

Background independent ERG for conformally reduced gravity

UK QFT Imperial 30/3/15

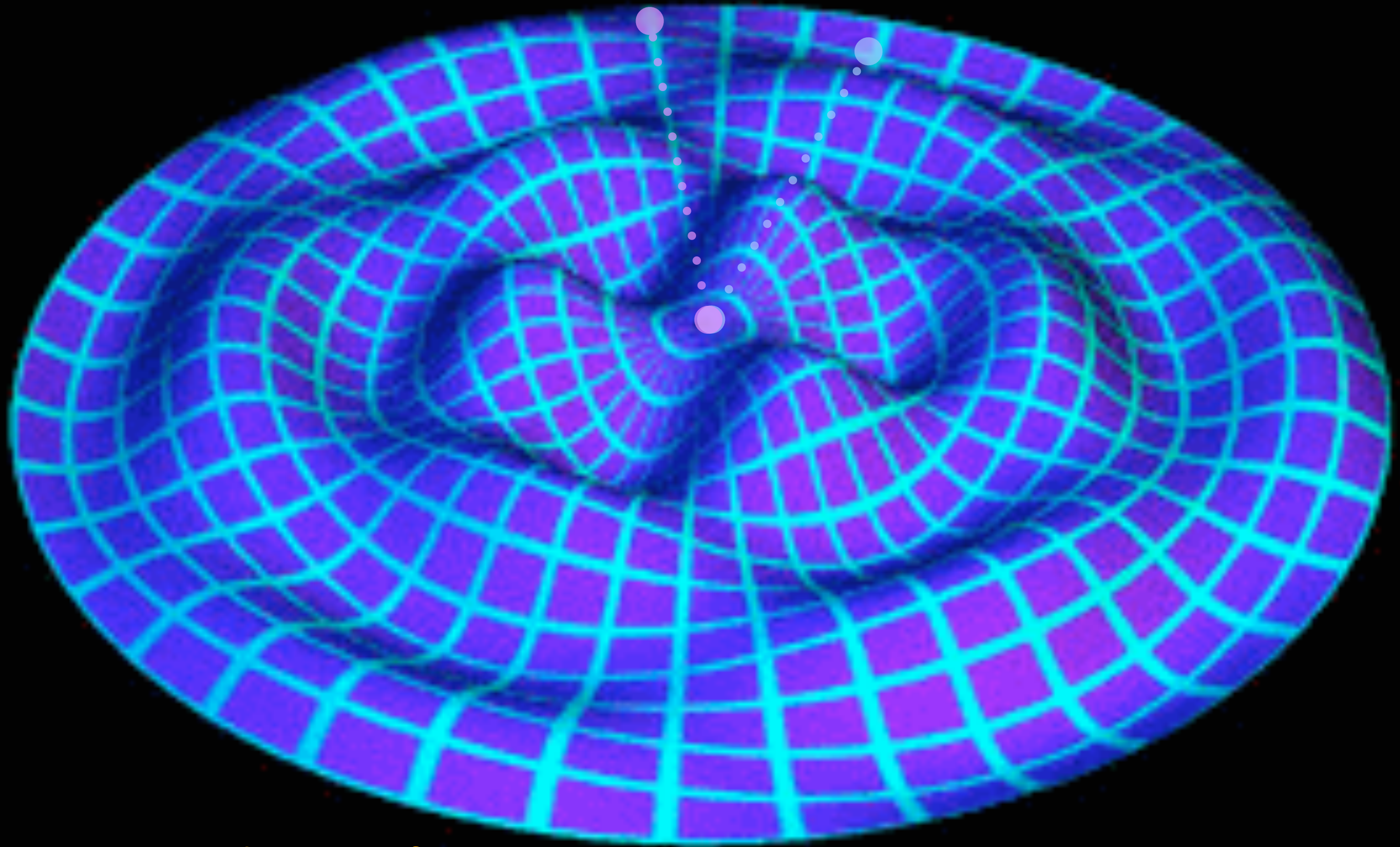
Tim Morris,

Physics & Astronomy,

University of Southampton, UK.

Juergen A. Dietz & TRM, arXiv:1502.0739.

What does scale $1/k$ mean?



Conformal field theory? How to separate scales?
What does scale mean? What is long wavelength?
How to form a continuum limit?

Use background field method:

$$\tilde{g}_{\mu\nu} = \bar{g}_{\mu\nu} + \tilde{h}_{\mu\nu}$$

cutoff:

$$R_k \sim r \left(-\frac{\bar{\nabla}^2}{k^2} \right)$$

- Now the scale k is defined through $\bar{g}_{\mu\nu}$ so the notion depends on the choice of background
- The RG itself depends on both $g_{\mu\nu}$ and $\bar{g}_{\mu\nu}$ so also is inherently background dependent, and necessarily bi-metric
- Is there a background independent notion of scale k ?

Conformally reduced gravity

$$\phi = \varphi + \chi$$

$$g_{\mu\nu} = f(\phi) \delta_{\mu\nu} \quad \text{and} \quad \bar{g}_{\mu\nu} = f(\chi) \delta_{\mu\nu}$$

- Remnant diffeomorphism invariance ...

$$x^\mu \mapsto x^\mu / \lambda, \quad f(\chi) \mapsto \lambda^2 f(\chi)$$

E.g. $S_k[\varphi, \chi] = \int d^d x f(\chi) \left(\frac{1}{2} K(\varphi, \chi) g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi + V(\varphi, \chi) \right)$

Maintain background covariance:

- φ at $\mathcal{O}(\partial^2)$ but slowly varying χ

- Conformal factor problem $\chi \mapsto \chi + \ln \lambda$, $\phi \mapsto \phi + \ln \lambda$ wrong sign kinetic term

- K has absorbed the factor of $1/8\pi G$ and \sqrt{g} (absorbed) $S_k[\varphi, \bar{g}] = \frac{1}{2} \int d^d x \sqrt{\bar{g}} (K(\varphi, \chi) g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi + V(\varphi, \chi))$

Break quantum covariance: fixes factors of $f(\chi)$

$$\chi \mapsto \chi, \quad \phi \mapsto \phi + \ln \lambda, \quad \varphi \mapsto \varphi + \ln \lambda$$

Flow equation

$$\partial_t \Gamma_k[\varphi, \chi] = \frac{1}{2} \int_x \int_y \Delta(x, y) \partial_t R_k(y, x)$$

$$\int_x \equiv \int d^d x f(x)^{\frac{d}{2}}$$

$$f(x) \equiv f(\chi(x))$$

$$\Delta(x, y) = \left[f(x)^{-\frac{d}{2}} f(y)^{-\frac{d}{2}} \frac{\delta^2 \Gamma_k}{\delta \varphi(x) \delta \varphi(y)} + R_k(x, y) \right]^{-1}$$

$$t = \ln(k/\mu)$$

Broken split Ward Identity (msWI)

$$\varphi(x) \mapsto \varphi(x) + \varepsilon(x)$$

$$\chi(x) \mapsto \chi(x) - \varepsilon(x)$$

$$\phi = \varphi + \chi$$

$$\int_w f^{-\frac{d}{2}} \left(\frac{\delta \Gamma_k}{\delta \chi} - \frac{\delta \Gamma_k}{\delta \varphi} \right) \varepsilon = \frac{1}{2} \int_x \int_y \Delta(x, y) \left[\frac{d}{2} \partial_\chi \ln f(y) R_k(y, x) + \partial_\chi R_k(y, x) \right] \varepsilon(y)$$

Sketch of derivative expansion...

Write trace in terms of differential operators...

$$E.g. \quad R_{k,x} = R_k \left(-\bar{\nabla}_x^2 \right) = R_k \left(-f(x)^{-1} \partial_x^2 \right)$$

$$\int d^d x \Delta_x \mathcal{K}_x \delta(x - x') \Big|_{x'=x} = \int d^d x \frac{d^d p}{(2\pi)^d} \underbrace{\left\{ e^{-ip \cdot x} \Delta_x e^{ip \cdot x} \right\}}_{\substack{\uparrow \\ Q(p,x)}} \mathcal{K} \left(f^{-1} p^2 \right)$$

$$Q(p, x) \equiv e^{-ip \cdot x} \left[\Gamma^{(2)} + R \right]^{-1} e^{ip \cdot x}$$

satisfies differential equation:

$$\underbrace{e^{-ip \cdot x} R \left\{ e^{ip \cdot x} Q \right\}}_{\substack{\uparrow \\ R(p^2/f)Q - 2i\partial_{p^2} R(p^2/f) p^\mu \partial_\mu Q - \partial_{p^2} R(p^2/f) \partial^2 Q - 2\partial_{p^2}^2 R(p^2/f) p^\mu p^\nu \partial_\mu \partial_\nu Q + \mathcal{O}(\partial^3)}} = 1 - e^{-ip \cdot x} \Gamma^{(2)} \left\{ e^{ip \cdot x} Q \right\}$$

$$R(p^2/f)Q - 2i\partial_{p^2} R(p^2/f) p^\mu \partial_\mu Q - \partial_{p^2} R(p^2/f) \partial^2 Q - 2\partial_{p^2}^2 R(p^2/f) p^\mu p^\nu \partial_\mu \partial_\nu Q + \mathcal{O}(\partial^3)$$

iterate

Flow equations, msWIs

$$\partial_t V(\varphi, \chi) = f(\chi)^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} \frac{\dot{R}(p^2/f)}{\partial_\varphi^2 V - K p^2/f + R(p^2/f)}$$

$$\partial_\chi V - \partial_\varphi V + \frac{d}{2} \partial_\chi \ln f V = f(\chi)^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} \frac{\partial_\chi R + \frac{d}{2} \partial_\chi \ln f R}{\partial_\varphi^2 V - K p^2/f + R(p^2/f)}$$

$$f^{-1} \partial_t K(\varphi, \chi) = 2f^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} P \dot{R}$$

$$f^{-1} \left\{ \partial_\chi K - \partial_\varphi K + \frac{d-2}{2} \partial_\chi \ln f K \right\} = 2f^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} P \left[\partial_\chi R + \frac{d}{2} \partial_\chi \ln f R \right]$$

$$P = -\frac{1}{2} \frac{\partial_\varphi^2 K}{f} Q_0^2 + \frac{\partial_\varphi K}{f} \left(2\partial_\varphi^3 V - \frac{2d+1}{d} \frac{\partial_\varphi K}{f} p^2 \right) Q_0^3$$

$$- \left[\left\{ \frac{4+d}{d} \frac{\partial_\varphi K}{f} p^2 - \partial_\varphi^3 V \right\} \left(\partial_{p^2} R - \frac{K}{f} \right) + \frac{2}{d} p^2 \partial_{p^2}^2 R \left(\frac{\partial_\varphi K}{f} p^2 - \partial_\varphi^3 V \right) \right] \left(\partial_\varphi^3 V - \frac{\partial_\varphi K}{f} p^2 \right) Q_0^4$$

$$- \frac{4}{d} p^2 \left(\partial_{p^2} R - \frac{K}{f} \right)^2 \left(\partial_\varphi^3 V - \frac{\partial_\varphi K}{f} p^2 \right)^2 Q_0^5 \quad Q_0 = \frac{1}{\partial_\varphi^2 V - K p^2/f + R}$$

Want momentum dependence so that:

$$\dot{R}(p^2/f) \propto \partial_x R + \frac{d}{2} \partial_x \ln f R$$

Cutoff

dimensionless cutoff profile

$$R(p^2/f) = -k^{d-\eta-\frac{d}{2}d_f} r\left(\frac{p^2}{k^{2-d_f} f}\right)$$

$$\implies \partial_x \ln f \dot{R} = (2 - d_f) \left[\partial_x R + \frac{d}{2} \partial_x \ln f R \right] - \eta \partial_x \ln f R$$

• $\eta = 0$ ✓

$$\eta \partial_x \ln f R \propto \partial_x R + \frac{d}{2} \partial_x \ln f R \implies f \frac{\partial R}{\partial f} \propto R$$

• $r(z) = \frac{1}{z^n}$ ✓

Linear PDEs

$$\alpha = 2 \left(1 - \frac{\eta}{d + 2n} \right) - d_f$$

$$\partial_{\chi} \ln f \partial_t V = \alpha \left(\frac{d}{2} \partial_{\chi} \ln f V + \partial_{\chi} V - \partial_{\varphi} V \right)$$

$$\partial_{\chi} \ln f \partial_t K = \alpha \left(\partial_{\chi} K - \partial_{\varphi} K + \frac{d-2}{2} \partial_{\chi} \ln f K \right)$$

Solutions

$$V(k, \varphi, \chi) = f(\chi)^{-\frac{d}{2}} \tilde{V}(\tilde{k}, \phi), \quad K(k, \varphi, \chi) = f(\chi)^{-\frac{d}{2}+1} \tilde{K}(\tilde{k}, \phi), \quad \tilde{k} = k f(\chi)^{\frac{1}{\alpha}}$$

Dimensionless variables

$$\hat{k} = \tilde{k}^{\frac{1}{1+d_f/\alpha}} = k^{\frac{1}{1+d_f/\alpha}} f(\chi)^{\frac{1}{\alpha+d_f}} \quad [\hat{k}] = 1$$

$$\tilde{V} = \hat{k}^d \hat{V}_{\hat{k}}(\hat{\phi}), \quad \tilde{K} = \hat{k}^{d-2-\eta} \hat{K}_{\hat{k}}(\hat{\phi}), \quad \phi = \hat{k}^{\frac{\eta}{2}} \hat{\phi}, \quad p = \hat{k} \hat{p}$$

msWIs, flow equations

$$\partial_t V(\varphi, \chi) = f(\chi)^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} \frac{\dot{R}(p^2/f)}{\partial_\varphi^2 V - K p^2/f + R(p^2/f)}$$

$$\partial_\chi V - \partial_\varphi V + \frac{d}{2} \partial_\chi \ln f V = f(\chi)^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} \frac{\partial_\chi R + \frac{d}{2} \partial_\chi \ln f R}{\partial_\varphi^2 V - K p^2/f + R(p^2/f)}$$

$$f^{-1} \partial_t K(\varphi, \chi) = 2f^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} P \dot{R}$$

$$f^{-1} \left\{ \partial_\chi K - \partial_\varphi K + \frac{d-2}{2} \partial_\chi \ln f K \right\} = 2f^{-\frac{d}{2}} \int_0^\infty dp p^{d-1} P \left[\partial_\chi R + \frac{d}{2} \partial_\chi \ln f R \right]$$

$$P = -\frac{1}{2} \frac{\partial_\varphi^2 K}{f} Q_0^2 + \frac{\partial_\varphi K}{f} \left(2\partial_\varphi^3 V - \frac{2d+1}{d} \frac{\partial_\varphi K}{f} p^2 \right) Q_0^3$$

$$- \left[\left\{ \frac{4+d}{d} \frac{\partial_\varphi K}{f} p^2 - \partial_\varphi^3 V \right\} \left(\partial_{p^2} R - \frac{K}{f} \right) + \frac{2}{d} p^2 \partial_{p^2}^2 R \left(\frac{\partial_\varphi K}{f} p^2 - \partial_\varphi^3 V \right) \right] \left(\partial_\varphi^3 V - \frac{\partial_\varphi K}{f} p^2 \right) Q_0^4$$

$$- \frac{4}{d} p^2 \left(\partial_{p^2} R - \frac{K}{f} \right)^2 \left(\partial_\varphi^3 V - \frac{\partial_\varphi K}{f} p^2 \right)^2 Q_0^5 \quad Q_0 = \frac{1}{\partial_\varphi^2 V - K p^2/f + R}$$

Background Independent flow equations

$$\partial_{\hat{t}} \hat{V} + d\hat{V} - \frac{\eta}{2} \hat{\phi} \hat{V}' = - (d - \eta + 2n) \int_0^{\infty} d\hat{p} \hat{p}^{d-1} \hat{Q}_0 r(\hat{p}^2)$$

• No χ : $\hat{V} \equiv \hat{V}_{\hat{k}}(\hat{\phi}), \quad \hat{K} \equiv \hat{K}_{\hat{k}}(\hat{\phi})!$

$$\partial_{\hat{t}} \hat{K} + (d - 2 - \eta) \hat{K} - \frac{\eta}{2} \hat{\phi} \hat{K}' = -2(d - \eta + 2n) \int_0^{\infty} d\hat{p} \hat{p}^{d-1} \hat{P} r(\hat{p}^2)$$

• No f ! $\hat{t} = \ln(\hat{k}/\mu)$

$$\hat{P} = -\frac{1}{2} \hat{K}'' \hat{Q}_0^2 + \hat{K}' \left(2\hat{V}'''' - \frac{2d+1}{d} \hat{K}' \hat{p}^2 \right) \hat{Q}_0^3$$

$$+ \left[\left\{ \frac{4+d}{d} \hat{K}' \hat{p}^2 - \hat{V}'''' \right\} \left(r'(\hat{p}^2) + \hat{K} \right) + \frac{2}{d} \hat{p}^2 r''(\hat{p}^2) \left(\hat{K}' \hat{p}^2 - \hat{V}'''' \right) \right] \left(\hat{V}'''' - \hat{K}' \hat{p}^2 \right) \hat{Q}_0^4$$

$$- \frac{4}{d} \hat{p}^2 \left(r'(\hat{p}^2) + \hat{K} \right)^2 \left(\hat{V}'''' - \hat{K}' \hat{p}^2 \right)^2 \hat{Q}_0^5$$

$$\hat{Q}_0 = \frac{1}{\hat{V}'' - \hat{K} \hat{p}^2 - r(\hat{p}^2)}$$

Two notions of fixed point

$$V(\varphi, \chi) = k^{d - \frac{d}{2}d_f} \bar{V}_k(\bar{\varphi}, \bar{\chi}), \quad \varphi = k^{\eta/2} \bar{\varphi}, \quad \chi = k^{\eta/2} \bar{\chi}$$

$$= f(\chi)^{-\frac{d}{2}} \hat{k}^d \hat{V}_{\hat{k}}(\hat{\phi})$$

$$\partial_t \bar{V} = (k^{-d_f} f)^{\frac{d\eta}{2(d-\eta+2n)}} \left\{ \frac{1}{\alpha + d_f} \left(\alpha + \frac{\eta}{2} \chi \partial_\chi \ln f \right) \partial_{\hat{t}} \hat{V}(\hat{\phi}) \right.$$

$$\left. + \left(\frac{\eta}{2} \chi \partial_\chi \ln f - d_f \right) \left(\frac{d\eta}{2(d-\eta+2n)} \hat{V} - \frac{\eta}{2(\alpha + d_f)} \hat{\phi} \hat{V}' \right) \right\}$$

If $\eta \chi \partial_\chi \ln f$ varies with χ then

$$\partial_t \bar{V} = 0 \quad \implies \quad \partial_{\hat{t}} \hat{V} = 0 \quad \text{and}$$

$$\frac{d\eta}{2(d-\eta+2n)} \hat{V} - \frac{\eta}{2(\alpha + d_f)} \hat{\phi} \hat{V}' = 0 \quad \implies \quad \hat{V} \propto \hat{\phi}^{\frac{2d}{d+2n}}$$



Two notions of fixed point

$$V(\varphi, \chi) = k^{d - \frac{d}{2} d_f} \bar{V}_k(\bar{\varphi}, \bar{\chi}), \quad \varphi = k^{\eta/2} \bar{\varphi}, \quad \chi = k^{\eta/2} \bar{\chi}$$

$$= f(\chi)^{-\frac{d}{2}} \hat{k}^d \hat{V}_{\hat{k}}(\hat{\phi})$$

$$\partial_t \bar{V} = (k^{-d_f} f)^{\frac{d\eta}{2(d-\eta+2n)}} \left\{ \frac{1}{\alpha + d_f} \left(\alpha + \frac{\eta}{2} \chi \partial_\chi \ln f \right) \partial_{\hat{t}} \hat{V}(\hat{\phi}) \right.$$

$$\left. + \left(\frac{\eta}{2} \chi \partial_\chi \ln f - d_f \right) \left(\frac{d\eta}{2(d-\eta+2n)} \hat{V} - \frac{\eta}{2(\alpha + d_f)} \hat{\phi} \hat{V}' \right) \right\}$$

If $\eta \chi \partial_\chi \ln f$ varies with χ then $\partial_t \bar{V} \neq 0$ (!)

$\partial_t \bar{V} = 0 \implies \partial_{\hat{t}} \hat{V} = 0$ **k-fixed points exist if:** **\hat{k} -fixed points coincide if:**

$$\frac{d\eta}{2(d-\eta+2n)} \hat{V} - \frac{\eta}{2(\alpha + d_f)} \hat{\phi} \hat{V}' = 0 \implies \hat{V} d_{\hat{\phi}} \hat{\phi}^{\frac{2d}{d+2n}} = 0$$

- $f(\chi) \propto \chi^\gamma$
- ~~$d_f = \gamma \frac{\eta}{2}$~~

Background independence is incompatible
with k -fixed points!

$$f(\phi) = k_0^{d_f} F \left(\frac{\phi}{k_0^{\eta/2}} \right)$$

\implies

$$\bar{f}(\bar{\chi}) = F \left(\bar{\chi} e^{\frac{\eta}{2}(t-t_0)} \right) e^{d_f(t_0-t)}$$

k -fixed points exist if: \hat{k} -fixed points coincide if:

- $\eta = 0$

- $d_f = 0$

- $f(\chi) \propto \chi^\gamma$

- $d_f = \gamma \frac{\eta}{2}$

In conformally truncated gravity...

- Background independent notion of scale \hat{k} exists
- which is also independent of the form of $f(\phi)$
- Generally k -fixed points are **incompatible** with background independence (i.e. the msWI)

...exceptions depending on form (running) of $f(\chi)$

\hat{k} -fixed points exist independently of all this...

...at a simpler deeper level

