

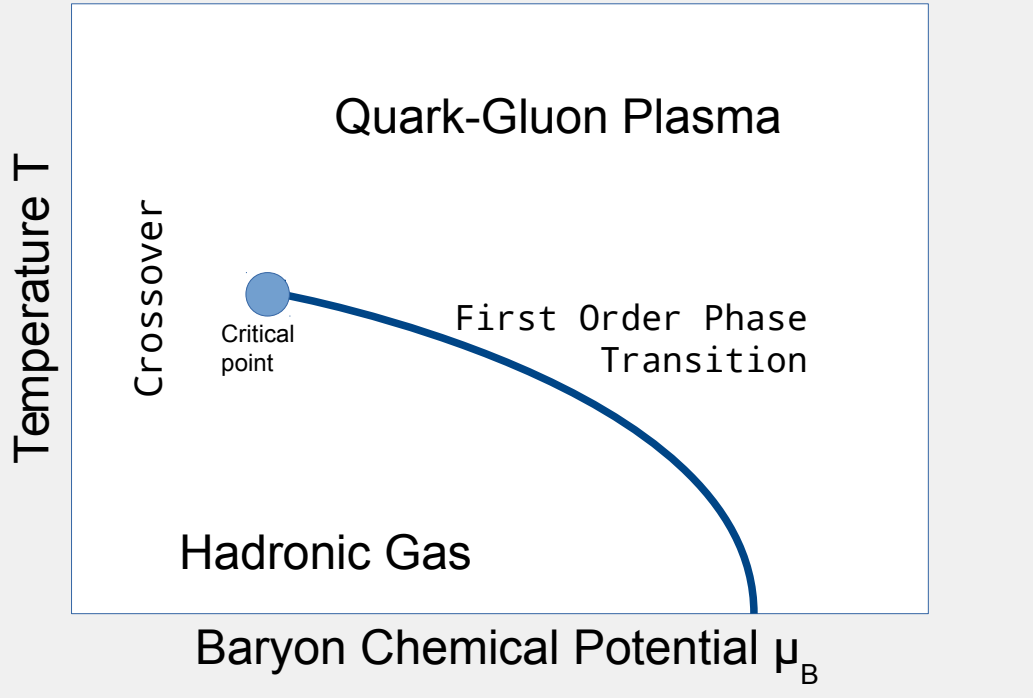
Strongly intensive fluctuations in nucleus-nucleus collisions at high energy in Monte Carlo model with string fusion

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Saint Petersburg State University

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Santiago de Compostela

QCD phase diagram and search for critical point

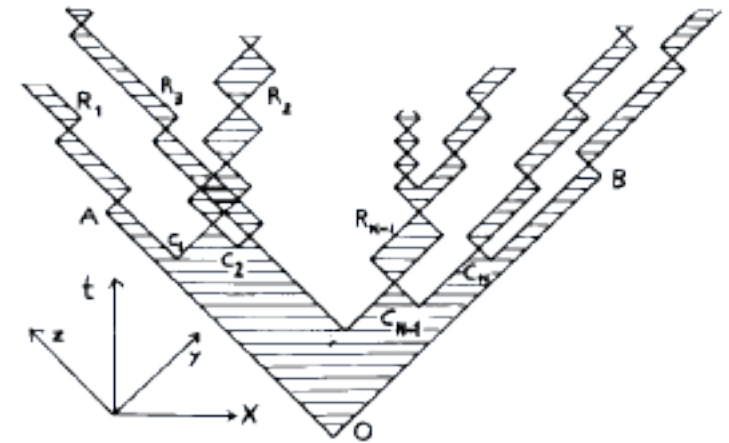
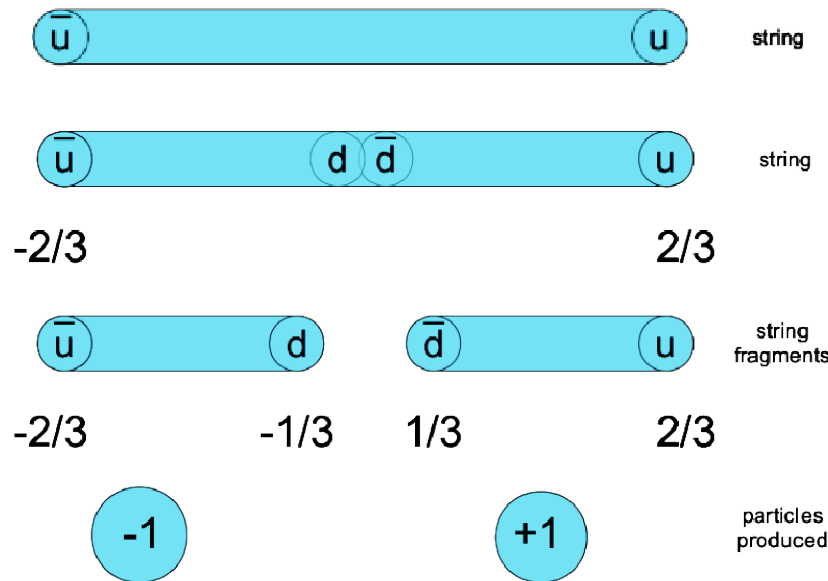


$$\frac{\bar{p}}{p} = \frac{e^{-(E+\mu_B)/T}}{e^{-(E-\mu_B)/T}} = e^{-(2\mu_B)/T}$$

HADES, GSI	2.3 – 2.7 GeV	p+p, Au+Au, Ar+KCl, C+C
NA61, SPS, CERN	6.3 - 17.3 GeV	p+p, Be+Be, p+Pb, Ar+Ca, Xe+La, Pb+Pb, ...
CBM, FAIR, GSI	2.7 - 8.3 GeV	p, Ca, Au
RHIC BES	5 - 200 GeV	Au+Au
NICA, JINR	3 - 11 GeV	from p to Au

String models

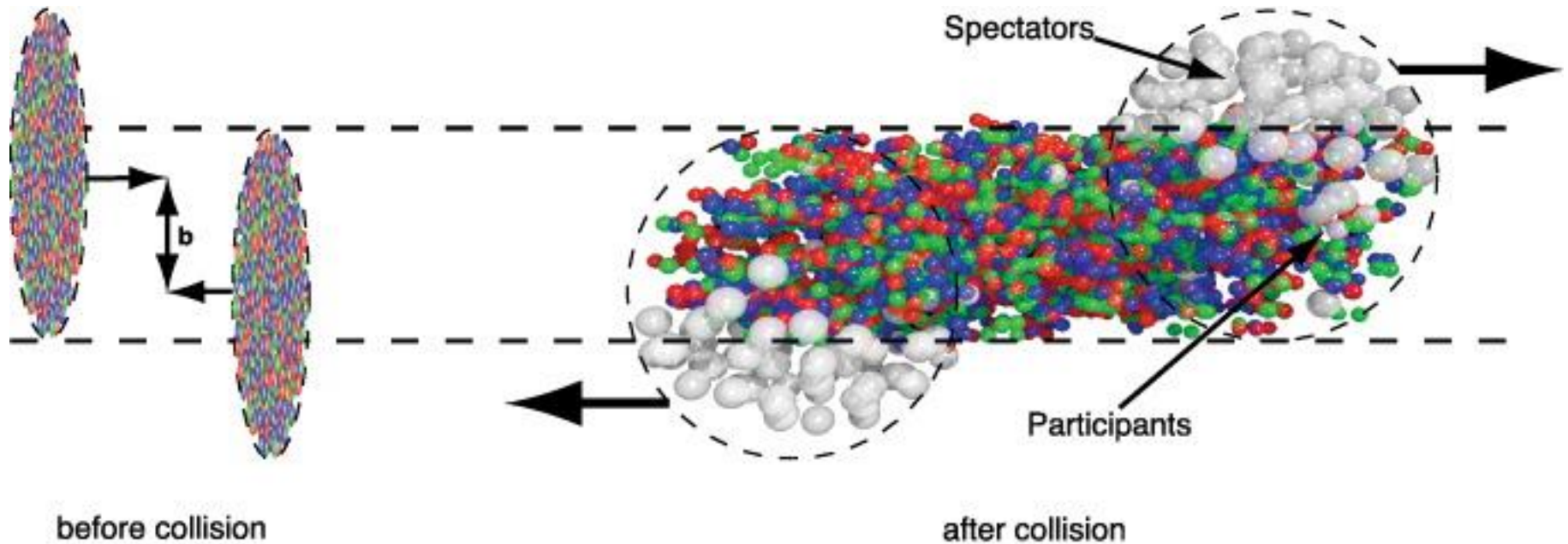
- The soft QCD processes is not described by usual perturbation theory
- The model of quark-gluon strings, stretched between projectile and target partons
 - semiphenomenological approach to the multiparticle production



X. Artru and G. Mennessier, Nucl Phys B 70 (1973) "String Model and Multiproduction",

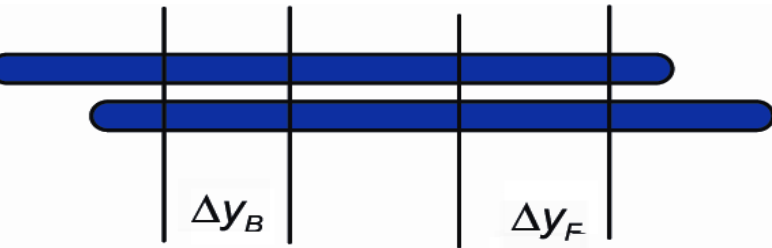
- Correlations play crucial role:
 - causality requires appearance of long-range correlations – if they exist – at the very early stages between particles detected in separated rapidity intervals

Centrality



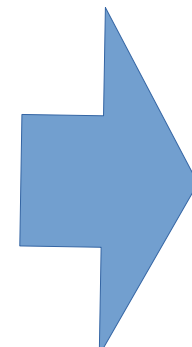
- Nucleon-participants N_{part} - nucleons collided at least once
- Nucleon-spectators N_{spect} - nucleons, which didn't interact
- Number of nucleon-nucleon collisions N_{coll}
- Multiplicity of charge particles N_{ch}

Long-range correlations and fluctuations



Select 2 variables in windows
(or in one window):

- n – number of charged particles in the window
- $p_t = \frac{1}{n} \sum_{i=1}^n p_{ti}$ – mean (in the event!) transverse momentum of charged particles in the given window
- $P_T = \sum_{i=1}^n p_{ti}$ – total transverse momentum of charged particles in the given window



Several types of correlations and fluctuations:

n-n , pt-n, pt-pt

long-range/one window, etc

correlation coefficient

$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

scaled variance

$$\omega = \frac{\langle n^2 \rangle - \langle n \rangle^2}{\langle n \rangle}$$

String fusion

$$Q^2(n) = \left(\sum_{i=1}^n \vec{Q}_i(1) \right)^2 = \sum_{i=1}^n Q_i^2(1) + \sum_{i \neq j} \vec{Q}_i(1) \cdot \vec{Q}_j(1)$$

$$\langle Q^2(n) \rangle = nQ^2(1)$$

overlaps

SFM

$$C = \{S_1, S_2, \dots\}$$

S_k – area covered k-times



$$\langle \mu \rangle_k = \mu_0 \sqrt{k} \frac{S_k}{\sigma_0}$$

$$\langle p_t^2 \rangle_k = p_0^2 \sqrt{k}$$

$$\langle p_t \rangle_k = p_0 \sqrt[4]{k}$$

S_k – area, where k strings are overlapping, σ_0 single string transverse area, μ_0 and p_0 – mean multiplicity and transverse momentum from one string

String fusion mechanism predicts:

- decrease of multiplicity
- increase of p_T
- growth of p_T with multiplicity in pp, pA and AA collisions
- growth of strange particle yields
- results are in a good agreement with the experiment

M. A. Braun, C. Pajares, Nucl. Phys. B 390 (1993) 542.

M. A. Braun, R. S. Kolevatov, C. Pajares, V. V. Vechernin, Eur. Phys. J. C 32 (2004) 535.

N.S. Amelin, N. Armesto, C. Pajares, D. Sousa, Eur.Phys.J.C22:149-163 (2001), arXiv:hep-ph/0103060

G. Ferreiro and C Pajares J. Phys. G: Nucl. Part. Phys. 23 1961 (1997)

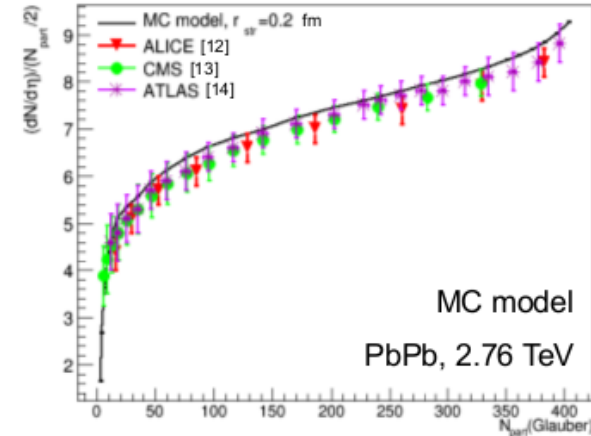
Monte-Carlo model

- Partonic dipole-based picture of nucleons interaction.
- Energy and angular momentum conservation in the initial state of a nucleon.
- The probability of dipoles are defined by their transverse coordinates [7-8]:

$$f = \frac{\alpha_s^2}{2} \ln^2 \frac{|\vec{r}_1 - \vec{r}'_1| |\vec{r}_2 - \vec{r}'_2|}{|\vec{r}_1 - \vec{r}'_2| |\vec{r}_2 - \vec{r}'_1|}$$

Multiplicity and transverse momentum are obtained in the approach of colour strings, stretched between projectile and target partons

- The interaction of strings is realized in the accordance with the **string fusion** model
- Multiplicity from one string is distributed according to Poisson



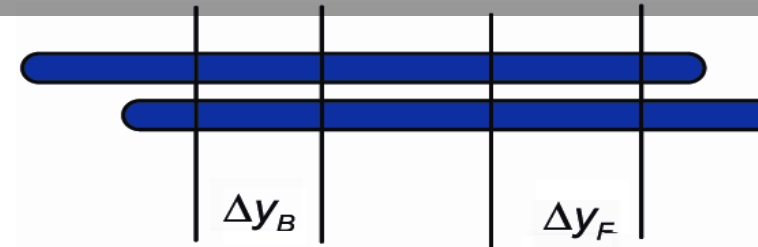
V. N. Kovalenko.. Phys. Atom. Nucl. 76, 1189 (2013), arXiv:1211.6209 [hep-ph]

V. Kovalenko, V. Vechernin., PoS (Baldin ISHEPP XXI) 077, arXiv:1212.2590 [nucl-th], 2012

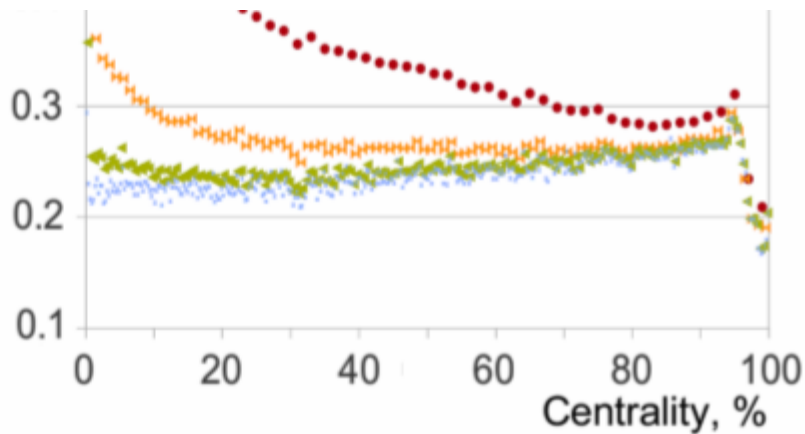
Long-range correlations and fluctuations

n-n
correlation
coefficient

$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

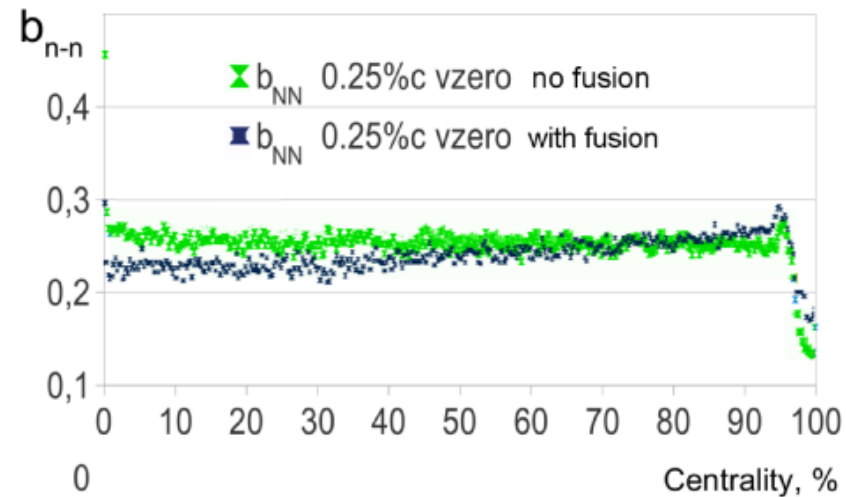


b_{n-n}



- ✕ 0.25% c vzero
- ◀ 0.5% c vzero
- ✎ 1% c vzero
- 2% c vzero

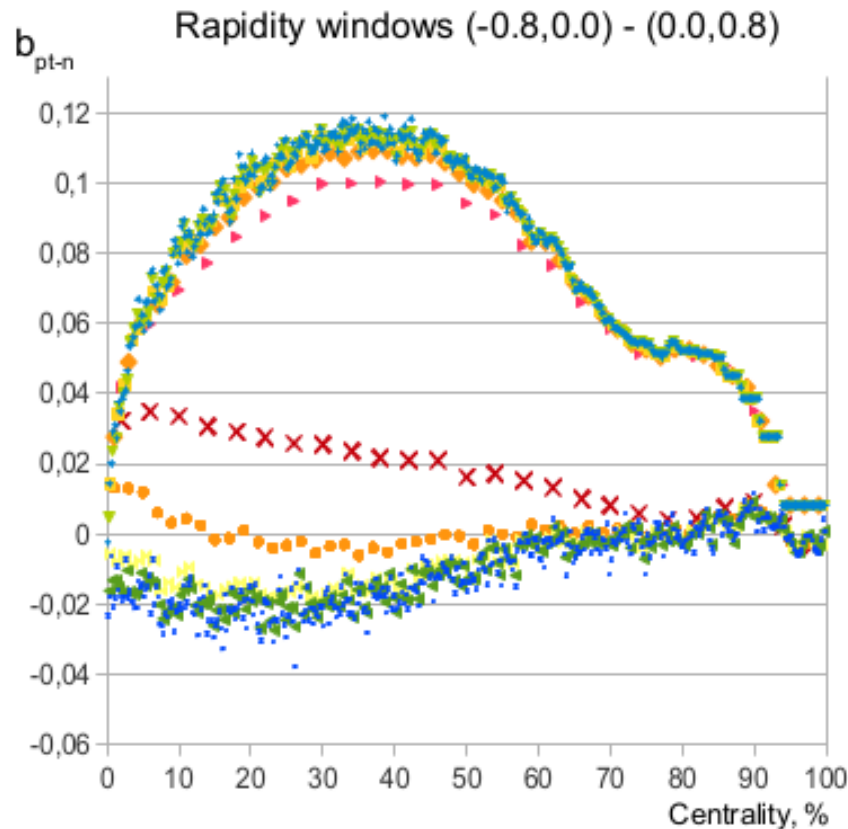
Pb-Pb, 2.76 TeV



Long-range correlations and fluctuations

pt-n
correlation
coefficient

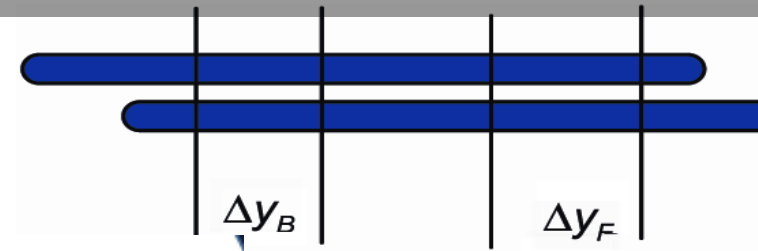
$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$



PbPb,
2.76 TeV,
MC model

- + b_{ptN} 0.25% c Npart
- ▼ b_{ptN} 0.5% c Npart
- ▲ b_{ptN} 1.5% c Npart
- b_{ptN} 1% c Npart
- ◆ b_{ptN} 2% c Npart
- ▶ b_{ptN} 4% c Npart
- ⊠ b_{ptN} 0.25% c vzero
- ◀ b_{ptN} 0.5% c vzero
- ◀ b_{ptN} 1% c vzero
- b_{ptN} 2% c vzero
- × b_{ptN} 4% c vzero

Pb-Pb, 2.76 TeV

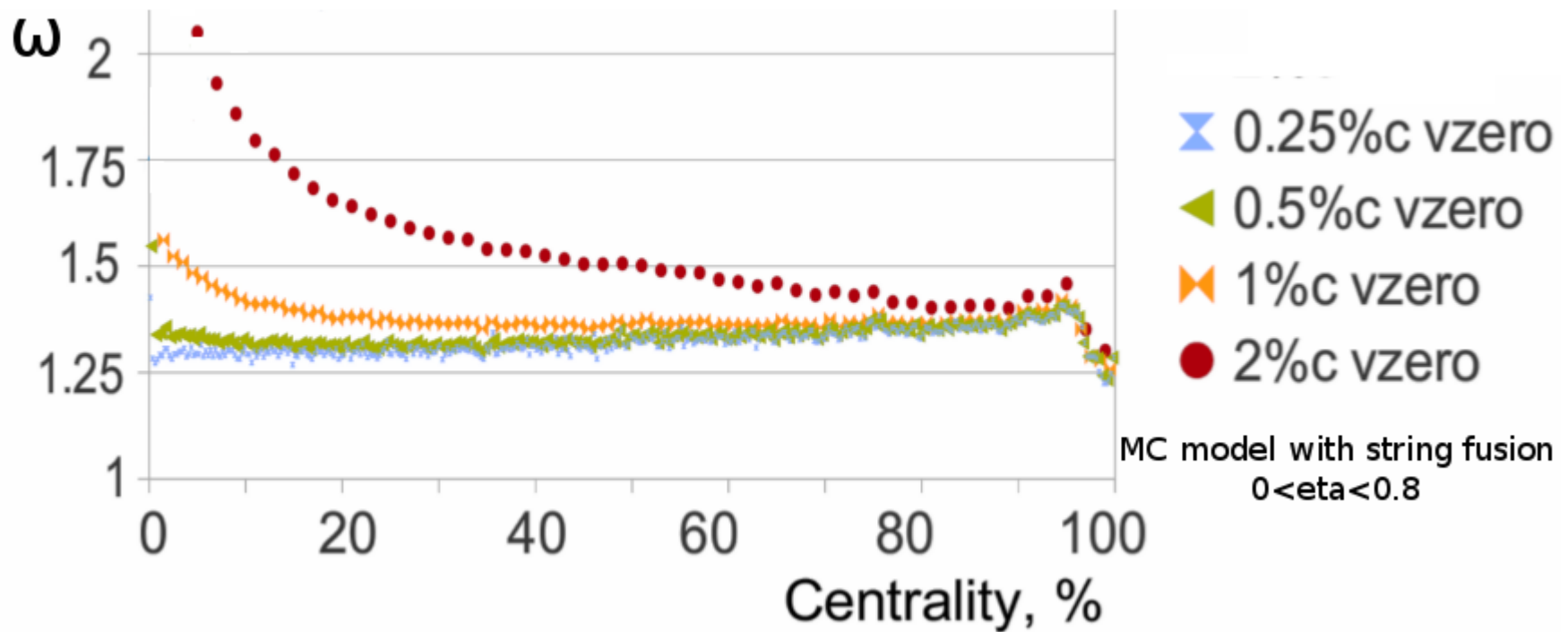


Long-range correlations and fluctuations

scaled variance $\omega = \frac{\langle n^2 \rangle - \langle n \rangle^2}{\langle n \rangle}$



Pb-Pb, 2.76 TeV



V. Kovalenko, V. Vechernin. EPJ Web of Conferences 66, 04015 (2014), arXiv:1308.6618 [nucl-th]

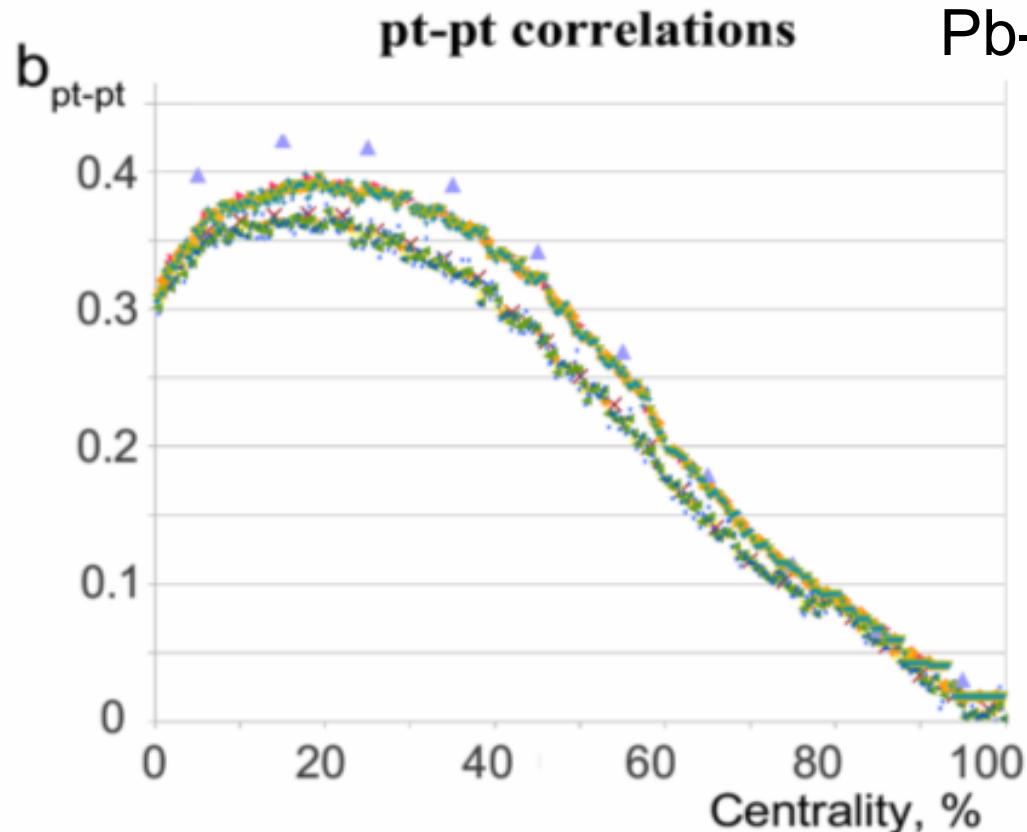
mean pt-pt correlations

correlation coefficient

$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

$$B, F, \rightarrow p_t = \frac{1}{n} \sum_{i=1}^n p_{ti}$$

- + 0.25% c N_{part}
- ▼ 0.5% c N_{part}
- ▲ 1.5% c N_{part}
- 1% c N_{part}
- ◆ 2% c N_{part}
- ▶ 4% c N_{part}
- ▲ 10% c N_{part}
- ✕ 0.25% c v_{zero}
- ◀ 0.5% c v_{zero}
- ✦ 1% c v_{zero}
- 2% c v_{zero}
- ✕ 4% c v_{zero}



V. Kovalenko, V. Vechernin. EPJ Web of Conferences 66, 04015 (2014), arXiv:1308.6618 [nucl-th]

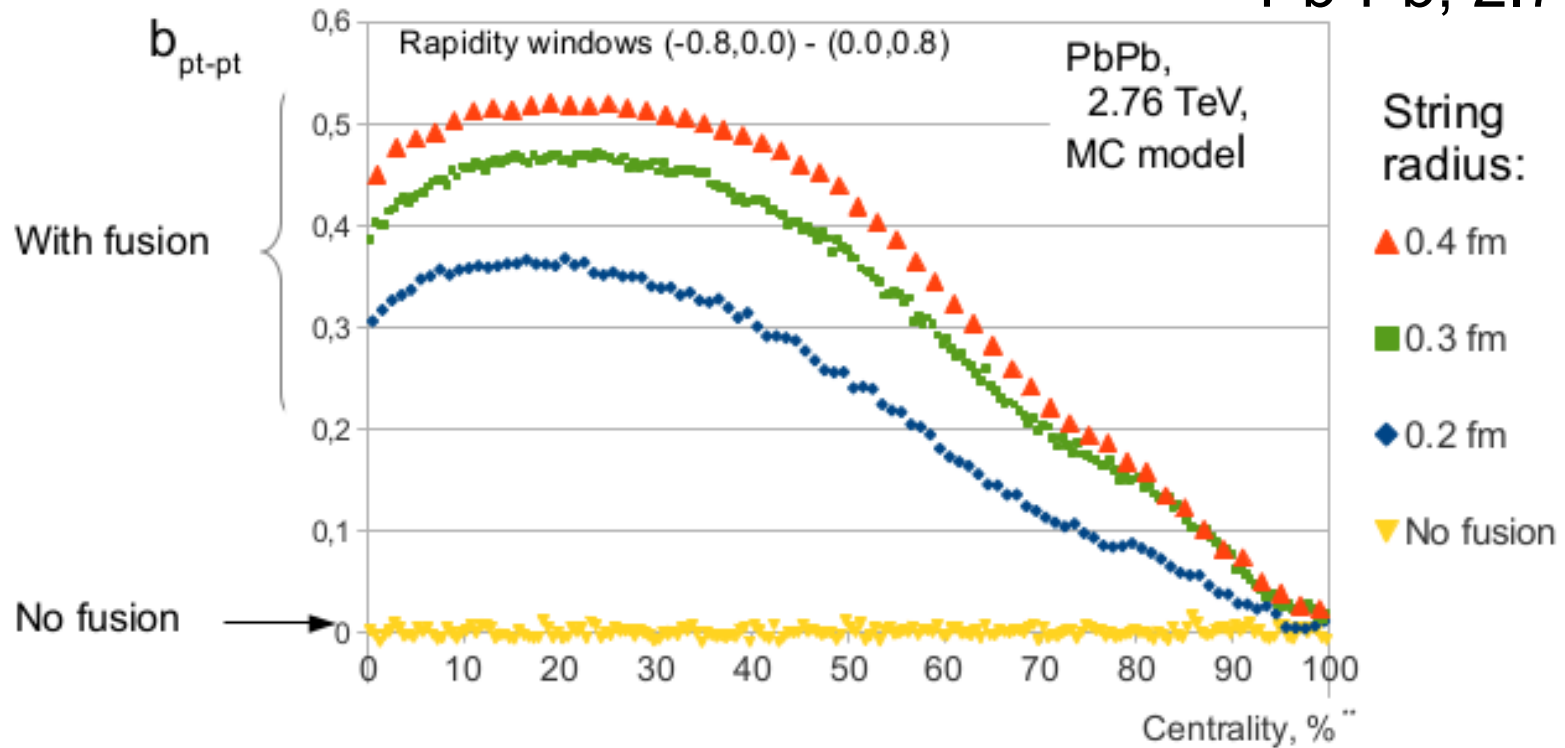
mean pt-pt correlations

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$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

$$B, F, \rightarrow p_t = \frac{1}{n} \sum_{i=1}^n p_{ti}$$

Pb-Pb, 2.76 TeV



V. Kovalenko, V. Vechernin. EPJ Web of Conferences 66, 04015 (2014), arXiv:1308.6618 [nucl-th]

Strongly intensive variables

Let A and B be two *extensive* quantities. Combining $\langle A \rangle$, $\langle A^2 \rangle$, $\langle B \rangle$, $\langle B^2 \rangle$, $\langle AB \rangle$ one can get:

$$\Delta[A, B] = \frac{1}{C_\Delta} (\langle B \rangle_\omega [A] - \langle A \rangle_\omega [B])$$

$$\Sigma[A, B] = \frac{1}{C_\Sigma} (\langle B \rangle_\omega [A] + \langle A \rangle_\omega [B] - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle))$$

M. Gorenstein, M. Gazdzicki, arXiv:1101.4865 [nucl-th]

Independent particle model.

In IPM we have for A and B for an event with n particles:

$$A = \alpha_1 + \cdots + \alpha_n$$

$$B = \beta_1 + \cdots + \beta_n$$

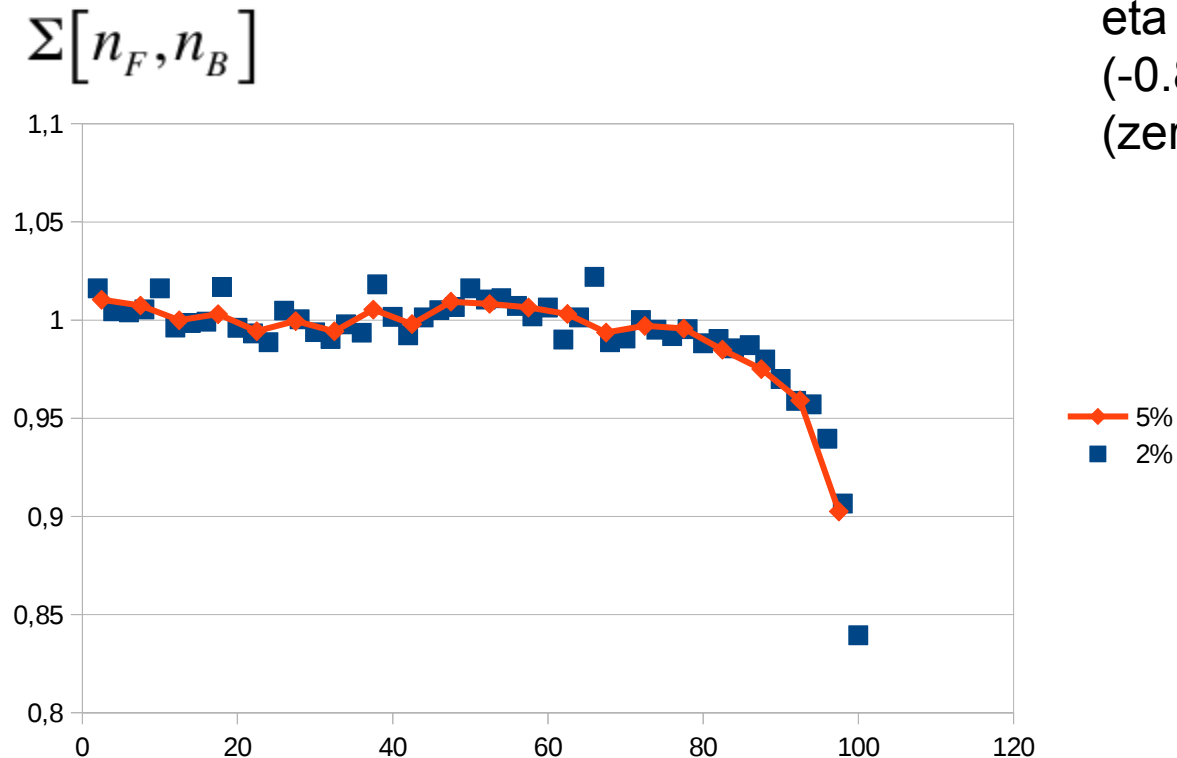
where for α_i and β_i we may introduce moments: $\bar{\alpha}$, $\overline{\alpha^2}$, $\bar{\beta}$, $\overline{\beta^2}$, $\overline{\alpha\beta}$
Particles are "independent" and "identical":

$$P(\alpha_1, \beta_1, \cdots, \alpha_n, \beta_n) = P_{mult}(n) * P_{part}(\alpha_1, \beta_1) * \cdots * P_{part}(\alpha_n, \beta_n)$$

Long-range SI multiplicity fluctuations

$$\Sigma[n_F, n_B] = \frac{\langle n_B \rangle \omega[n_F] + \langle n_F \rangle \omega[n_B] - 2(\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle)}{\langle n_B \rangle + \langle n_F \rangle}$$

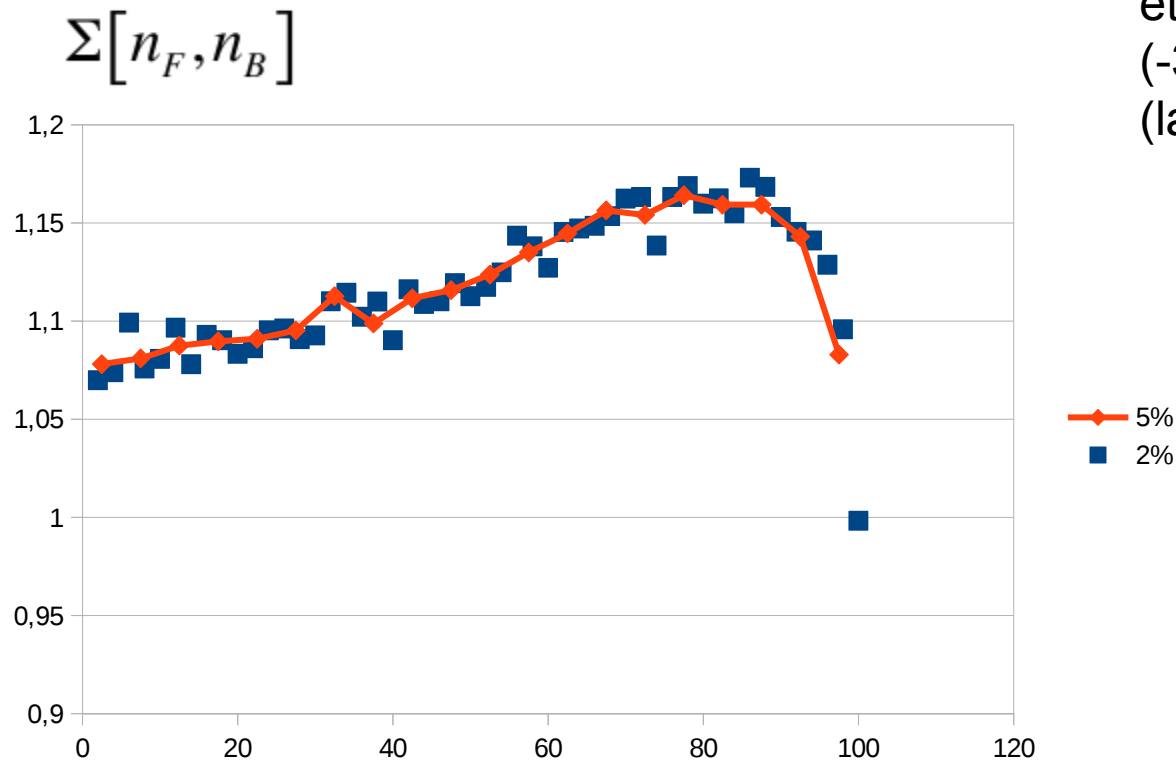
Pb-Pb, 2.76 TeV
eta windows:
(-0.8, 0), (0, 0.8)
(zero gap)



Long-range SI multiplicity fluctuations

$$\Sigma[n_F, n_B] = \frac{\langle n_B \rangle \omega[n_F] + \langle n_F \rangle \omega[n_B] - 2(\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle)}{\langle n_B \rangle + \langle n_F \rangle}$$

Pb-Pb, 2.76 TeV
eta windows:
(-3.8, -3), (3, 3.8)
(large gap)

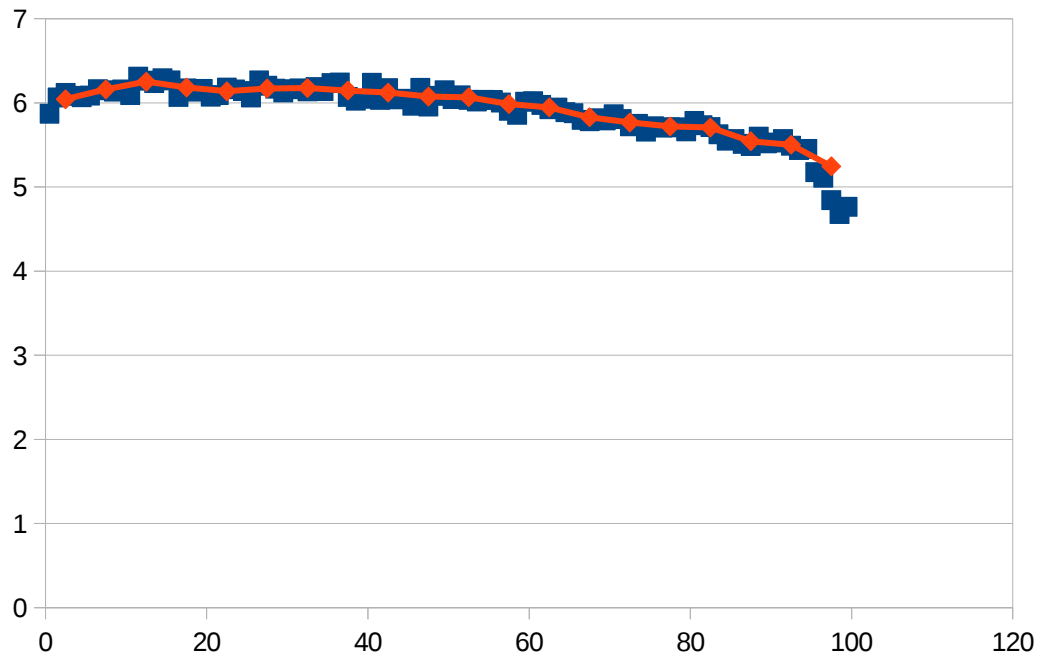


SI Pt-multiplicity fluctuations

$$\Sigma [A, B] = \frac{1}{C_{\Sigma}} (\langle B \rangle \omega [A] + \langle A \rangle \omega [B] - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle))$$

$$A = P_T = \sum_{i=1}^N p_{T_i}, \quad B = N \quad \longrightarrow \quad C_{\Delta} = C_{\Sigma} = \langle N \rangle \omega [p_T]$$

$$\Sigma [P_T \ N]$$

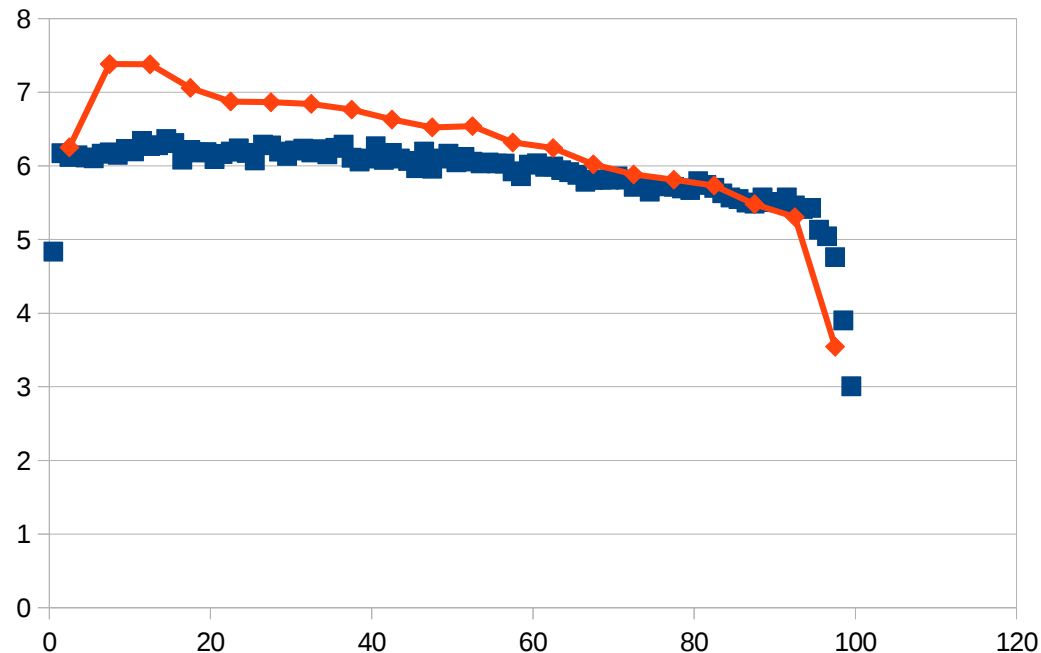


SI Pt-multiplicity fluctuations

$$\Delta[A, B] = \frac{1}{C_{\Delta}} (\langle B \rangle_{\omega} [A] - \langle A \rangle_{\omega} [B])$$

$$A = P_T = \sum_{i=1}^N p_{T_i}, \quad B = N \quad \longrightarrow \quad C_{\Delta} = C_{\Sigma} = \langle N \rangle_{\omega} [p_T]$$

$\Delta[P_T, N]$

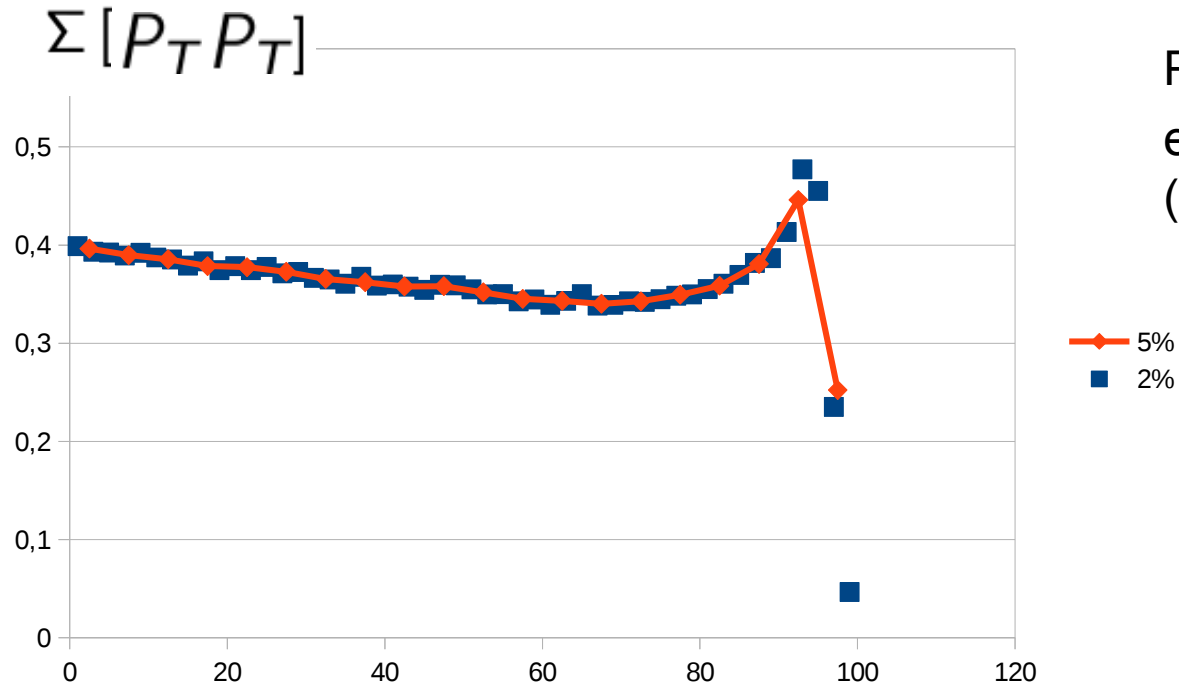


Pb-Pb, 2.76 TeV
eta window:
(-0.8, 0.8)

5%
2%

SI Pt-multiplicity fluctuations

$$\Sigma [A, B] = \frac{1}{C_{\Sigma}} (\langle B \rangle_{\omega} [A] + \langle A \rangle_{\omega} [B] - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle))$$



Conclusions and outlook

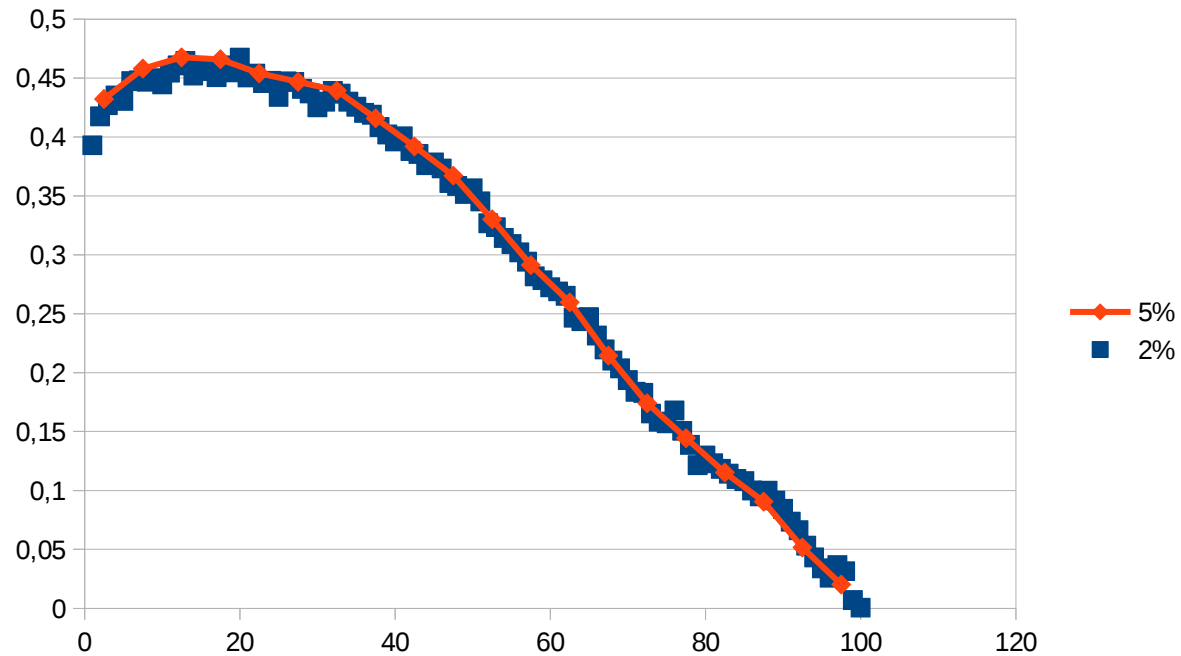
- Centrality determination and a width of centrality class influences the value of fluctuational and correlational observables
- Introducing of intensive variables helps to overcome the dependence on centrality determination issues
- With SI values, n-n and pt-n long-range correlation variables obtained, which don't depend on the centrality class width; another pt-pt long-range correlational observable introduced
- The set of the observables could be extended using particle species, charge and azimuthal angle – large area for strongly intensive quantities studies

Thank you!

Backup

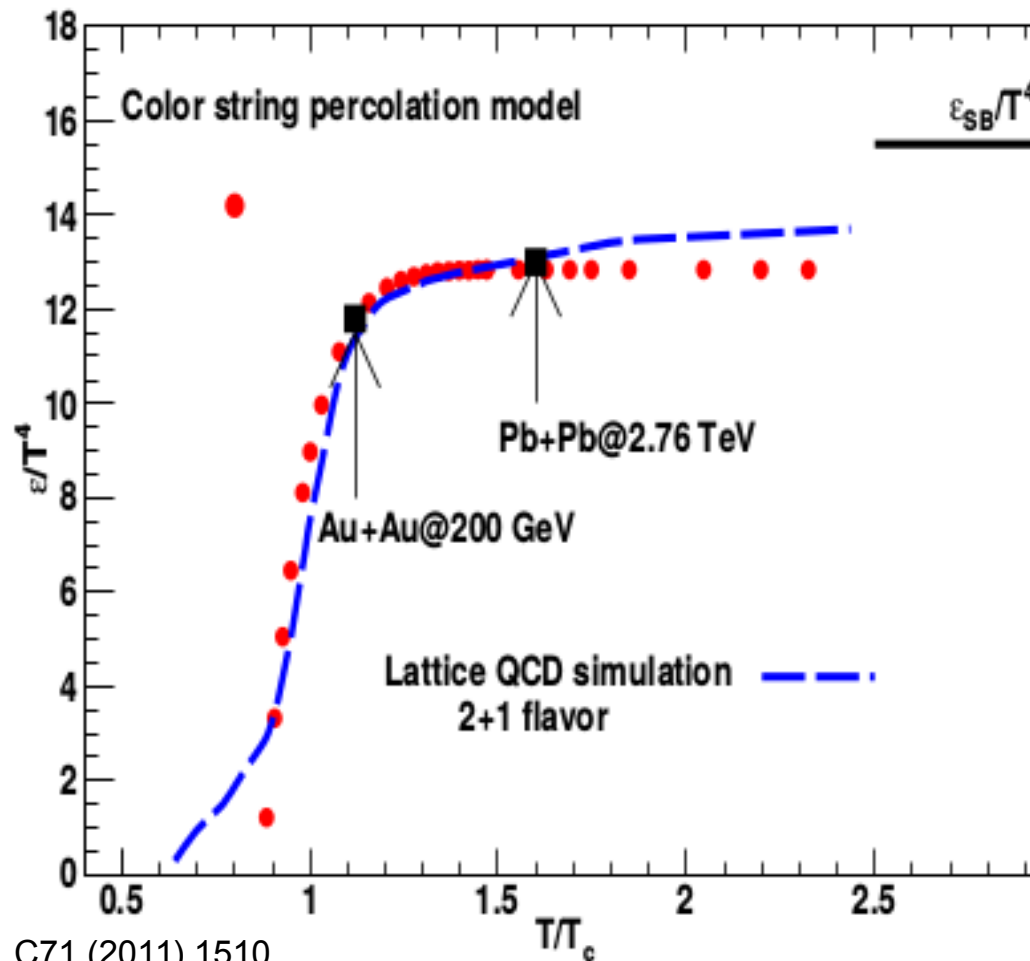
Mean pt-pt correlations

$$b = \frac{\langle pt_F pt_B \rangle - \langle pt_F \rangle \langle pt_B \rangle}{\langle pt_F^2 \rangle - \langle pt_F \rangle^2}$$



String fusion

In the recent papers it was shown that the equation of state of QGP (ϵ/T^4 as a function of T) at zero chemical potential, obtained in the colour string percolation model is in a good agreement with the lattice results.



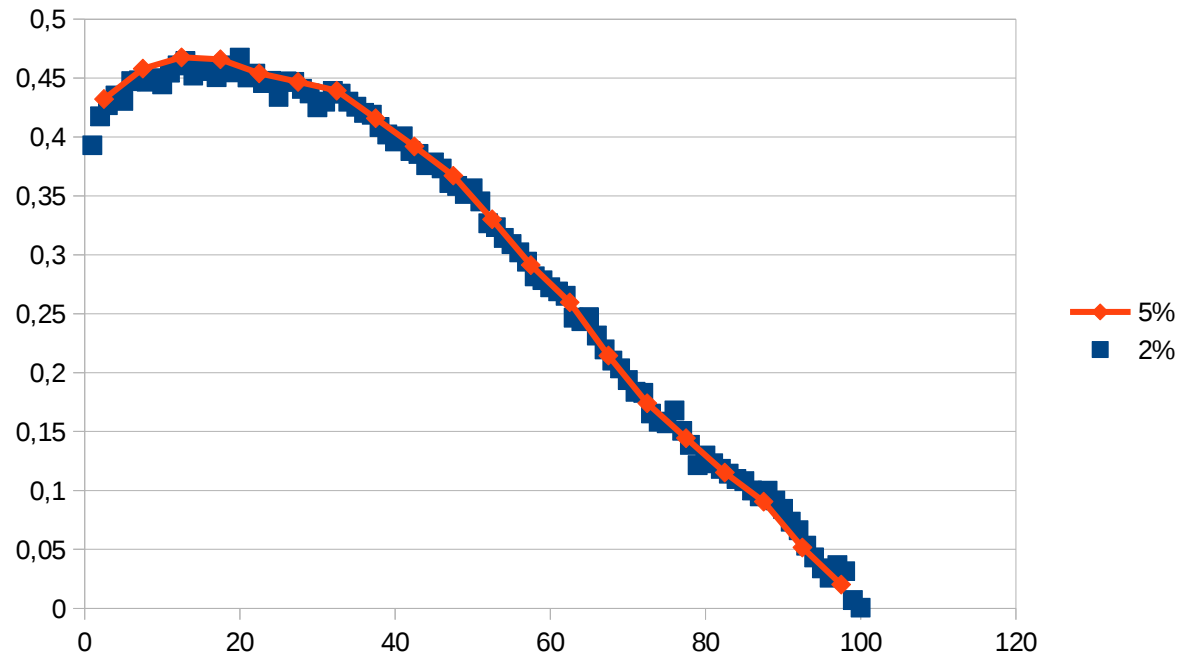
R.P. Scharenberg, B.K. Srivastava, A.S. Hirsch Eur.Phys.J. C71 (2011) 1510

J. Dias de Deus, C. Pajares, Phys.Lett. B642 (2006) 455-458

Brijesh K Srivastava, EP J Web of Conferences 70, 00032 (2014)

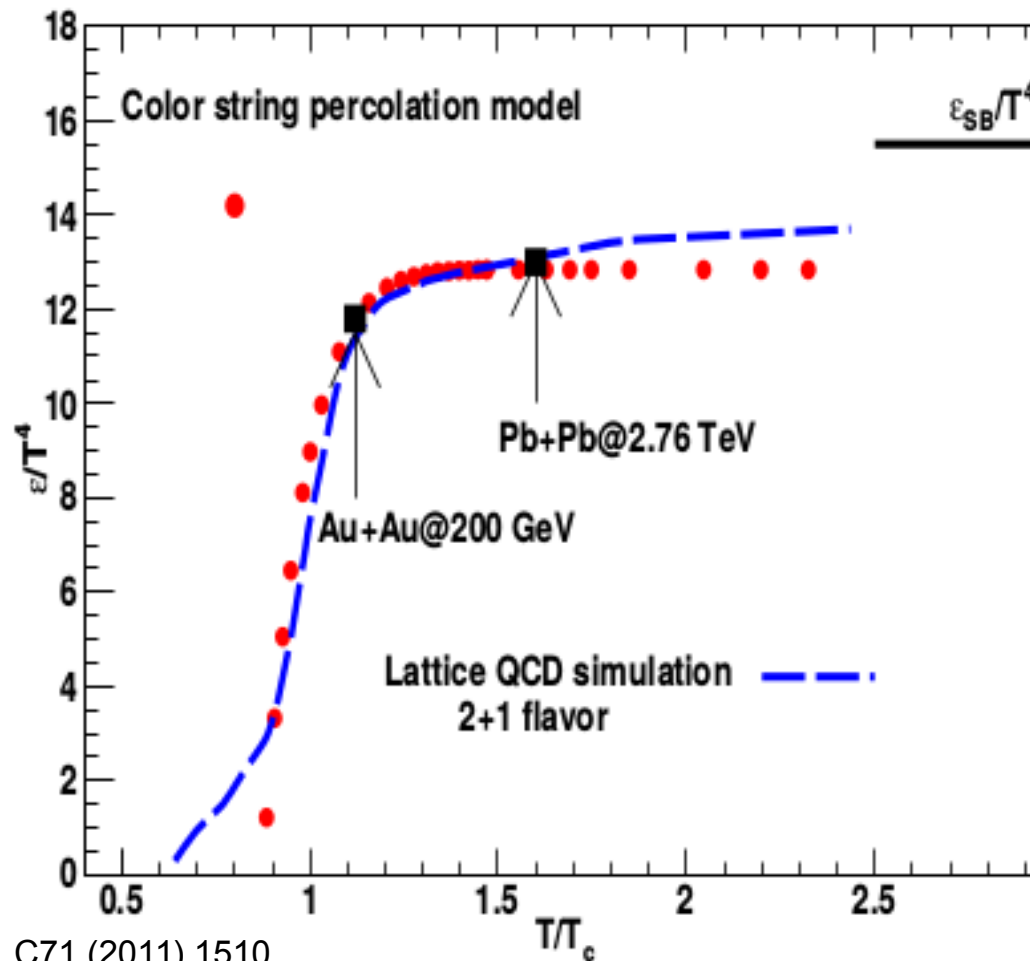
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$$b = \frac{\langle pt_F pt_B \rangle - \langle pt_F \rangle \langle pt_B \rangle}{\langle pt_F^2 \rangle - \langle pt_F \rangle^2}$$



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