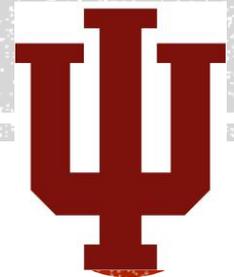
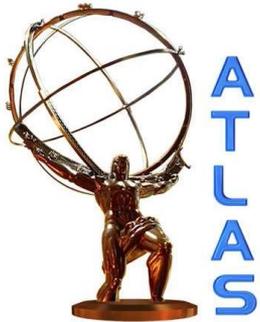


# QUARKONIA PRODUCTION WITH THE ATLAS EXPERIMENT

PHENO2015  
MAY 4 – 6, 2015



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Indiana University

(on behalf of the ATLAS Collaboration)

# Results in Talk

In this talk I will focus on two different methods of studying quarkonium at ATLAS.

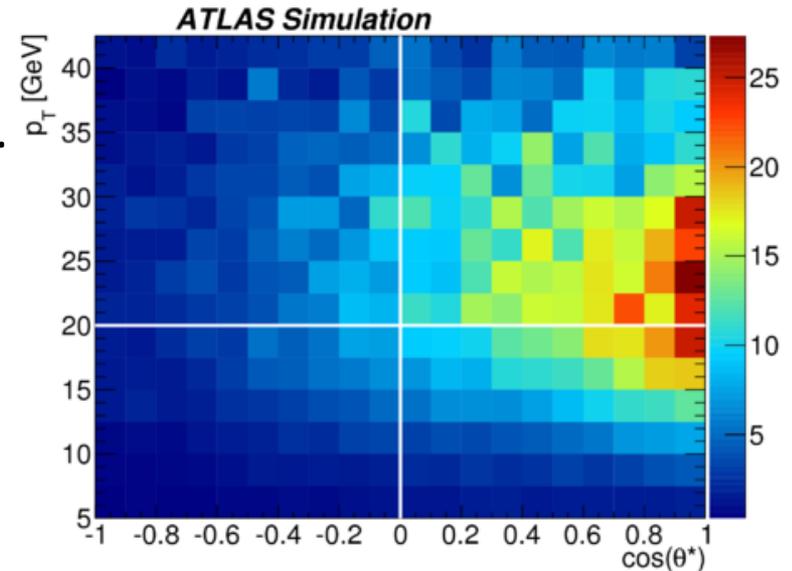
- The first is trying to understand exotic quarkonium through searches for the bottomonium counterpart for  $X(3872)$ .
  - “Search for the  $X_b$  and other hidden-beauty states in the  $\pi^+\pi^-\Upsilon(1S)$  channel at ATLAS”: [arXiv:1410.4409 \[hep-ex\]](#), Phys. Lett. B740 (2015) 199-217
- The second is through studying rare production models of quarkonium.
  - “Observation and measurements of the production of prompt and non-prompt  $J/\psi$  mesons in association with a Z boson in pp collisions at  $\sqrt{s} = 8$  TeV with the ATLAS detector”: [arXiv:1412.6428 \[hep-ex\]](#), **submitted to EPJC**.

# $X(3872)$ and $X_b$

- $X(3872)$  is the first and best-studied hidden-charm state in the last decade. First observed by Belle in  $B^\pm \rightarrow K^\pm X(\rightarrow \pi^+\pi^-J/\psi)$  (arXiv:hepex/0309032v2).
- CMS has measured the product of the cross section of prompt  $X(3872)$  times  $\text{BF}(X(3872) \rightarrow \pi^+\pi^-J/\psi)$  to be  $(6.56 \pm 0.29 \pm 0.65)\%$  of the value for the product of the cross section of  $\psi(2S)$  and  $\text{BF}(\psi(2S) \rightarrow \pi^+\pi^-J/\psi)$  (arXiv:1302.3968).
- Due to the mass, small width, and  $J^{PC} = 1^{++}$  (LHCb: arXiv:1504.0633) it is unlikely to be a conventional quarkonium state.
  - $D^0\bar{D}^{*0}$  molecular models (arXiv:hep-ph/0311229) and tetraquark states (arXiv:hep-ph/0412098) are popular.
- Heavy-quark symmetry suggests the presence of hidden-beauty partner  $X_b$ . We will look at the analogous decay to  $\pi^+\pi^-\Upsilon(1S)$ .
  - CMS found no signal and set an upper limit on the cross section times  $\text{BF}(X_b \rightarrow \pi^+\pi^-\Upsilon(1S))$  at values between 0.9-5.4% of the  $\Upsilon(2S)$  rate (arXiv:1309.0250).
  - Using this decay we also search for  $\Upsilon(1^3D_J)$ ,  $\Upsilon(10860)$ , and  $\Upsilon(11020)$ .

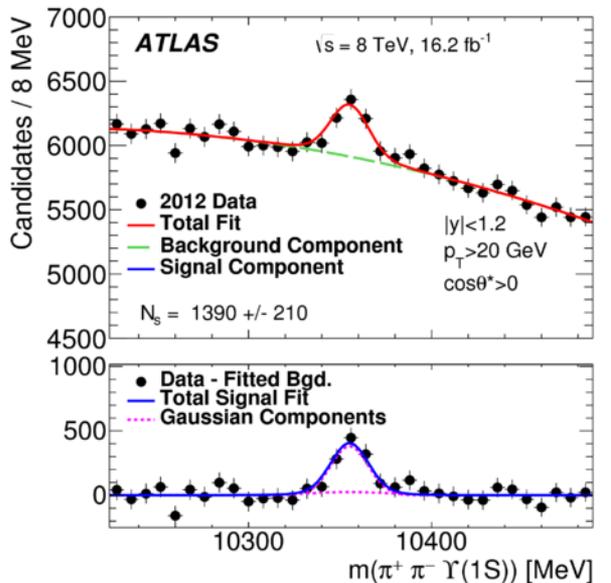
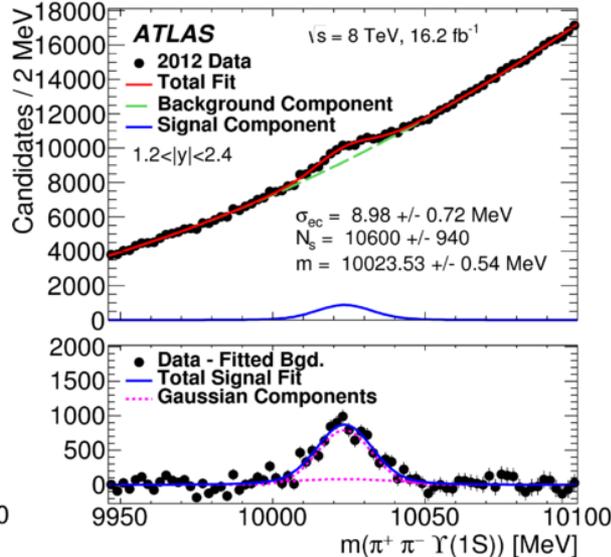
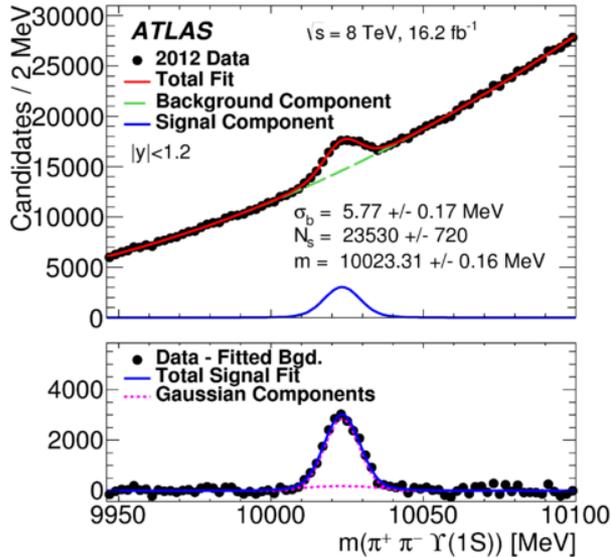
$$X_b \rightarrow \pi^+\pi^-\Upsilon(1S)$$

- 16.2 fb<sup>-1</sup> of  $\sqrt{s} = 8$  TeV ATLAS Data.
- Each event passes a trigger requiring two opp. sign muons with  $p_T > 4$  GeV and  $8 \text{ GeV} < M(\mu\mu) < 12 \text{ GeV}$ .
  - $\Upsilon(1S)$ :  $|\eta^\mu| < 2.3$ ,  $|M(\mu\mu) - M(\Upsilon(1S))| < 350 \text{ MeV}$ .
  - $\pi$ : opp. charge,  $|\eta^\pi| < 2.5$ ,  $p_T > 400 \text{ MeV}$ .
  - $\pi^+\pi^-\Upsilon(1S)$ :  $M(\pi^+\pi^-\Upsilon(1S)) < 11.2 \text{ GeV}$ ,  $p_T > 5 \text{ GeV}$ ,  $|\eta| < 2.4$ .



- ATLAS measurements of  $\Upsilon(nS)$  production cross sections at  $\sqrt{s} = 7$  TeV (in the di-muon channel) are reweighted using Pythia 8 MC to provide predictions for  $\Upsilon(nS)$  cross sections at  $\sqrt{s} = 8$  TeV.
- The data are split into 8 kinematic bins and the signal is extracted via three splitting functions and their complements.
  - Bins are at  $|\eta| = 1.2$ ,  $p_T = 20 \text{ GeV}$ , and  $\text{Cos}(\vartheta^*) = 0$ . Determined by the  $S/\sqrt{B}$  significance.

# $\Upsilon(2S)$ and $\Upsilon(3S)$ Test



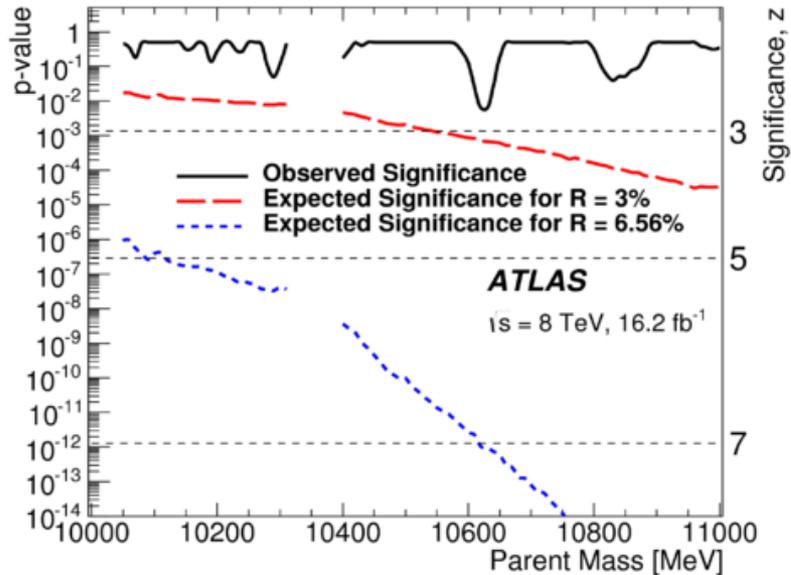
- The  $\Upsilon(2S)$  and  $\Upsilon(3S)$  peaks are fit to validate the measurement techniques.
- The signal shape parameters are fixed to the simulated values.

- The expected yield (calculated in the backup) is compared to the total fitted yield.

- $\Upsilon(2S): N_{2S} = 34,300 \pm 800.$  ✓  
  - $N_{2S}^{expect.} = 33,300 \pm 2,500.$

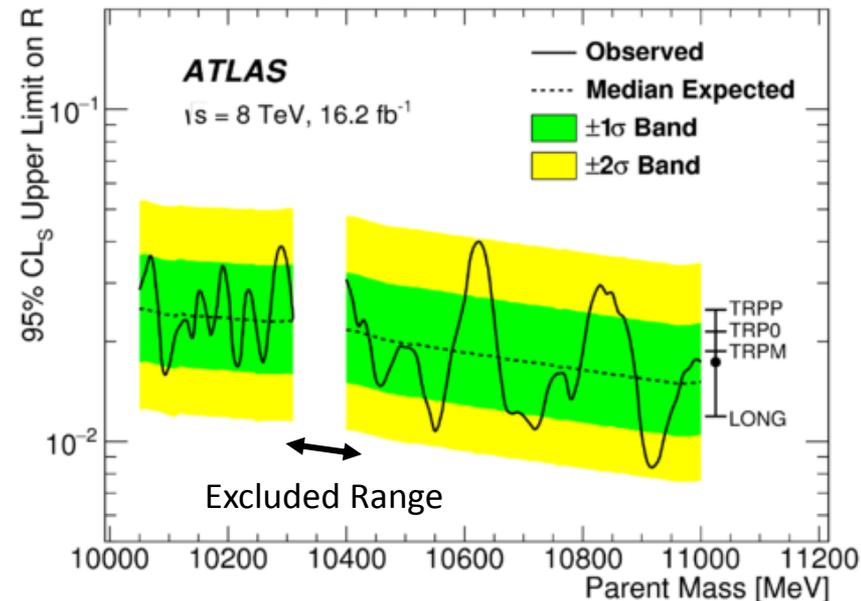
- $\Upsilon(3S): N_{3S} = 11,600 \pm 1,300.$  ✓  
  - $N_{3S}^{expect.} = 11,400 \pm 1,500.$

# Results for $X_b$ search



- The data is fit every 10 MeV excluding the range around  $\Upsilon(2S)$  and  $\Upsilon(3S)$ . The range is:
  - 10.05–10.31 and 10.40–11.00 GeV.
  - **No excess with a local significance of  $3\sigma$  or greater is observed.**
- $$N = N_{2S} * R * \frac{A}{A_{2S}} * \frac{\epsilon}{\epsilon_{2S}}$$
- $R \equiv (\sigma\text{BF})/(\sigma\text{BF})_{2S}$ .

- **Upper limits on  $R = 0.8\text{-}4.0\%$  at 95% CL.**
  - **Excludes analogous  $R = 6.56\%$  from  $X(3872)$ .**
- Upper limit is more restrictive than the CMS results above  $\sim 10.1$  GeV.
- If  $X_b$  exists in this range:
  - The BF or production cross section relative to  $\Upsilon(2S)$  is smaller than  $X(3872)$  relative to  $\psi(2S)$
  - There are no strong isospin-violating effects present as in  $X(3872) \rightarrow \pi^+\pi^-J/\psi$ .
    - If so  $X_b$  would have more prominent decays to  $\pi^+\pi^-\chi_{b1}$  or  $\pi^+\pi^-\pi^0\Upsilon(1S)$ .



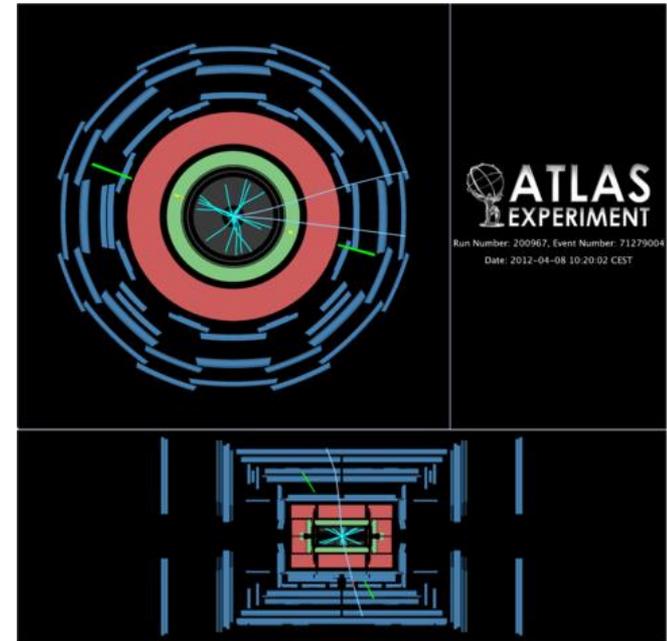
# $\Upsilon(1^3D_J)$ , $\Upsilon(10860)$ , and $\Upsilon(11020)$

- To fit the  $\Upsilon(1^3D_J)$  triplet two extra peaks are added.
  - CLEO (arXiv:hep-ex/0404021) and BaBar (arXiv:1004.0175) measured an average  $M(\Upsilon(1^3D_2)) = (10163.7 \pm 1.4)$  MeV.
  - Averaging over multiple models gives 10156, 10164, and 10170 MeV.
  - **Assuming a common lineshape model for  $J=0,1,2$ , no significant signals are observed.**
  - Assuming  $J = 2$  dominates or the mass splitting is larger than the resolution and  $BF(\Upsilon(1^3D_2) \rightarrow \pi^+\pi^-\Upsilon(1S))$  (arXiv:1004.0175) gives:
    - **$\sigma(\Upsilon(1^3D_2))/\sigma(\Upsilon(2S)) \leq 0.55$ .**
- Broader range for  $\Upsilon(10860)$ , and  $\Upsilon(11020)$ .
  - Assuming the world average for the mass and width, significances of  $0.6\sigma$  and  $0.3\sigma$  are found for  $\Upsilon(10860)$  and  $\Upsilon(11020)$ .
  - If the mass and width are calculated in a grid with  $m \pm 20$  MeV and  $\Gamma \pm \Delta\Gamma$ :
    - **$\Upsilon(10860)$ : Significance of  $1.1\sigma$  at  $m = 10856$  MeV and  $\Gamma = 55$  MeV.**
    - **$\Upsilon(11020)$ : Significance of  $0.6\sigma$  at  $m = 11039$  MeV and  $\Gamma = 95$  MeV.**

# Z+J/ψ

- The first measurement of Z+J/ψ.
  - 20.3 fb<sup>-1</sup> of  $\sqrt{s} = 8$  TeV ATLAS Data.
  - Look at both di-electron and di-muon decay modes of Z boson.
- Interests:
  - Spin-alignment of charmonium and distinguish between color octet and singlet models.
  - Improve perturbative convergence of calculations.
  - Heavy flavor production in association with a Z boson.
  - ZZ\* and a background to Z→l+l'J/ψ.
  - Rare Higgs decay, charm couplings, and CP properties.
  - Double Parton Scattering.

Z(→e<sup>+</sup>e<sup>-</sup>) + J/ψ



## Z boson selection

$$p_T(\text{trigger lepton}) > 25 \text{ GeV}, \quad p_T(\text{sub-leading lepton}) > 15 \text{ GeV}$$

$$|\eta(\text{lepton from } Z)| < 2.5$$

$$|m(Z) - 91.1876 \text{ GeV}| < 10 \text{ GeV}$$

## J/ψ selection

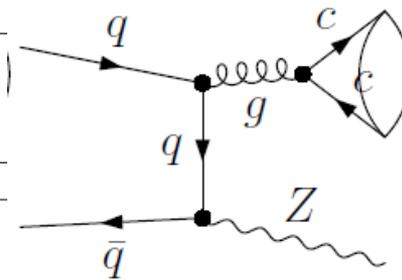
$$2.6 < m(J/\psi) < 3.6 \text{ GeV}$$

$$8.5 < p_T(J/\psi) < 100 \text{ GeV}, \quad |y(J/\psi)| < 2.1$$

$$p_T(\text{leading muon}) > 4.0 \text{ GeV}, \quad |\eta(\text{leading muon})| < 2.5$$

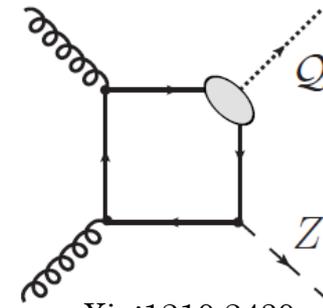
$$\left\{ \begin{array}{l} \text{either } p_T(\text{sub-leading muon}) > 2.5 \text{ GeV}, \quad 1.3 \leq |\eta(\text{sub-leading muon})| < 2.5 \\ \text{or } p_T(\text{sub-leading muon}) > 3.5 \text{ GeV}, \quad |\eta(\text{sub-leading muon})| < 1.3 \end{array} \right\}$$

COM



arXiv:1102.0398

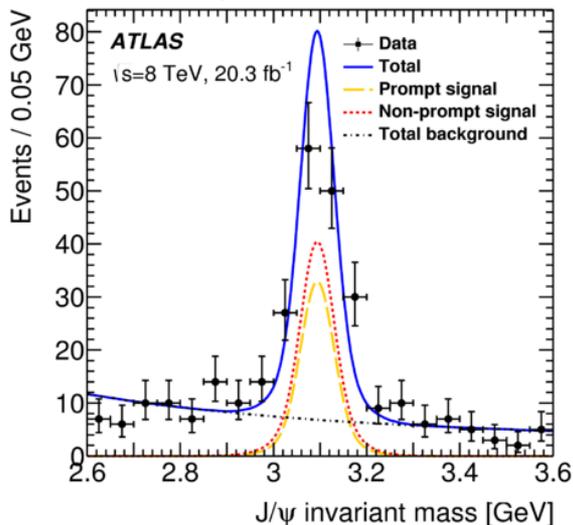
CSM



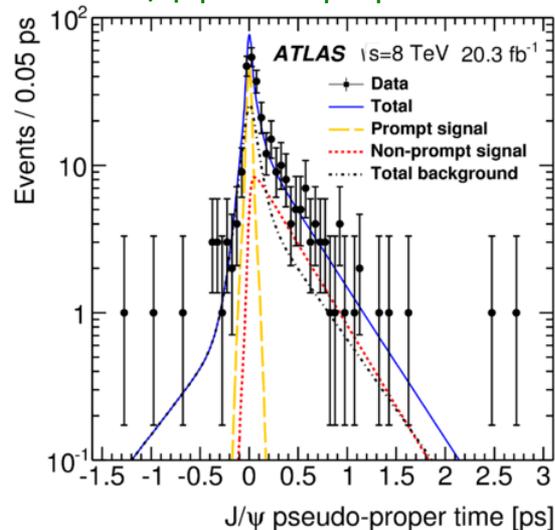
arXiv:1210.2430

# Extracting Signal

J/ψ invariant mass

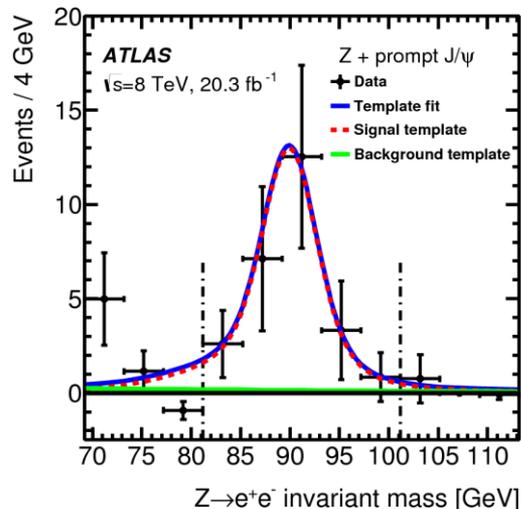


J/ψ pseudo-proper time

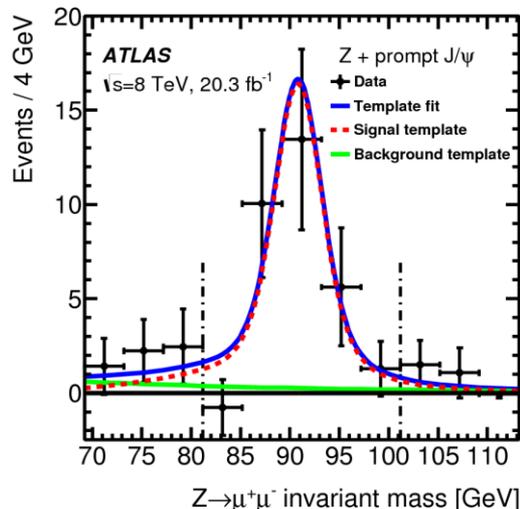


- $Z \rightarrow \tau\tau$ ,  $W \rightarrow l\nu$ ,  $t\bar{t}$ , single top and diboson backgrounds are modeled by MC.
- Multijet and misidentified leptons are modelled by reversing isolation req.
- Pileup background events are estimated (backup).

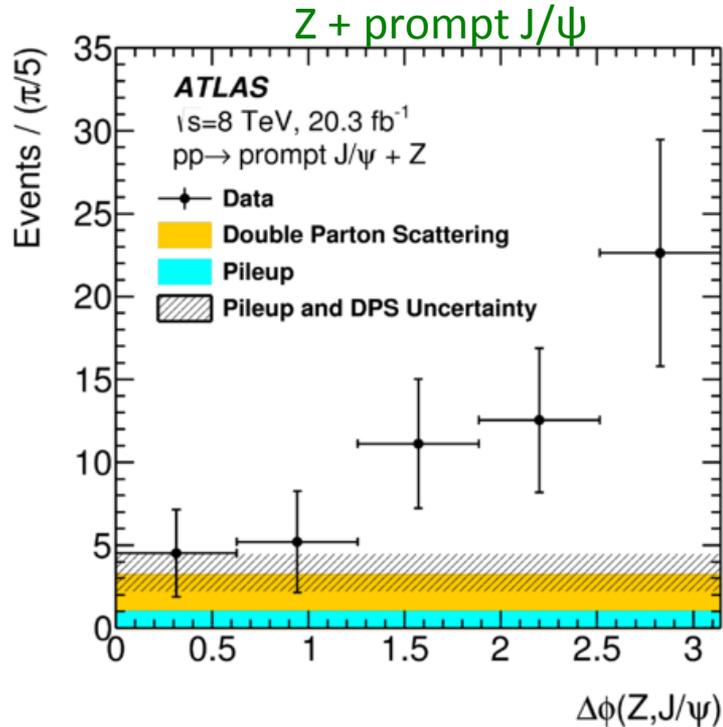
$Z(\rightarrow e^+e^-) + \text{Prompt } J/\psi$



$Z(\rightarrow \mu^+\mu^-) + \text{Prompt } J/\psi$

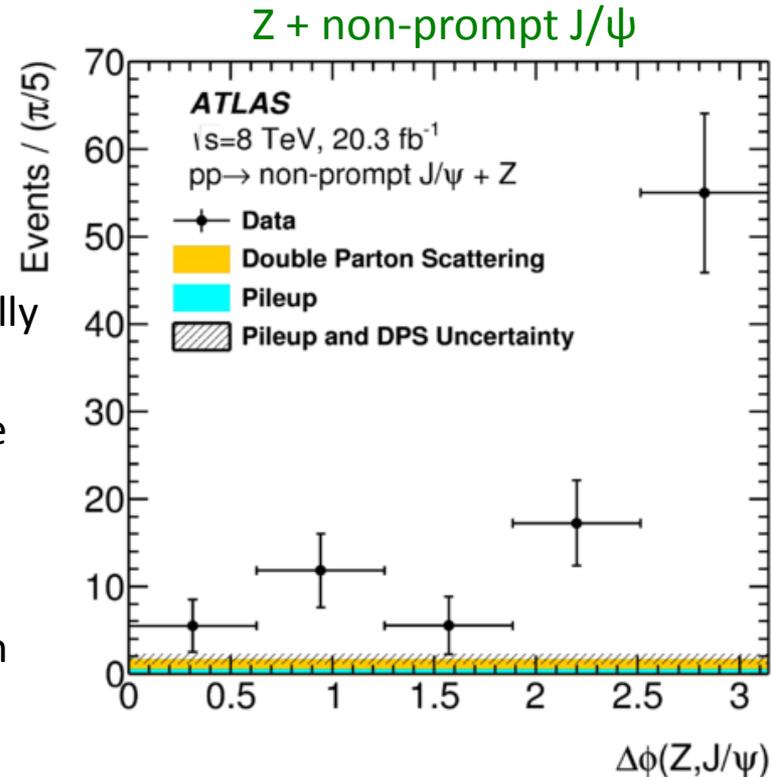


# Double Parton Scattering



- **Double Parton Scattering (DPS):** two parton-parton interactions in a single hadron-hadron collision.
- $$\sigma_{DPS} = \frac{\sigma_Z \sigma_{J/\psi}}{\sigma_{Eff}}$$
  - $\sigma_{Eff} = 15 \pm 3(\text{stat})_{-3}^{+5}(\text{sys}) \text{ mb}$  taken from ATLAS W+2j results (arXiv:1301.6872[hep-ex]) for DPS estimates.
- DPS events are distributed uniformly in  $\Delta\phi(Z, J/\psi)$ .

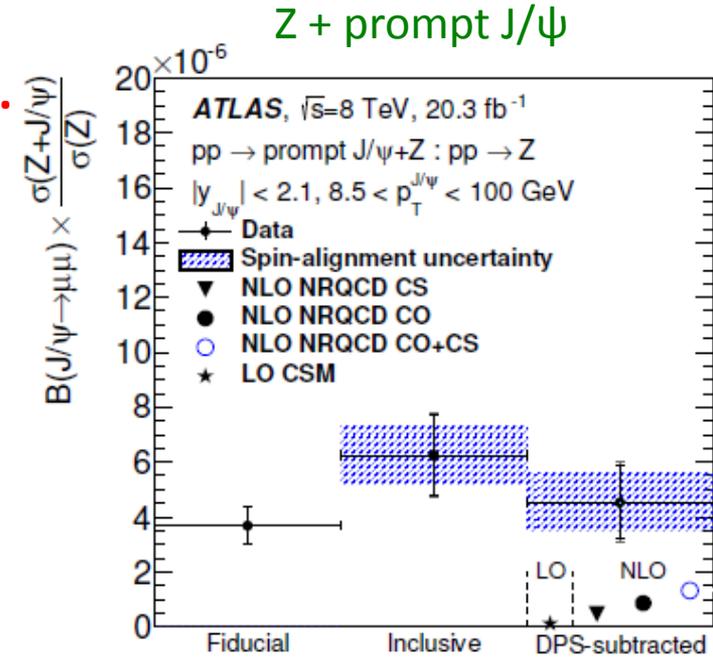
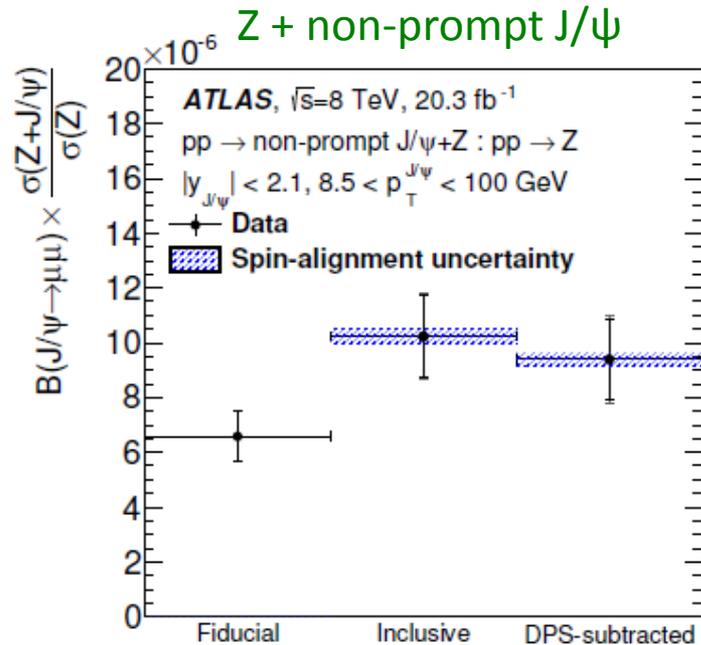
- **Single Parton Scattering (SPS)** events preferentially produced back-to-back, peaking at  $\Delta\phi = \pi$ .
- We set a conservative lower limit on the effective cross-section (upper limit on DPS production) by assuming  $|\Delta\phi| < \pi/5$  is purely DPS.
  - $\sigma_{Eff} > 5.3 \text{ mb}$  ( $3.7 \text{ mb}$ ) at 68% (95%) CL.
- Results are measured in the fiducial volume, then the inclusive volume, and finally after DPS subtraction.



# Results: Z+J/ψ

- First measurement of Z+J/ψ. Null hypothesis excluded at 5σ for prompt and 9σ for non-prompt.

- $$R_{Z+J/\psi} = \text{BF}(J/\psi \rightarrow \mu^+ \mu^-) \frac{\sigma(\text{pp} \rightarrow Z+J/\psi)}{\sigma(\text{pp} \rightarrow Z)}$$
  - $$R_{Z+J/\psi}^{\text{fid}} = (36.8 \pm 6.7 \pm 2.5) \times 10^{-7} \text{ (prompt).}$$
  - $$R_{Z+J/\psi}^{\text{fid}} = (65.8 \pm 9.2 \pm 4.2) \times 10^{-7} \text{ (non-prompt).}$$

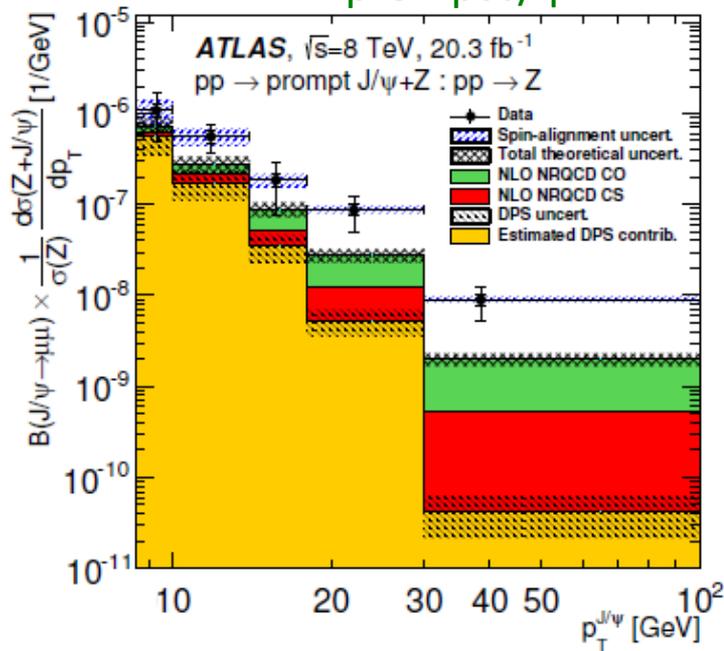


- $$f_{\text{DPS}} = (29 \pm 9)\% \text{ prompt}, (8 \pm 2)\% \text{ non-prompt.}$$

- Theory underestimates SPS production.
  - Does not include feed-down from excited charmonium states, but only accounts for about 20%.

# Differential Cross Section Ratio

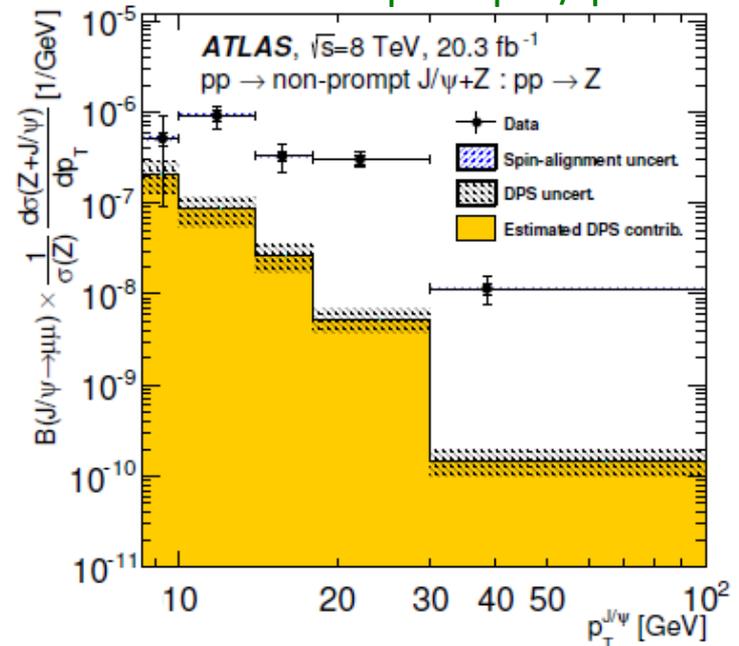
## Z + prompt J/ψ



- Prompt differential cross section ratio is compared to NLO SPS predictions and DPS expectations.
- Theory predicts  $CO > 2 * CS$  with CO increasing with  $p_T$ .
- DPS + NLO SPS predictions underestimate the data by a factor of 4-5 for  $p_T^{J/\psi} > 18$  GeV.

- Significant SPS contribution to Z + non-prompt J/ψ is expected from Z + b-jet production.
- The data presents the opportunity to test Z + b-jet production at low  $p_T$ .

## Z + non-prompt J/ψ



NLO Theory: arXiv:1102.0398 [hep-ph].

# Summary

- Search for  $X_b$  and other hidden beauty states ([arXiv:1410.4409 \[hep-ex\]](#)):
  - No evidence of new states for masses 10.05–10.31 GeV and 10.40–11.00 GeV.
  - **Upper limit set on  $R = (\sigma_{BF})/(\sigma_{BF})_{2S}$  ranging from 0.8 – 4.0%. The  $X_b$  limit is the most stringent in this channel to-date.**
    - Excludes the analogous  $R = 6.56\%$  of  $X(3872)$  for all masses.
  - No significant signal for the  $\Upsilon(1^3D_j)$  triplet,  $\Upsilon(10860)$ , or  $\Upsilon(11020)$ .
  - **95% CLs upper limit of 0.55 is set on  $\sigma(\Upsilon(1^3D_2))/\sigma(\Upsilon(2S))$ .**
- Z+J/ $\psi$  production ([arXiv:1412.6428 \[hep-ex\]](#)):
  - **Measured for the first time. Background-only rejection at  $5\sigma$  for prompt and  $9\sigma$  for non-prompt.**
  - DPS plays a large role ( $\sim 29\%$  for prompt J/ $\psi$ ). A lower limit on the effective cross section (upper limit on the rate of DPS) is measured:
    - **$\sigma_{Eff} > 5.3$  mb (3.7 mb) at 68% (95%) CL.**
  - NLO SPS + DPS predictions underestimate the Z + prompt J/ $\psi$  data.
  - Large Z + non-prompt J/ $\psi$  SPS component is expected from Z+b-jet.



# BACKUP SLIDES

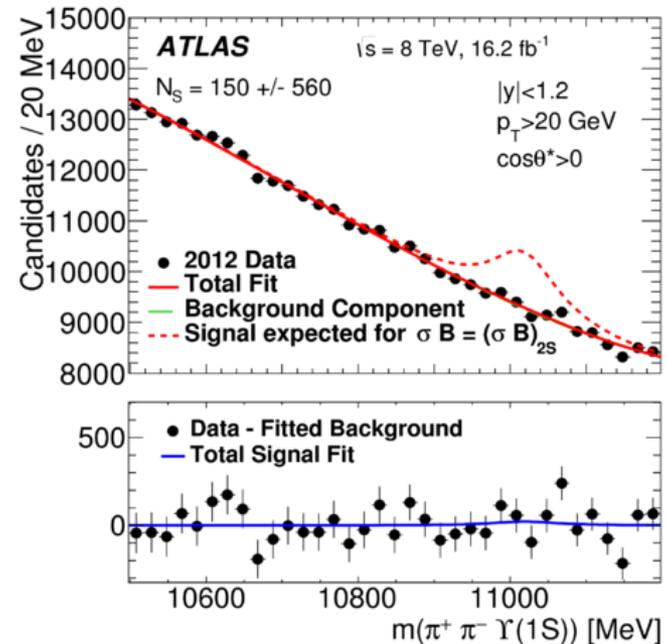
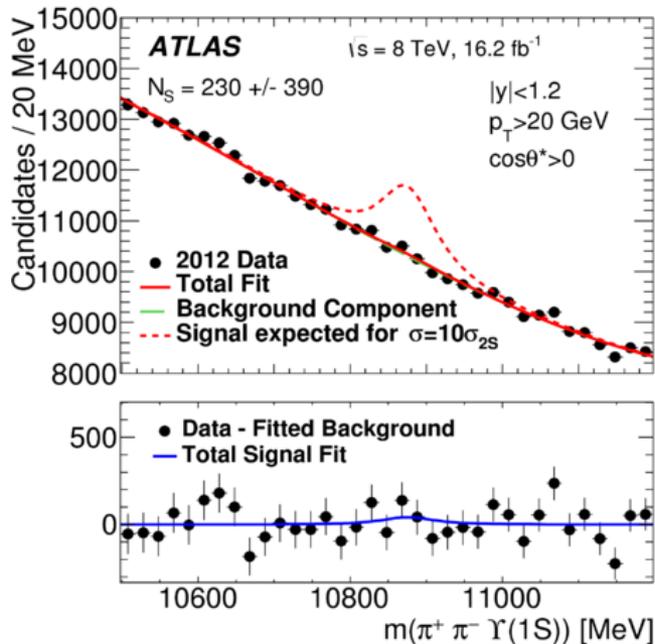
# Splitting Functions

- Splitting functions are used to calculate the signal yield in any particular ( $|y|$ ,  $p_T$ ,  $\cos\vartheta^*$ ) bin.
  - $S_{|y|} = 0.606 \pm 0.004$  for  $|y| < 1.2$ .
  - $S_{p_T}^{b(ec)}(m) = a/(1+e^{-b(m-c)})$  for  $p_T < 20$  GeV.
  - $S_{\cos\vartheta^*}(m) = a + bm + cm^2$  for  $\cos\vartheta^* < 0$ .
- Multiplying the splitting functions and their complements gives the signal yield. For example in the bin ( $|y| < 1.2$ ,  $p_T > 20$  GeV,  $\cos\vartheta^* < 0$ ) the function is:
  - $S_{|y|} \cdot (1 - S_{p_T}^b(m)) \cdot S_{\cos\vartheta^*}(m)$ .

# Calculation of Expected $\Upsilon(2S)$

- $N_2^{exp} = (\sigma\text{BF})_{2S} \cdot L \cdot A \cdot \epsilon$ 
  - $(\sigma\text{BF})_{2S} \equiv$  The product of the cross section and branching fraction estimated from the extended cross section measurement using world averaged values for  $\text{BF}(\Upsilon(1S) \rightarrow \mu^+\mu^-)$  and  $\text{BF}(\Upsilon(2S) \rightarrow \pi^+\pi^-\Upsilon(1S))$ .
  - $L \equiv$  Luminosity.
  - $A \equiv$  The acceptance calculated from the shape of the doubly differential cross section.  $A = (1.442 \pm 0.004)\%$  assuming the CLEO dipion mass spectrum (arXiv:hep-ex/9802024) and isotropic signal decays.
  - $\epsilon \equiv$  The reconstruction efficiency for decays within the acceptance is taken from the  $\Upsilon(2S)$  simulation,  $\epsilon = 0.283 \pm 0.002$ .

# $\Upsilon(10860)$ and $\Upsilon(11020)$



- The  $\Upsilon(10860)$  (left) and  $\Upsilon(11020)$  (right). simultaneous fits projected onto the most sensitive bin.
- The dashed line is a strong signal with  $\sigma = 10\sigma_{2S}$  for  $\Upsilon(10860)$  and  $\sigma B = (\sigma B)_{2S}$  for  $\Upsilon(11020)$ .
- No attempt is made to explain the large partial width of  $\Upsilon(10860)$  compared to  $\Upsilon(3S)$  and  $\Upsilon(4S)$ .
- No strong signal is observed for either state.

# Systematic Uncertainty: Search for $X_b$

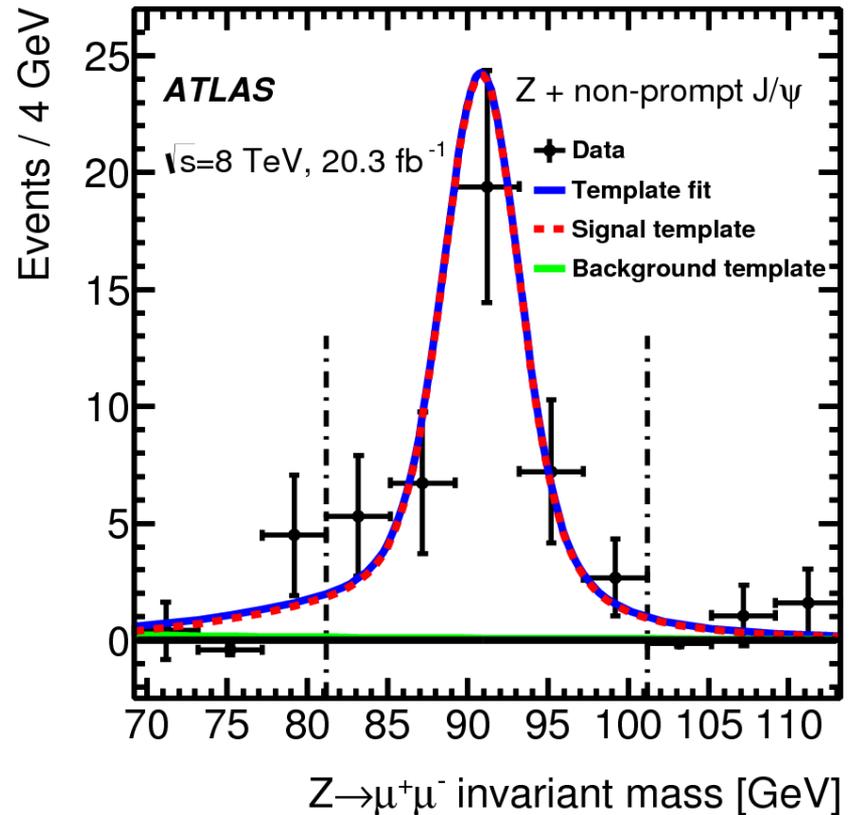
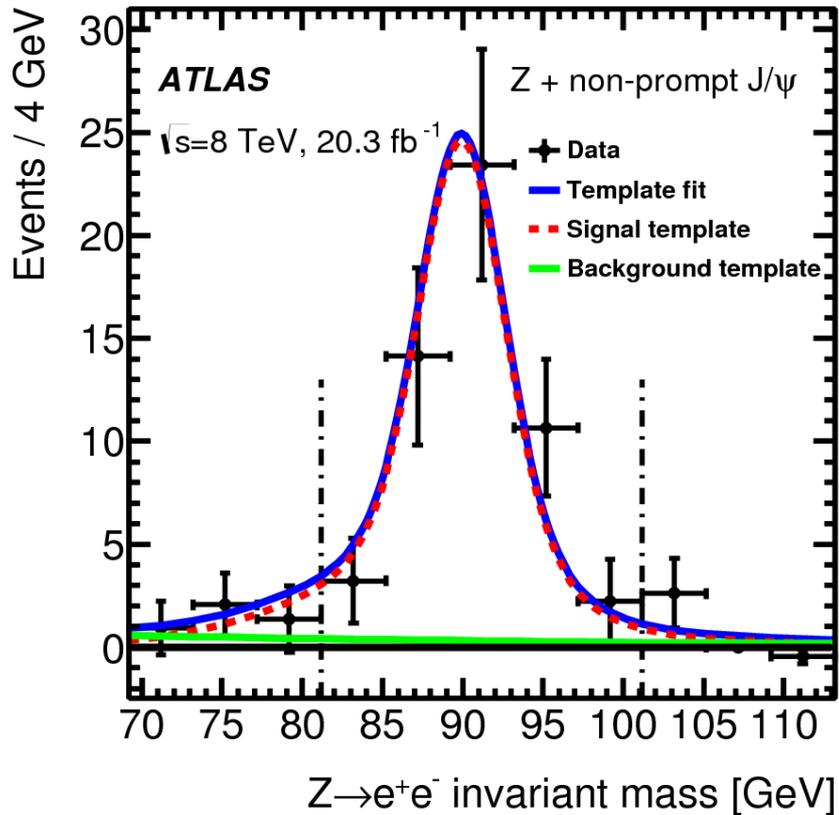
Signal shape parameters

Signal bin splitting functions

	$\sigma_b$ [%]	$\sigma_{ec}$ [%]	$f_b$ [%]	$f_{ec}$ [%]	$r_b$ [%]	$r_{ec}$ [%]	$S_{ y }$ [%]	$S_{p_T}^b$ [%]	$S_{p_T}^{ec}$ [%]	$S_{\cos\theta^*}^{(1)}$ [%]	$S_{\cos\theta^*}^{(2)}$ [%]	$S_{\cos\theta^*}^{(3)}$ [%]	$S_{\cos\theta^*}^{(4)}$ [%]
Extracting $f, r$			0.5	1.1	1.2	1.4							
Extrapolating $\sigma$	0.1	0.2											
Data/MC difference in $\sigma$	1.9	4.2											
$ y $ scale factors							5.8						
Production weighting							0.3	8.4	7.0	0.9	2.8	2.1	3.4
Bin splittings: fit							0.2	0.5	0.8	2.4	4.2	2.8	6.0
Bin splittings: parameterisation							1.8	1.0	1.2	0.2	0.2	0.4	0.2
$m_{\pi^+\pi^-}$ shape							0.2	8.0	11.5	34.7	16.2	15.9	15.0
<b>Total</b>	2.0	4.2	0.5	1.1	1.2	1.4	6.1	11.6	13.6	34.8	17.0	16.3	16.6

	$N_{2S}$ [%]	$\epsilon/\epsilon_{2S}$ [%]	$\mathcal{A}/\mathcal{A}_{2S}$ [%]	$\epsilon/\epsilon_{2S} \cdot \mathcal{A}/\mathcal{A}_{2S}$ [%]
$N_{2S}$ yield	2.3			
$\epsilon$ vs. $m$ : fit		1.0		
$\epsilon$ vs. $m$ : parameterisation		0.5		
Production weighting		1.0		
Acceptance Extrapolation			11.7	
$m_{\pi^+\pi^-}$ shape				17.3
<b>Total</b>	2.3	1.5	11.7	17.3

# Z + Non-prompt J/ $\psi$ Mass fits



# Z + J/ψ Fit Results

Process	$ y_{J/\psi}  < 1.0$	$1.0 <  y_{J/\psi}  < 2.1$	Total	
			Events found	From pileup
Prompt signal	$24 \pm 6 \pm 2$	$32 \pm 8 \pm 5$	$56 \pm 10 \pm 5$	$5.2^{+1.8}_{-1.3}$
Non-prompt signal	$54 \pm 9 \pm 3$	$41 \pm 8 \pm 7$	$95 \pm 12 \pm 8$	$2.7^{+0.9}_{-0.6}$
Background	$61 \pm 11 \pm 6$	$77 \pm 13 \pm 7$	$138 \pm 17 \pm 9$	

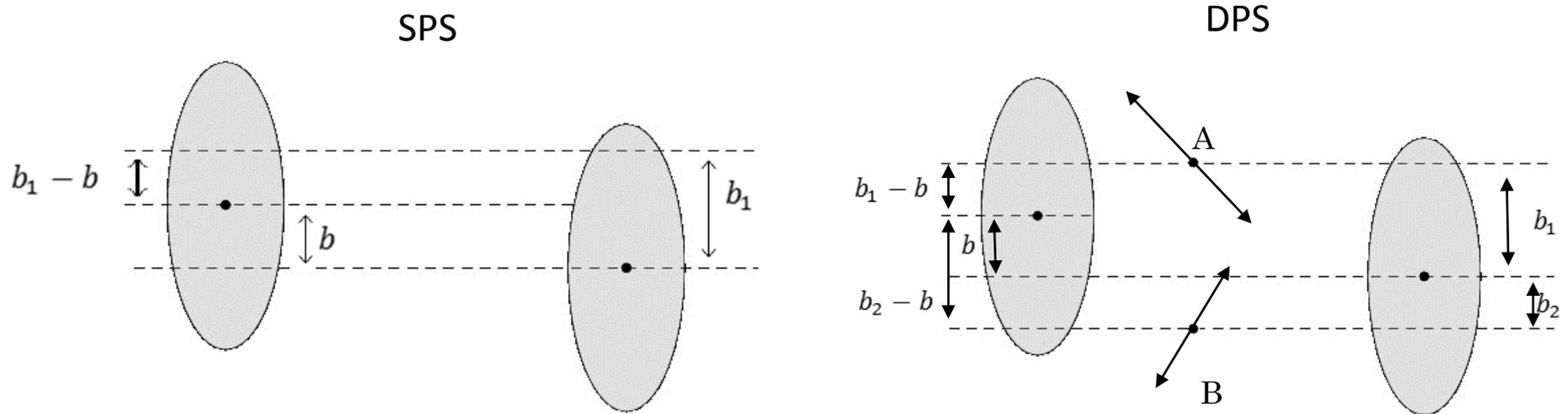
# “Pile-up”: Double Interactions

- Cut on  $|d_z| > 10$  mm between J/ $\psi$  and Z boson limits DI background without biasing the signal.
  - Need to calculate number of DI events that make it into cut.
- $$N_{DI} = N_{\text{extra}} N_Z P_{J/\psi}^{ij} = N_{\text{extra}} N_Z \frac{\sigma_{J/\psi}^{ij}}{\sigma_{\text{inel}}}$$
  - $N_{\text{extra}} = 2.3 \pm 0.2$ . The number of additional vertices within 10 mm of a randomly selected vertex.
  - $N_Z \equiv$  number of Z candidates.
  - $P_{J/\psi}^{ij} \equiv$  the probability for a J/ $\psi$  to be produced at a given pileup vertex.
  - $\sigma_{\text{inel}} = 73$  mb (arXiv:1101.2185 [hep-ex]).
- $N_{DI} = 5.2^{+1.8}_{-1.3}$  (prompt),  $2.7^{+0.9}_{-0.6}$  (non-prompt).

# Double Parton Scattering: Features

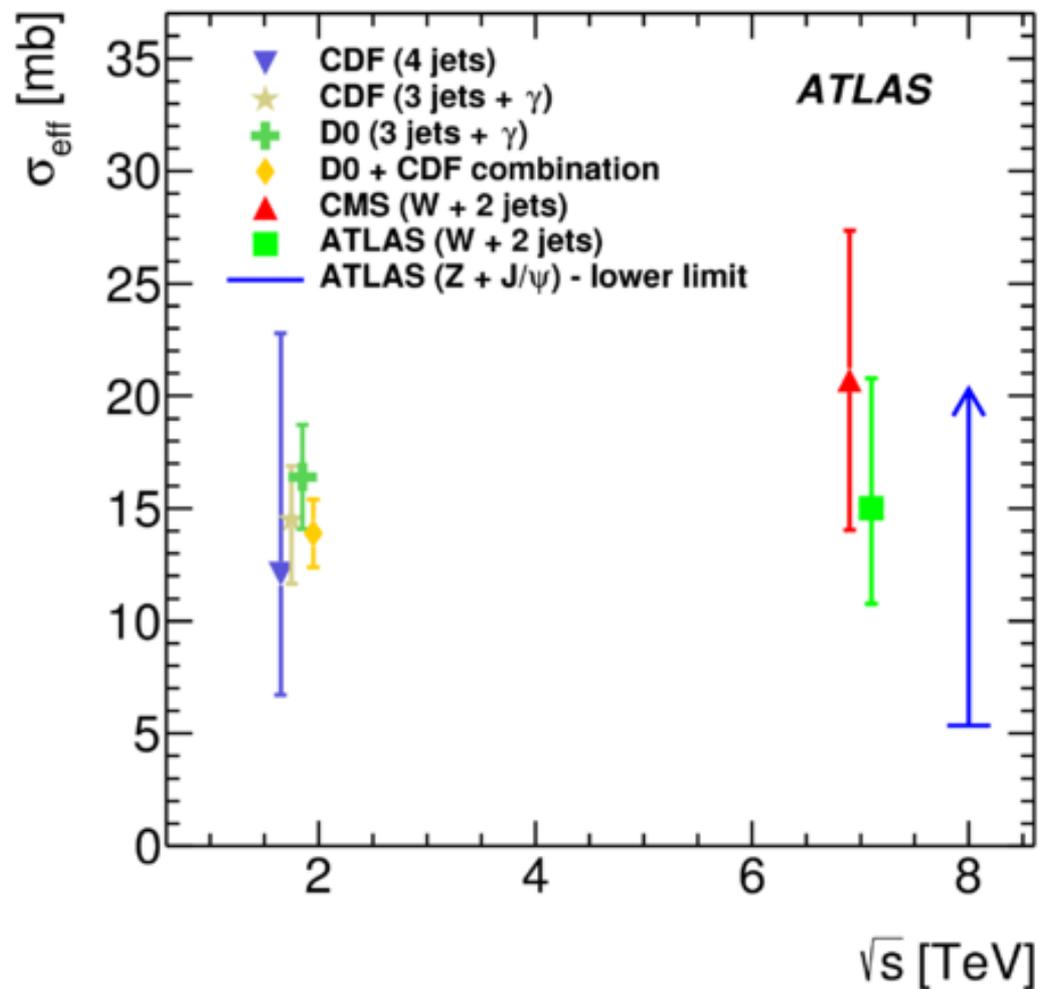
- Double Parton Scattering (DPS) requires large c.m. energies and low values of incoming fractional momenta ( $x_F$ ). This is possible to achieve at a non-negligible rate at the LHC.
- Because it is dependent on the transverse distance between interactions regions, the DPS cross section decreases quickly as a function of transverse energy.
- Assuming that the two processes ( $\sigma_A, \sigma_B$ ) are independent of each other, DPS cross section can be written as: 
$$\sigma_{DPS}^{(A,B)} = \frac{m}{2} \frac{\sigma_A \sigma_B}{\sigma_{Eff}}$$
  - $m = 1$  if A and B are indistinguishable and  $m = 2$  if they are distinguishable
  - $\sigma_{Eff} = \left[ \int \left[ \int f(b_1) f(\mathbf{b}_1 - \mathbf{b}) d^2 b_1 \right]^2 d^2 b \right]^{-1}$  where  $f(b)$  parton density in the transverse plane and is assumed to be a universal function, same for both protons.
  - $\frac{\sigma_B}{\sigma_{Eff}}$  is the probability for scattering B to occur given scattering A has already occurred.
  - $\sigma_{Eff}$  measures the size in impact parameter space of the incident hadron's partonic core.
- $\sigma_{Eff} \sim \frac{1}{4} \sigma_{Inel}$ .
  - If the effective cross section was equal to the inelastic cross section, it would imply uncorrelated scatterings.
  - This result indicates a correlation (“clumpiness”) in the hadron structure.
- A constant value of  $\sigma_{Eff}$  has been able to describe results in different kinematical regions. CDF has also tested the dependence of  $\sigma_{Eff}$  on  $x_F$  and had compatible results with being independent of  $x_F$ .

# Double Parton Scattering



- $b_1, b_2$  are the transverse positions of the two partons.
- $b$  is the impact parameter (distance between the centers of the colliding hadrons).
- $pp \rightarrow A \oplus pp \rightarrow B$ .
- A and B with a clean signature, and large rates.
- Example:  $A = \gamma + j$ ,  $B = j + j$ .

# Effective Cross Section



# Systematic Uncertainty: Z+J/ $\psi$ Production

Source	Prompt		Non-prompt	
	$ y_{J/\psi}  < 1.0$	$1.0 <  y_{J/\psi}  < 2.1$	$ y_{J/\psi}  < 1.0$	$1.0 <  y_{J/\psi}  < 2.1$
Fit procedure	3%	3%	4%	8%
Z boson kinematics	1%	1%	1%	1%
$\mu_{J/\psi}$ efficiency	1%	1%	1%	1%
Vertex separation	7%	16%	2%	15%