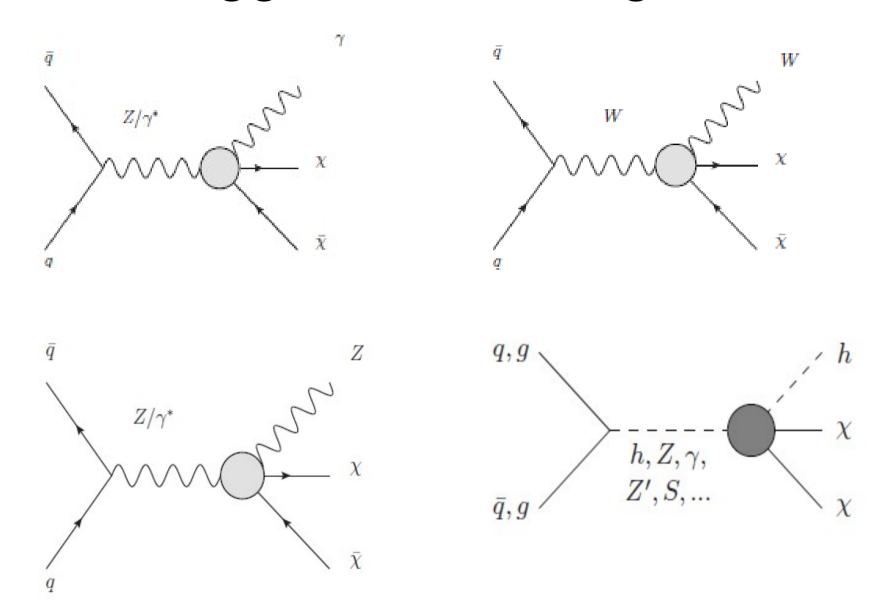
# Mono-Higgs & EW Gauge Boson



arXiv:1212.3352 arXiv:1307.5064 arXiv:1403.6734

## Why?

Low Background

Explore Interesting Portals

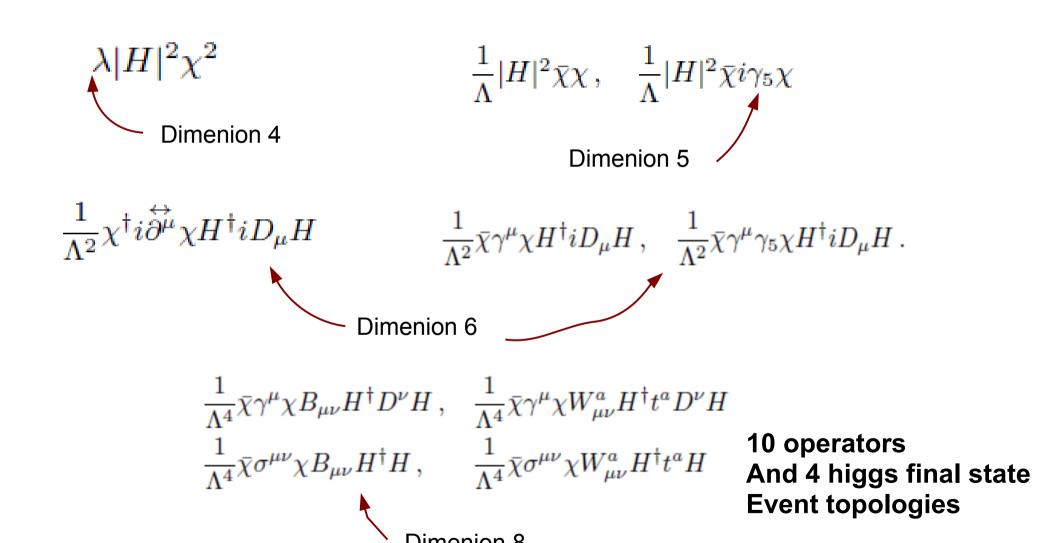
- Complimentary Indirect Experiments looking for annihilation (e.g. line searches)
- Many of these type of couplings hard for Direct Detection, leaving room for collider discovery

# Dimension 6 and 7 Operators

$$\begin{split} \mathcal{L} &= \frac{1}{\Lambda_{C1,2}^3} \; \bar{\chi} \chi \; \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \; + \\ &\frac{1}{\Lambda_{C3,4}^3} \; \bar{\chi} \chi \; \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^{\tilde{i}} \end{split} \qquad \text{Fermionic DM}$$

$$\mathcal{L} = \frac{1}{\Lambda_{C5,6}^3} \, \bar{\chi} \gamma^5 \chi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \frac{1}{\Lambda_{C7,8}^3} \, \bar{\chi} \gamma^5 \chi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B1,2}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ + \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ + \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^2} \, \bar{\phi} \phi \, \sum_i k_i F_i^{\mu\nu} F_i^i \, + \\ \mathcal{L} = \frac{1}{\Lambda_{B3,4}^$$

# Mono-Higgs A Menagerie of Models



#### Extra Slides

#### Operators of Effective Dim 5

$$\frac{1}{\Lambda_5^3} \, \bar{\chi} \chi \, (D_\mu H)^\dagger D^\mu H$$

Inserting Higgs vevs and expanding the covariant derivative yields

$$\frac{m_W^2}{\Lambda_5^3} \; \bar{\chi} \chi \; W^{+\mu} W_{\mu}^- + \frac{m_Z^2}{2\Lambda_5^3} \; \bar{\chi} \chi \; Z^{\mu} Z_{\mu}$$

Notice a single parameter determines ratio of DM couplings to W's and Z's.  $\gamma\gamma$  a  $\gamma$ Z couplings genrated at higher at loop order.

## A Note on Gauge Invariance

The general gauge invariant operators are given by

$$L = \frac{1}{\Lambda_7^3} \,\bar{\chi}\chi \,\sum_i k_i F_i^{\mu\nu} F_{\mu\nu}^i$$

3 independent couplings for 3 SM gauge groups

$$g_{gg} = \frac{k_3}{\Lambda_7^3}$$
 One independent coupling for SU(3)

$$g_{WW} = \frac{2\kappa_2}{s_w^2 \Lambda_7^3}$$

$$g_{ZZ} = \frac{1}{4s_w^2 \Lambda_7^3} \left( \frac{k_1 s_w^2}{c_w^2} + \frac{k_2 c_w^2}{s_w^2} \right)$$

$$g_{\gamma\gamma} = \frac{1}{4c_w^2} \frac{\kappa_1 + \kappa_2}{\Lambda_7^3}$$

$$g_{Z\gamma} = \frac{1}{2s_w c_w \Lambda_7^3} \left( \frac{k_2}{s_w^2} - \frac{k_1}{c_w^2} \right)$$

Related by Gauge invariance, determined by 2 coefficients