# Dark matter explained through two distinct ideas related to Higgs

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In collaboration with S. Chakdar, K. Ghosh and S. Nandi, arXiv:1410.7331 [hep-ph] (Submitted to PRD) K.S. Babu, S. Chakdar and R.N Mohapatra, Phys. Rev. D 91, 075020, arXiv:1412.7745

Phenomenology Symposium 2015, Pittsburgh 5th May 2015

#### **Outline of Talk**

- Dark matter through multiple higgs signals at the ILC
  - Motivation
  - Model and formalism
  - Phenomenology
  - Results
- Warm dark matter in two higgs doublet models
  - Model and formalism
  - Cosmological details
  - Allowed parameter space and results
  - X-ray anamoly explanation
- Conclusions

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- In first work, we look into  $e^+e^-$  Collider Phenomenology for  $\delta m \sim 100$  MeV between two Higgs bosons

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- Also assume post-inflationary reheating in two universes different and parallel universe colder than our universe. This makes possible to maintain the successful prediction of BBN

# Running of the strong coupling constant

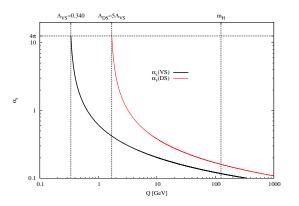


Figure: The running of the strong coupling constant in the visible sector and dark sector shows  $\alpha_S^{DS}(Q=m_H=125~{\rm GeV})=1.4~\alpha_S^{VS}(Q=m_H=125~{\rm GeV}).$ 

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- The recoil mass to the *Z*-boson is the invariant mass of the decay products against which the *Z*-boson recoils:

$$M_{recoil}^2 = (\sqrt{s} - E_Z)^2 - |\vec{p}_Z|^2 = s + M_Z^2 - 2E_Z\sqrt{s}$$

where  $M_Z$ =Z-boson mass as reconstructed from the decay products,  $E_Z$  is the corresponding energy and  $\sqrt{s}=250$  GeV

## Results in direct Higgs decays and Recoil to the Z-boson in ILC

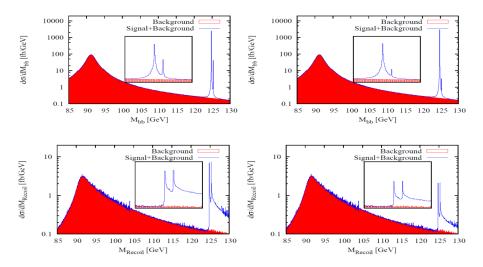


Figure: Invariant mass distributions of  $b\bar{b}$  pairs(top panel) and recoil mass distribution for invisible Higgs decays(bottom panel) where  $m_{h_{SM}}=124.75~{\rm GeV}$  and  $m_{h_{DH}}=125.25~{\rm GeV}$ . 123  ${\rm GeV} < m_{bb} < 127~{\rm GeV}$  regions magnified in insets. 6/14

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- $\bullet$  Consider proposed ILC as it has advatages of clean environment, precise measurements and large number of Higgs boson production. In ILC Higgs masses with  $\delta m \sim 100$  MeV will be possible
- Find that for  $\sqrt{s}=250$  GeV ILC with 500 MeV mass splitting we can see two clear mass peaks when the Higges decay to  $b\bar{b}$  or invisibly.

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- ullet In this work we focus on alt possibility with 2HDMs consisting of warm dark matter candidate in the form of neutral scalar boson having mass  $\sim$  keV
- This scenario is completely consistent with known observations and would have distinct signatures at colliders as well as in cosmology and astrophysics

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- With tiny breaking of  $Z_2$ ,  $\sigma \to \gamma \gamma$  can occur with lifetime longer than age of the universe which can explain reported anomaly in the X-ray spectrum for  $m_{\sigma} = 7.1 \text{ keV}$

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$$m_h^2 = 4\lambda_1 v^2, \quad m_\sigma^2 = m_2^2 + (\lambda_3 + \lambda_4 + \lambda_5) v^2;$$
  
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•  $m_A^2=m_\sigma^2-2\lambda_5 v^2$  and  $m_{H^\pm}^2=m_\sigma^2-(\lambda_4+\lambda_5)v^2$ ,  $(m_\sigma\sim \text{keV},~m_A~\text{and}~m_{H^\pm}~\text{can be}\sim \text{few hundred GeV})$ 

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- At very high temperature N was in thermal equilibrium. 2 and 3 body decays of N will come into equilibrium at  $T_d$  at which N would begin to decay. If  $T_d$  in range 150 1 MeV, decay products as do the photons gain entropy

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m of}~\sigma\sim 150~{
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- For DM  $\sim$  keV,  $\Omega_{\sigma}$  is a factor of 34 larger than the observed value of 0.265.
- The decay of N involved in the seesaw mechanism dilutes the abundance of  $\sigma$  and make the scenario consistent.
- At very high temperature N was in thermal equilibrium. 2 and 3 body decays of N will come into equilibrium at  $T_d$  at which N would begin to decay. If  $T_d$  in range 150 1 MeV, decay products as do the photons gain entropy
- Since  $\sigma$  froze out at  $T \approx 150$  MeV, N decay causes temp of photons to increase relative to that of  $\sigma$ . Thus a dilution in abundance of  $\sigma$  is realized.

- $\sigma$  remains in thermal equilibrium in early universe down to  $T_f$  through weak interaction processes. Freeze-out occurs at  $T_{f,\sigma}\approx 150$  MeV through the scattering  $\gamma\gamma\to\sigma\sigma$
- ullet We obtain the abundance of  $\sigma$  today as

$$\Omega_\sigma = 9.02 \left( rac{17.25}{
m g_{eff}} 
ight) \left( rac{m_\sigma}{1 \, {
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- $T_d > 1$  MeV, so BBN is not affected. The desired range for the lifetime of N is thus  $\tau_N = (10^{-4} 1)$  sec

# Allowed parameter space in $M_N - Y_N$ plane

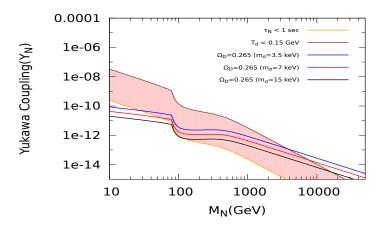


Figure : Allowed parameter space in  $M_N-Y_N$  plane in shown. The shaded region corresponds to decay temperature  $T_d$  of N in the range 150 MeV-1 MeV. The three solid curves generate the correct dark matter density  $\Omega_D$  for three values of the WDM mass :  $m_\sigma=\{3.5,7,15\}$  keV

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- This is matched to partial lifetime  $\Gamma^{-1}(\sigma \to \gamma \gamma) = (4 \times 10^{27} 4 \times 10^{28})$  sec
- Once  $\sigma$  develops a vev u, it also mixes with SM Higgs field h, but this effect is subleading for the decay  $\sigma \to \gamma \gamma$

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- The model can also explain the anomalous X-ray signal reported by different groups in the extra-galactic spectrum.
- The charged scalar and the pseudoscalar in the model cannot be much heavier than a few hundred GeV so could be detectable at the colliders

# Thank you!

# Backup Slide: Symmetry breaking of Parallel universe model

- Our gauge symmetry is  $SU(4)_C \times SU(2)_L \times SU(2)_R$
- SU(4) color symmetry is spontaneously broken to  $SU(3)_C \times U(1)_{B-L}$  in the usual Pati-Salam way using the Higgs fields (15,1,1) at a scale  $V_c$ , where  $V_c > 2300 \, TeV$  (Valencia, Willenbrock, 1994)
- $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$  can be broken to the SM using the Higgs representations (1,3,1)+(1,1,3) at a scale  $V_{LR}$ , where  $V_{LR}>2.5\,TeV$
- Finally the remaining symmetry is broken to the  $U(1)_{EM}$  using the Higgs bi-doublet (1,2,2) and (15,2,2). Similar Higgs representations are used to break the symmetry in the parallel universe to U'(1)
- $SU(4)_C \otimes SU(2)_L \otimes SU(2)_R \rightarrow SU(3)_C \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L} \rightarrow SU(3)_C \otimes SU(2)_L \otimes U(1)_Y \rightarrow U(1)_{EM}$

# Backup Slide: Details of higgs sector of Parallel universe model

- A study of the Higgs potential shows a parameter space where only one neutral Higgs in the bi-doublet remains light and becomes very similar to the SM Higgs in our universe. Similar is true in the parallel universe
- The quartic Higgs interactions  $\lambda(H_{VS}^{\dagger}H_{VS})(H_{DS}^{\dagger}H_{DS})$  leads to mixing between the two light SM like Higgs fields:

$$\mathcal{L}_{Scalar} \supset m_{VS}^2 h_1^2 + m_{DS}^2 h_2^2 + 2\lambda v_{VS} v_{DS} h_1 h_2$$

from which two mass eigenstates and mixing can be calculated

• The two physical light Higgs states are defined as,

$$h_1^{(p)} = \cos\theta \ h_1 + \sin\theta \ h_2$$
  
$$h_2^{(p)} = -\sin\theta \ h_1 + \cos\theta \ h_2$$

• The masses and the mixing angle of these physical states are given by,

$$\begin{array}{lcl} m_{h_1^{(p)},h_2^{(p)}}^2 & = & \frac{1}{2}[(m_{VS}^2+m_{DS}^2)\mp\sqrt{(m_{VS}^2-m_{DS}^2)^2+4\lambda^2v_{VS}^2v_{DS}^2}]\\ \tan & = & \frac{2\lambda~v_{VS}~v_{DS}}{m_{DS}^2-m_{VS}^2}. \end{array}$$