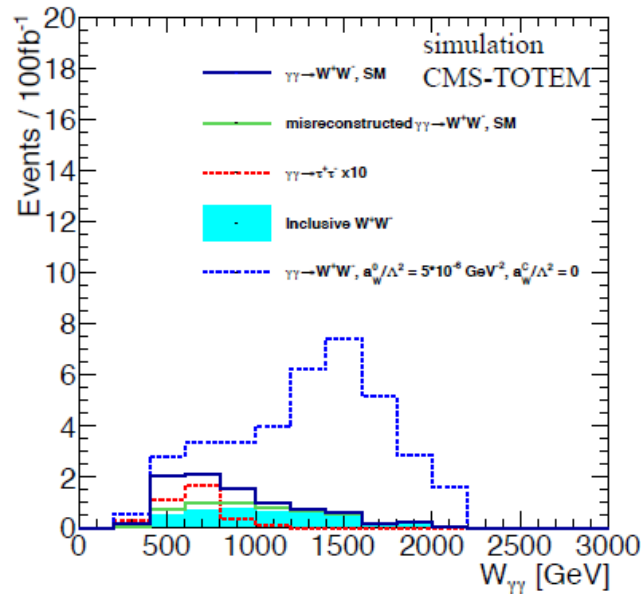
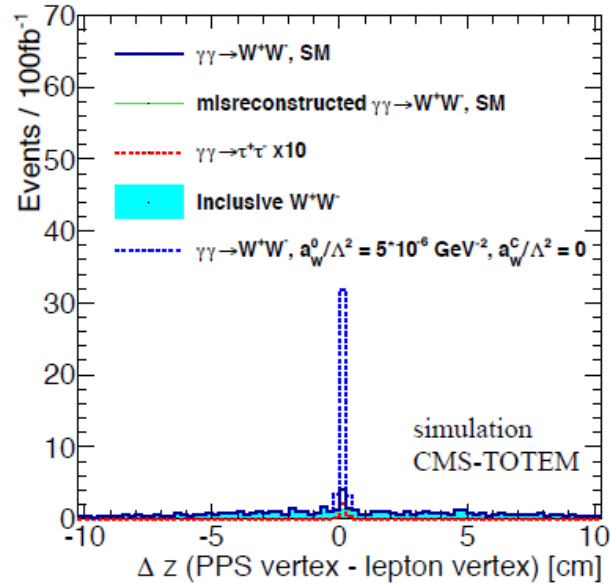


**Exclusive dijet production**

For these studies events generated at  $\sqrt{s}=13$  TeV

# Exclusive WW production (TDR study)

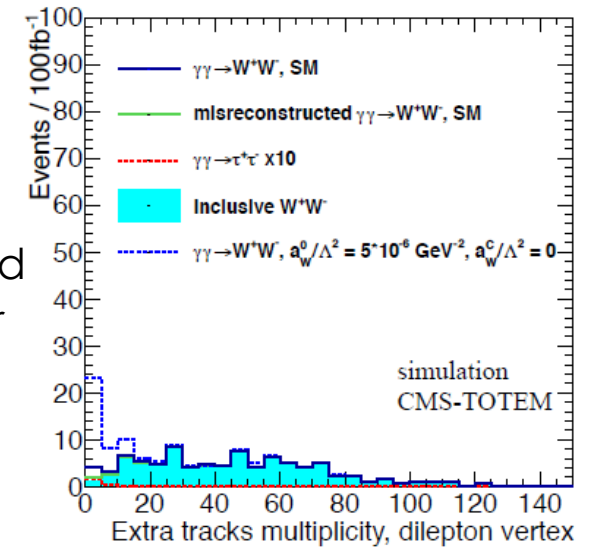
Study only  $e\mu$  final state



**Proton timing** is a powerful tool **to reject background**

Timing resolution of **10 ps** is assumed

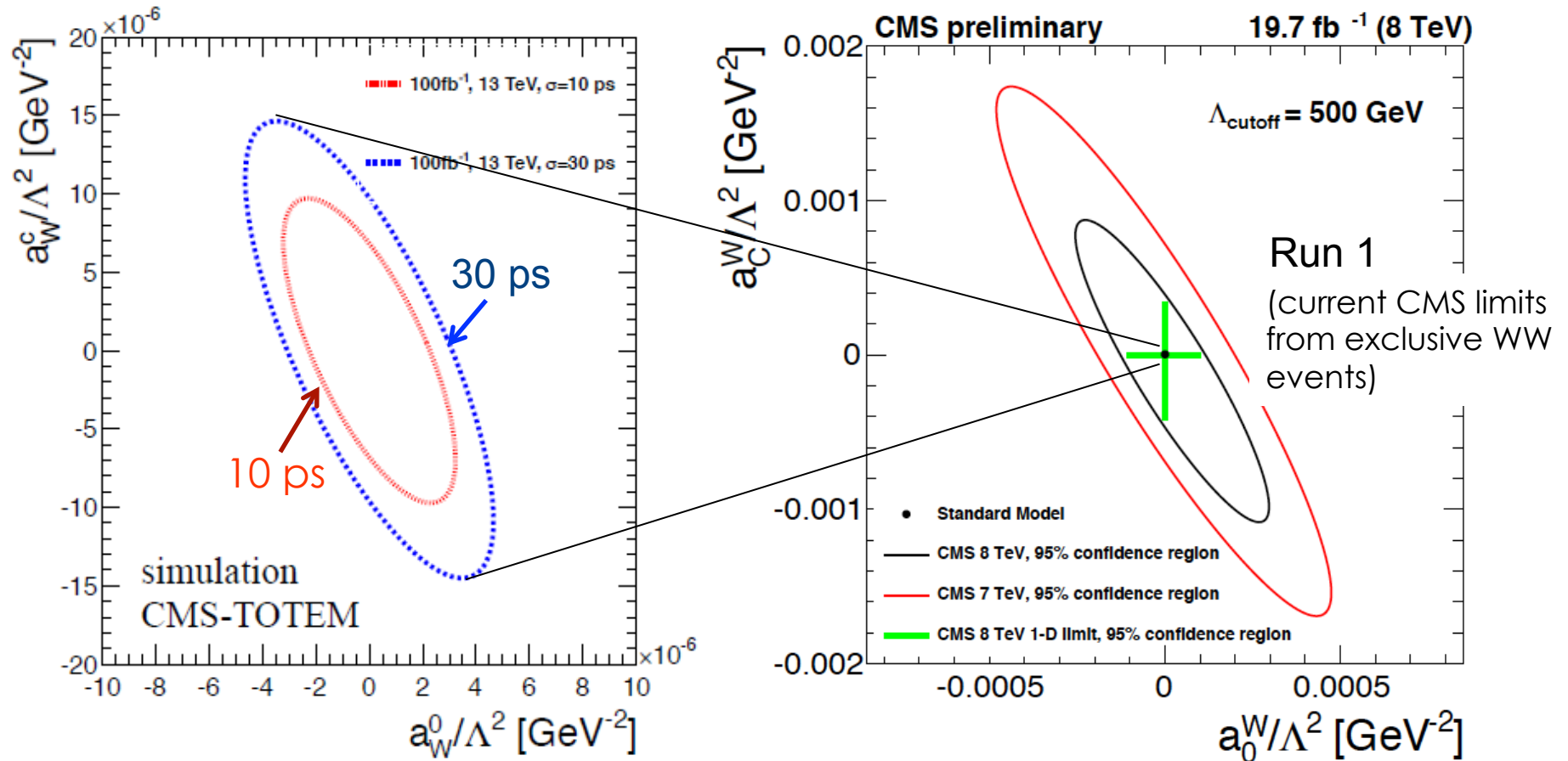
Track multiplicity associated to the dilepton vertex after the timing selection cut



Tail of the  $W_{\gamma\gamma}$  distribution ( $W_{\gamma\gamma} = M_X = \sqrt{s \cdot \xi_1 \cdot \xi_2} > 1 \text{ TeV}$ ), where SM contribution is expected to be small, provides a very **clear separation of AQC events**

# AQGC expected limits with CT-PPS

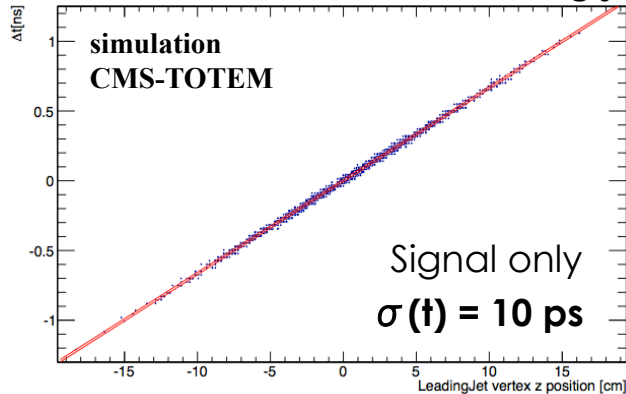
Expected two dimensional limits in the  $a_W^0 / \Lambda^2$ ,  $a_C^W / \Lambda^2$  parameter space @95%CL:



# Exclusive dijet production (TDR study)

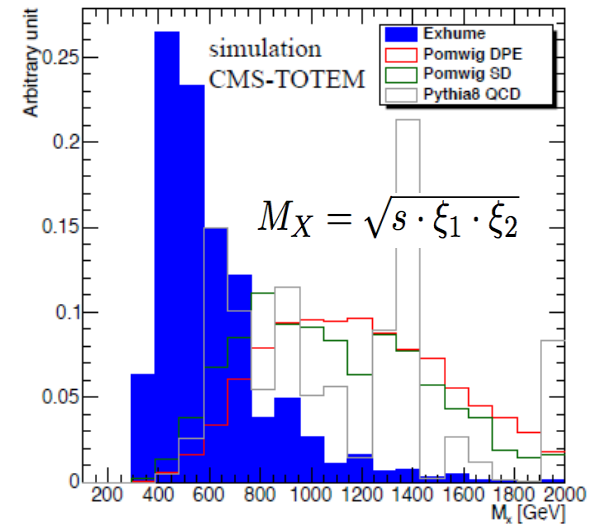
Require 2 jets with  $p_T > 100$  GeV and  $|\eta| < 2$

## PPS vertex vs vertex of leading jet



**Proton timing** as powerful discriminant for exclusive states:

- retain  $\sim 66\%$  ( $1\sigma$ ) of signal
- rejects  $\sim 70\%$  of QCD bkg



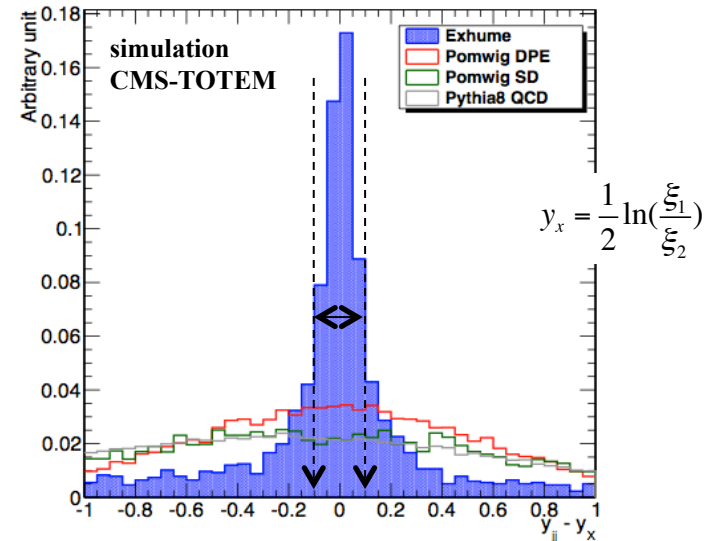
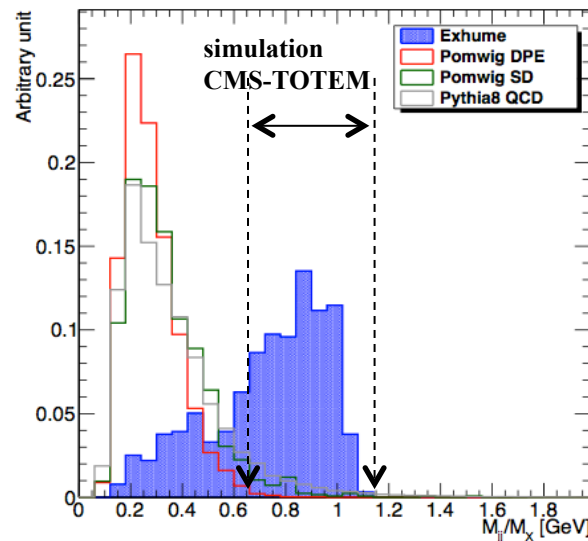
Exclusivity cuts

## Rapidity difference:

$y_{jj}$  (central detector)  $- y_x$  (PPS)

## Dijet mass fraction:

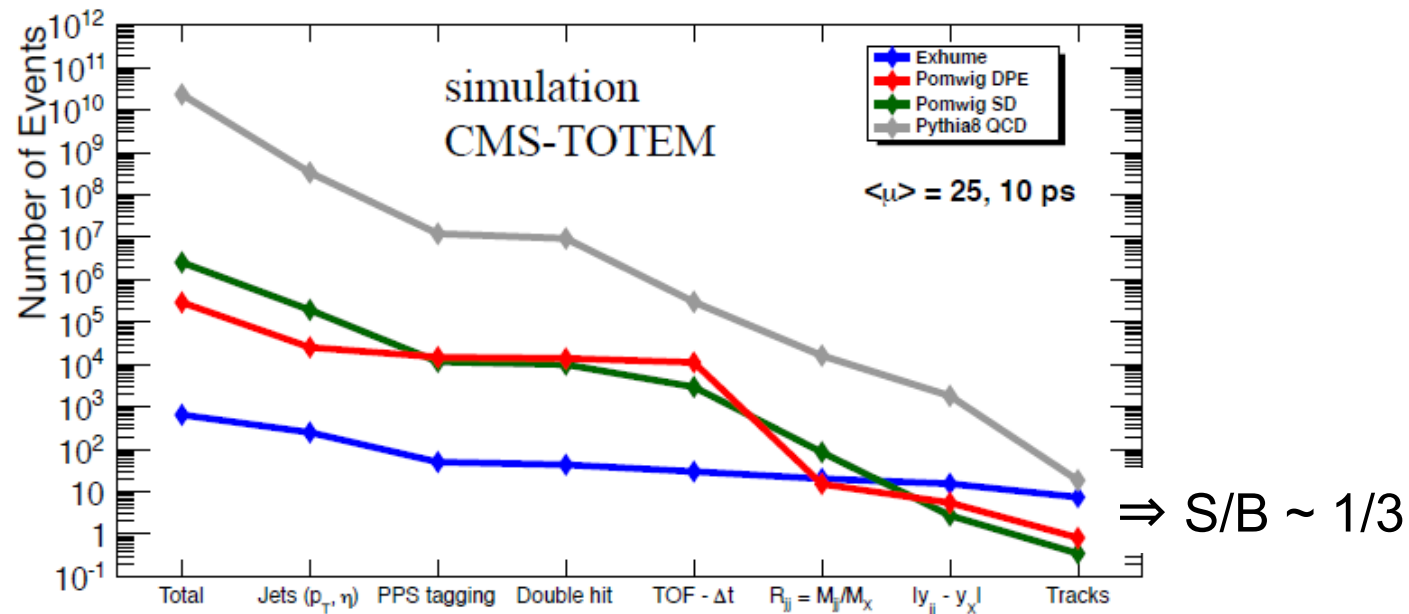
$$\frac{M_{jj} \text{ (central detector)}}{M_X \text{ (PPS)}}$$





# Expected dijet yields (1 fb<sup>-1</sup> - Pileup=25)

Event yields for signal and background processes as a function of the cuts applied.



	Exclusive dijets	DPE	SD	Inclusive dijets	S:B
<b>pileup <math>\mu = 25</math></b>					
$M_X \leq 500 \text{ GeV}$	$4.0 \pm 0.2$	$0.2 \pm 0.1$	$0 \pm 1$	$1 \pm 1$	3:1
$500 < M_X \leq 800 \text{ GeV}$	$3.1 \pm 0.2$	$0.3 \pm 0.1$	$0 \pm 1$	$15 \pm 1$	1:5
$M_X > 800 \text{ GeV}$	$0.3 \pm 0.1$	$0.3 \pm 0.1$	$1 \pm 1$	$4 \pm 1$	1:18

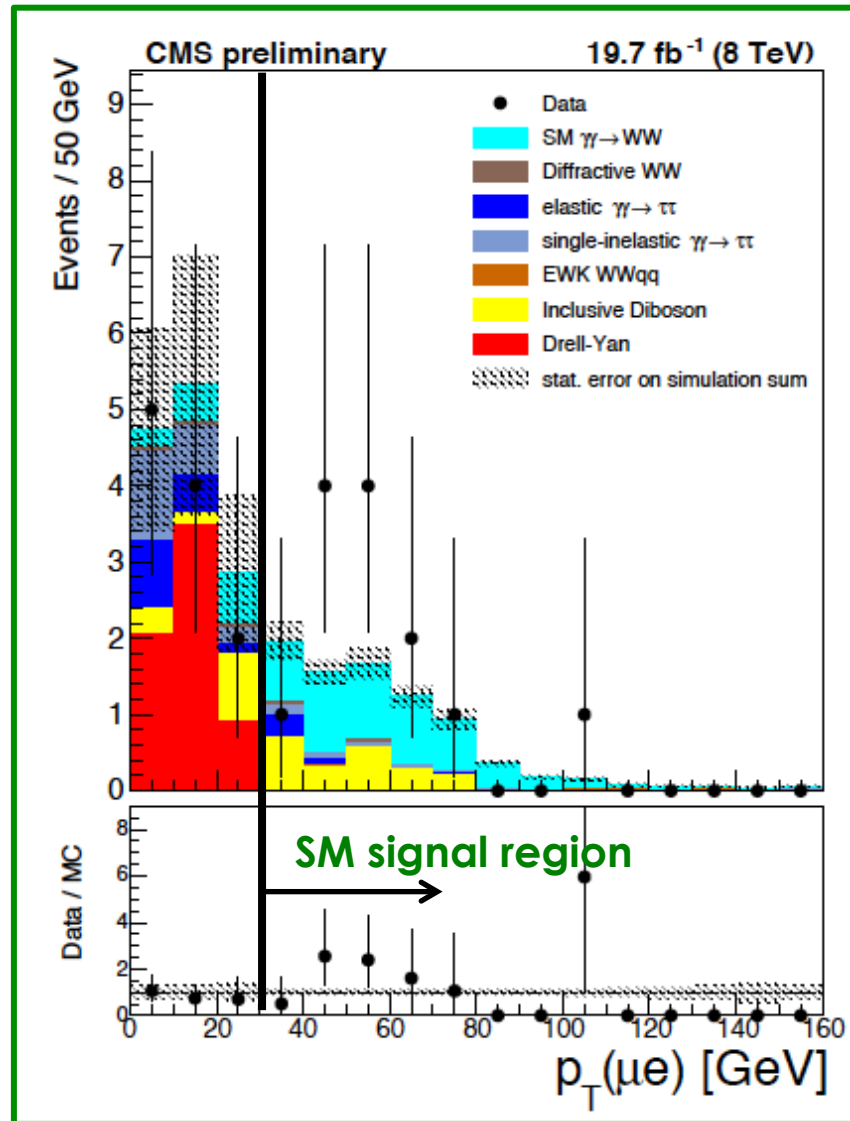
# Summary

- Central Exclusive Production allows to study a **variety of physics topics**.
- **New CMS results for  $\gamma\gamma \rightarrow WW$  exclusive production at 8 TeV** presented:
  - 13 events are observed; their properties consistent with the SM signal plus background prediction
  - Cross section measurement
  - More stringent limits on AQGCs
- Prospects for CEP measurements with CT-PPS discussed:  
**precision proton tracking and timing detectors** in the very forward region to study CEP in **standard LHC running at high luminosity**  
Project planning: **detector installation in 2016**  
**aim at accumulating 100 fb<sup>-1</sup> of data before LHC LS2**

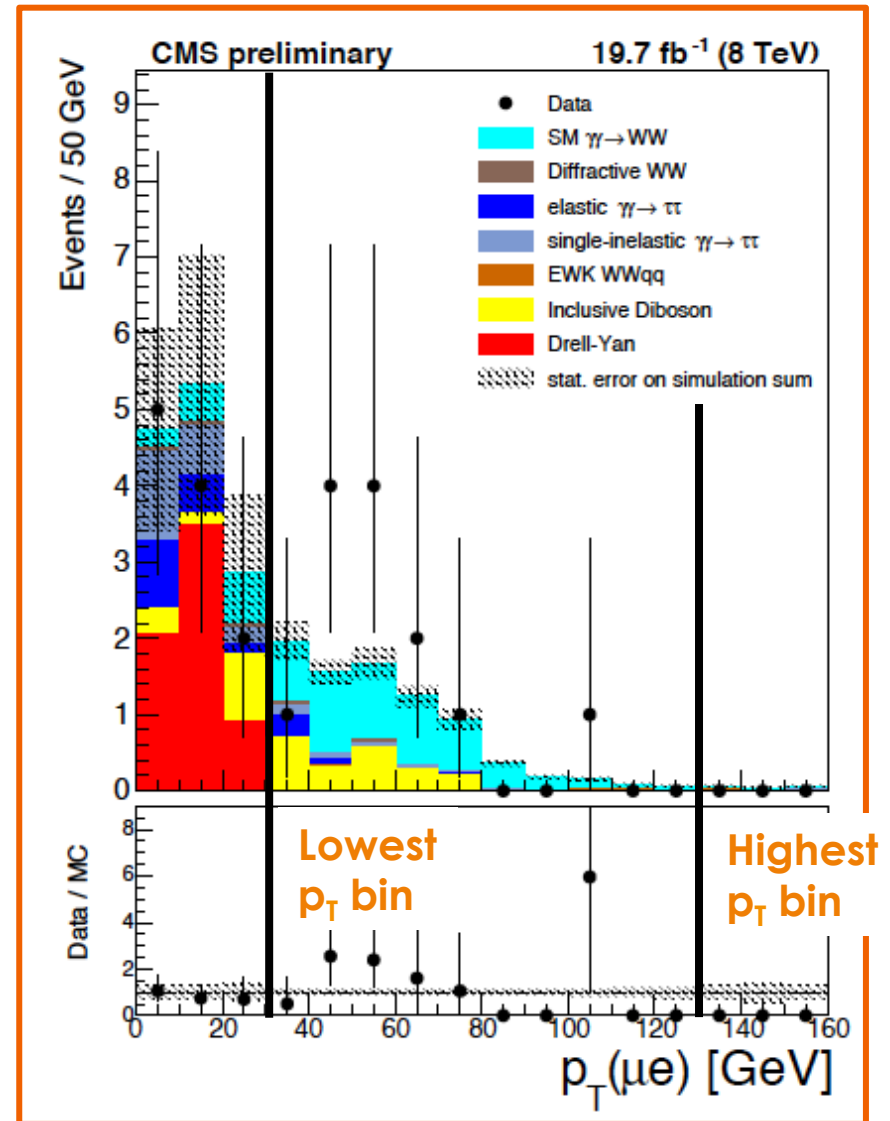
# Backup

# Signal region

SM signal



AQGCs search



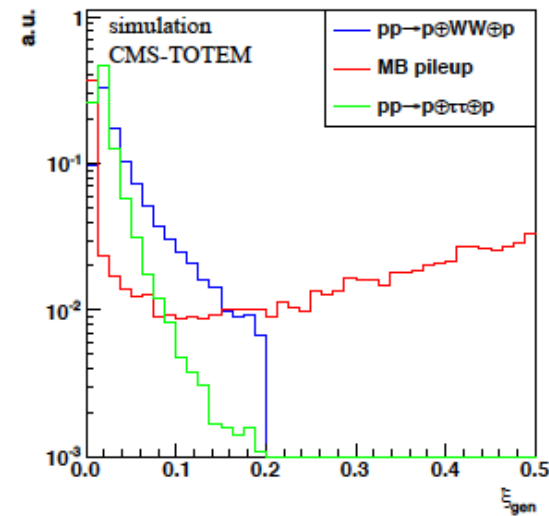
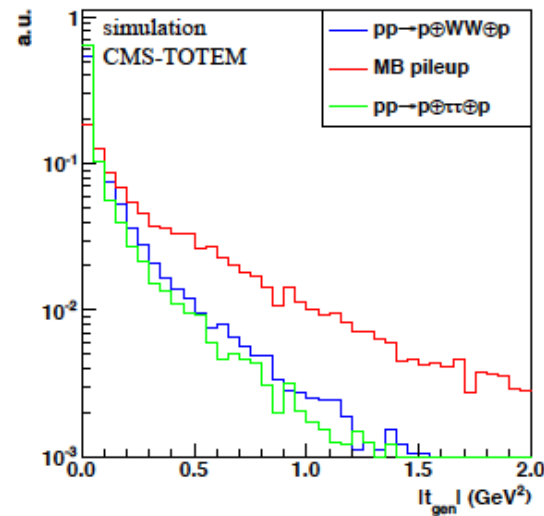
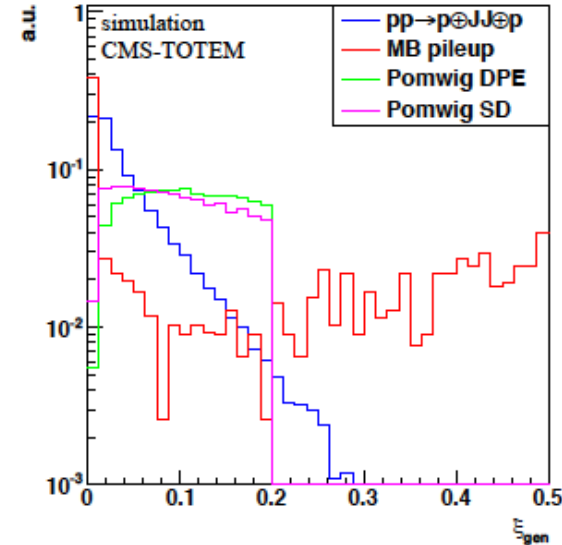
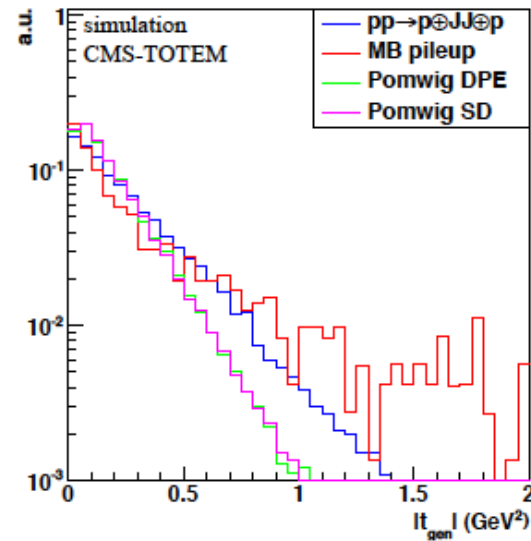
Muon-electron transverse momentum for events which pass all the selection criteria

# CT-PPS MC simulation ( $\sqrt{s}=13$ TeV )

- **$\gamma\gamma \rightarrow WW$** 
  - ✓ Events generated with FPMC (V1.0)
  - ✓ HERWIG (v6.5) used to simulate the WW pair decay
  - ✓ Events generated in the region  $0 < |t| < 4 \text{ GeV}^2$  and  $0.01 < \xi < 0.2$
- **$pp \rightarrow pJp$** 
  - ✓ Events generated with ExHuME (V1.3.2, but GSP 5%) interfaced with PYTHIA (v6.426)
- **Backgrounds**
  - ✓ SD+DPE events generated with POMWIG (v2.0) interfaced with HERWIG (v6.521)
  - ✓ Multijet QCD events simulated with PYTHIA (v8.175)
  - ✓ Inclusive WW events simulated with PYTHIA (v6.426)
  - ✓ Exclusive  $\gamma\gamma \rightarrow \tau\tau$  generated with FPMC
- **Pileup** incorporated by simulating additional interactions with an average pileup multiplicity of  $\mu=50$ . Minimum bias events generated with
  - ✓ PYTHIA (v6.175) for exclusive WW analysis
  - ✓ PYTHIA (v8.4) for exclusive dijets analysismixed to signal events
- **Particle transport** from the IP to the CT-PPS detector location performed with HECTOR

# Generator-level distributions of $|t|$ and $\xi$

Signal and background events in the exclusive dijet (top) and exclusive WW (bottom) samples.



# Generator-level distributions of $|t|$ and $\xi$

Signal and background events in the exclusive dijet and exclusive WW samples.

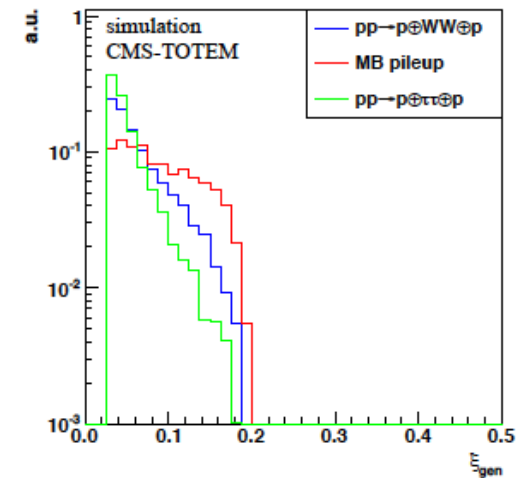
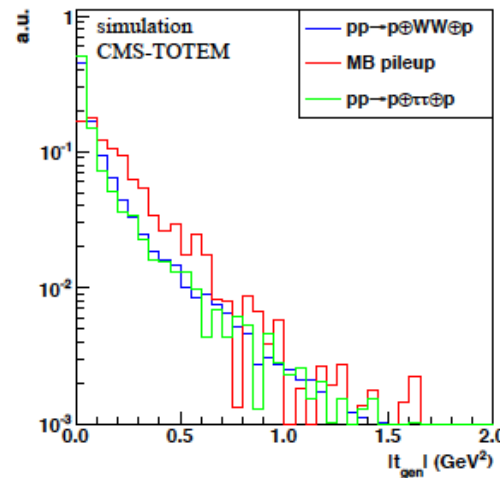
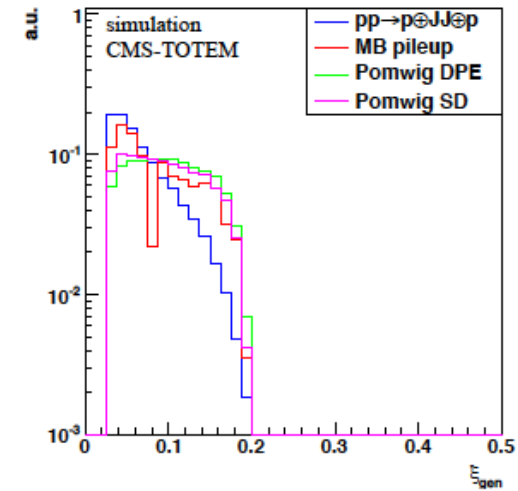
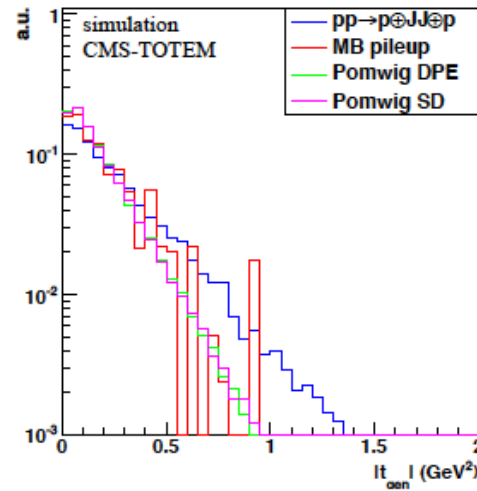
Coincidence of hits in the 2 tracking stations required.

Tracking detectors located at a distance of  $15 \sigma$  from the beam center.

Smearing effects due to

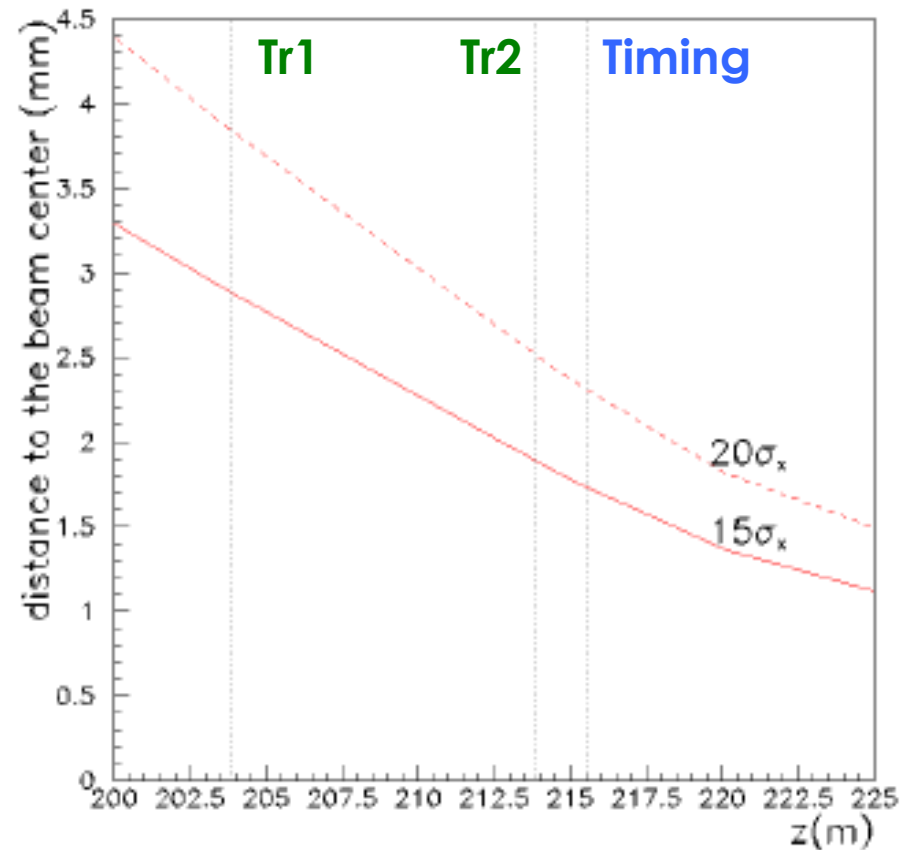
- vertex position,
- beam energy dispersion
- crossing angle

accounted for.





# Horizontal beam size



**Amplitude function @ IP ( $\beta^*$ ):** 0.6 m

**Emittance ( $\epsilon$ ):**  $5.4 \cdot 10^{-10}$  m

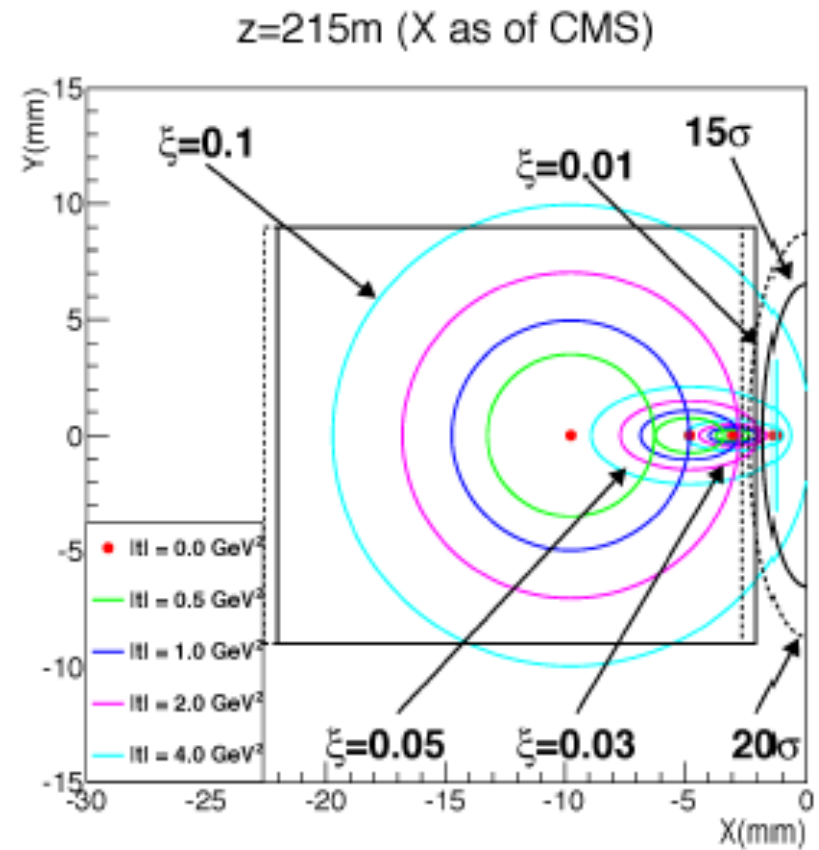
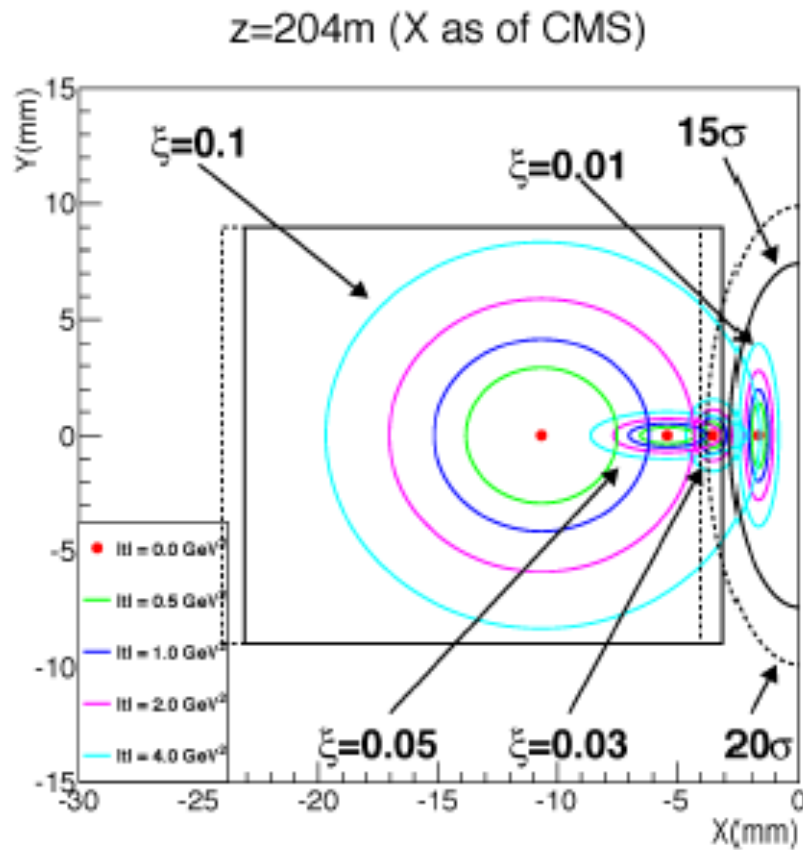
**Crossing angle:**  $142.5 \mu$  rad (horizontal plane)

**Vertex resolution of  $\sigma^*_{x,y}$ :**  $15 \mu$  m

**Angular beam divergence @ IP :**  $30 \mu$  rad

# Detector acceptance

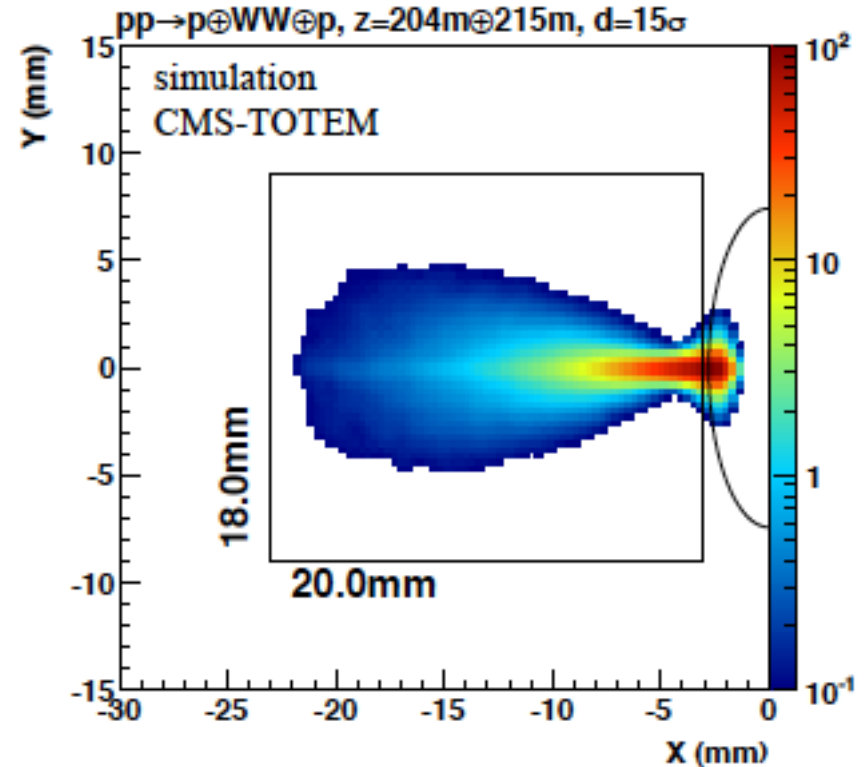
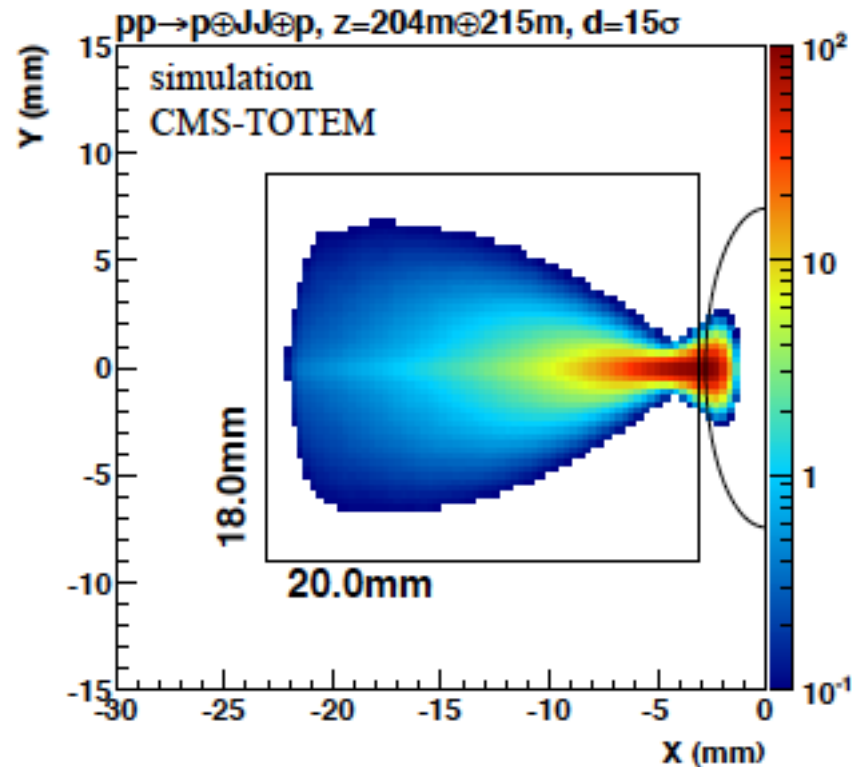
“Particle gun” sample: protons generated at fixed  $(|t|, \xi)$  transported from IP to CT-PPS detector location with HECTOR.



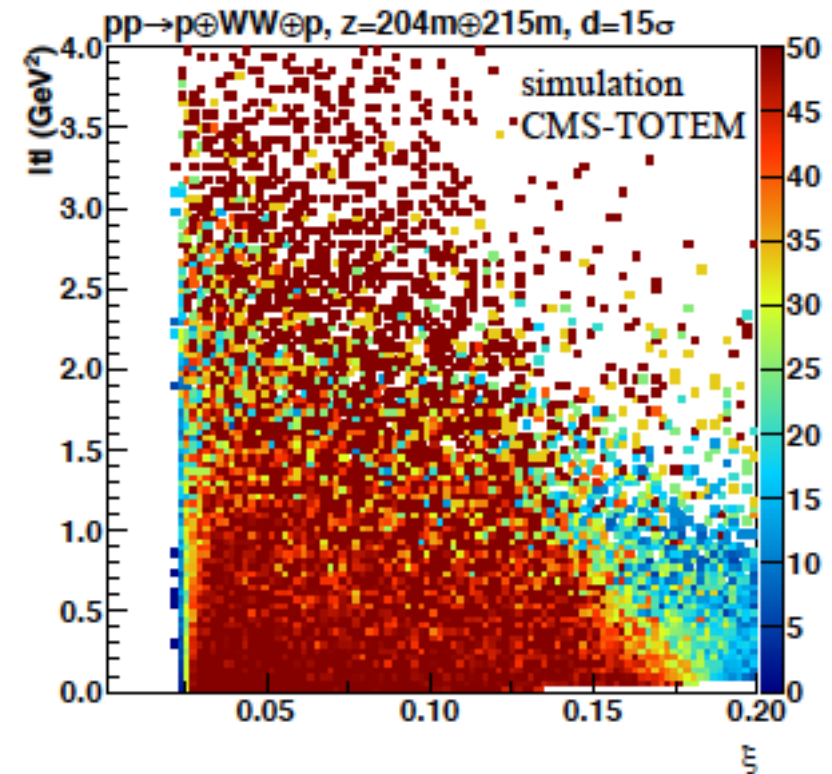
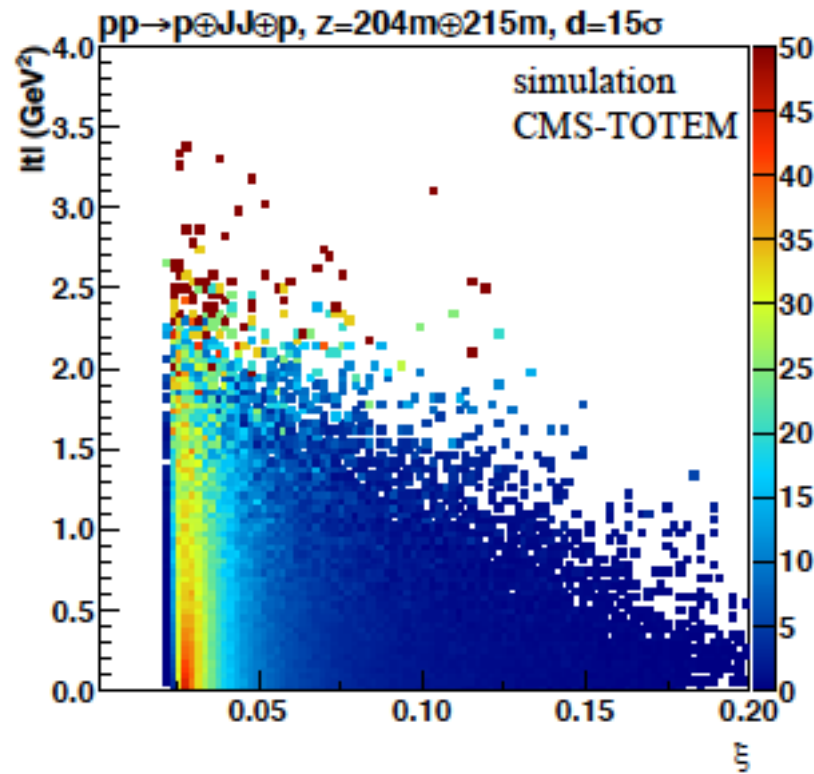
# Hit distribution

Hit distributions for centrally produced exclusive dijet and WW events for the tracking detectors located at  $z=204$  m.

Beam energy, vertex smearing and detector resolution are accounted for.



# Detector acceptance: $|t|$ and $\xi$



Coincidence of hits in the tracking stations at  $z=204$  m and  $z=215$  m on both sides of the IP required

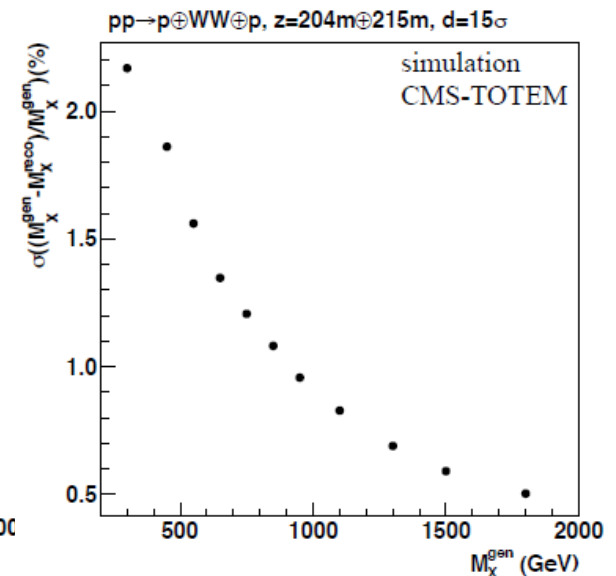
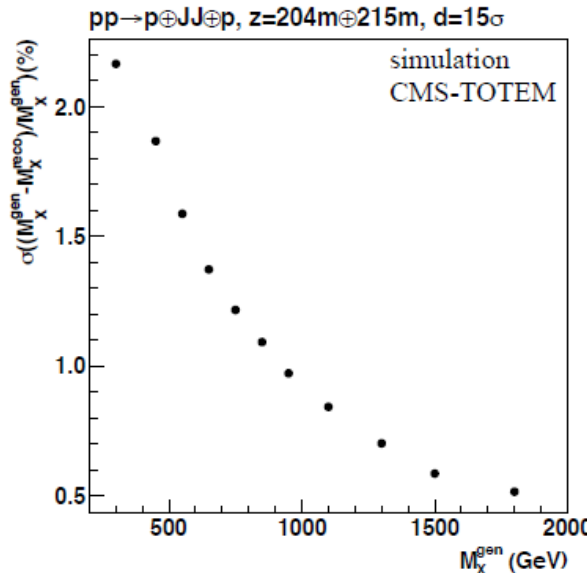
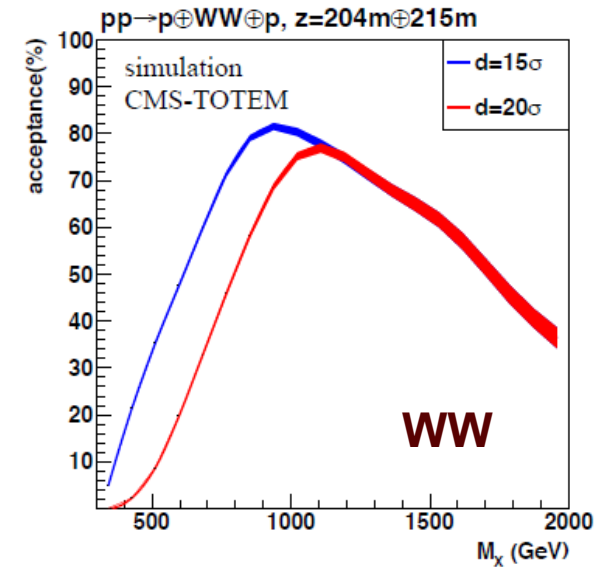
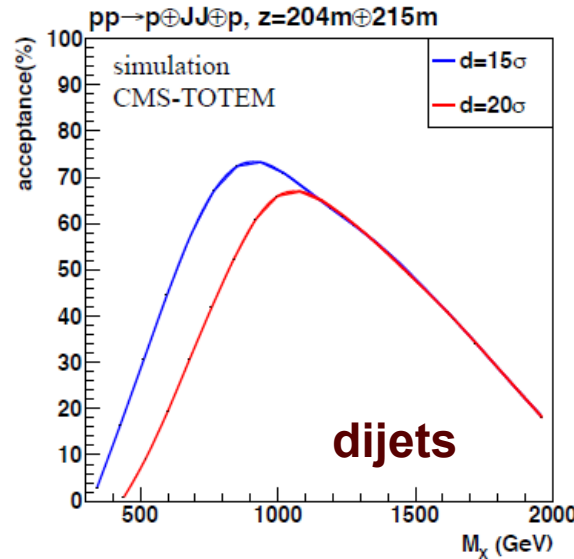
# $M_X$ acceptance and resolution

- With detectors at  $15\sigma$  from the beam centre, PPS selects exclusive systems in the 350-1700 GeV range ( $\epsilon > 5\%$ )



**Search for exclusive (light) Higgs production is not possible**

- At  $15\sigma$  acceptance larger wrt  $20\sigma$  for lower mass range
- $\gamma\gamma$  central system acceptance greater than that for gluon-gluon events because of smaller  $|t|$
- Mass resolution  $\sim 1.5\%$  at 500 GeV



# MC simulation: CMS $\gamma\gamma \rightarrow WW$ analysis ( $\sqrt{s}=8$ TeV )

- **Signal**

- ✓ SM and anomalous signal samples: MADGRAPH(v2.0.0), based on the equivalent photon approximation (EPA).

- **Backgrounds**

- ✓ Inclusive WW, ttbar, W+jets processes: MADGRAPH; yields normalized to the NLO cross section predictions obtained with MCFM
- ✓ Inclusive Drell-Yan process: POWHEG(v1.0)
- ✓ Exclusive two photon processes:  $\gamma\gamma \rightarrow ll$  : LPAIR
- ✓ WW production from single diffractive interactions: POMPYT (v2.6.1); 100% gap survival probability is assumed

MADGRAPH and POWHEG are matched to parton showers from PYTHIA(v6.4.26) Z2\* tune

- **Control samples**

- ✓ The elastic and proton dissociation  $\gamma\gamma \rightarrow l^+l^-$  samples: LPAIR

# Systematic uncertainties affecting signal

	Uncertainty
Proton dissociation factor	10.5%
0 extra tracks Efficiency Correction	5.0%
Trigger and lepton ID	2.4%
Luminosity	2.6%
Total	12.1%

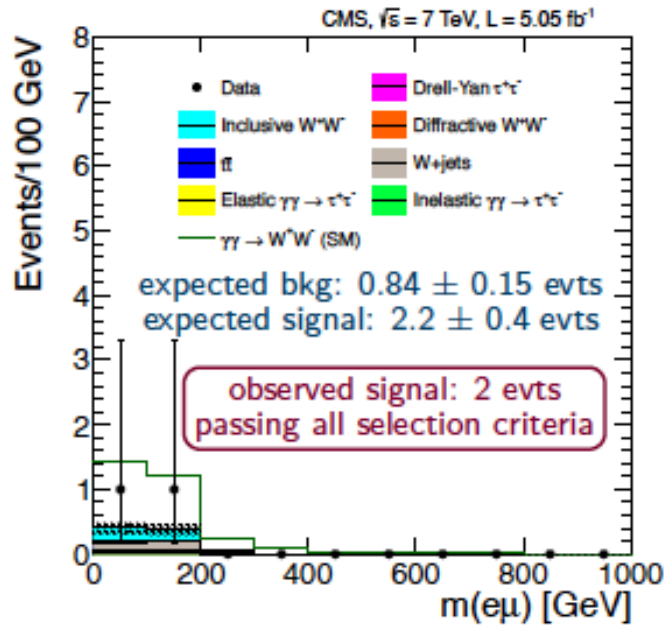
Proton dissociation factor:

- 9.2% → statistical uncertainty based on the combination of the  $\mu + \mu -$  and  $e + e -$  channels.
- 5.0% → difference between the matrix element prediction of LPAIR used for  $\gamma \gamma \rightarrow l^+ l^-$  and the equivalent photon approximation used to generate signal events

Systematic uncertainties considered for the **background** estimation include:

- trigger and lepton ID, luminosity, and simulation statistics
- an uncertainty of  $\pm 0.24$  events on the electroweak  $W^+ W^-$  background contribution is included, corresponding to the full difference between the background predictions of the MADGRAPH and PHANTOM generators.

# 7 TeV data results



2 events pass all the selection criteria, compared to the expectation of

- $2.2 \pm 0.4$  signal events
- $0.84 \pm 0.15$  background events

$$\sigma(pp \rightarrow p^{(*)}W^+W^-p^{(*)} \rightarrow p^{(*)}\mu^\pm e^\mp p^{(*)}) = 2.2^{+3.3}_{-2.0}$$

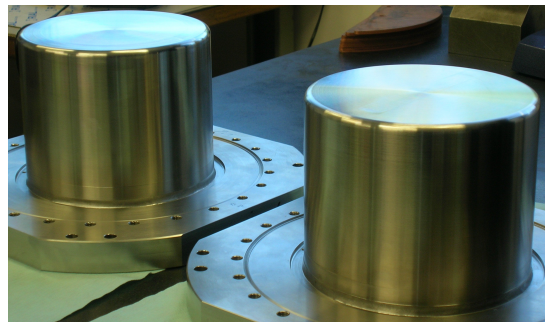
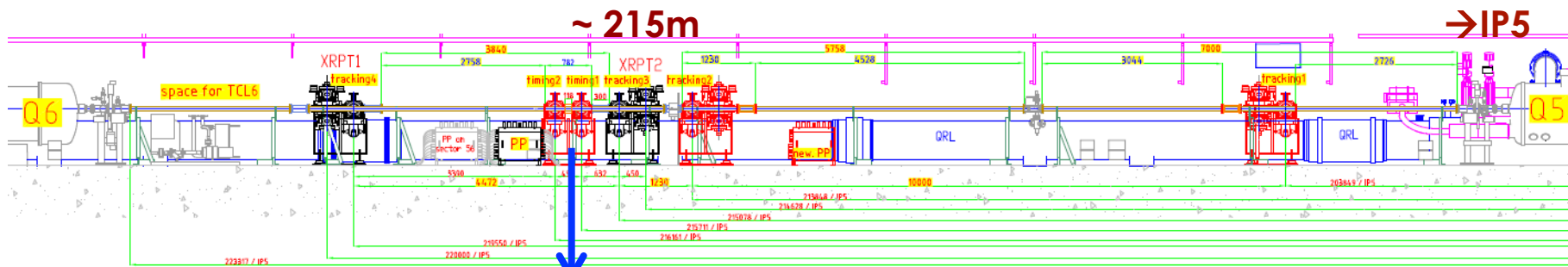


# Project planning

## Exploratory phase (2015-16)

- Prove the ability to operate detectors close to the beamline at high luminosity
  - Show that CT-PPS does not prevent the stable operation of the LHC beams and does not affect significantly the luminosity performance of the machine.
- In 2015:
  - Evaluate RPs in the 204-215 m region
  - Demonstrate the timing performance of the Quartic baseline
  - Use TOTEM silicon strip detectors at sustainable radiation intensity
  - Integrate the CT-PPS detectors into the CMS trigger/DAQ system.
- In 2016:
  - Evaluate the MBP option
  - Upgrade the tracking to pixel detectors
  - Upgrade the timing detectors if required/possible
- **Data Production phase**
  - Aim at accumulating  $100 \text{ fb}^{-1}$  of data before LHC LS2

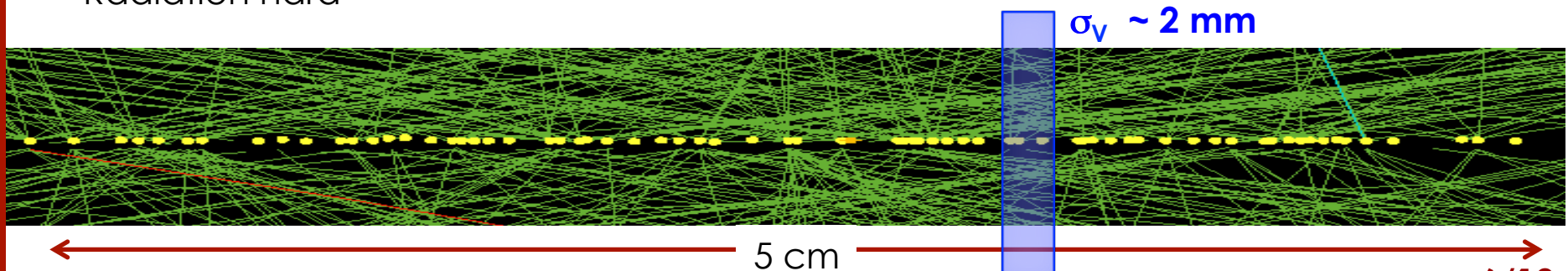
# Timing detectors for CT-PPS: requirements



2 new horizontal cylindrical pots

Timing measurement from both sides of CMS allows to determine the vertex of the protons and reject pile-up

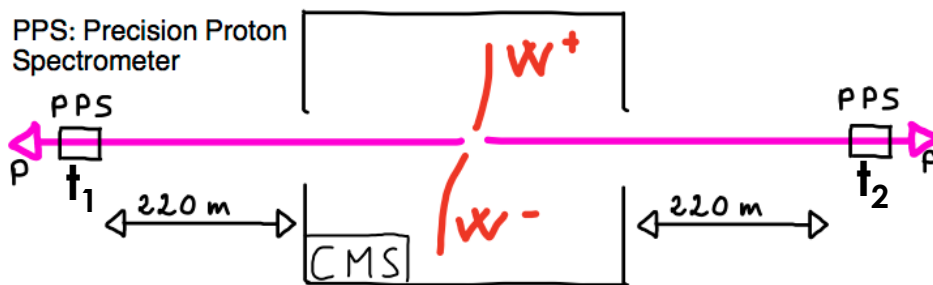
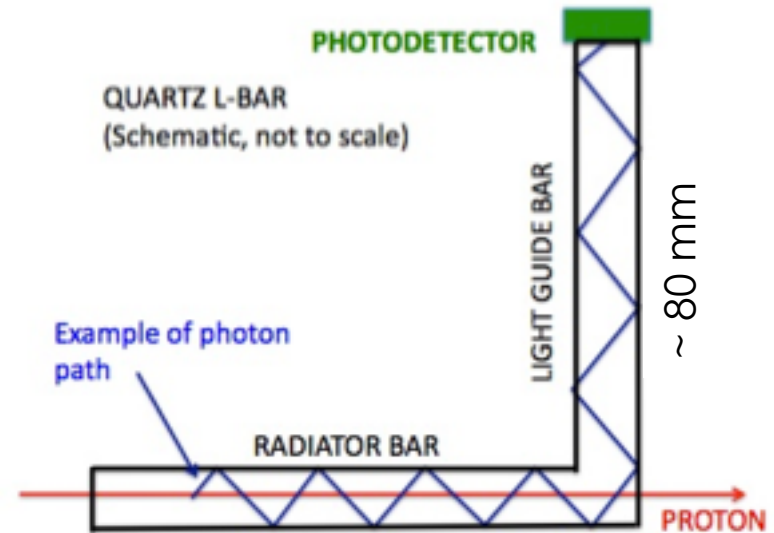
- Time resolution  $\sim 10$  ps  $\rightarrow$  Vertex z-by-timing:  $\sim 2$  mm
- Segmentation to cope with the high occupancy expected
- Edgeless ( $\sim 200$   $\mu$ m)
- Radiation hard



# L-bar QUARTIC for PPS

- All Cherenkov light is totally internally reflected along radiator bar
- 66% goes promptly along light guide to SiPM or segmented MCP-PMT.

$\sigma_T \sim 30$  ps for 30 mm bar  
Four-in-line  $\rightarrow$  15 ps

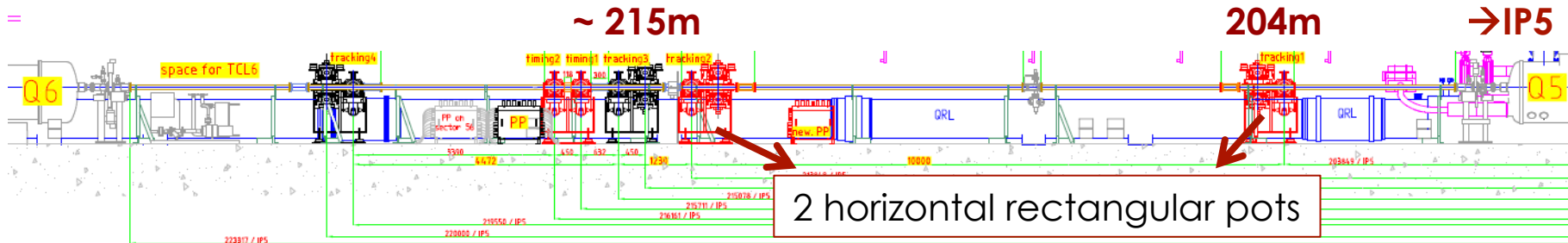


Time difference =  $\Delta t = t_1 - t_2$   
 $z(pp) = 0.5 c \Delta t$   
 If time resolution  $\sigma_T = 10$  ps  
 $\sigma_{\Delta T} = \sqrt{2} \sigma_T = 14.1$  ps  
 $\sigma_Z = 0.5 c 14.1$  ps  $\sim 2$  mm



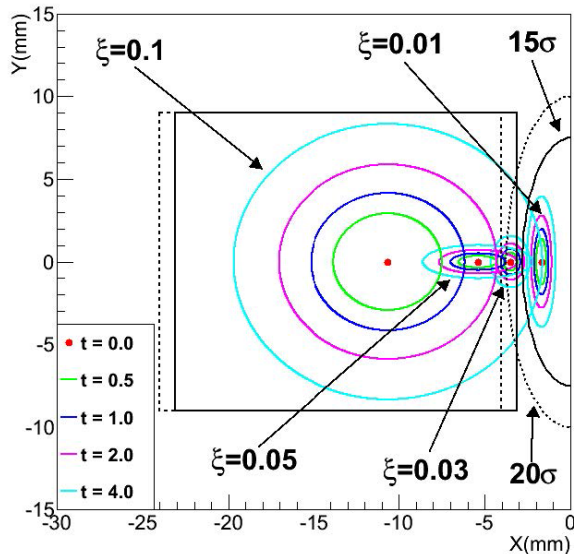
# Tracking system for CT-PPS

Two detector stations, ~ 10 meters apart, on each side of IP5



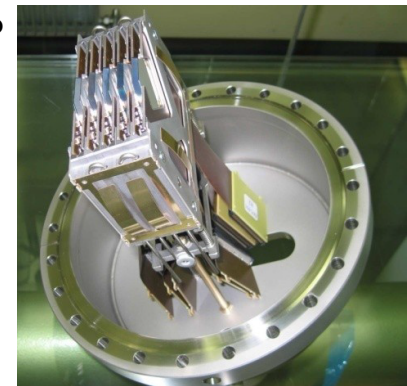
Position and angle, combined with the beam magnets, allow to determine the momentum of the scattered proton and in turn the missing mass:

- Position precision of 10  $\mu\text{m}$
  - Angular resolution of 1-2  $\mu\text{rad}$
- $\Rightarrow \Delta p/p \sim 2 \cdot 10^{-4}$   
Mass resolution:  $\sim 2 \text{ GeV}/c^2$



Position of scattered protons at 204 m, for fixed  $(\xi, t)$

- Slim edges of side facing beam: dead region  $\sim 100 \mu\text{m}$
- Tolerance to inhomogeneous irradiation: up to  $2 \cdot 10^{15} n_{eq}/\text{cm}^2$  close to the beam (for  $100 \text{ fb}^{-1}$ )
  - Should fit into existing RP

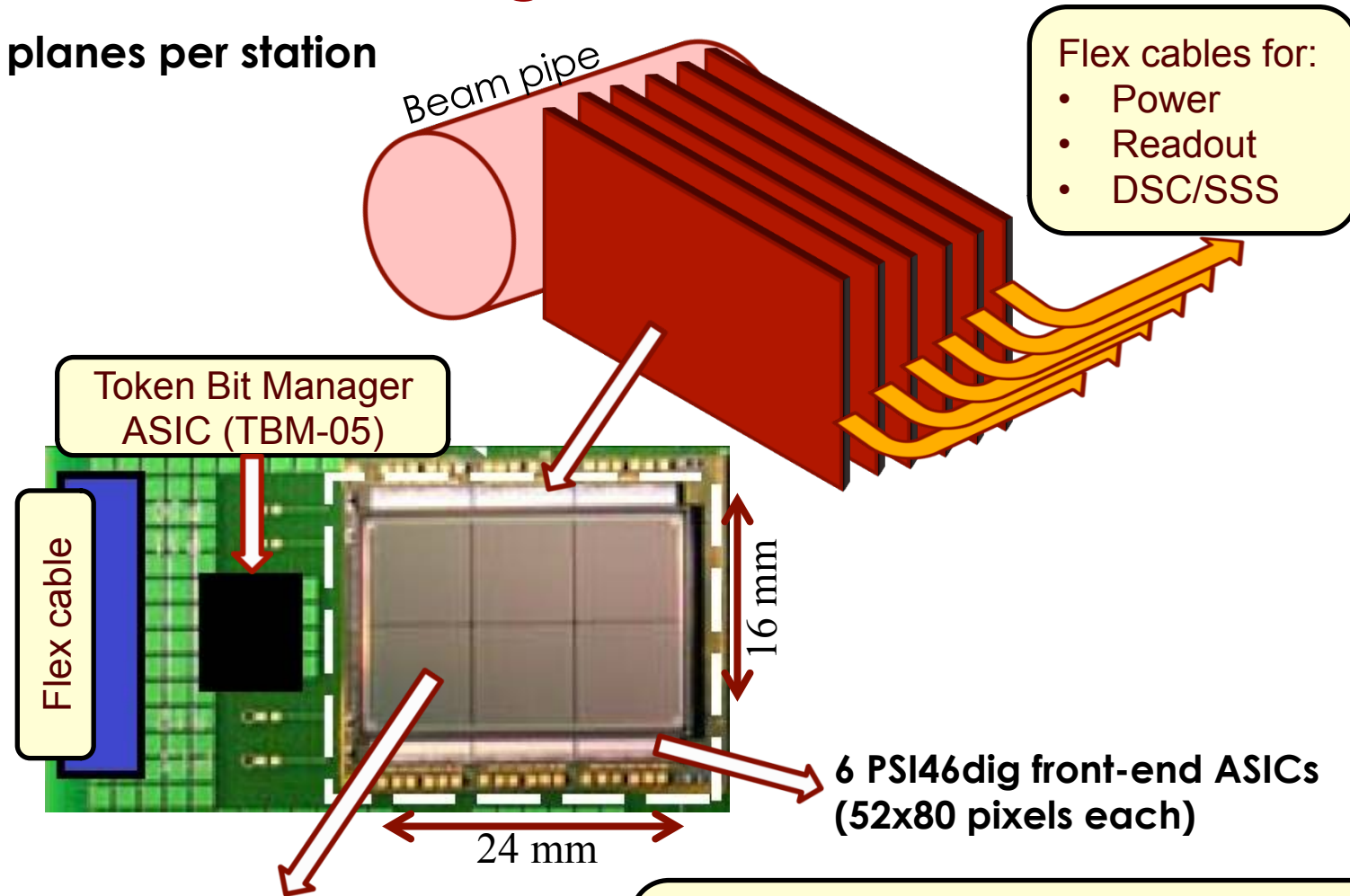


Chosen solution: 3D pixel silicon detectors



# Tracking stations

Six detector planes per station



- 16 x 24 mm<sup>2</sup> 3D pixel silicon sensor
- 150(y) x 100(x) μm<sup>2</sup> pixel pattern, same as CMS pixel sensors

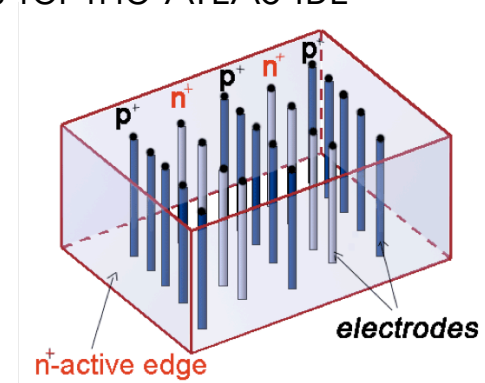
Same readout scheme as Phase-I upgrade of CMS forward pixel system  
→ Existing CMS DAQ components and software can be reused

Detector installation foreseen at the end 2016

# Why 3D sensors for CT-PPS?

After 15 years of R&D, 3D sensor technology has reached its maturity, as demonstrated by the successful fabrication of more than 300 3D sensors for the ATLAS IBL

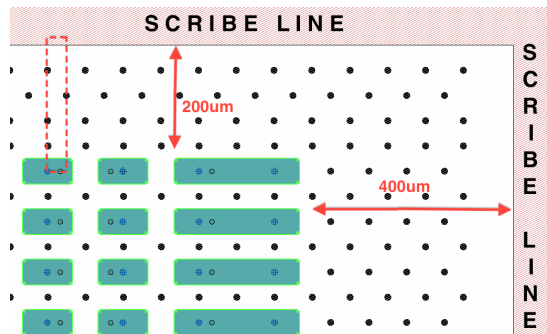
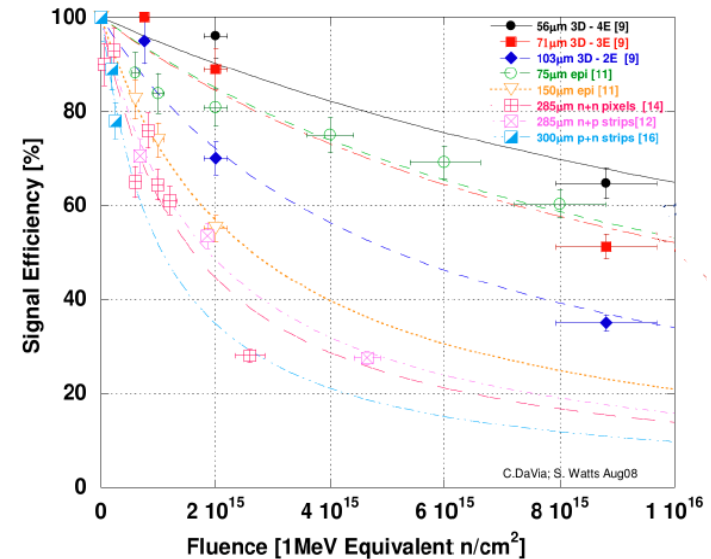
**3D sensors** consist of an array of columnar electrodes [ $r \sim 5\mu\text{m}$ ] of both doping types that penetrate in the silicon substrate perpendicularly to the surface.



➔ **Electrode distance and active substrate thickness are decoupled**

The close electrode distance provides several advantages w.r.t planar sensor

1. **low full depletion voltage** ( $\sim 10\text{ V}$ )
2. fast charge collection time
3. reduced charge-trapping probability and therefore **high radiation hardness**.



**Slim edges**, with dead area  $< 200\ \mu\text{m}$

**Active edges**, with dead area reduced to a few  $\mu\text{m}$

Layout details of the corner region of CMS 3D sensors produced by FBK

# Roman Pot spectrometer

