

Soft pomeron in AdS/QCD

Miguel S. Costa

CERN & Universidade do Porto

Work with A. Bayona, M. Djuric, R. Quevedo

EDS Blois 2015:

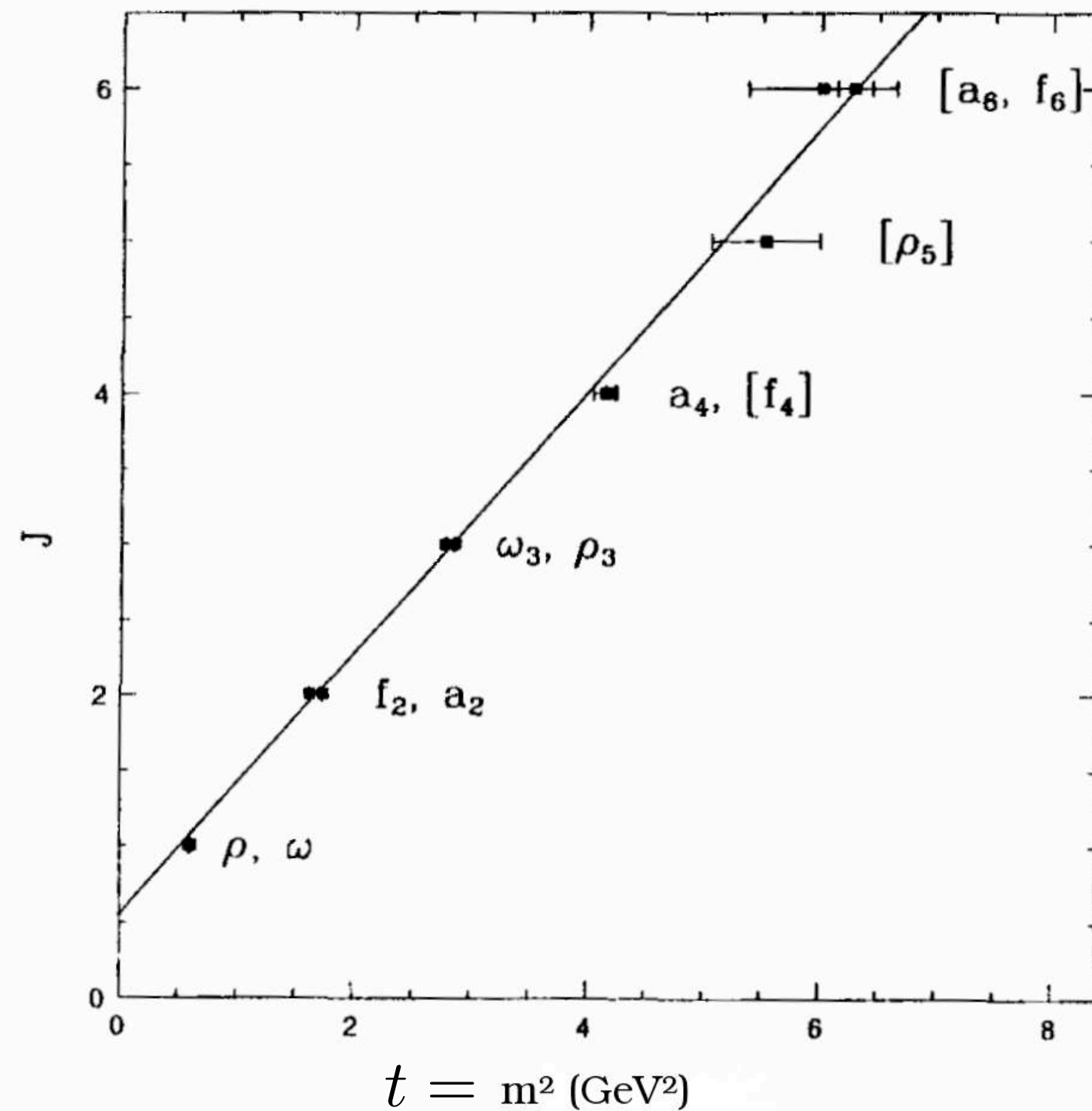
The 16th conference on Elastic and Diffractive scattering
Borgo, Corsica/France, June 2015

Regge behaviour in QCD

Regge behaviour in QCD

- Hadronic resonances fall in linear trajectories

$I = 1$ even parity mesons

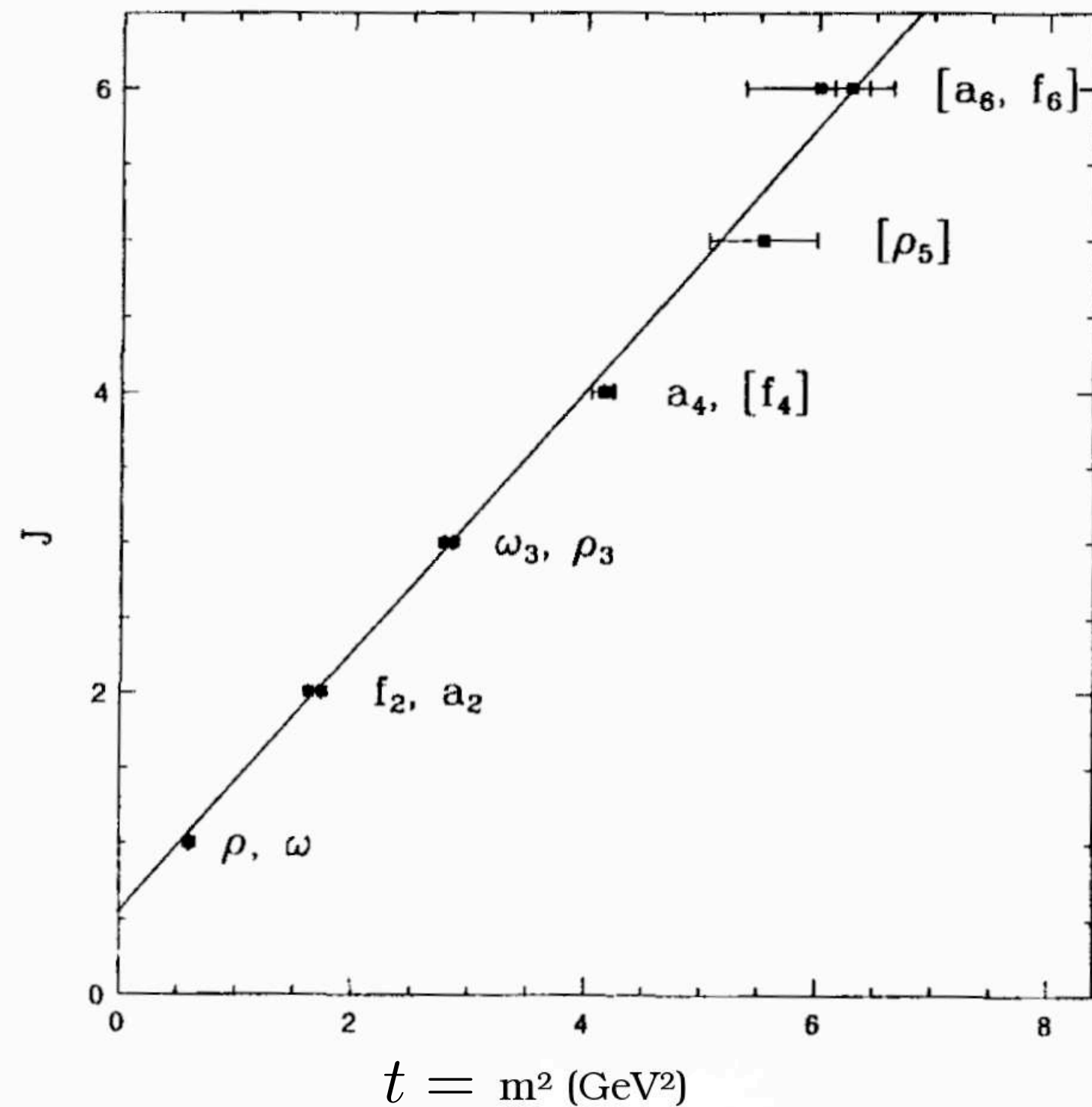


$$J = j(t) = j(0) + \alpha' t$$

Regge behaviour in QCD

- Hadronic resonances fall in linear trajectories

$I = 1$ even parity mesons

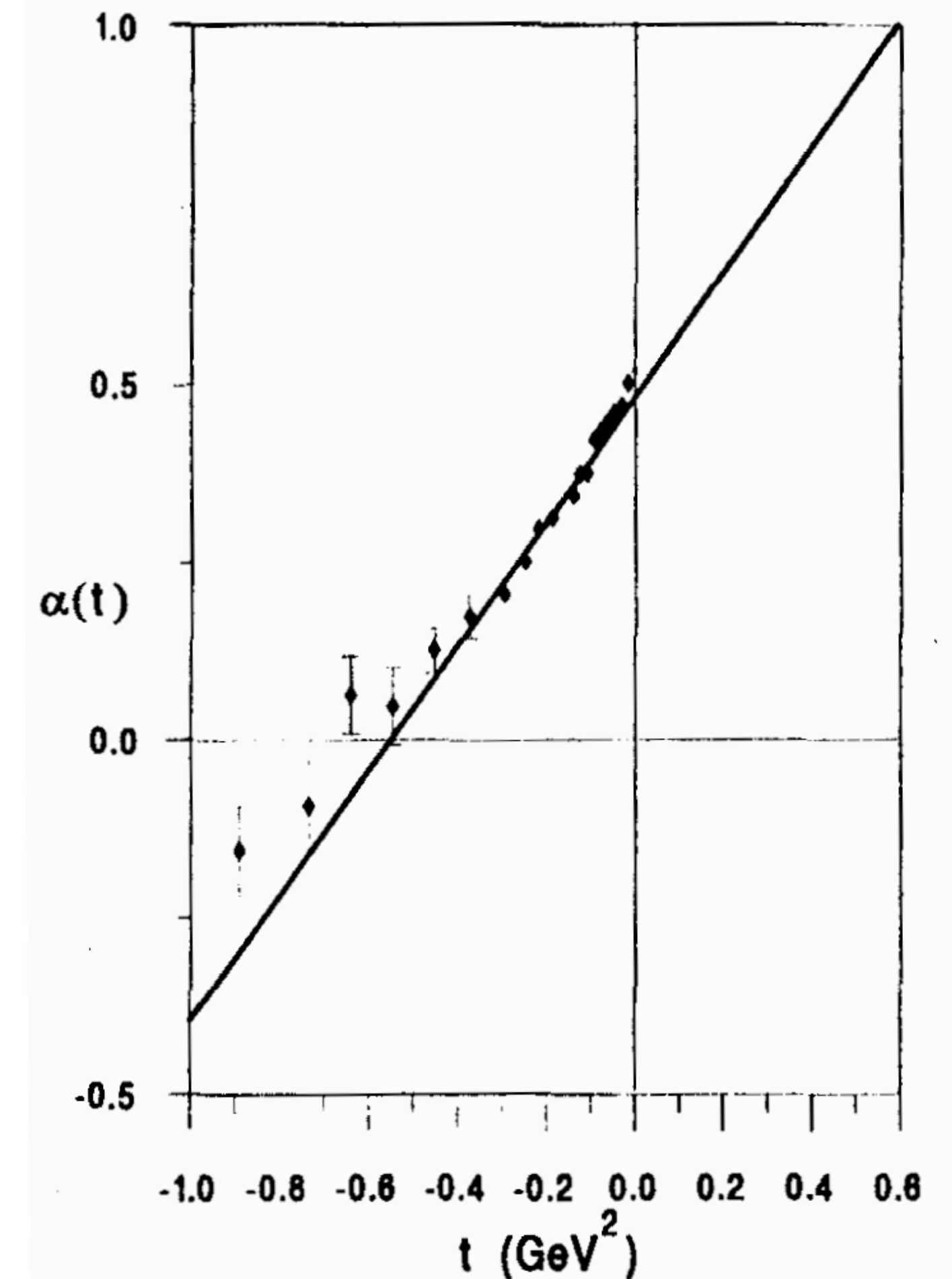
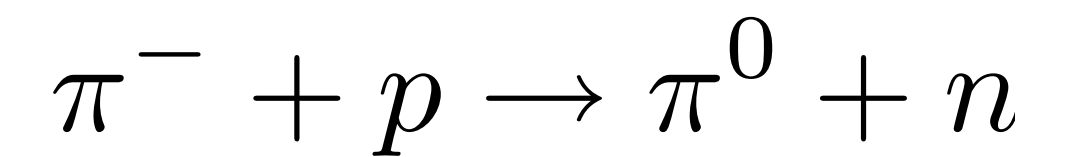


$$J = j(t) = j(0) + \alpha' t$$

$$A(s, t) \sim \beta(t) s^{j(t)} \quad (s \gg t)$$

Total cross section

$$\sigma \sim s^{j(0) - 1}$$

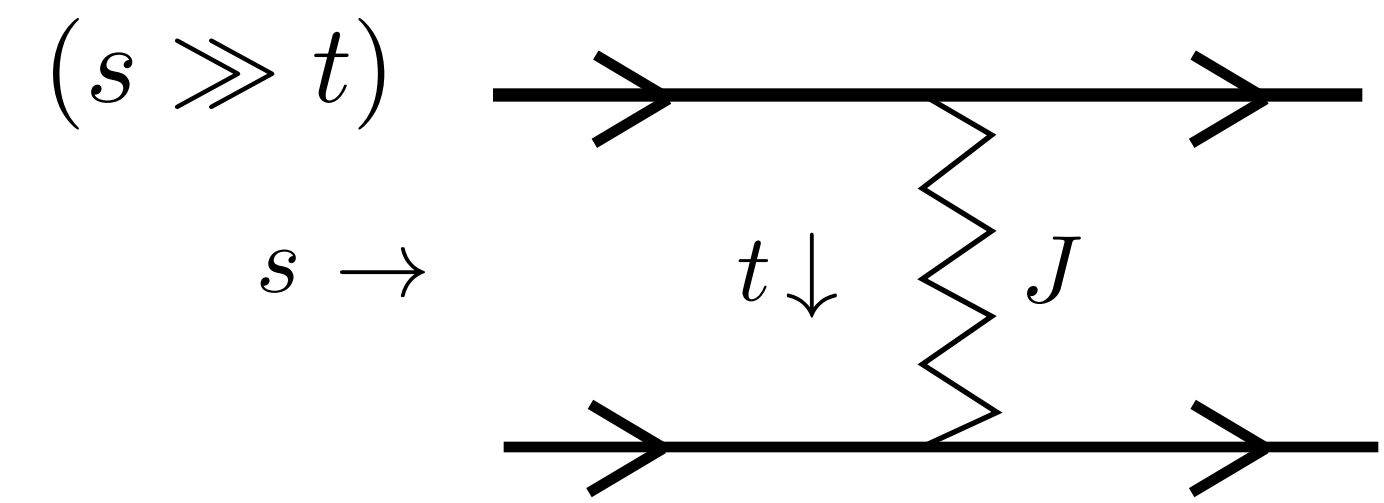


Regge theory

Regge theory

- Scattering dominated by t-channel exchange of a Regge trajectory

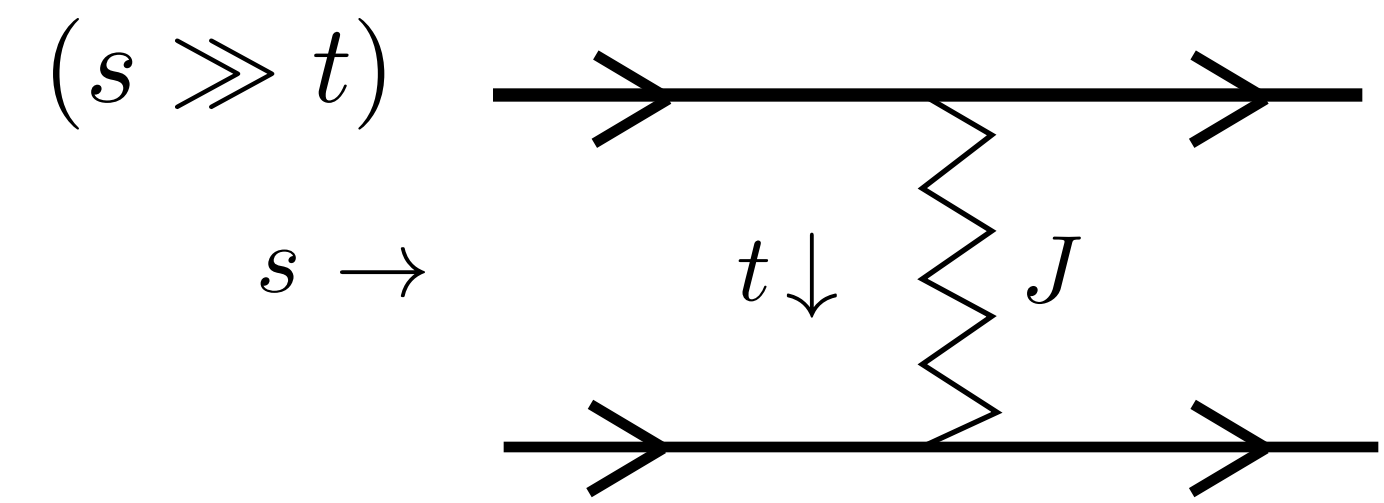
$$A(s, t) \approx \sum_J g_J \frac{s^J}{t - m^2(J)} \sim \sum_J g_J \frac{s^J}{J - j(t)}$$



Regge theory

- Scattering dominated by t-channel exchange of a Regge trajectory

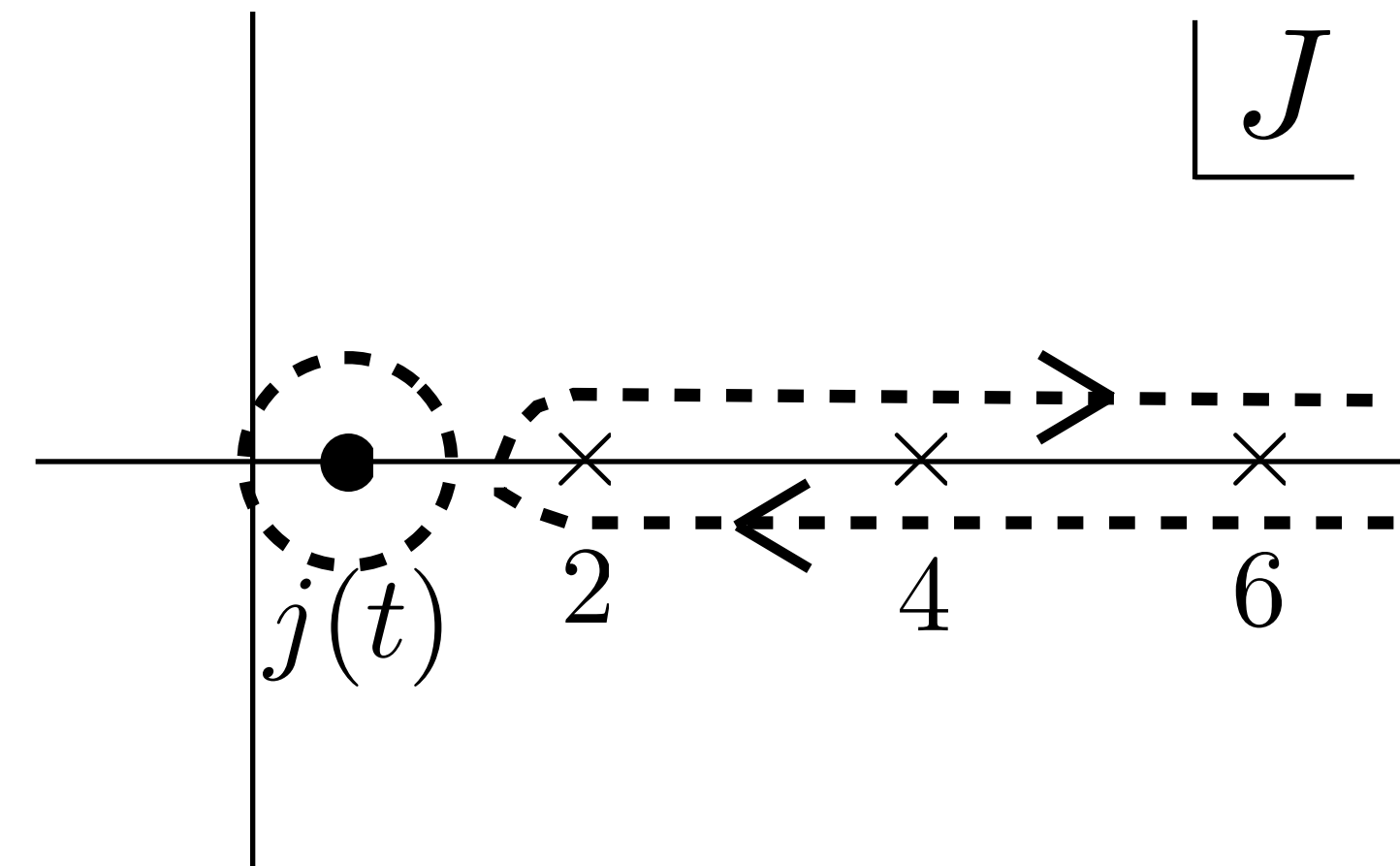
$$A(s, t) \approx \sum_J g_J \frac{s^J}{t - m^2(J)} \sim \sum_J g_J \frac{s^J}{J - j(t)}$$



- Sommerfeld-Watson transform:

$$\sum_J \rightarrow \int \frac{dJ}{\sin \pi J}$$

$$A(s, t) \sim \beta(t) s^{j(t)}$$



Soft Pomeron trajectory [Donnachie, Landshoff]

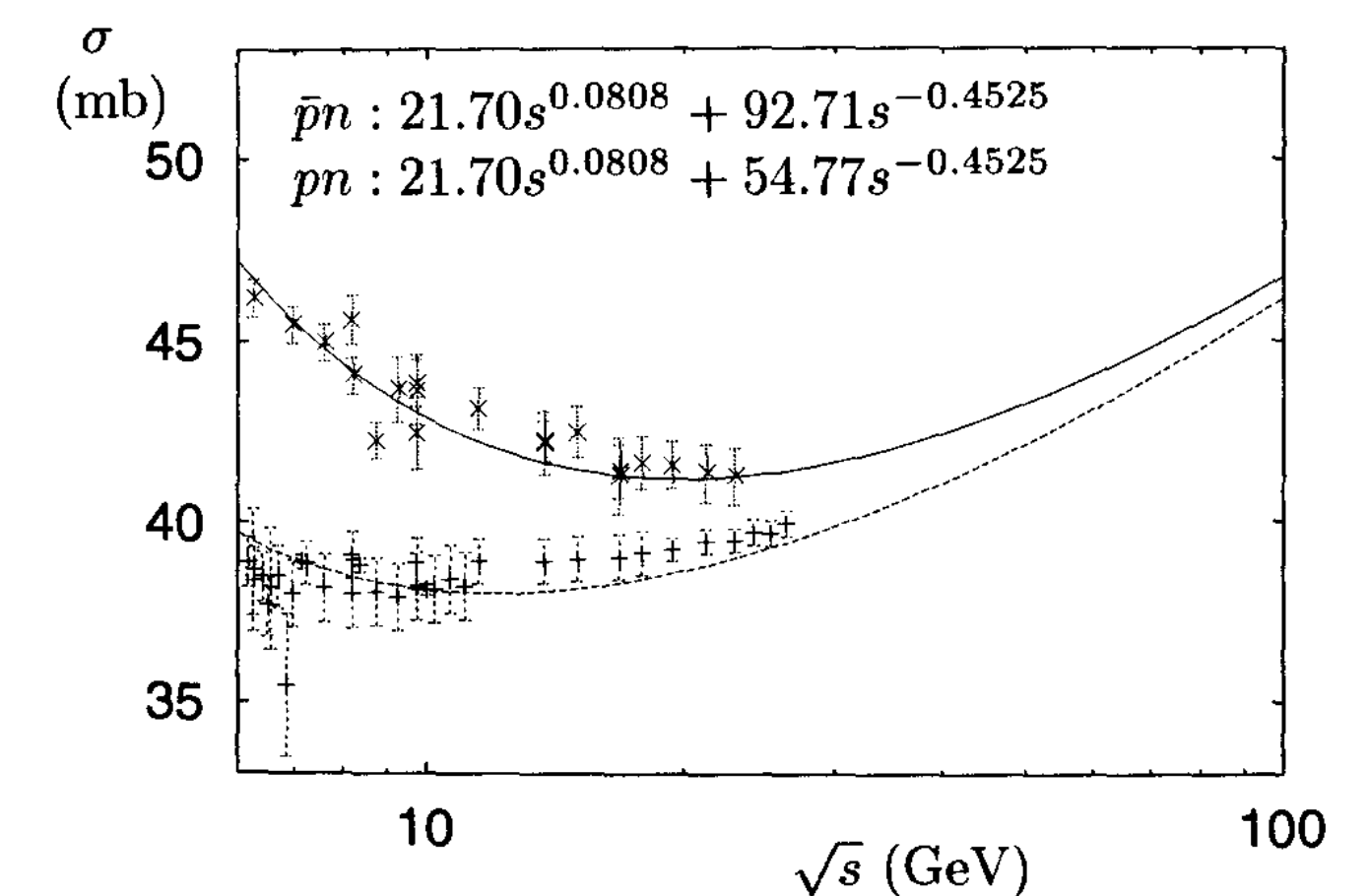
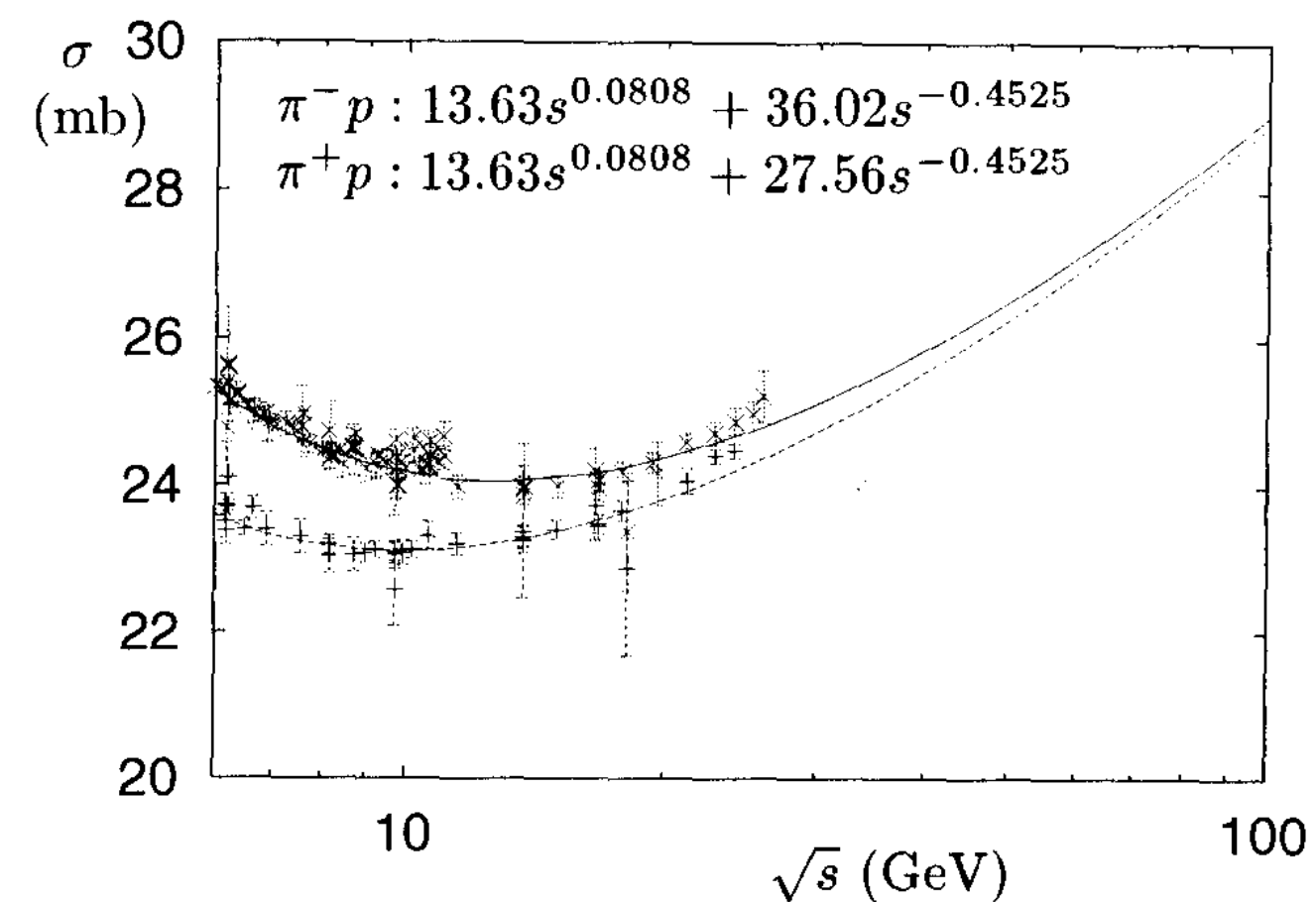
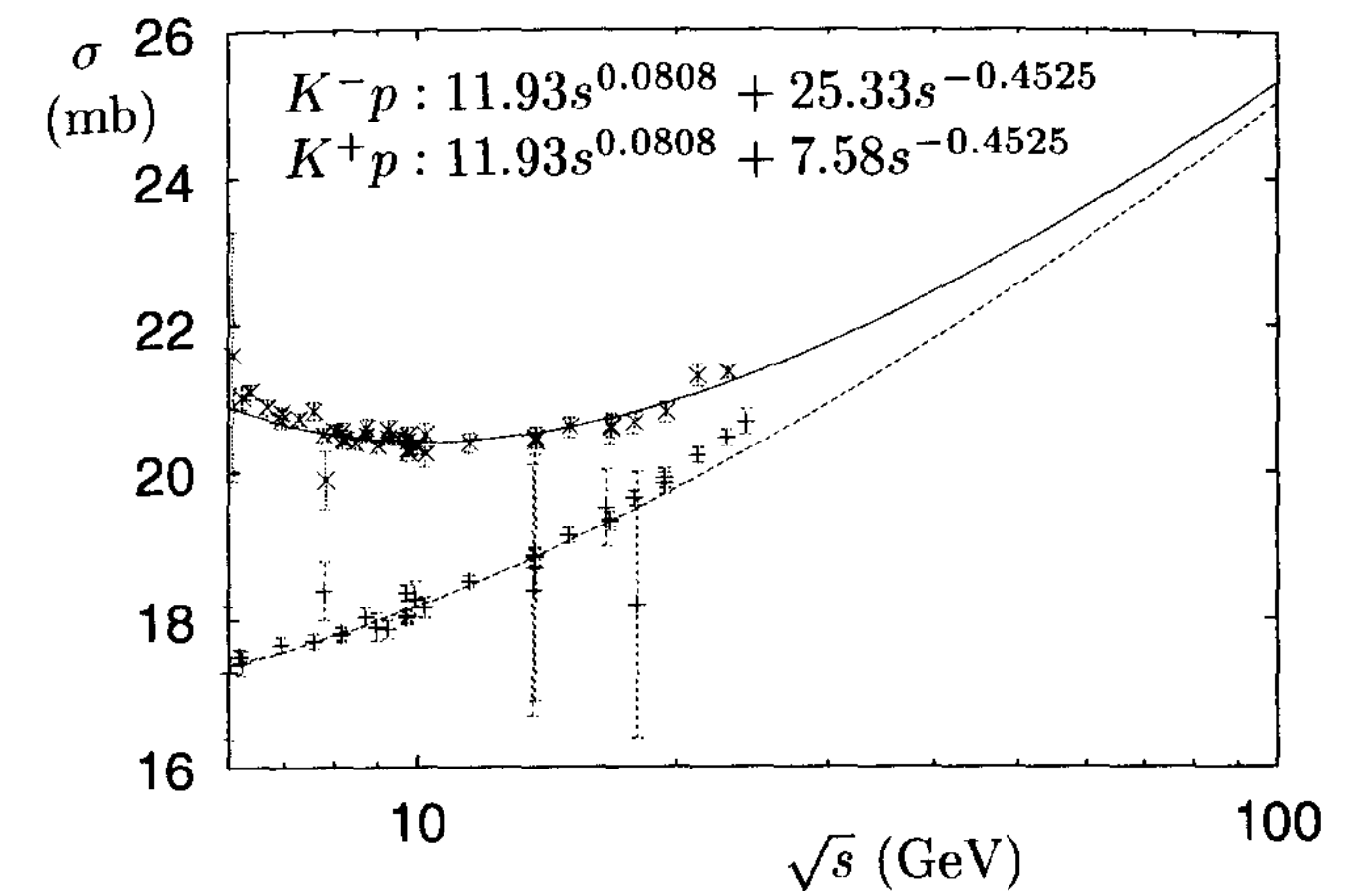
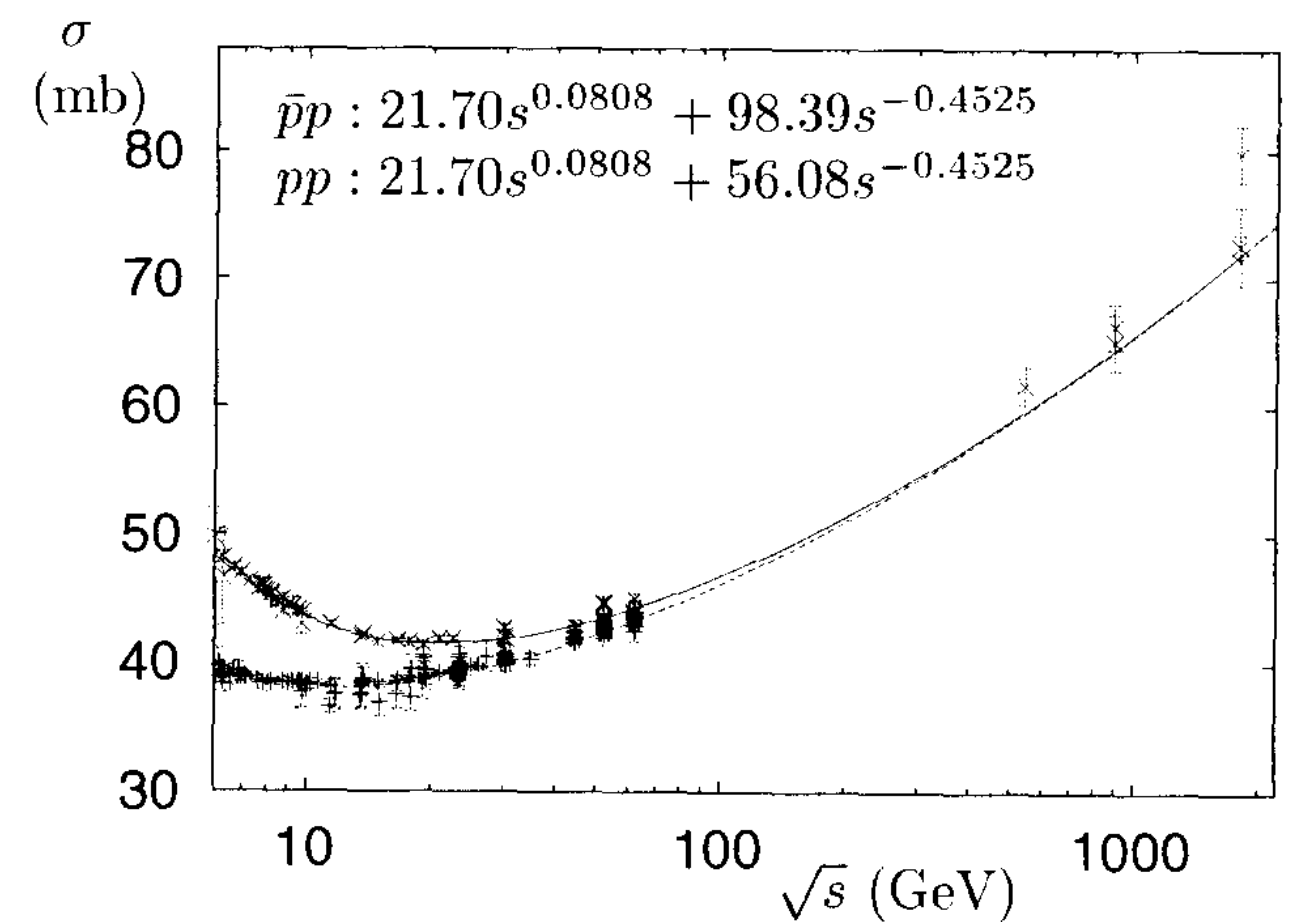
Soft Pomeron trajectory [Donnachie, Landshoff]

- Trajectory selected by exchanged quantum numbers. For elastic scattering these are the vacuum quantum numbers.

$$j_P(t) \approx 1.08 + 0.25t \quad (\text{GeV units})$$

Elastic cross sections in QCD

$$\sigma \sim s^{j_P(0)-1} \sim s^{0.08}$$



Soft Pomeron trajectory [Donnachie, Landshoff]

- Trajectory selected by exchanged quantum numbers. For elastic scattering these are the vacuum quantum numbers.

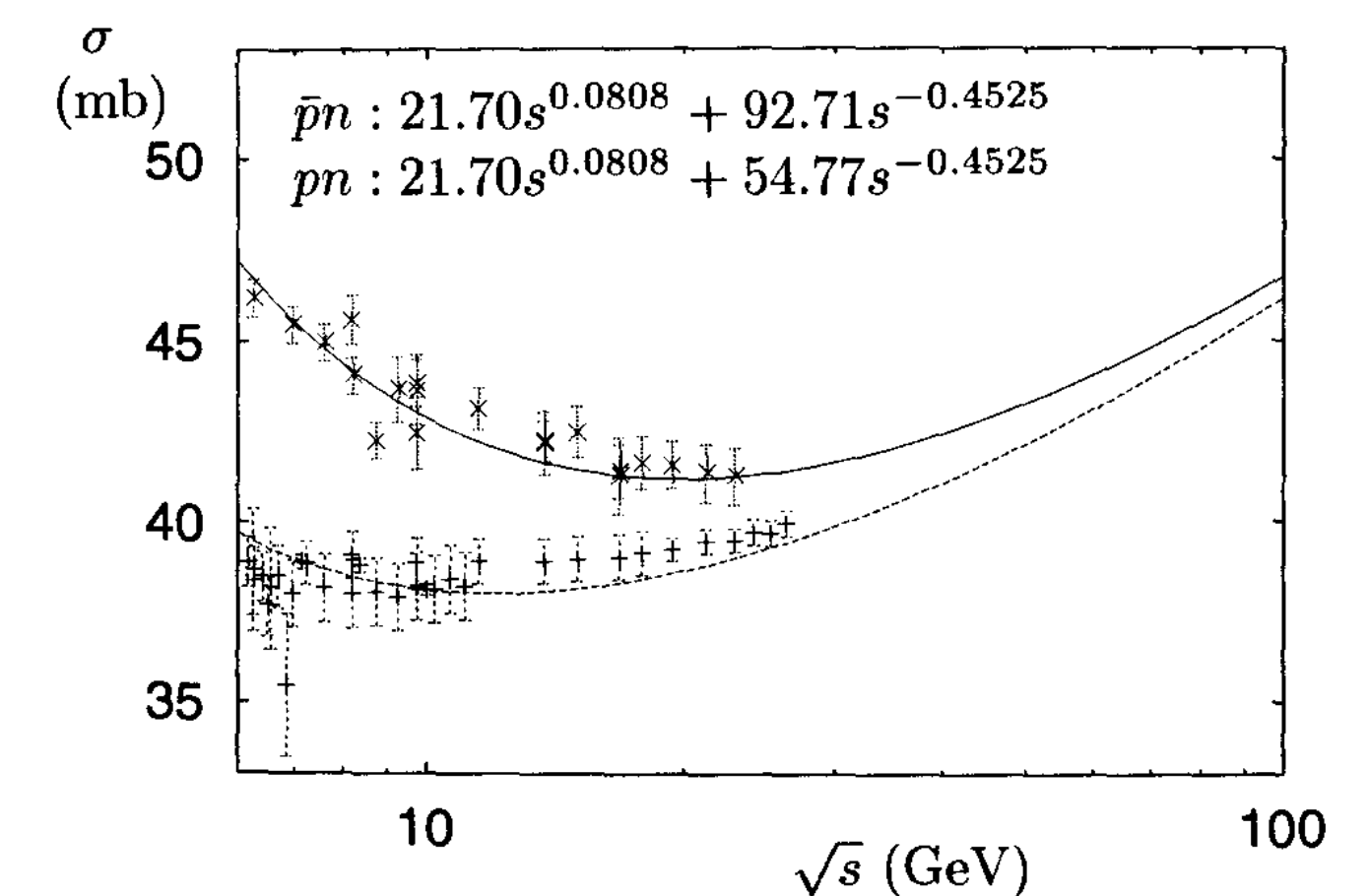
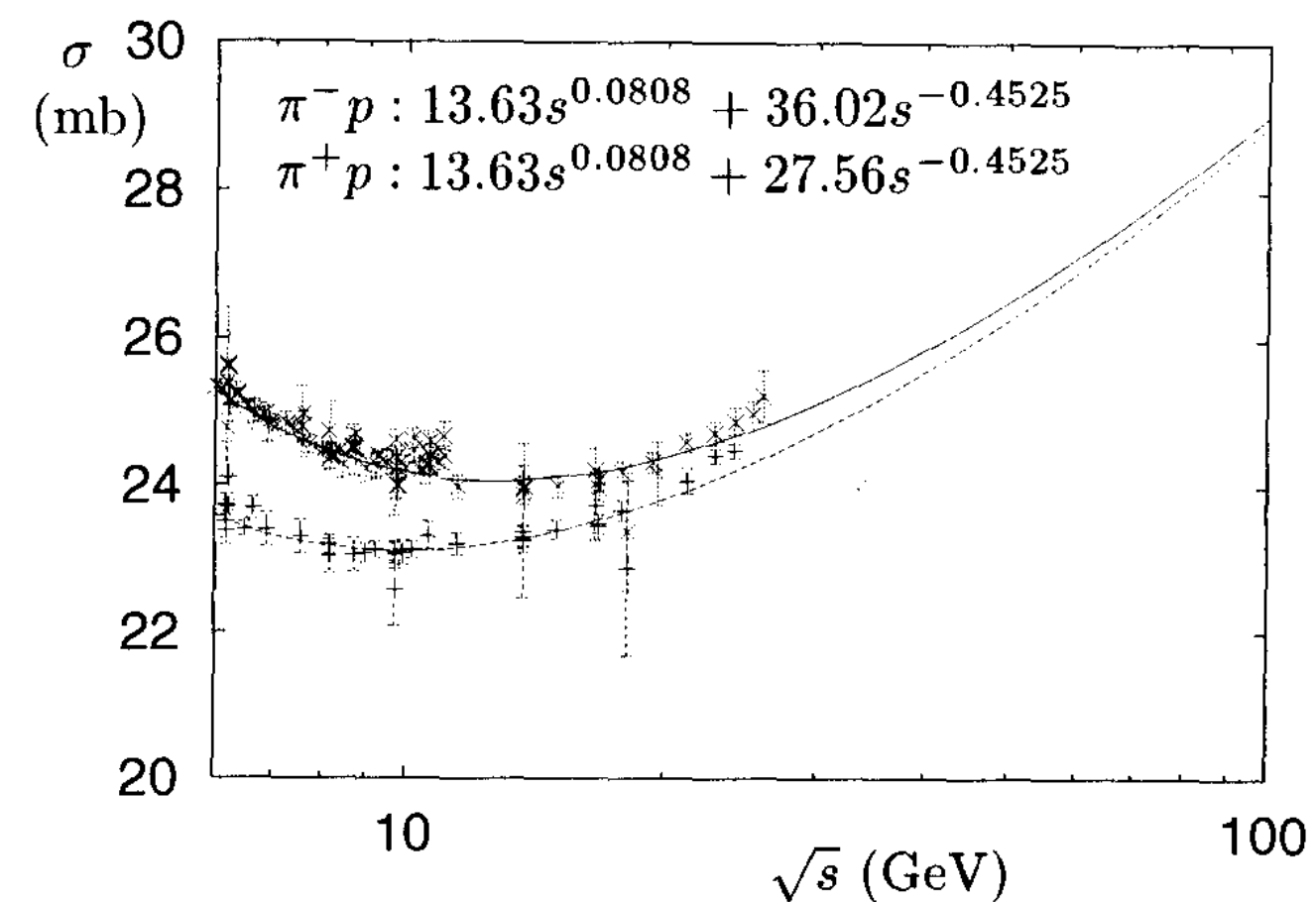
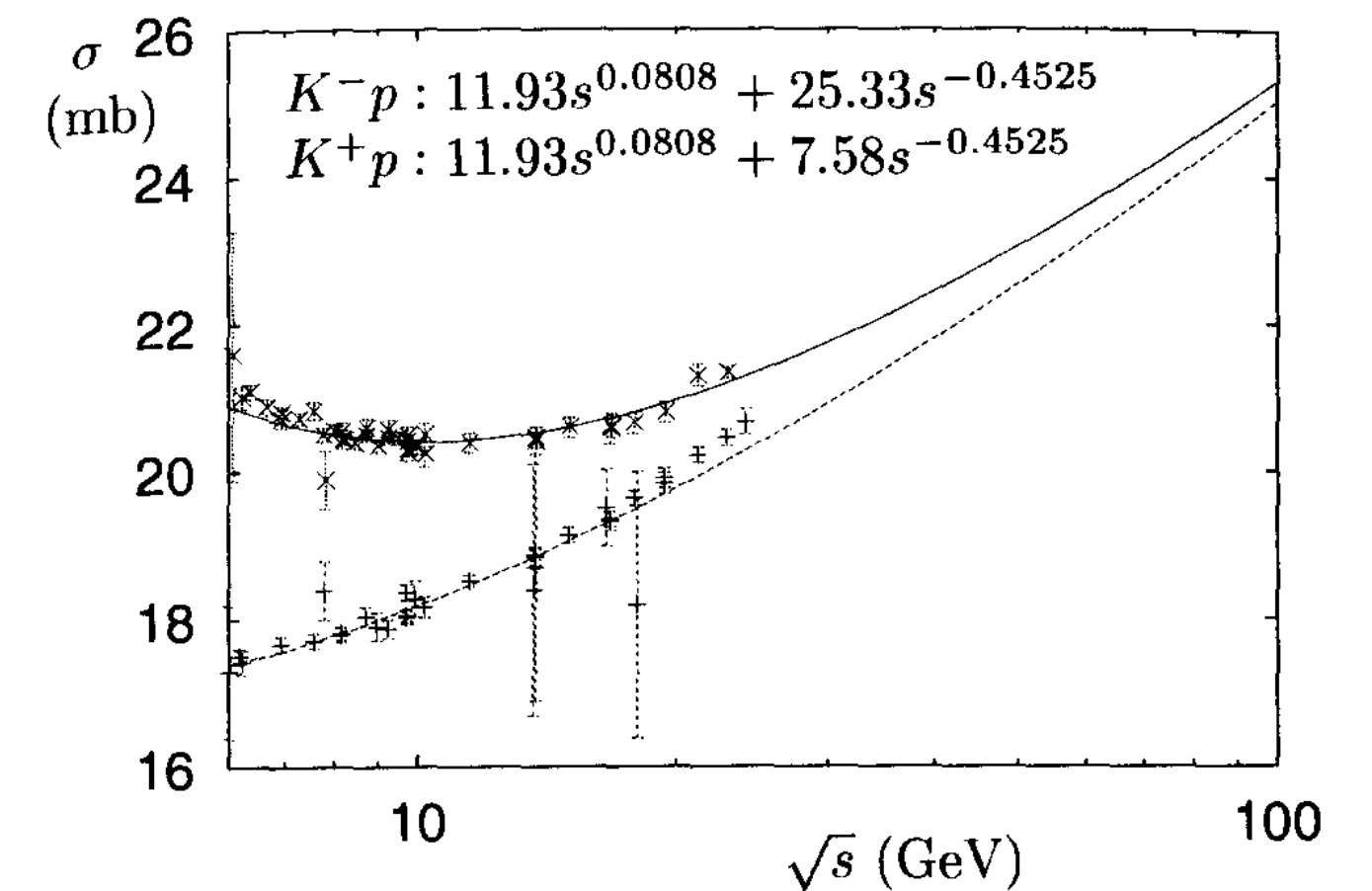
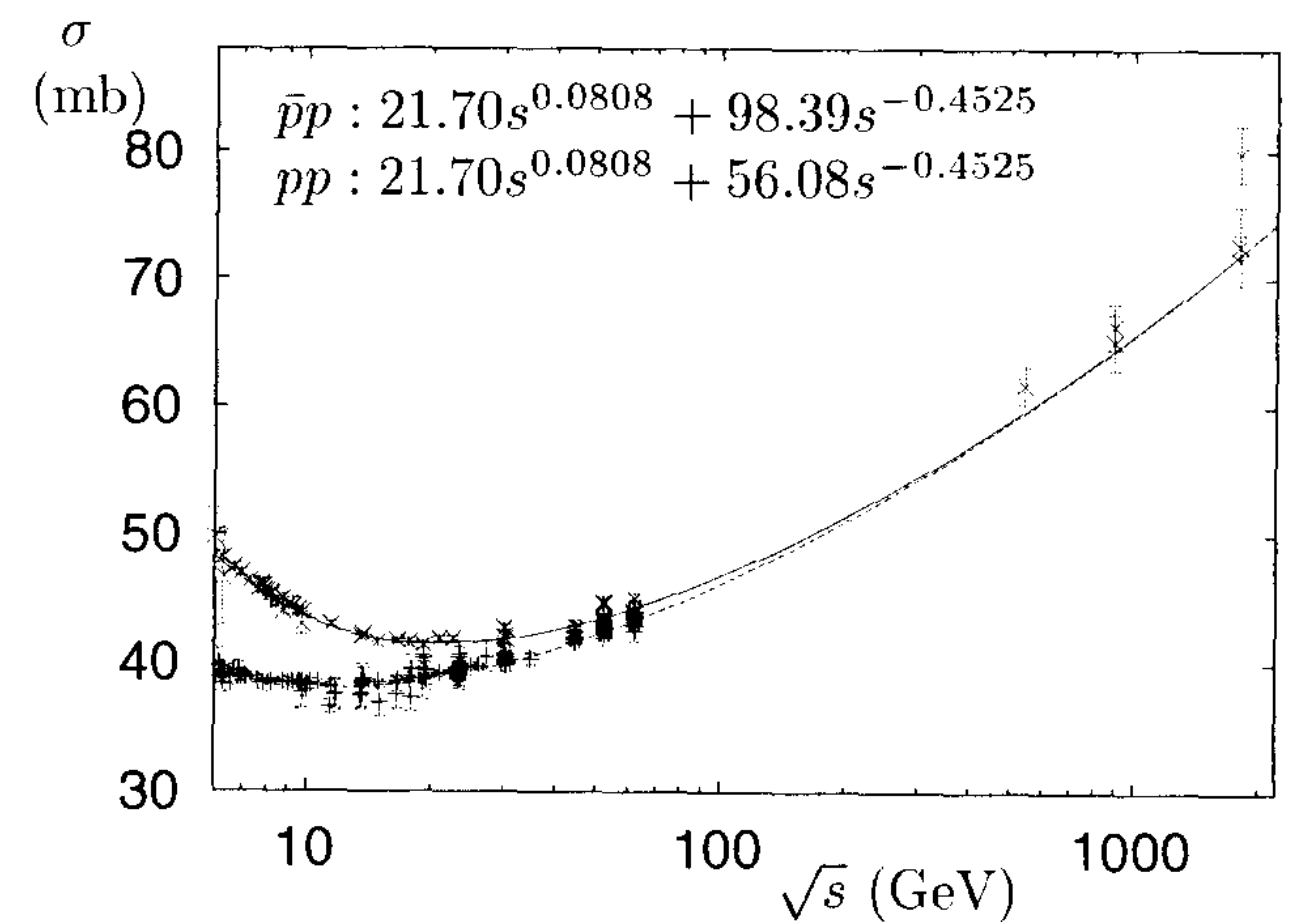
$$j_P(t) \approx 1.08 + 0.25t \quad (\text{GeV units})$$

Elastic cross sections in QCD

$$\sigma \sim s^{j_P(0)-1} \sim s^{0.08}$$

Exchange of even spin glueballs ($J \geq 2$)

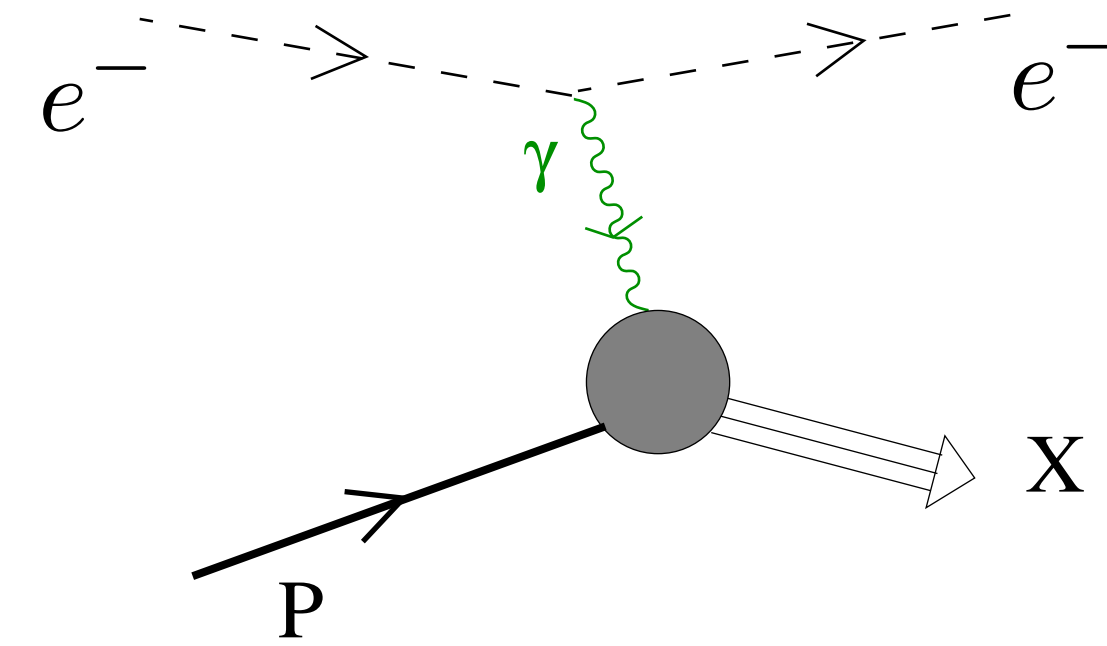
$$\mathcal{O}_J \sim F_{\alpha[\beta_1} D_{\beta_2} \dots D_{\beta_{J-1}} F_{\beta_J]}^{\alpha}$$



Deep Inelastic Scattering

Deep Inelastic Scattering

- Pomeron enters also in diffractive processes. For example DIS, where electron interacts with proton via exchange of off-shell photon



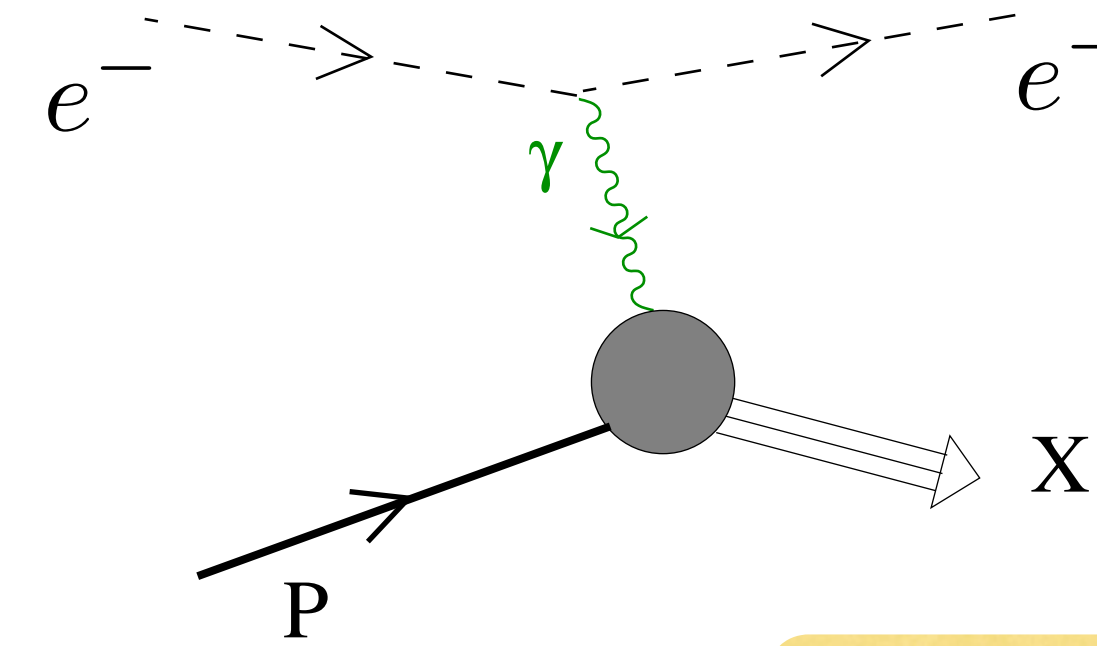
- Optical theorem

$$\sum_X \left| \begin{array}{c} \gamma \\ \nearrow \\ \text{---} \bullet \text{---} \\ \searrow \\ P \end{array} \begin{array}{c} \text{---} \text{---} \\ \longrightarrow \\ X \end{array} \right|^2 = \text{Im} \begin{array}{c} \gamma \\ \nearrow \\ \text{---} \bullet \text{---} \\ \searrow \\ P \end{array} \begin{array}{c} \gamma \\ \nearrow \\ \text{---} \bullet \text{---} \\ \searrow \\ P \end{array} \quad (t=0)$$

The equation states that the sum over all final states X of the squared magnitude of the transition amplitude for a proton (P) to transition to state X via the exchange of a virtual photon (γ), is equal to the imaginary part of the forward scattering amplitude (where the final state is also a proton P) at zero momentum transfer ($t=0$).

Deep Inelastic Scattering

- Pomeron enters also in diffractive processes. For example DIS, where electron interacts with proton via exchange of off-shell photon



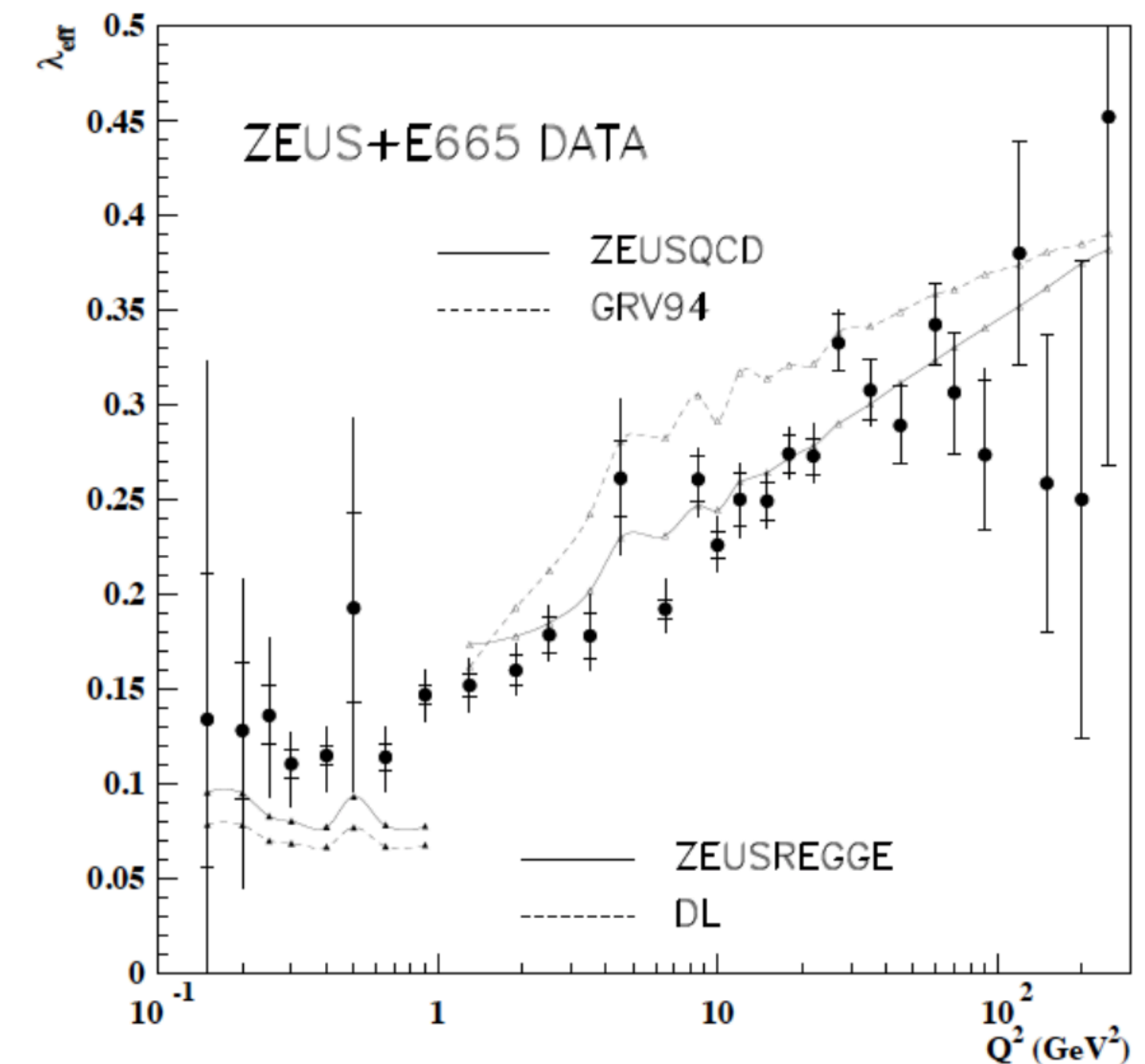
$$j_0 \sim 1.1 - 1.4$$

- Optical theorem

$$\sum_X \left| \text{Diagram} \right|^2 = \text{Im} \left(\text{Diagram} \right) \quad (t = 0)$$

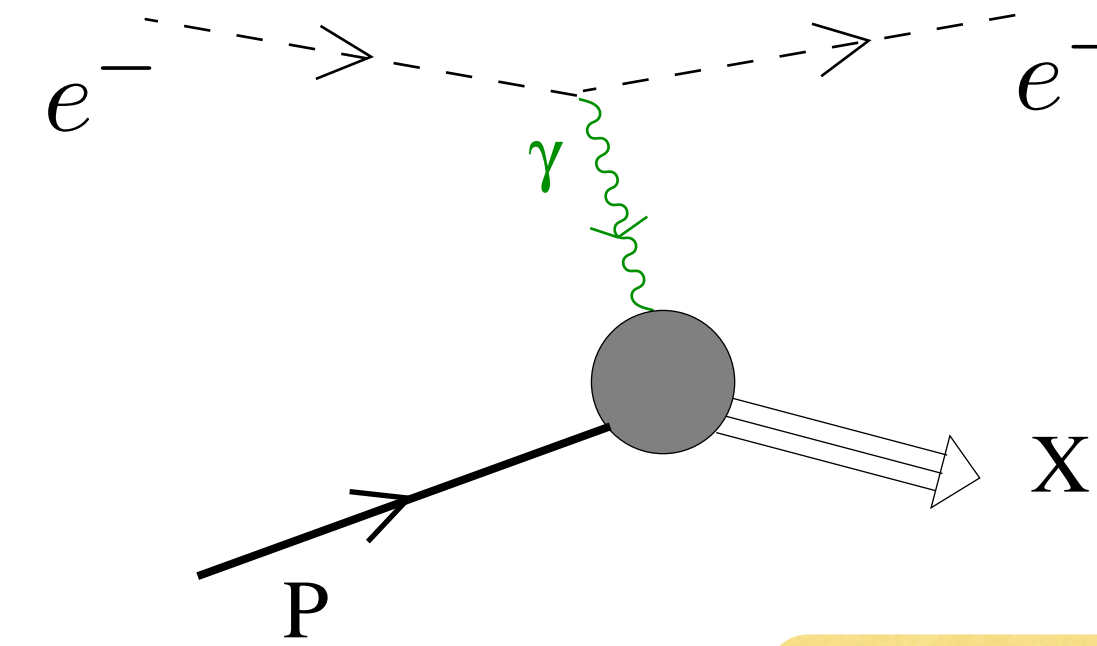
The diagram on the left shows a proton (P) interacting with a virtual photon (gamma) to produce a final state X. The diagram on the right shows a proton (P) interacting with a virtual photon (gamma) to produce another proton (P).

- Regge limit corresponds to low x ($s \sim Q^2/x$)



Deep Inelastic Scattering

- Pomeron enters also in diffractive processes. For example DIS, where electron interacts with proton via exchange of off-shell photon



$$j_0 \sim 1.1 - 1.4$$

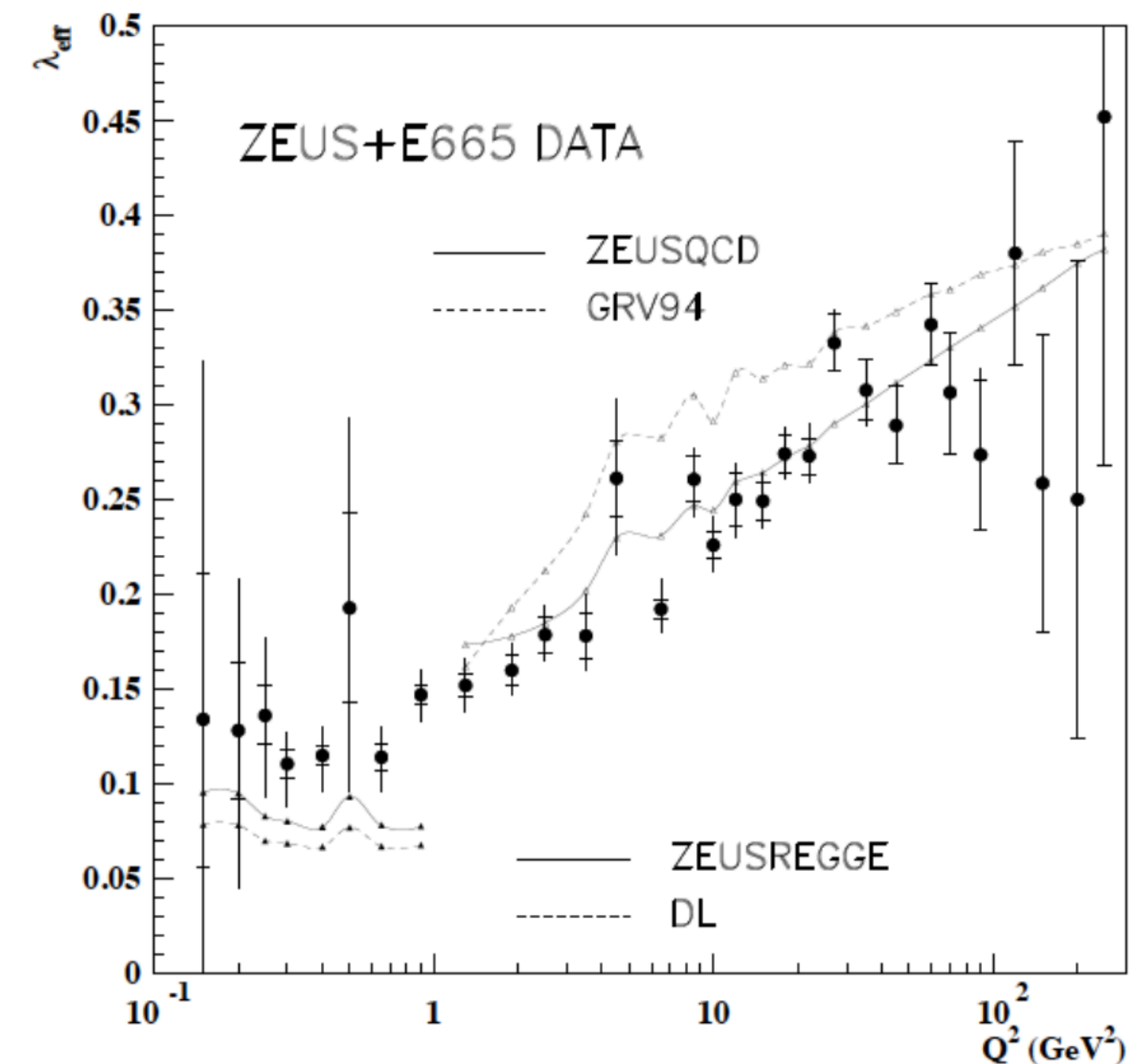
- Optical theorem

$$\sum_X \left| \text{Diagram} \right|^2 = \text{Im} \left(\text{Diagram} \right) \quad (t = 0)$$

The diagram on the left shows a proton (P) interacting with a virtual photon (gamma) to produce a final state X. The diagram on the right shows a proton (P) interacting with a virtual photon (gamma) to produce another proton (P).

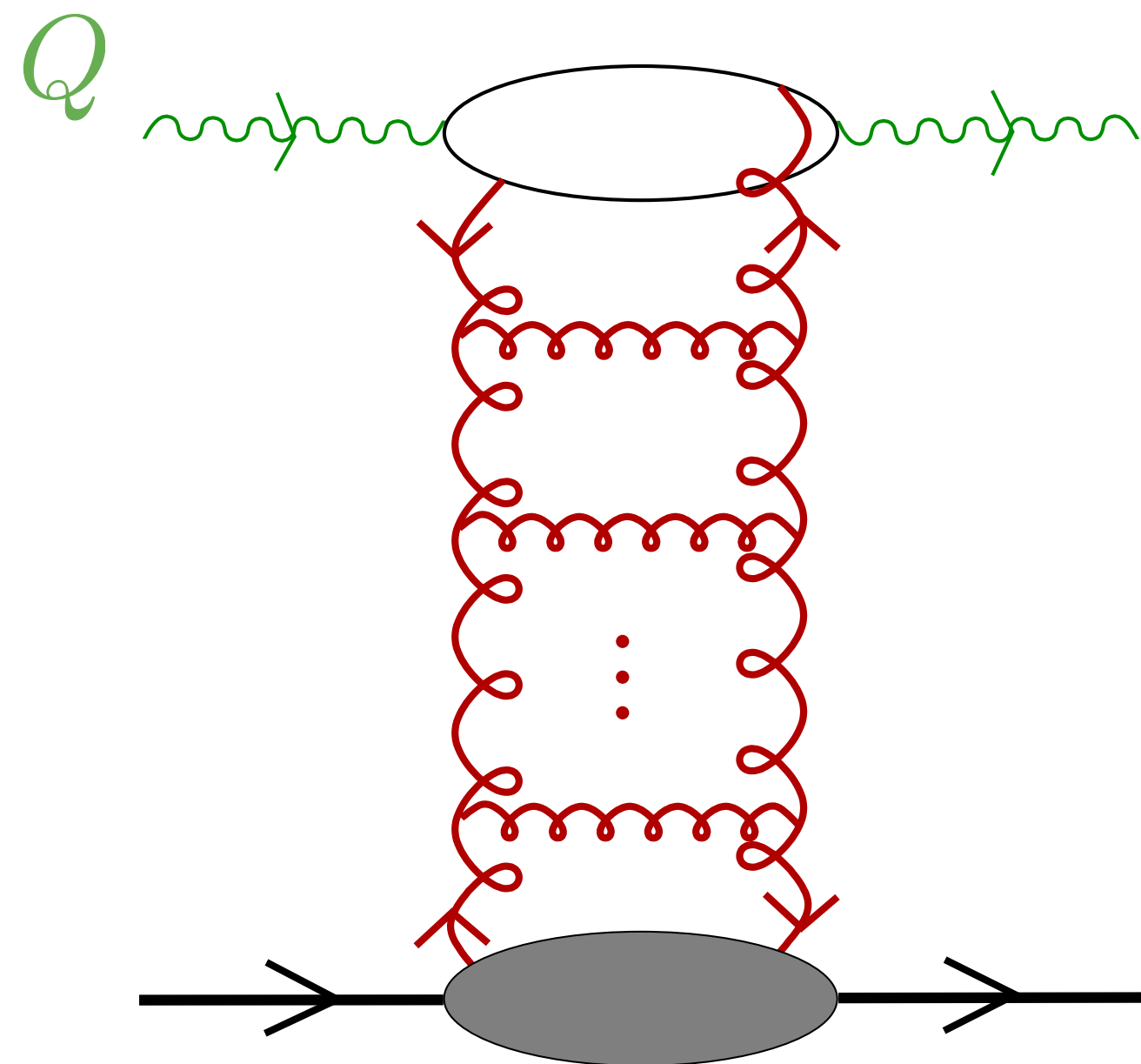
- Regge limit corresponds to low x ($s \sim Q^2/x$)

**Is it the same Regge trajectory?
One or two pomerons (soft and hard)?**



Hard Pomeron [BFKL - Balitsky, Fadin, Kuraev & Lipatov]

Hard Pomeron [BFKL - Balitsky, Fadin, Kuraev & Lipatov]



Two (reggeized) gluon exchange with ladder interactions

Resums $(\alpha_s \ln 1/x)^n$ contributions

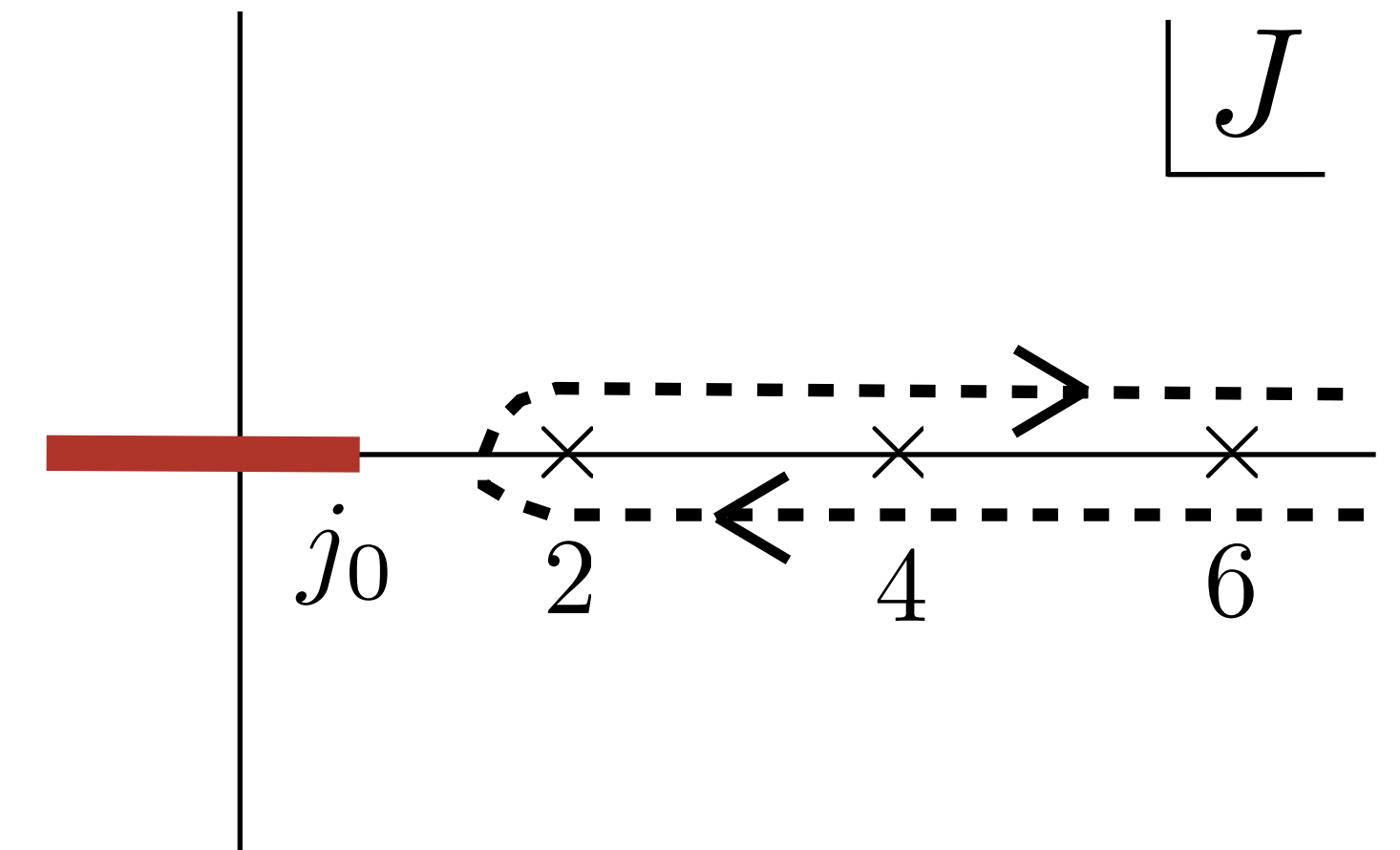
Valid for hard probes $Q \gg \Lambda_{QCD}$

Exhibits conformal symmetry, hard pomeron is a cut in J -plane starting at

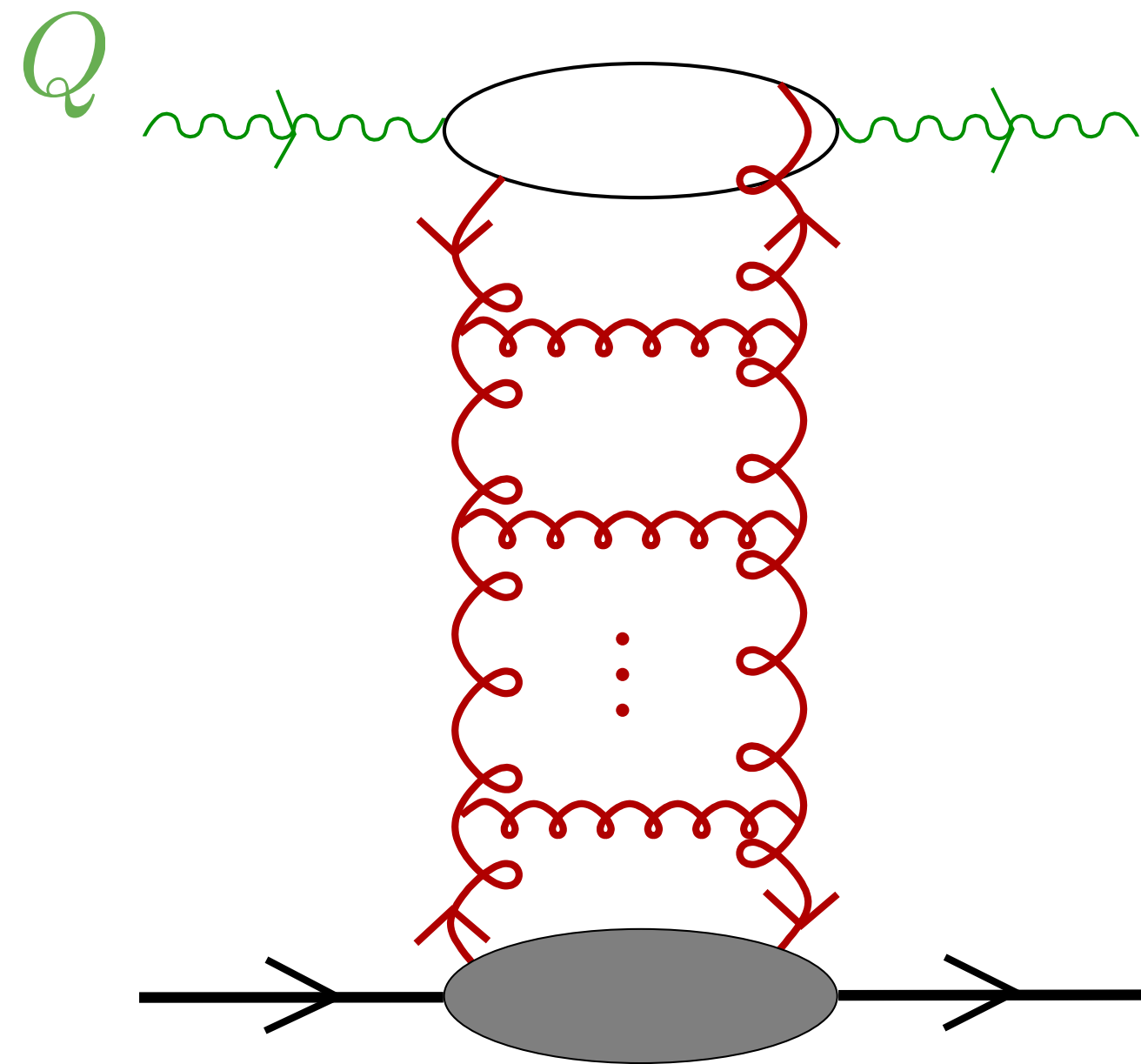
$$j_0 = 1 + \frac{12 \ln 2}{\pi} \alpha_s$$

$\frac{12 \ln 2}{\pi}$

↓
 2.65



Hard Pomeron [BFKL - Balitsky, Fadin, Kuraev & Lipatov]



Two (reggeized) gluon exchange with ladder interactions

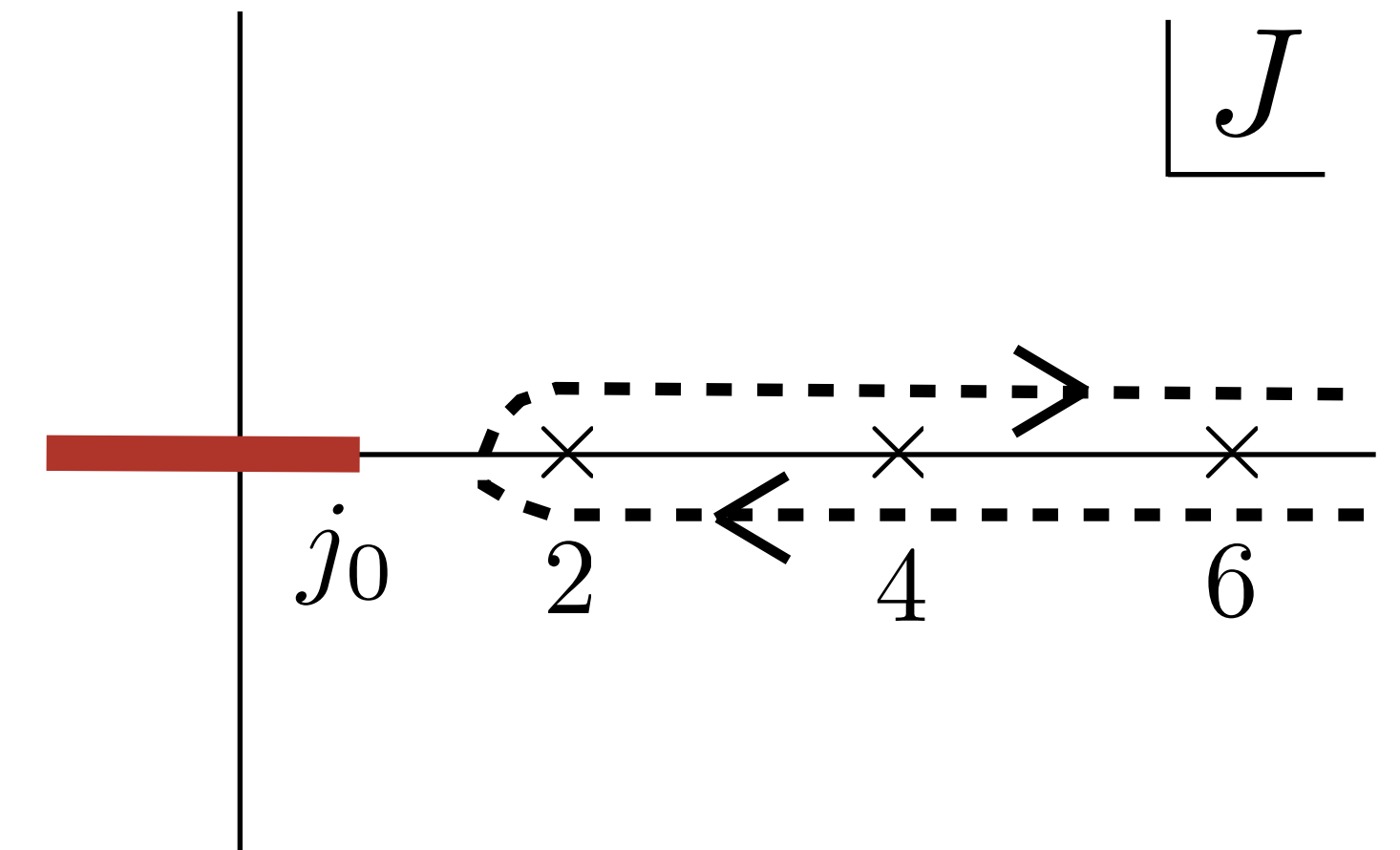
Resums $(\alpha_s \ln 1/x)^n$ contributions

Valid for hard probes $Q \gg \Lambda_{QCD}$

Exhibits conformal symmetry, hard pomeron is a cut in J-plane starting at

$$j_0 = 1 + \frac{12 \ln 2}{\pi} \alpha_s$$

2.65



- Breaking conformal symmetry, explains well DIS data outside the confining region $Q \sim \Lambda_{QCD}$

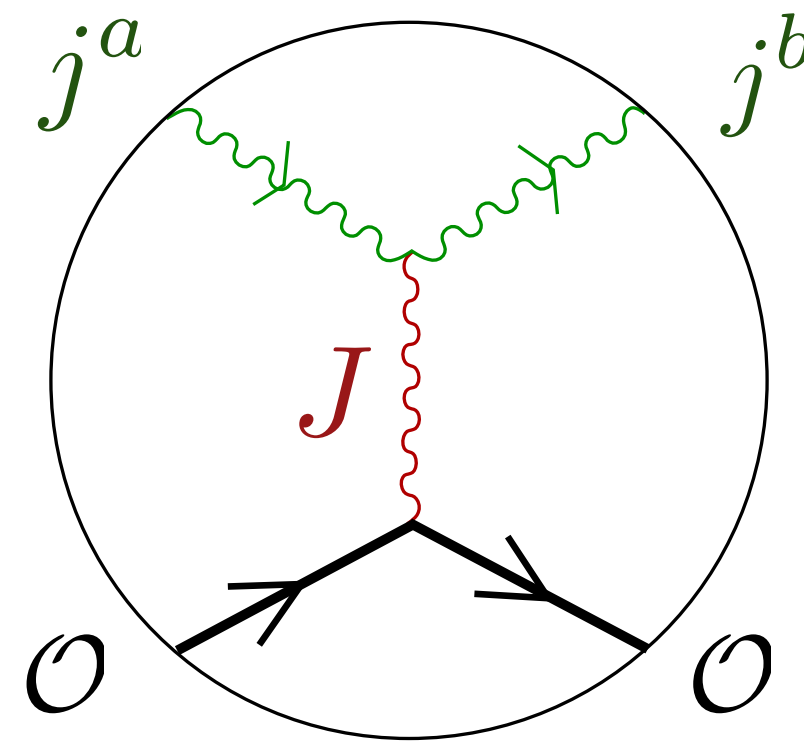
[Kowalski, Lipatov, Ross, Watt 10]

Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

- At strong coupling pomeron trajectory described by string theory graviton Regge trajectory in Anti-de Sitter space (large N, conformal theory $\mathcal{N} = 4$ SYM)

$$(D^2 - m^2) h_{a_1 \dots a_J} = 0$$

$$\text{with } m^2 = \Delta(\Delta - 4) - J$$



Again a cut in J-plane, starting at

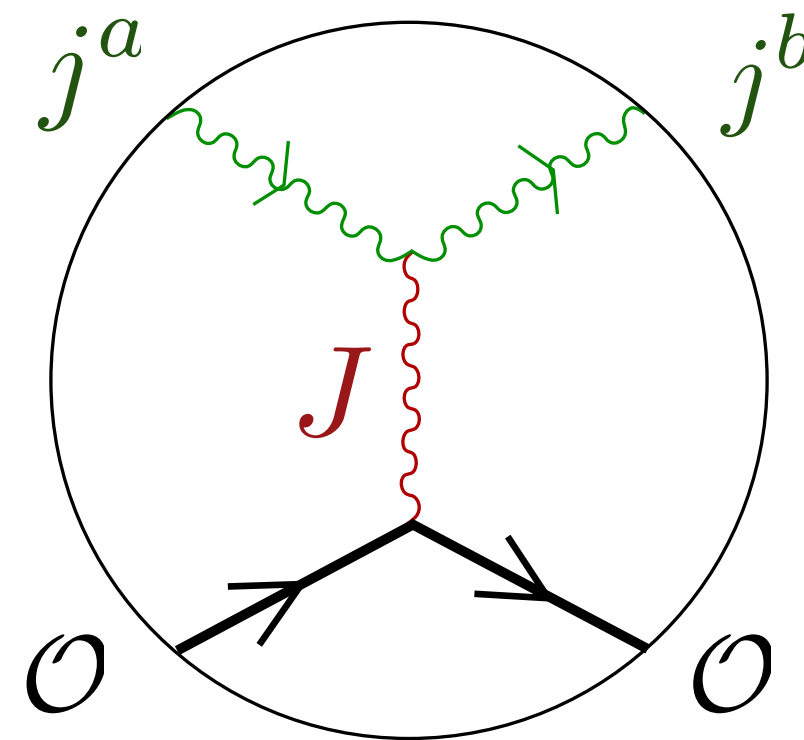
$$j_0 = 2 - \frac{2}{\sqrt{\lambda}} \quad \lambda = \frac{R^4}{\alpha'^2}$$

Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

- At strong coupling pomeron trajectory described by string theory graviton Regge trajectory in Anti-de Sitter space (large N, conformal theory $\mathcal{N} = 4$ SYM)

$$(D^2 - m^2) h_{a_1 \dots a_J} = 0$$

$$\text{with } m^2 = \Delta(\Delta - 4) - J$$



Again a cut in J-plane, starting at

$$j_0 = 2 - \frac{2}{\sqrt{\lambda}} \quad \lambda = \frac{R^4}{\alpha'^2}$$

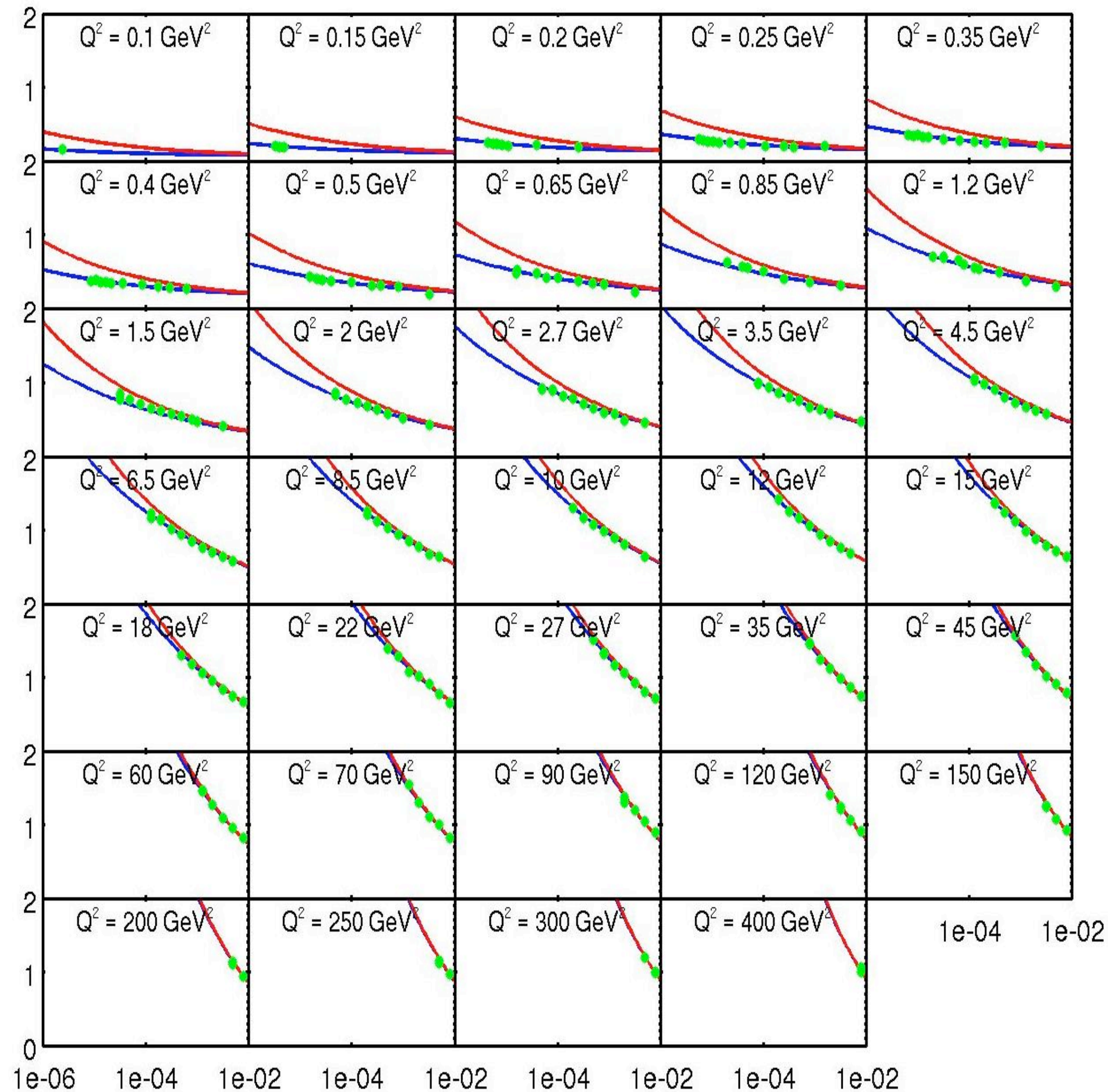
- Explains well low x data for DIS, DVCS, VMP including inside confining region $Q \sim \Lambda_{QCD}$

DIS - [Cornalba, MSC 08; Levin, Potashnikova 10; Brower, Djuric, Sarcevic, Tan 10]

DVSC - [MSC, Djuric 12]

VMP - [MSC, Djuric, Evans 13]

DIS - AdS Pomeron [Brower, Djuric, Sarcevic, Tan 10]



HERA combined data by H1 and ZEUS experiments [Aaron et al 10] with

$$0.10 < Q^2 < 400 \text{ GeV}^2, x < 10^{-2}$$

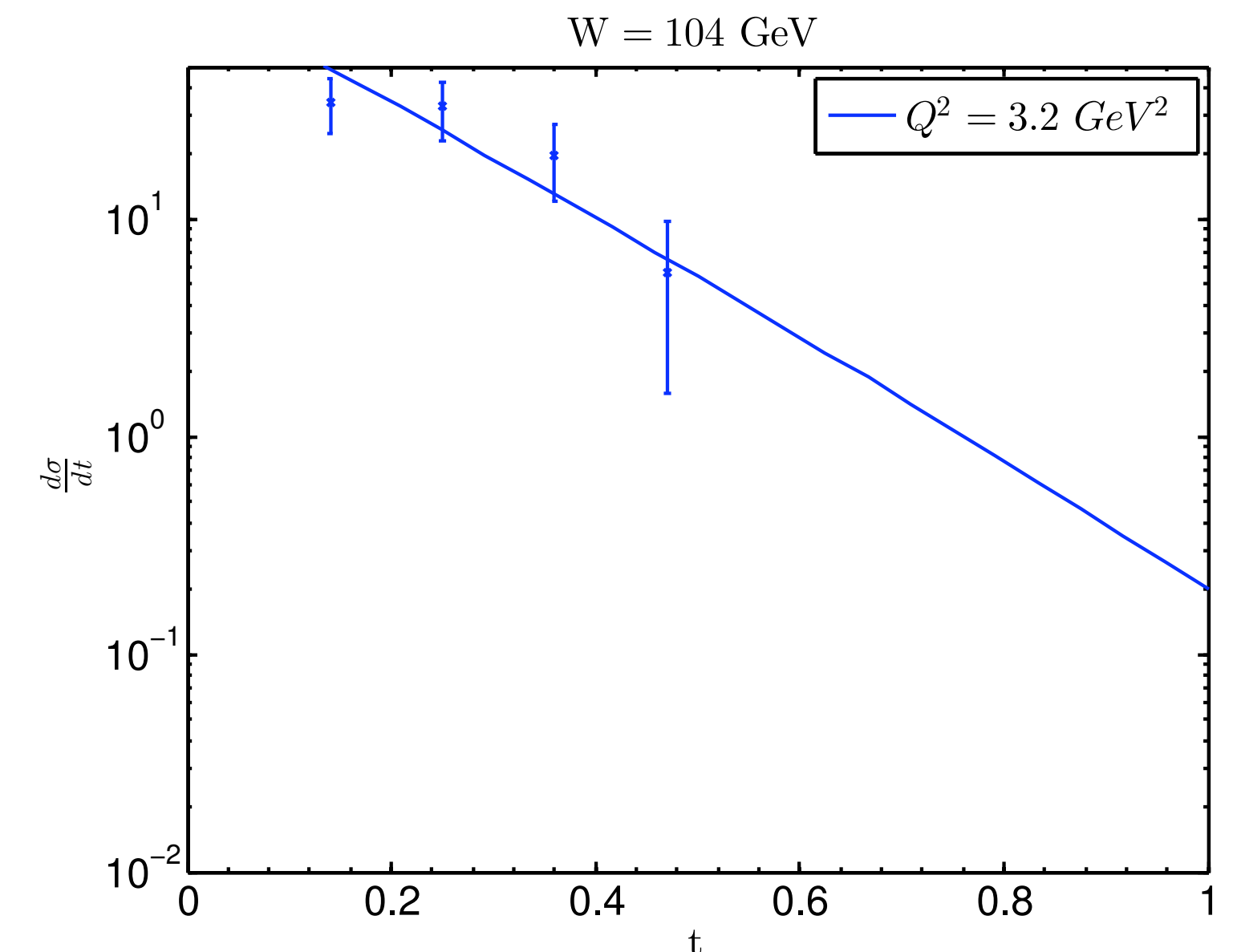
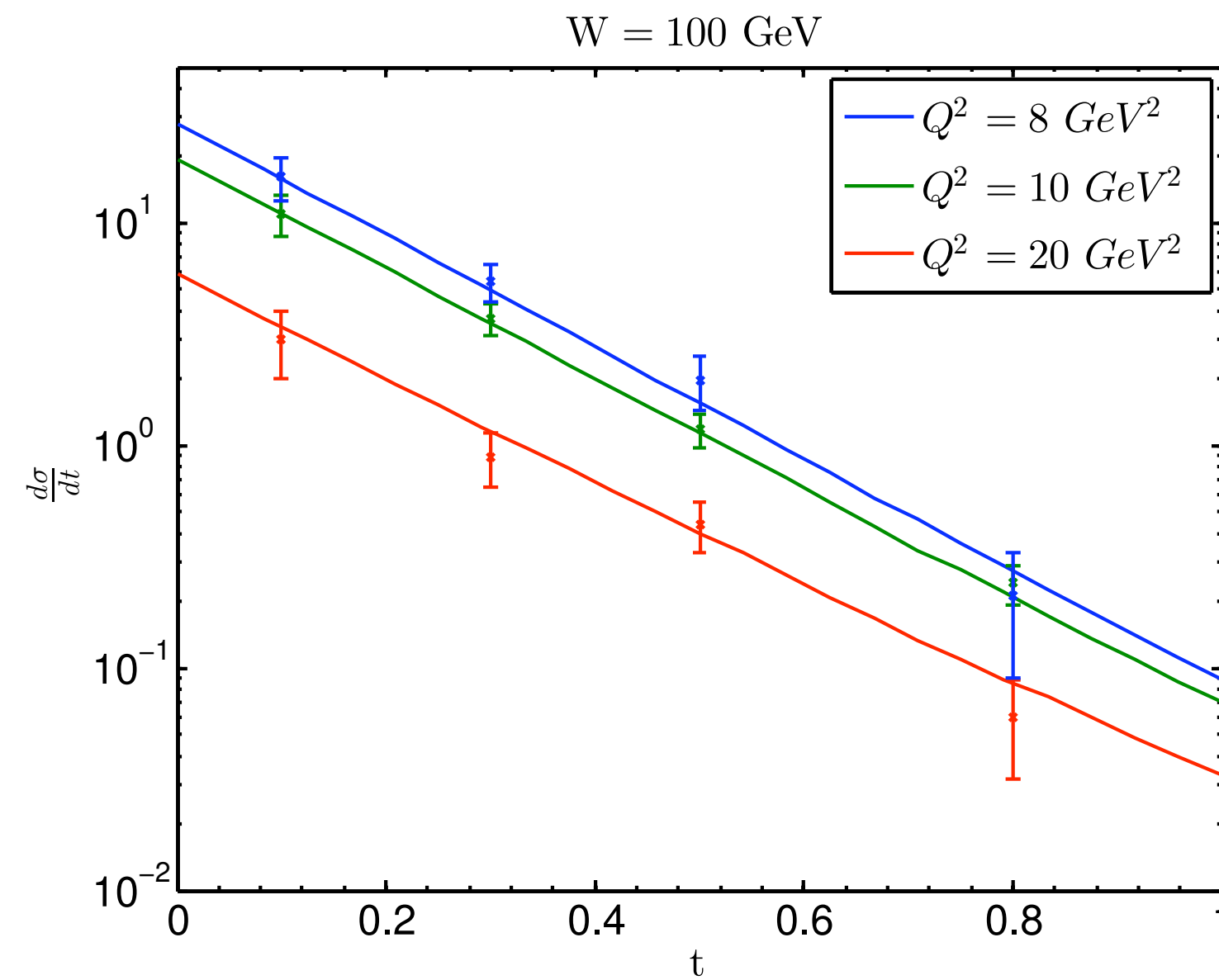
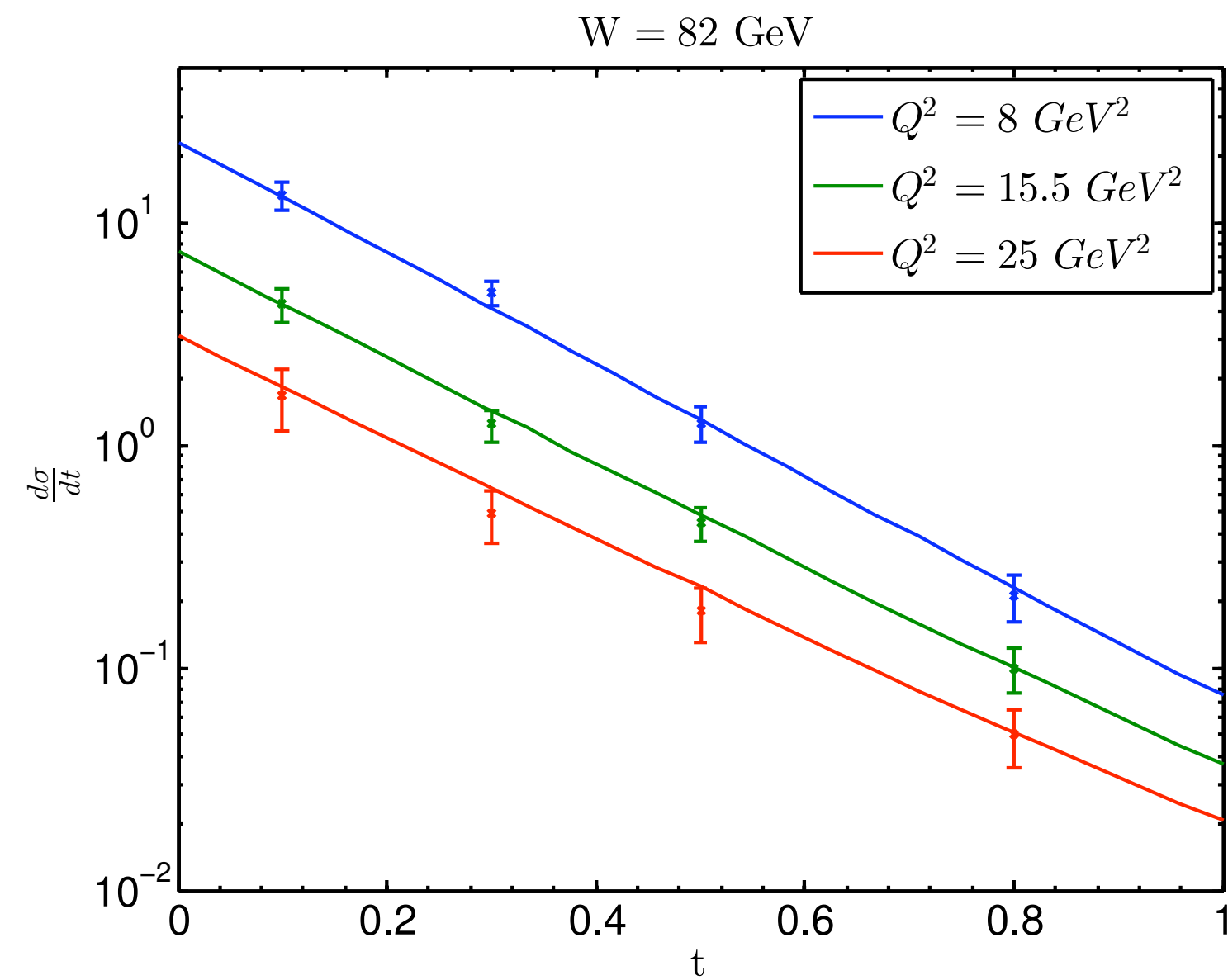
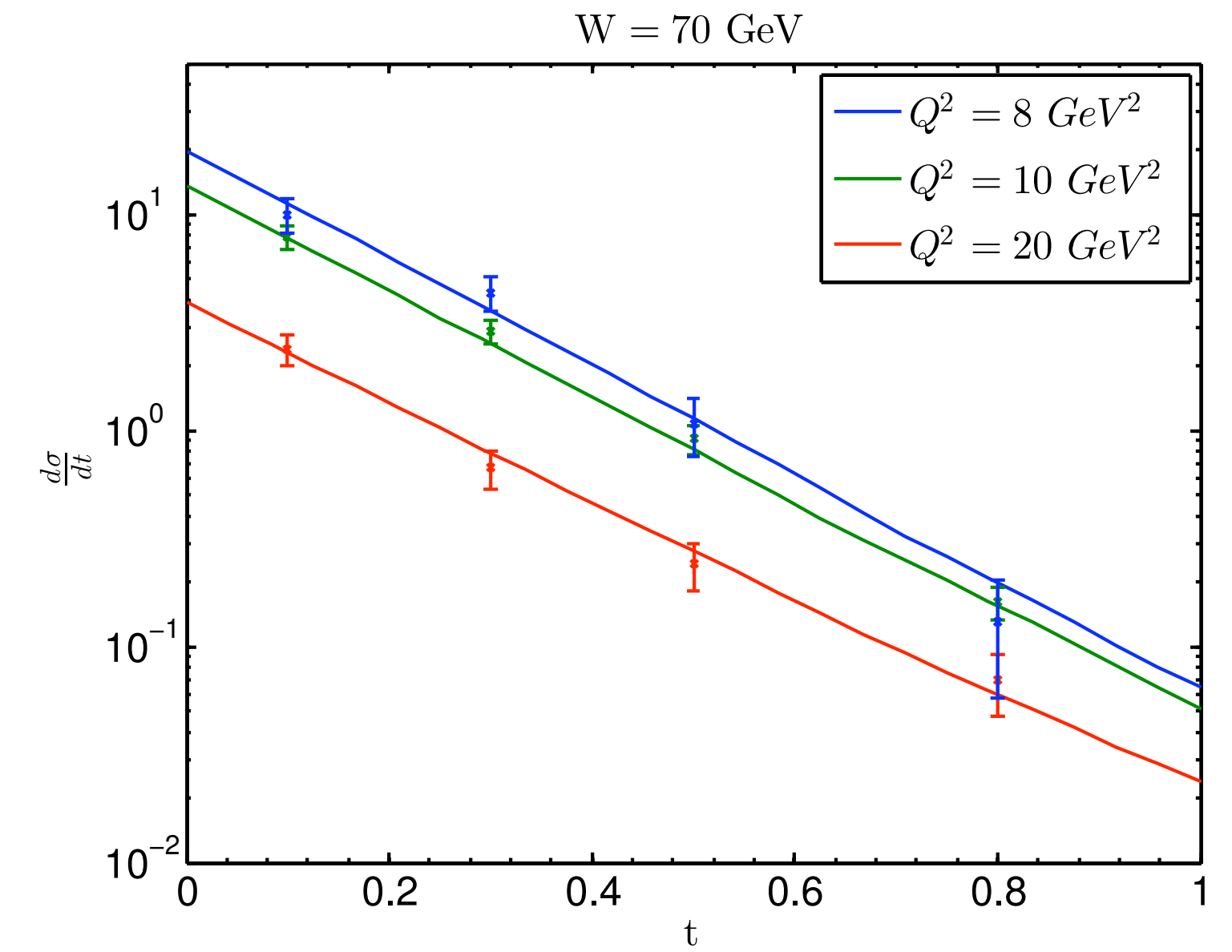
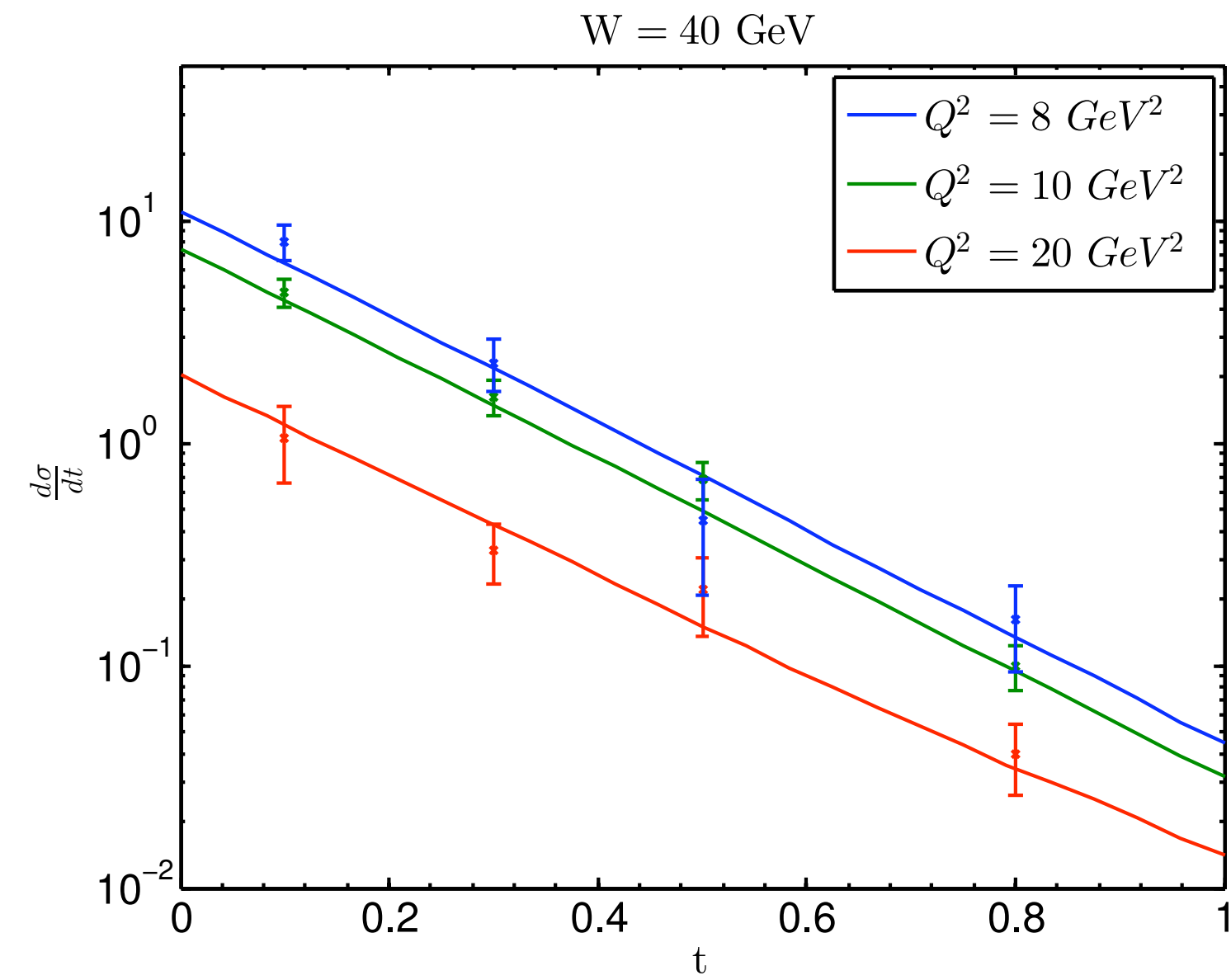
For hard wall model obtained excellent fit with (249 points)

$$\chi^2_{d.o.f.} = 1.07$$

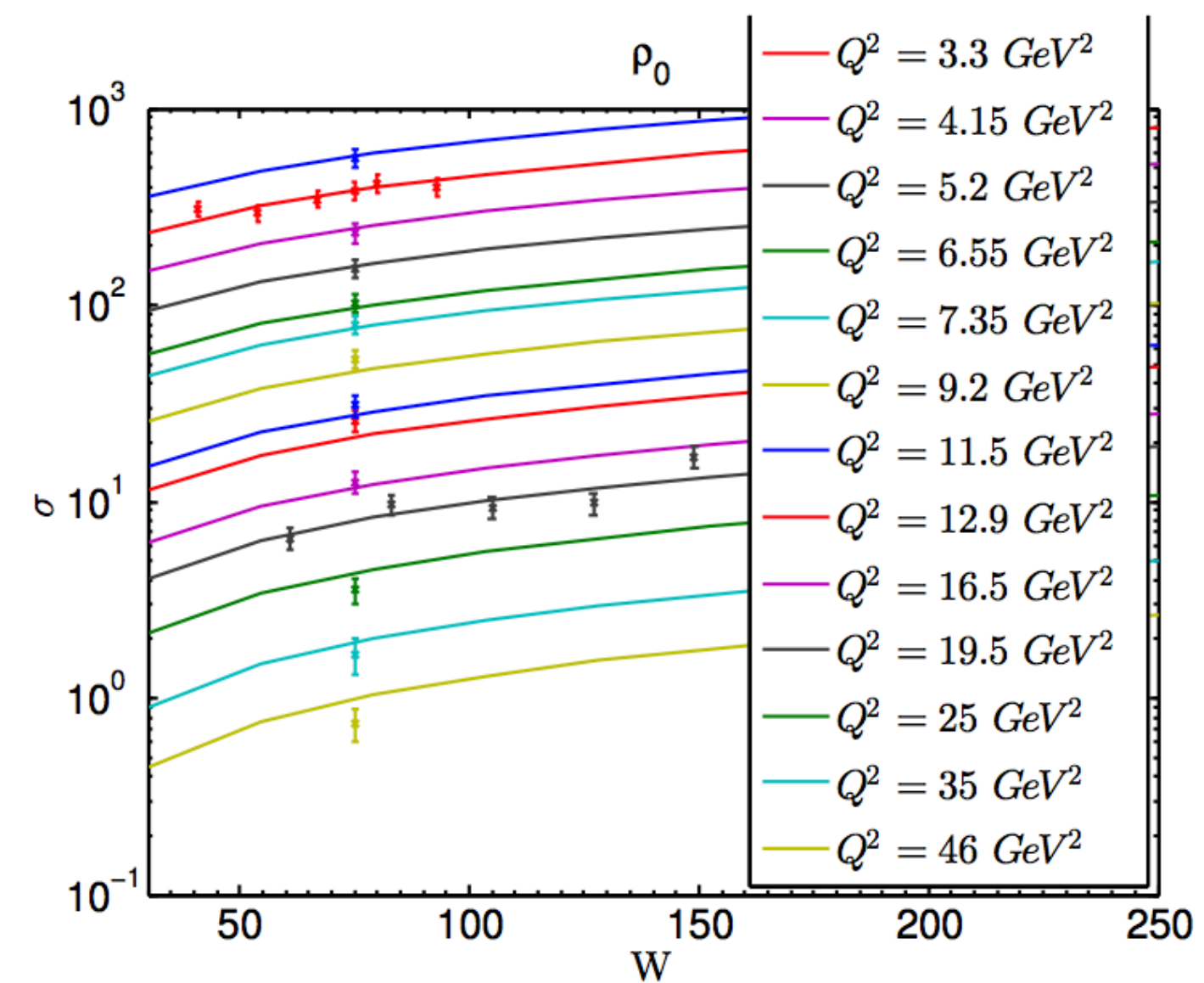
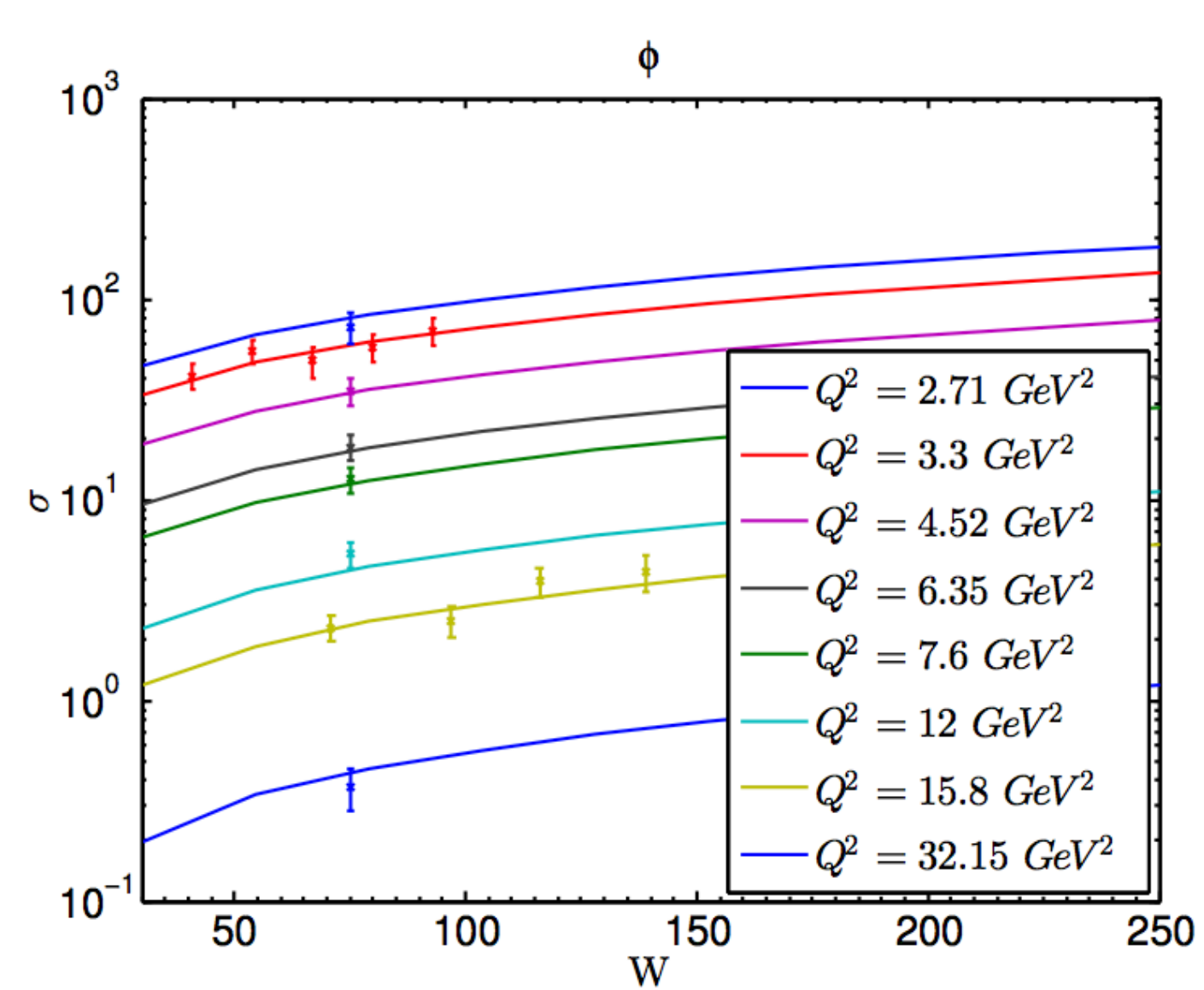
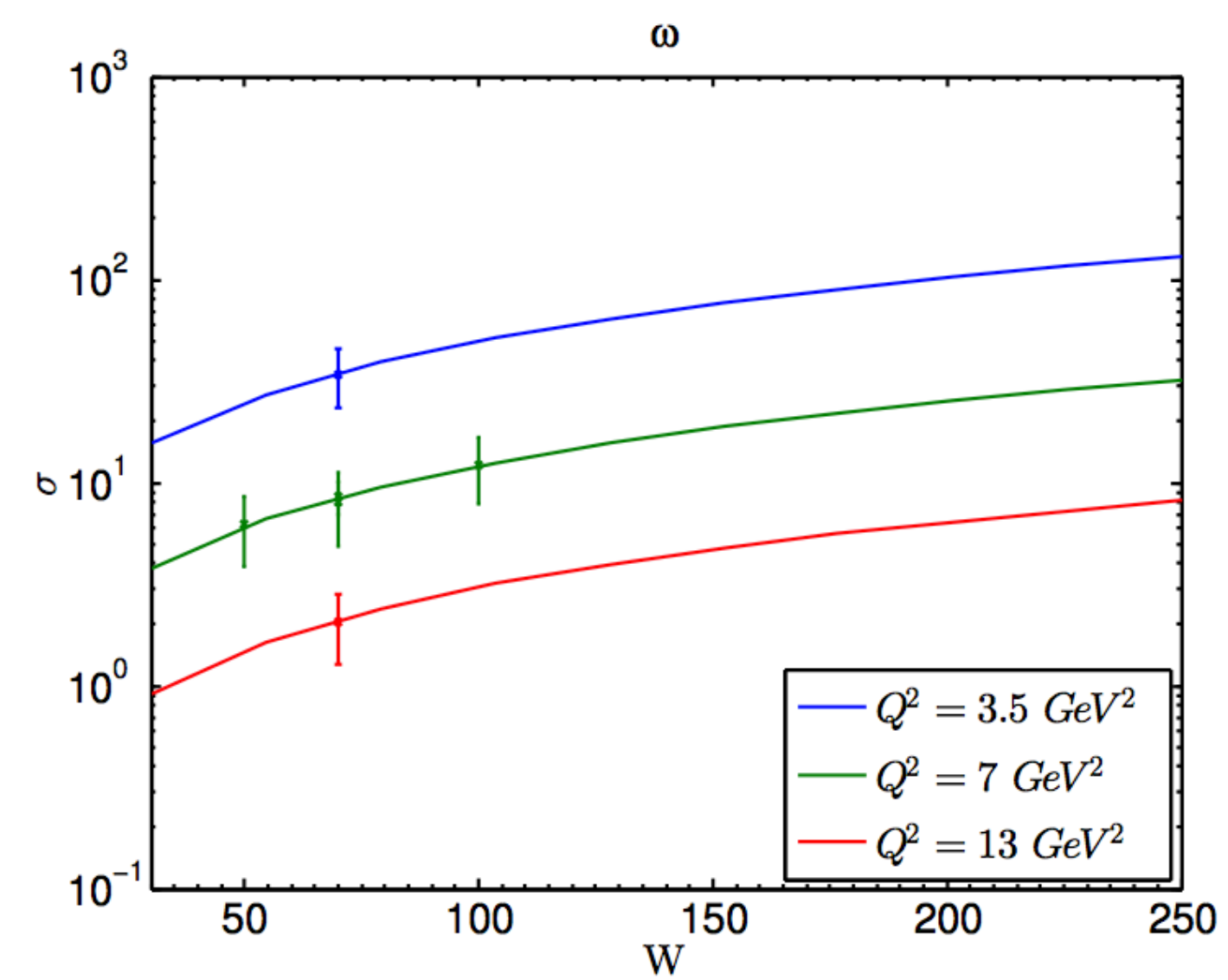
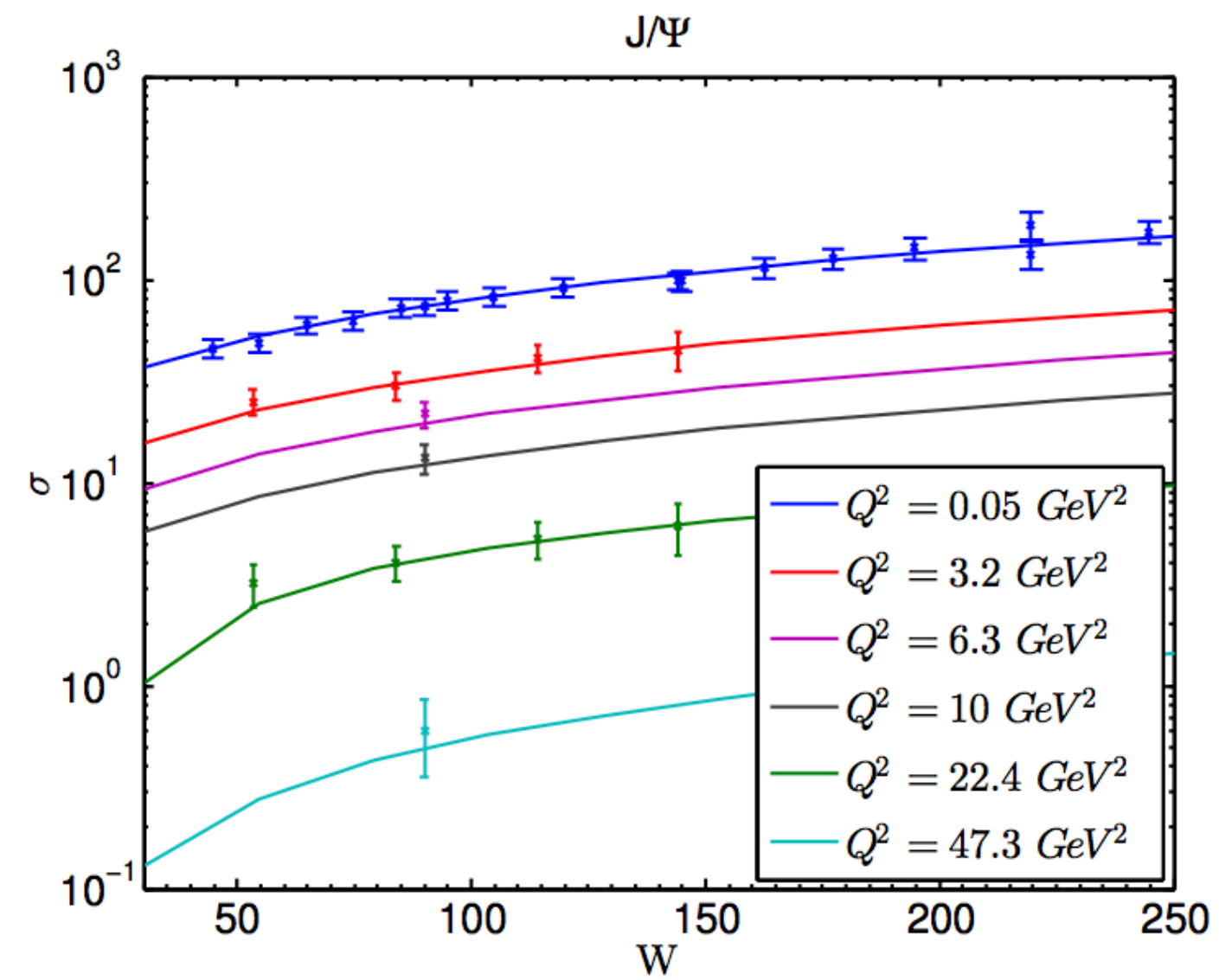
DVSC (differential cross section) [MSC, Djuric 12]

All data (52 points)

$$\chi_{d.o.f.}^2 = 0.51$$



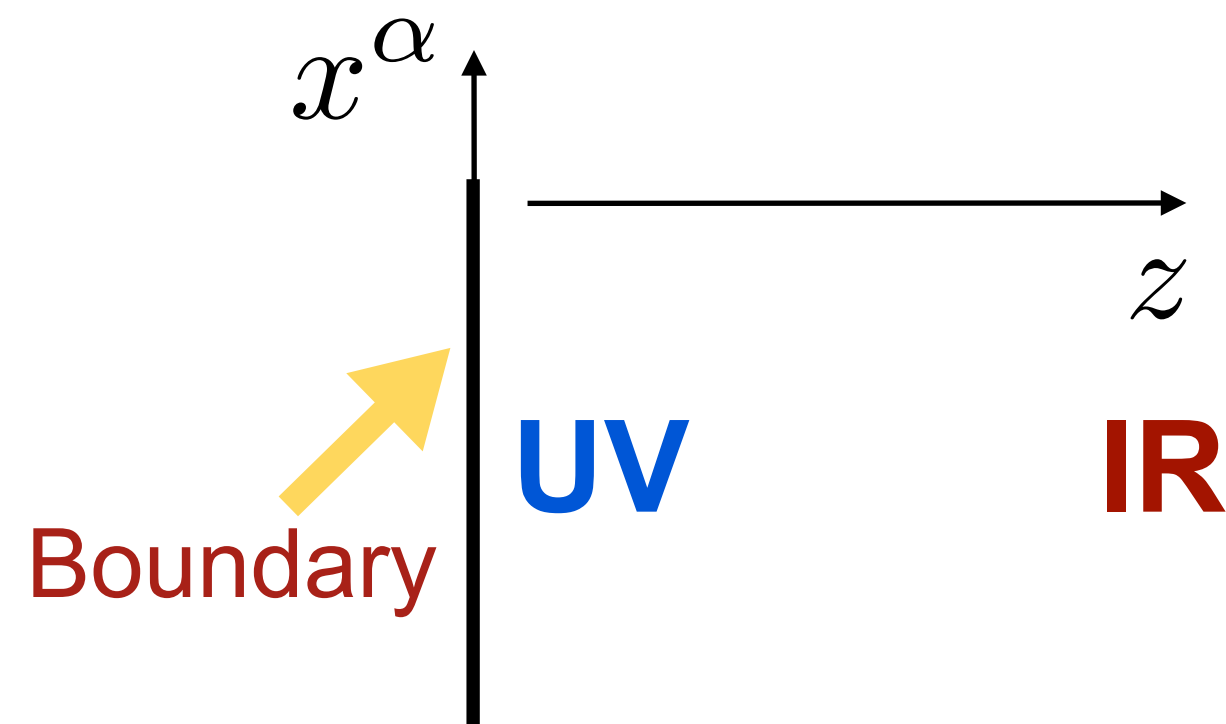
VMP (J/Ψ , ω , ϕ , ρ_0) [MSC, Djuric, Evans 13]



AdS/QCD [Gursoy, Kiritsis, Nitti 07]

- 5D dilaton-gravity phenomenological model constructed to reproduce QCD

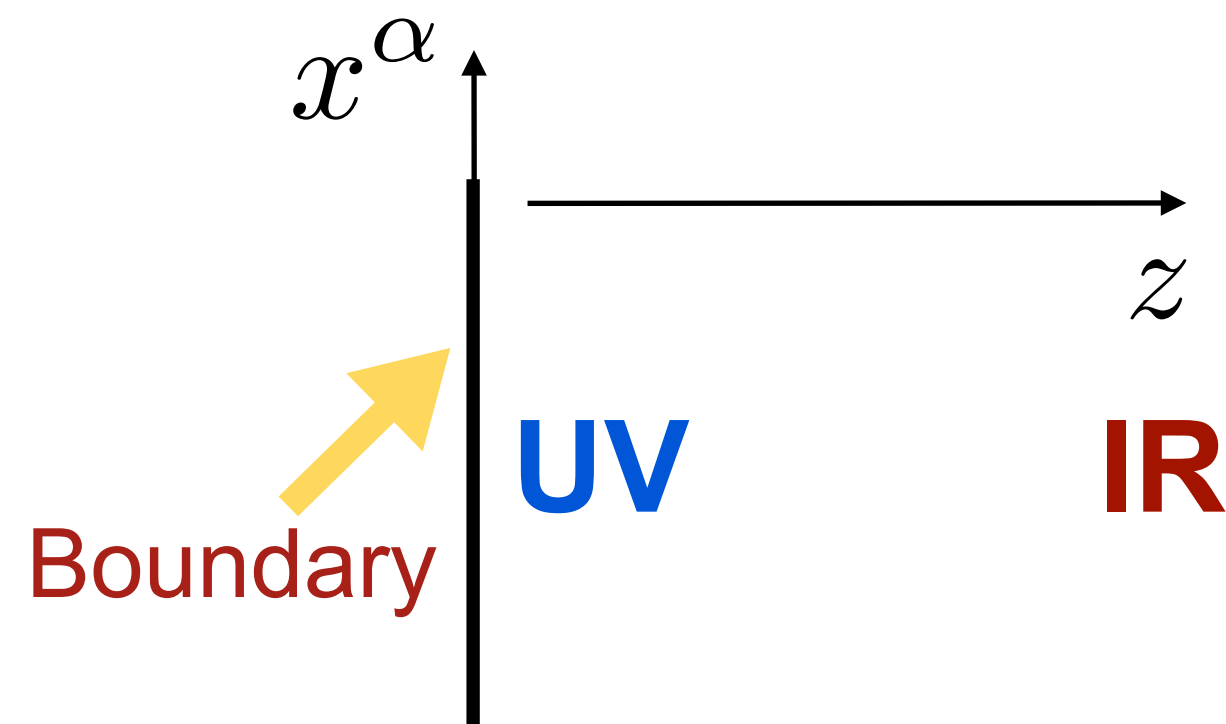
$$S = \frac{1}{2\kappa^2} \int d^5x \sqrt{-g} \left[R - \frac{4}{3} (\partial\phi)^2 + V(\phi) \right]$$



- 5D dilaton-gravity phenomenological model constructed to reproduce QCD

$$S = \frac{1}{2\kappa^2} \int d^5x \sqrt{-g} \left[R - \frac{4}{3} (\partial\phi)^2 + V(\phi) \right]$$

Judicious choice of potential with only 2 free parameters!



Constructed to match QCD perturbative beta function

Reproduces: - heavy quark-antiquark linear potential

- glueball spectrum from lattice simulations

- thermodynamic properties of QGP (bulk viscosity, drag force and jet quenching parameters)

Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

- Operators that contribute are the twist 2 operators

$$\mathcal{O}_J \sim F_{\alpha[\beta_1} D_{\beta_2} \cdots D_{\beta_{J-1}} F_{\beta_J]}^{\alpha}$$

Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

- Operators that contribute are the twist 2 operators

$$\mathcal{O}_J \sim F_{\alpha[\beta_1} D_{\beta_2} \cdots D_{\beta_{J-1}} F_{\beta_J]}^{\alpha}$$

- Dual to string theory spin J field in leading Regge trajectory

$$(D^2 - m^2) h_{a_1 \dots a_J} = 0$$

$$m^2 = \Delta(\Delta - 4) - J, \quad \Delta = \Delta(J)$$

Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

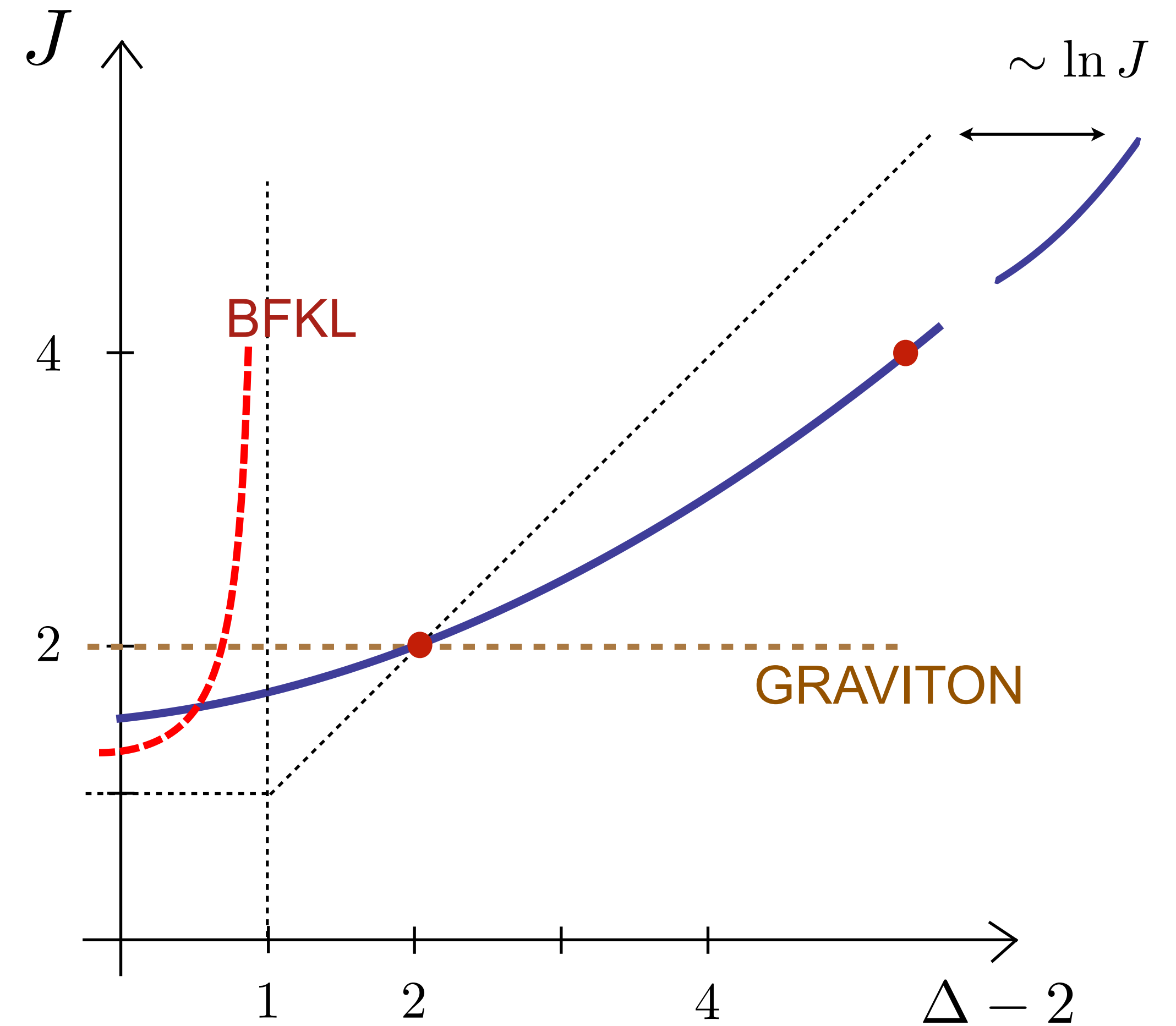
- Operators that contribute are the twist 2 operators

$$\mathcal{O}_J \sim F_{\alpha[\beta_1} D_{\beta_2} \dots D_{\beta_{J-1}} F_{\beta_J]}^{\alpha}$$

- Dual to string theory spin J field in leading Regge trajectory

$$(D^2 - m^2) h_{a_1 \dots a_J} = 0$$

$$m^2 = \Delta(\Delta - 4) - J, \quad \Delta = \Delta(J)$$



Graviton/Pomeron Regge trajectory [Brower, Polchinski, Strassler, Tan 06]

- Operators that contribute are the twist 2 operators

$$\mathcal{O}_J \sim F_{\alpha[\beta_1} D_{\beta_2} \dots D_{\beta_{J-1}} F_{\beta_J]}^{\alpha}$$

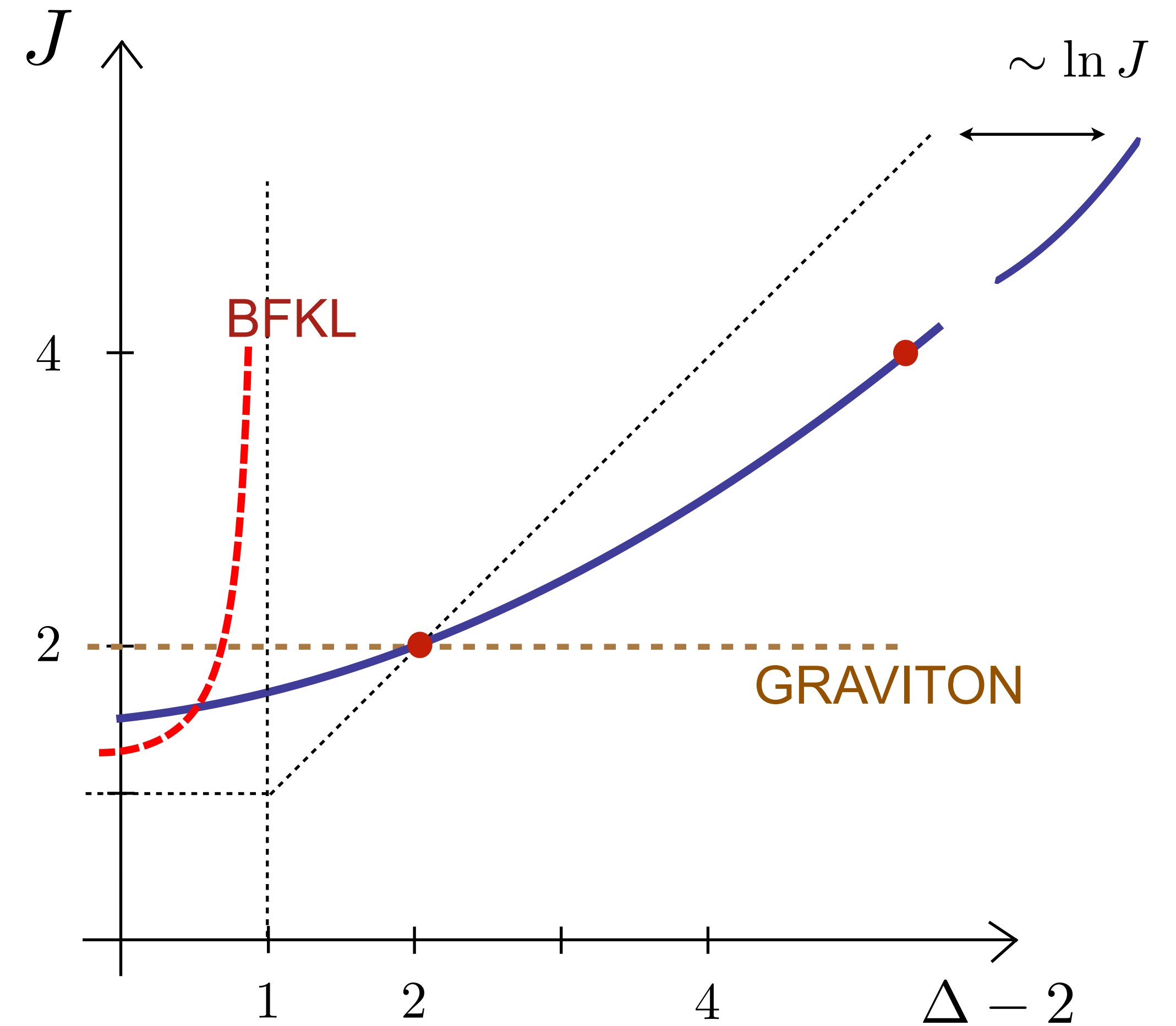
- Dual to string theory spin J field in leading Regge trajectory

$$(D^2 - m^2) h_{a_1 \dots a_J} = 0$$

$$m^2 = \Delta(\Delta - 4) - J, \quad \Delta = \Delta(J)$$

- Diffusion limit

$$J(\Delta) = J_0 + \mathcal{D}(\Delta - 2)^2 \Rightarrow m^2 = \frac{2}{\alpha'}(J - 2) - \frac{J}{L^2}$$



AdS Spin J field

AdS Spin J field

- Spin J field equation of motion (graviton $\Delta = 4$, $J = 2$)

$$\left[(D^2 - 2\partial^b \phi D_b - \Delta(\Delta - 4)) g_{a_1 b} g_{a_2 c} + J R_{a_1 b a_2 c} \right] h^{bc}_{a_3 \dots a_J} = 0$$

$\Delta = \Delta(J)$ e.g. in
perturbation theory
 $\Delta = 2 + J + \gamma_J$

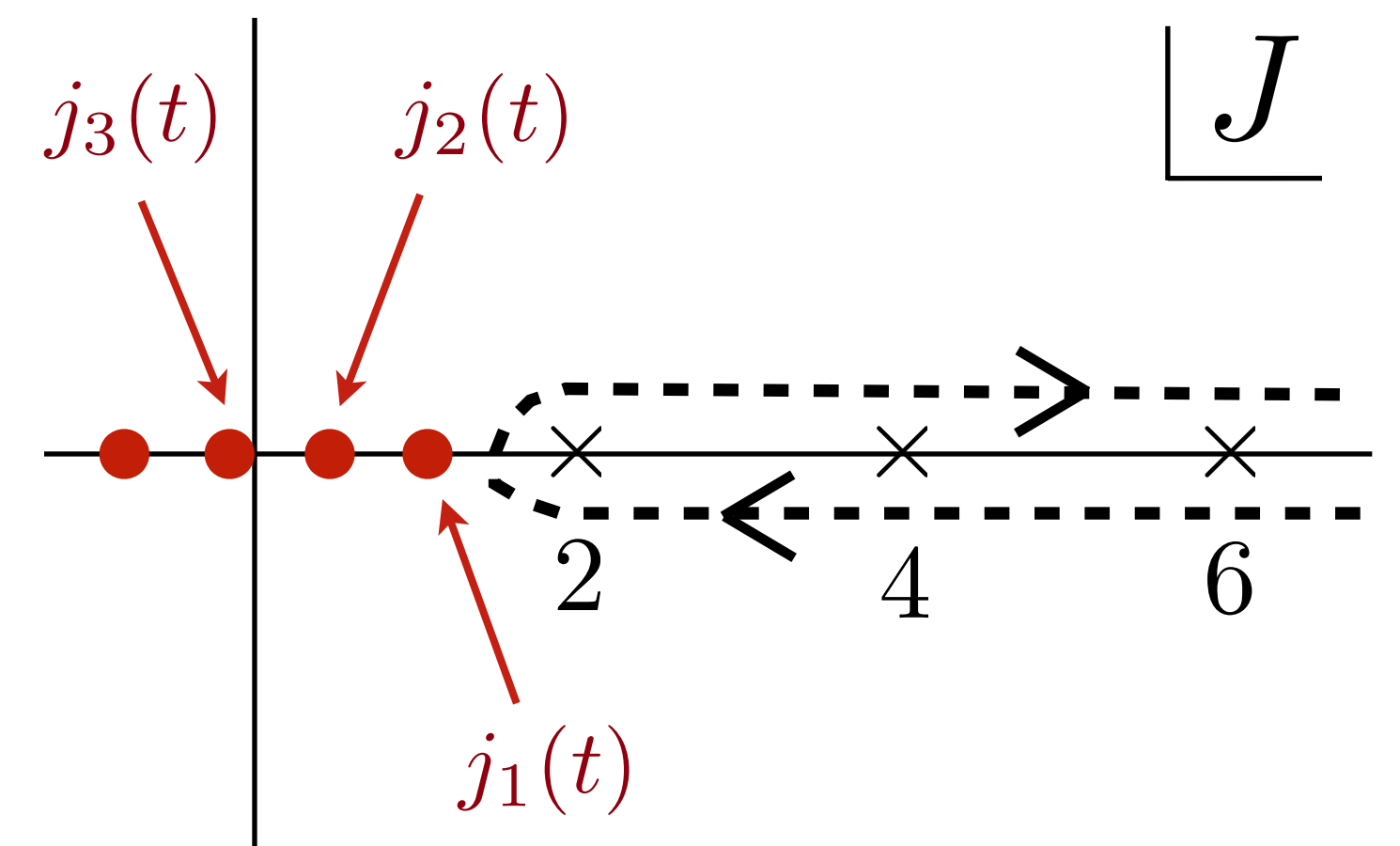
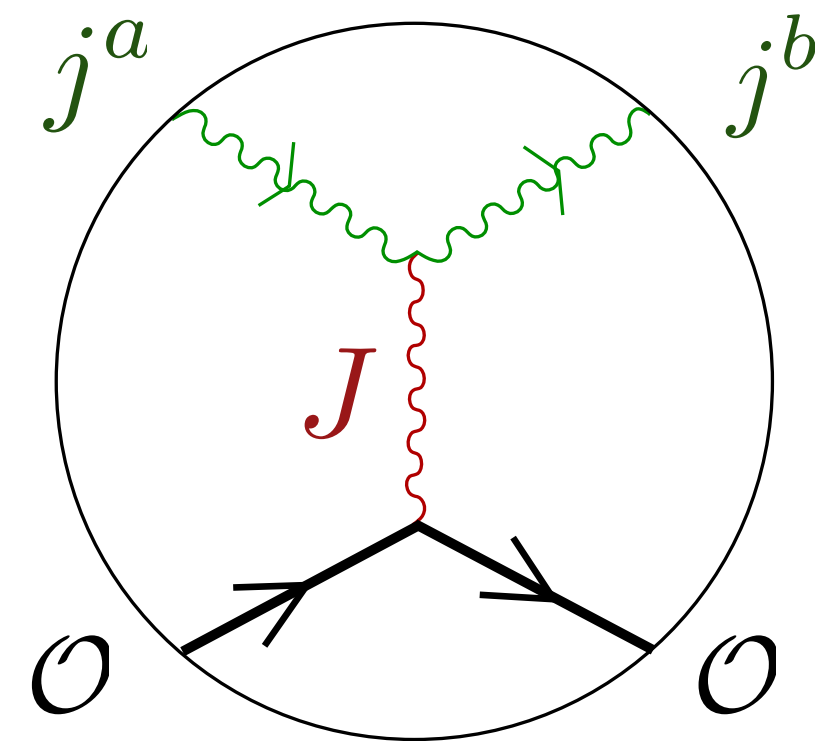
AdS Spin J field

- Spin J field equation of motion (graviton $\Delta = 4$, $J = 2$)

$$\left[(D^2 - 2\partial^b \phi D_b - \Delta(\Delta - 4)) g_{a_1 b} g_{a_2 c} + J R_{a_1 b a_2 c} \right] h^{bc}_{a_3 \dots a_J} = 0$$

$\Delta = \Delta(J)$ e.g. in
perturbation theory
 $\Delta = 2 + J + \gamma_J$

- Sum over spin J exchanges in 5D dual theory



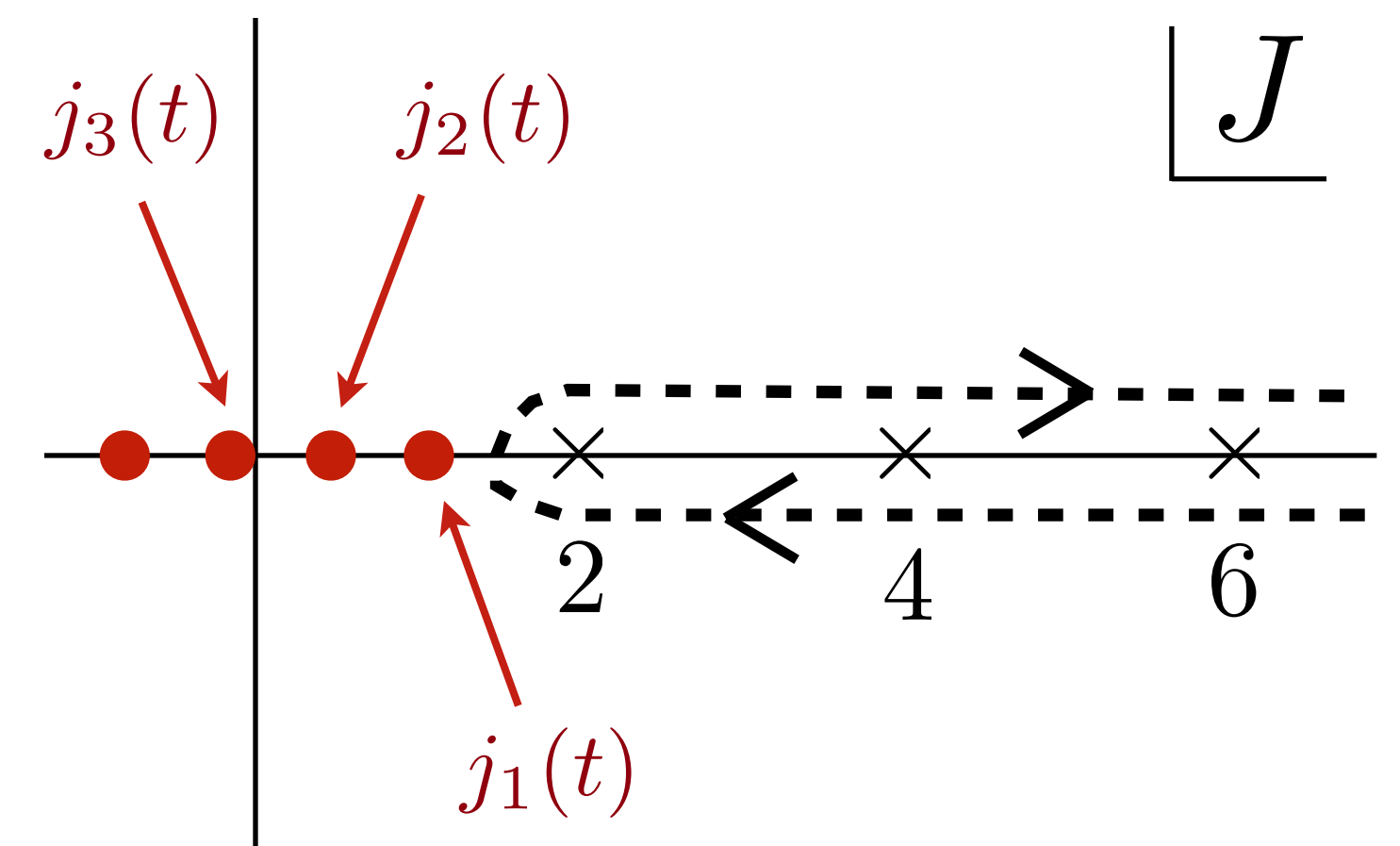
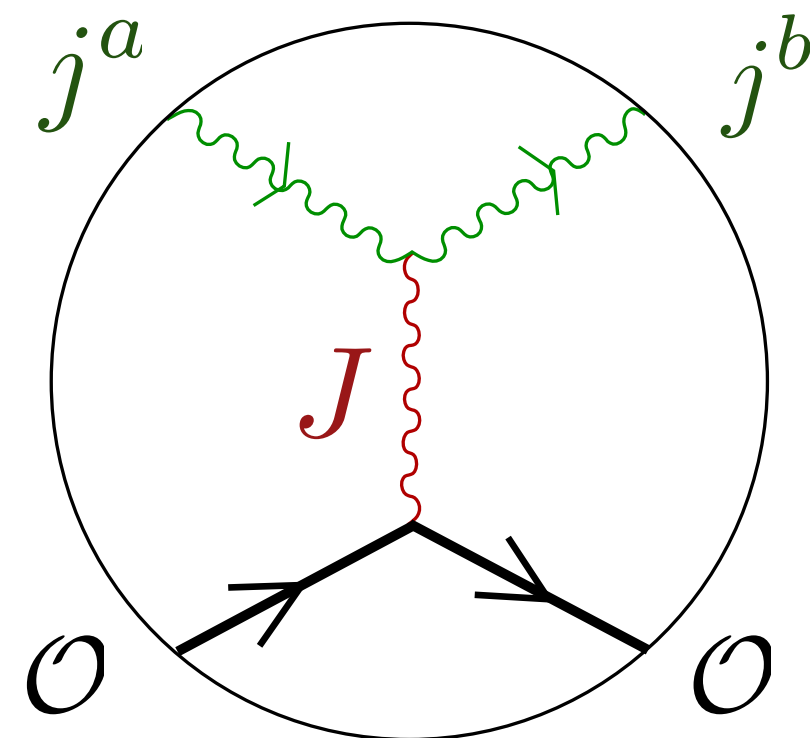
AdS Spin J field

- Spin J field equation of motion (graviton $\Delta = 4$, $J = 2$)

$$\left[(D^2 - 2\partial^b \phi D_b - \Delta(\Delta - 4)) g_{a_1 b} g_{a_2 c} + J R_{a_1 b a_2 c} \right] h^{bc}_{a_3 \dots a_J} = 0$$

$\Delta = \Delta(J)$ e.g. in
perturbation theory
 $\Delta = 2 + J + \gamma_J$

- Sum over spin J exchanges in 5D dual theory



- Analytically continue using diffusion approximation for $J < J_0$ ($\Delta \in \mathbb{C}$)

$$\Delta(\Delta - 4) = \frac{2}{\alpha'} (J - 2)$$

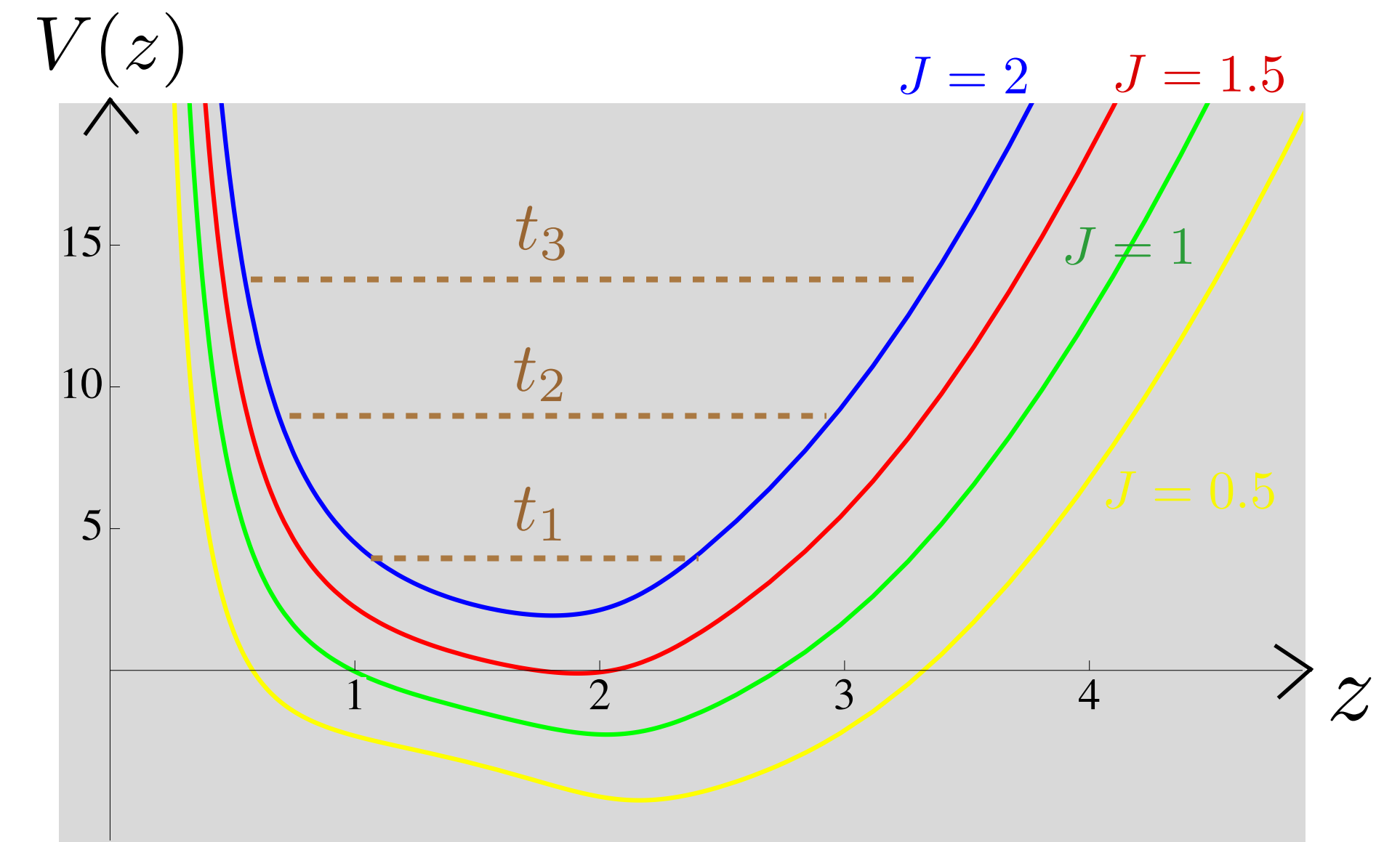
Pomeron Regge trajectories

Pomeron Regge trajectories

- Problem reduces to a J-dependent Schrodinger potential. Poles in J-plane

at $t = t_n(J) \Rightarrow J = j_n(t)$

$$T(s, t) \sim \sum_n \left(e^{-A(z) - A(z')} s \right)^{j_n(t)} \psi_n(z) \psi_n^*(z')$$



Pomeron Regge trajectories

- Problem reduces to a J-dependent Schrodinger potential. Poles in J-plane at

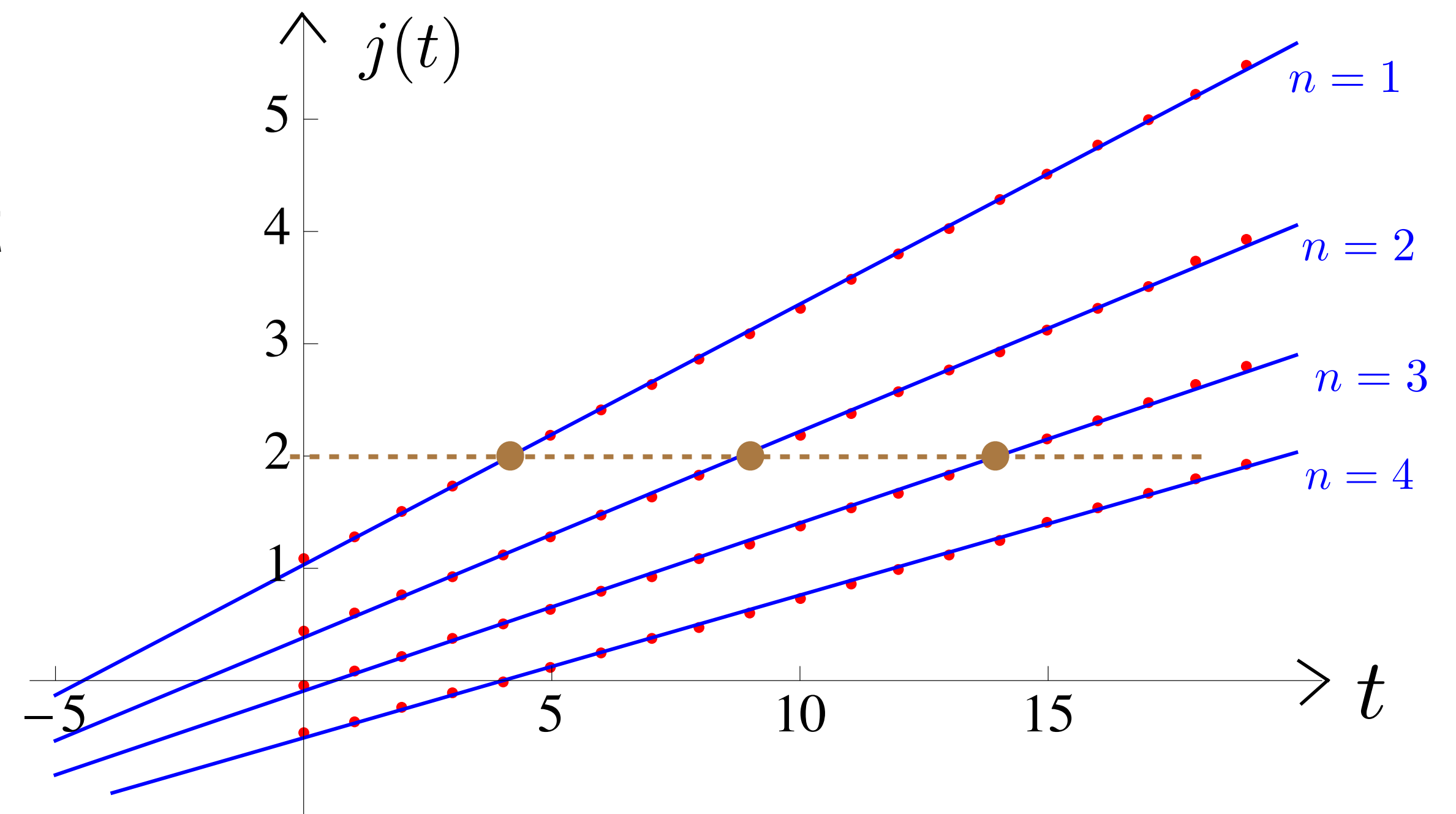
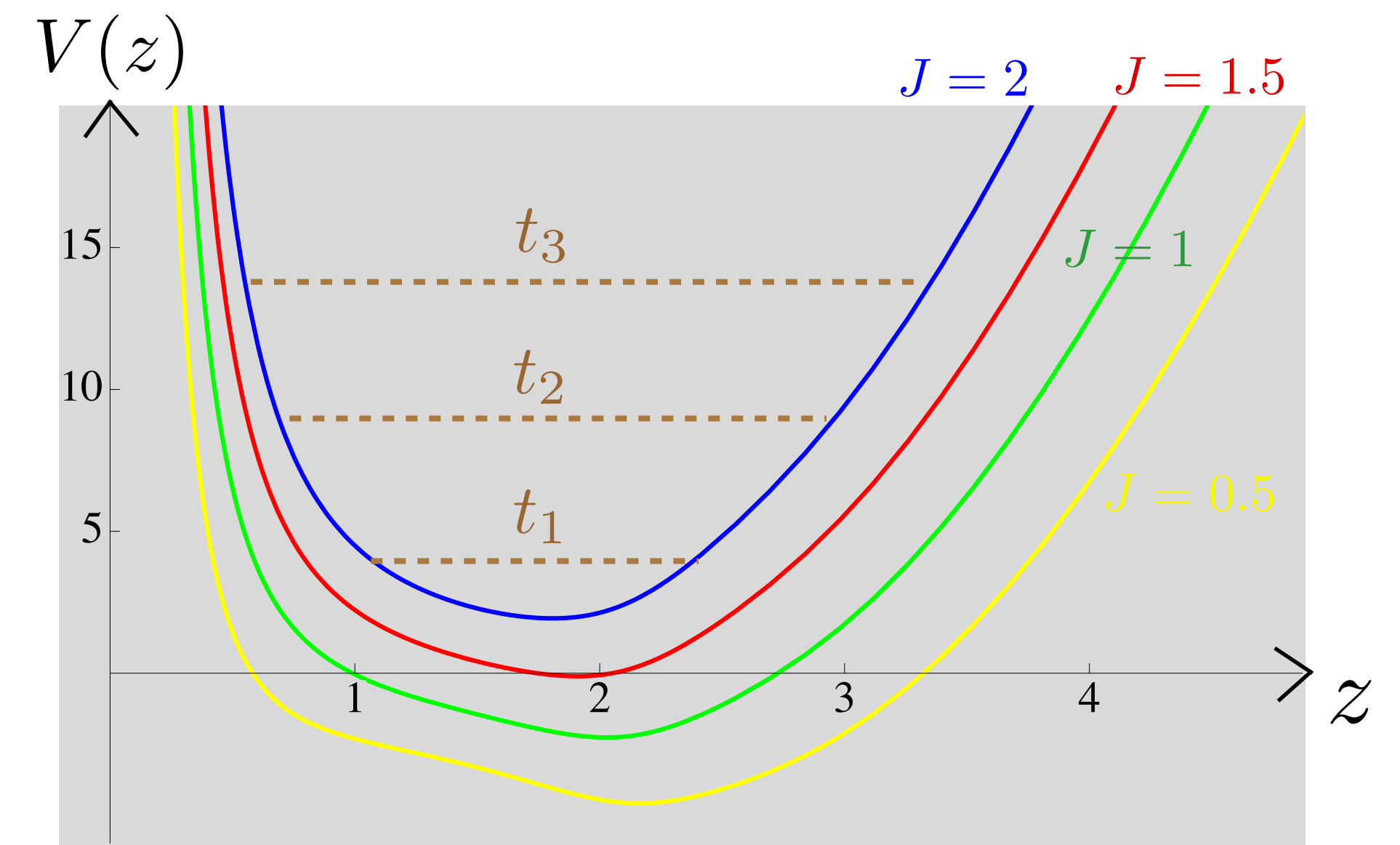
$$t = t_n(J) \Rightarrow J = j_n(t)$$

$$T(s, t) \sim \sum_n \left(e^{-A(z)-A(z')} s \right)^{j_n(t)} \psi_n(z) \psi_n^*(z')$$

- Obtained approximate linear Regge trajectory. One free parameter to fit soft pomeron intercept and slop. E.g.

$$j_1(t) \approx 1.08 + 0.22 t$$

Free parameter $l_s = 0.18$ consistent with value obtained for quark-antiquark potential



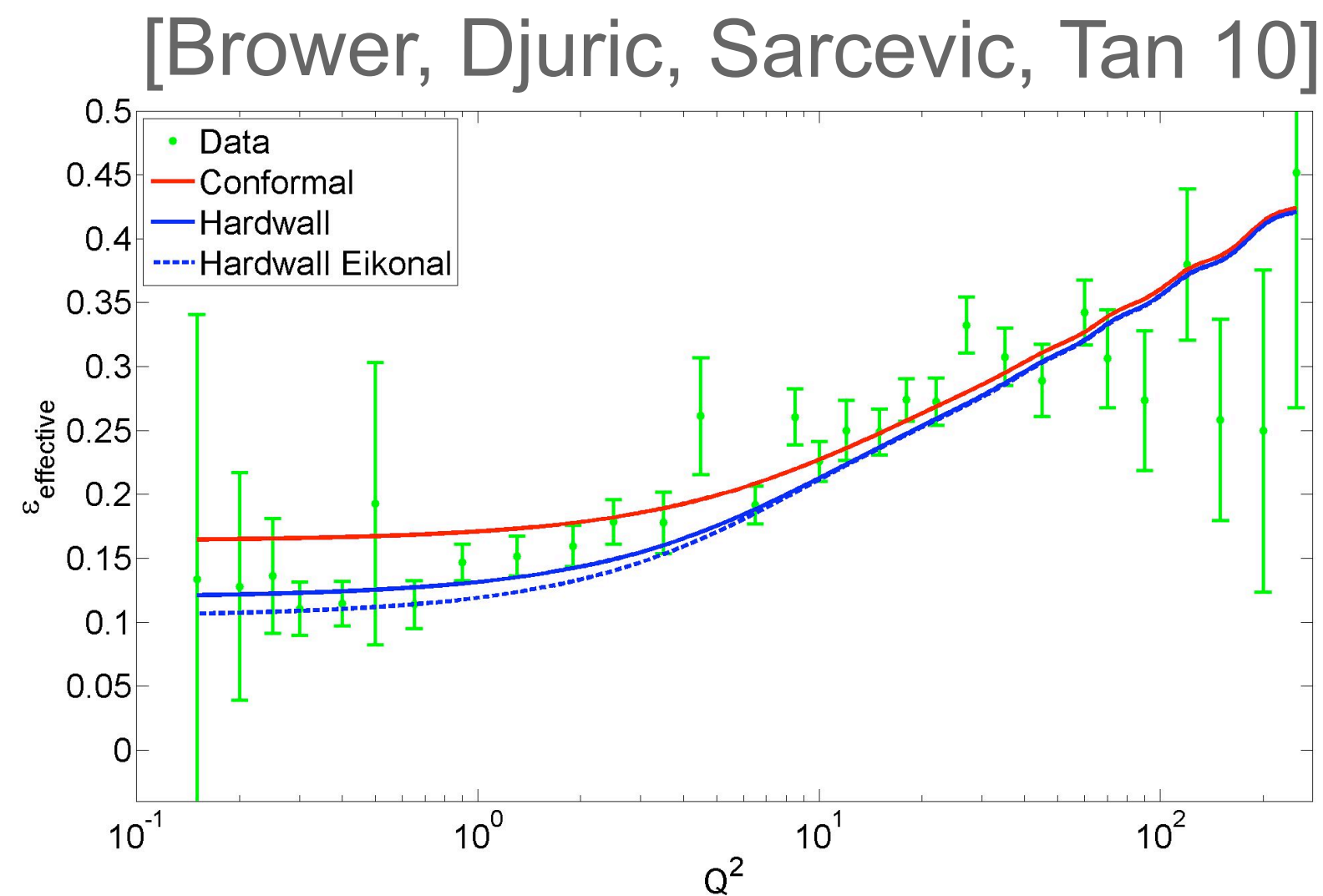
Concluding Remarks

Concluding Remarks

- AdS/QCD phenomenological model matches surprisingly well the intercept and slope of Donnachie-Landshoff pomeron

Concluding Remarks

- AdS/QCD phenomenological model matches surprisingly well the intercept and slope of Donnachie-Landshoff pomeron
- Connection with hard-pomeron. Understand running of effective exponent

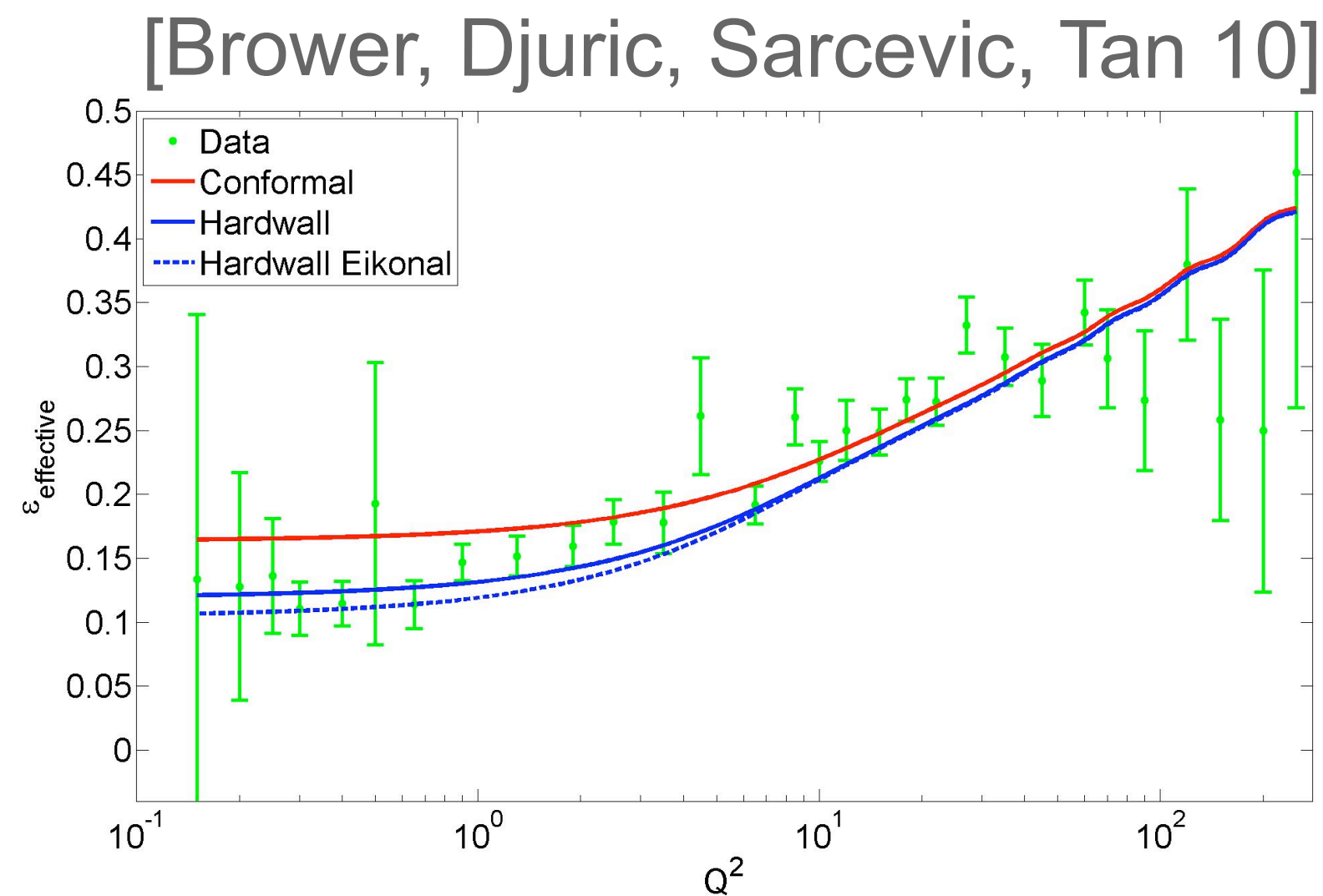


$$\sigma \sim \frac{1}{s} \text{Im } T(s, t = 0) \sim f(Q) \left(\frac{1}{x} \right)^{\epsilon_{eff}(Q)}$$

Add next sub-leading poles and study behavior with a varying probe scale Q

Concluding Remarks

- AdS/QCD phenomenological model matches surprisingly well the intercept and slope of Donnachie-Landshoff pomeron
- Connection with hard-pomeron. Understand running of effective exponent



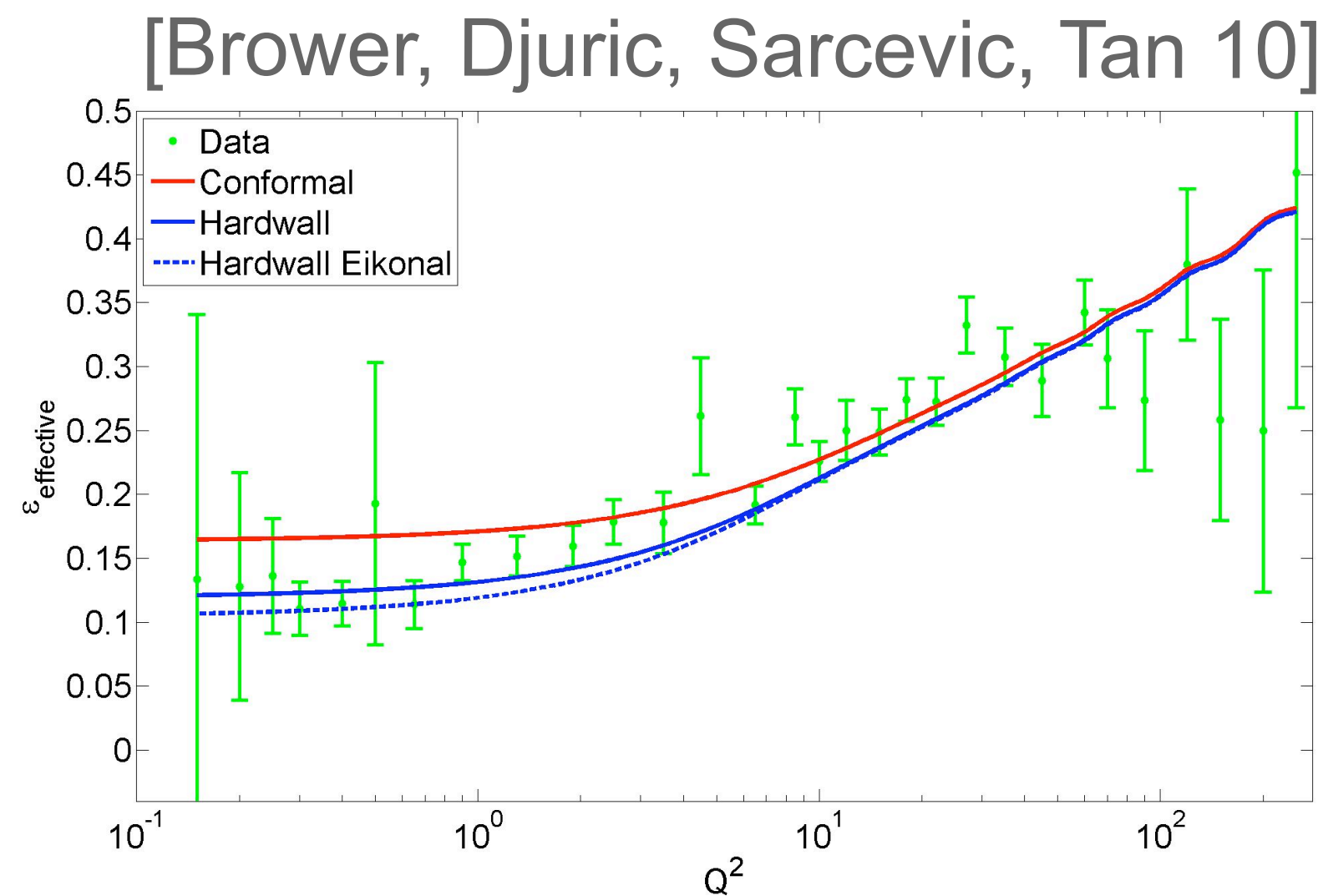
$$\sigma \sim \frac{1}{s} \text{Im } T(s, t = 0) \sim f(Q) \left(\frac{1}{x} \right)^{\epsilon_{eff}(Q)}$$

Add next sub-leading poles and study behavior with a varying probe scale Q

- Improve 5D equation for spin J field to reproduce anomalous dimensions

Concluding Remarks

- AdS/QCD phenomenological model matches surprisingly well the intercept and slope of Donnachie-Landshoff pomeron
- Connection with hard-pomeron. Understand running of effective exponent



$$\sigma \sim \frac{1}{s} \text{Im } T(s, t = 0) \sim f(Q) \left(\frac{1}{x} \right)^{\epsilon_{eff}(Q)}$$

Add next sub-leading poles and study behavior with a varying probe scale Q

- Improve 5D equation for spin J field to reproduce anomalous dimensions
- Make a careful analysis of data for DIS, DVSC and VMP. Goal:

unify soft and hard pomerons

THANK YOU