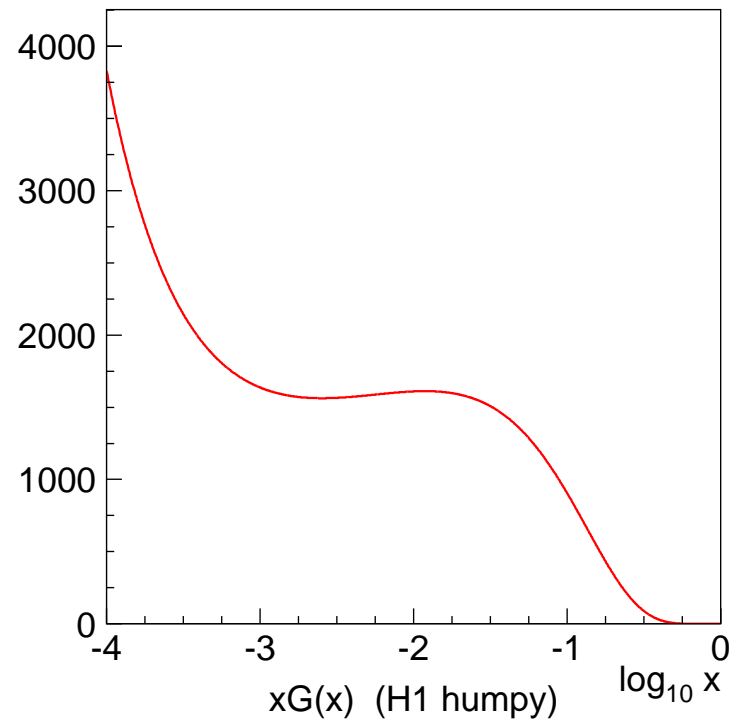
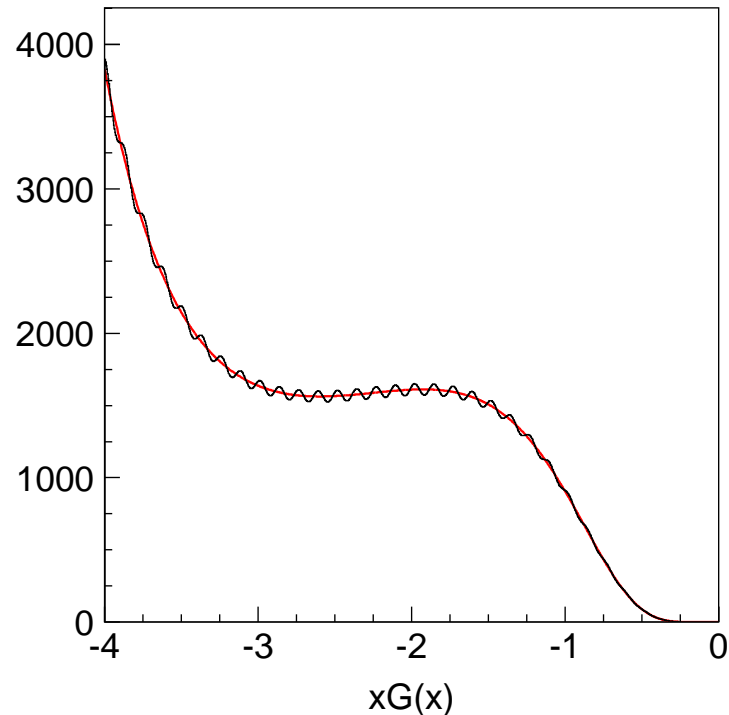


# Physical Constraints on PDFs



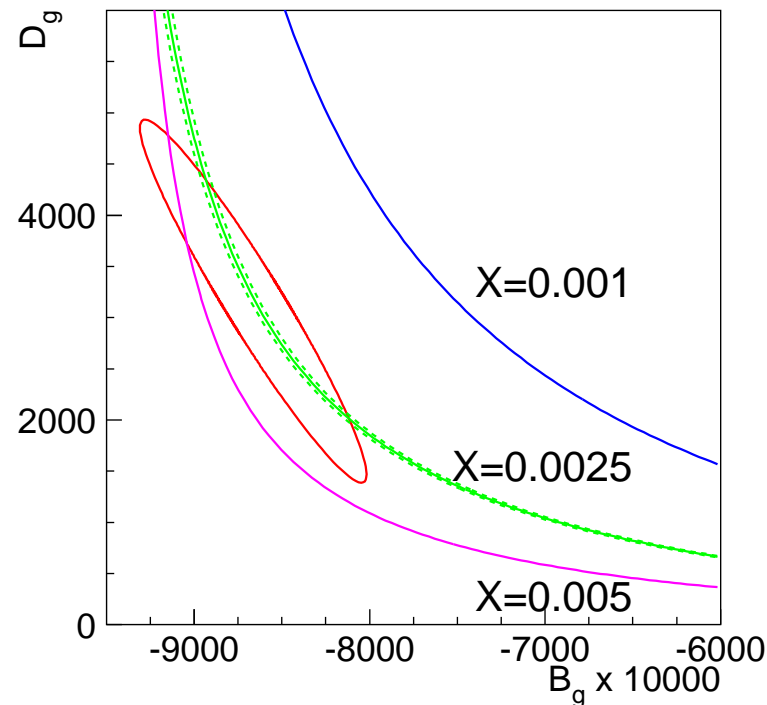
- H1 parameterization, fit to H1 data
- Humpty solution:  $xG(x) = Ax^B(1-x)^C(1+Dx)$ ,  
 $D = 3160 \pm 1773$ .
- Large correlation btw  $D$  and  $B$ :  $\rho = -95.5\%$ .

# Hypothetical unconstraint solution



- Blindly increasing number of parameters may lead to solutions with many minima.
- $F_2$  shows steady increasing rise vs  $x$  at  $Q^2 > 2 \text{ GeV}^2$ , large  $W^2 \approx Q^2(1 - x)/x$ .
- → Impossible for  $xS(x)$  in DIS scheme, probably can be proved impossible for  $xG(x)$ .

## “No extra minima” constraint



For  $xG(x) = Ax^B(1-x)^C(1+Dx)$  parameterization, require no extra minima for  $x < x_{min}$ :

$$\frac{dxG(x)}{dx} < 0 \quad \text{for } x < x_{min}$$

This imposes relation btw  $B, C, D$  which is for some  $x_{min}$  more restrictive than experimental precision.

## Example of working “clever” parameterization

A similar problem: determination of semileptonic form factor  $\hat{f}_+(t)$  for  $K \rightarrow \pi e \nu$  decays. Standard parameterization:

$$\hat{f}_+(t) = 1 + \lambda' t/m_\pi^2 + 0.5\lambda'' t^2/m_\pi^4 + \dots$$

does not allow to control phase space integral.

Z-mapping:  $t \rightarrow z$  (R. J. Hill, PRD **74** 096006 (2006)),

$$\hat{f}_+(t) = f_+^z(t_0) \frac{\phi(t_0, t_0, Q^2)}{\phi(t, t_0, Q^2)} \sum_{k=0}^{\infty} a_K(t_0, Q^2) z(t, t_0)^k.$$

For some choice of  $\phi(t, t_0, Q^2)$ ,

$$\sum_{k=0}^{\infty} a_k^2 \leq 170$$

from unitarity arguments. This provides better limit for  $k \geq 3$  than experimental data.

## Summary

- Can  $xF(x)$ ,  $(xF(x))'$ , etc be required to have no/finite amount of zeros for certain  $0 < x < x_{min}$  and  $Q^2 > Q_{min}^2$  kinematic domain and some factorization scheme ?
- If yes, what is the best parameterization to use this ?

Can theory be used to control higher order terms in PDF expansion ?

Ideal parameterization: “low orders” controlled from experiment, “higher orders” limited by theory.