

Detecting Higgsino Dark Matter at Colliders

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with H.Baer, V.Barger, D.Mickelson and X.Tata

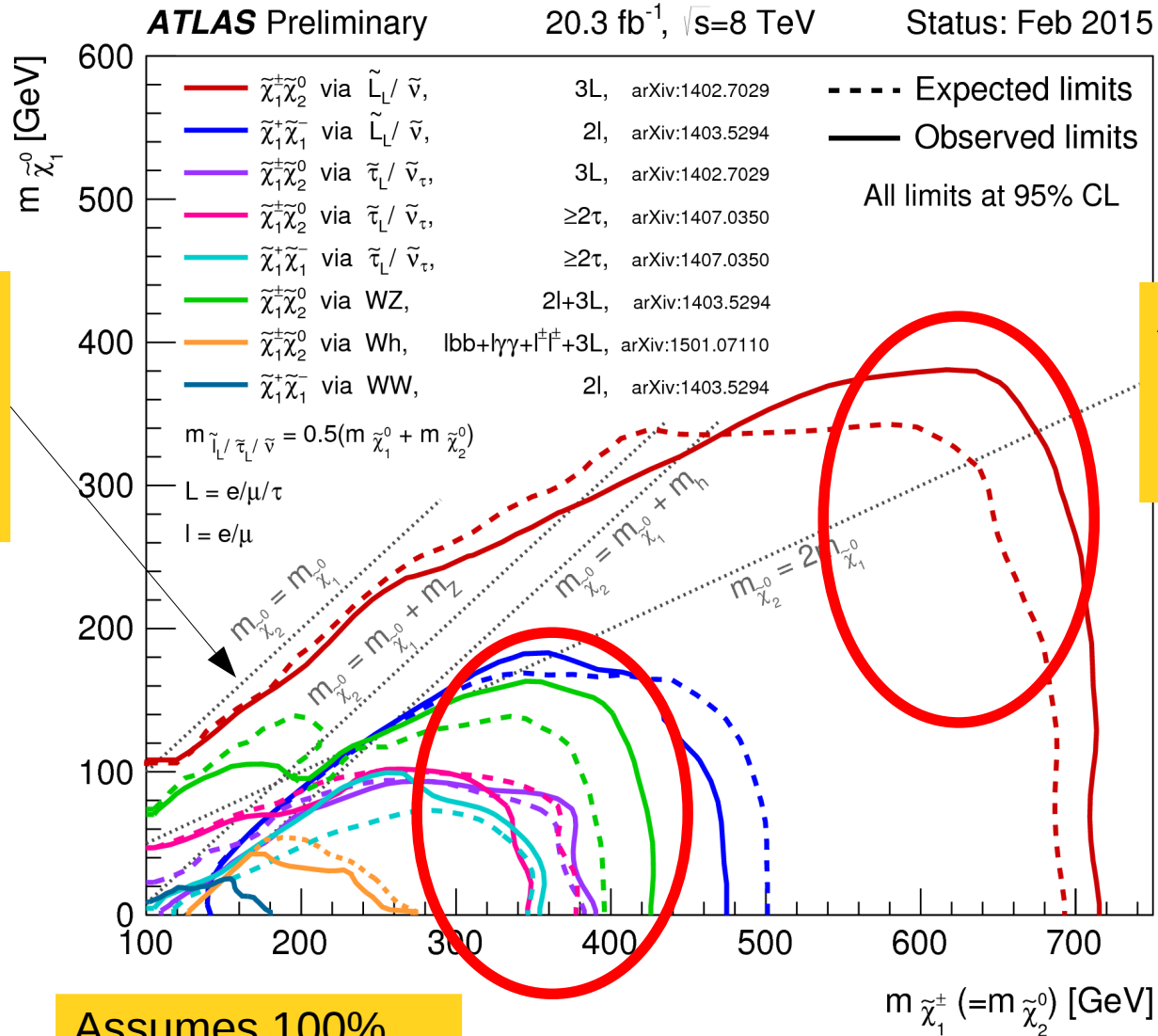
SUSY and Dark Matter

- SUSY has “natural” Dark Matter candidate, once R-parity is imposed: neutralino LSP
- Pure states have annih. cross section very large (Wino & Higgsino => right RD at $\sim 2\text{TeV}$) or very small (Bino)
- Many ways to achieve right RD with sub-TeV Higgsinos:
 - Non-thermal from decays of moduli
[Allahverdi, Dutta, Sinha]
 - Higgsinos + axions
 - Higgsinos/axinos/axions in SUSY KSVZ model
[Bae, Baer, Lessa]
 - ...
- Can we discover sub-TeV higgsino-like neutralino at colliders?

Motivation 2: SUSY and Naturalness

- In MSSM naturalness requires light higgsinos
- For $m_h \sim 125 \text{ GeV}$ and $\Delta_{\text{EW}} < 30$ (see talks by H.Baer and X.Tata):
 - ♦ $\mu \sim 100\text{-}300 \text{ GeV}$
 - ♦ $\text{stop}_1 \sim 1\text{-}2 \text{ TeV}$, $\text{stop}_2 = \text{sbottom}_1 \sim 2\text{-}4 \text{ TeV}$, highly mixed by large A_t
 - ♦ gluino $\sim 1\text{-}5 \text{ TeV}$
 - ♦ bino/wino - upto $\sim 8 \text{ TeV}$
 - ♦ 1st/2nd generation squarks $\sim 1\text{-}10 \text{ TeV}$
 - ♦ sleptons - upto $\sim 30 \text{ TeV}$
- Higgsinos are the only ones fully accessible at LHC.
Can we detect them?

Current EW-ino searches at LHC



Weak constraints on compressed spectrum

Assumes wino NLSP and light sleptons \rightarrow lepton rich final states.

Assumes 100% BF to WZ

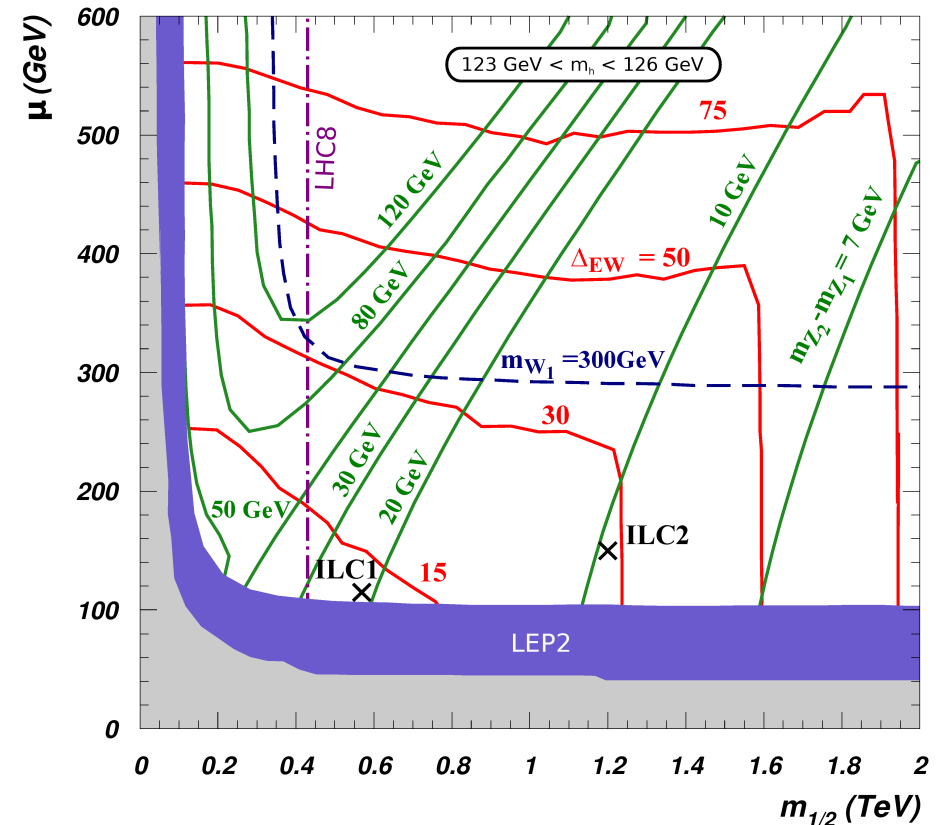
Higgsino mass splittings

- For $|\mu| \ll M_{1,2}$ lightest states $\tilde{Z}_1, \tilde{Z}_2, \tilde{W}_1^\pm$ are higgsino-like and have small mass gap $\leq 30 \text{ GeV}$

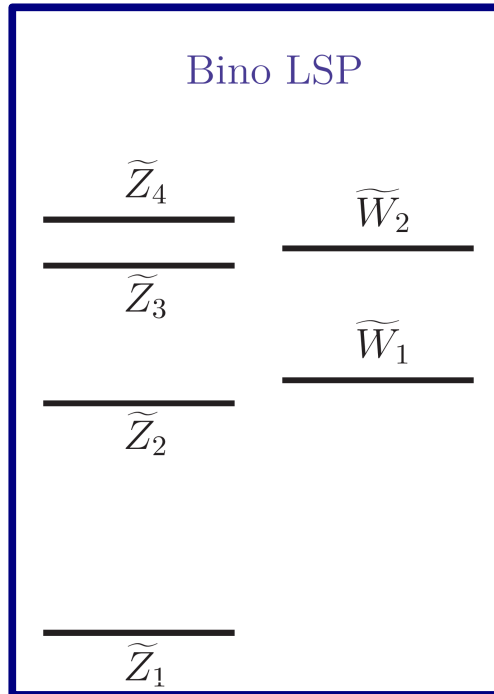
$$m_{\tilde{Z}_2} - m_{\tilde{Z}_1} \simeq M_Z^2 \left(\frac{s_W^2}{M_1} + \frac{c_W^2}{M_2} \right)$$

$$m_{\tilde{W}_1} - m_{\tilde{Z}_1} \simeq \frac{M_Z^2}{2} \left(\frac{s_W^2}{M_1} + \frac{c_W^2}{M_2} \right)$$

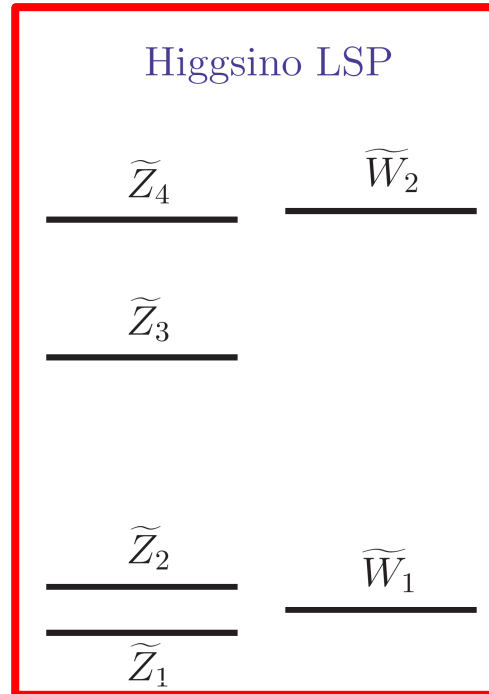
NUHM2: $m_0=5 \text{ TeV}, \tan\beta=15, A_0=-1.6m_0, m_A=1\text{TeV}, m_t=173.2 \text{ GeV}$



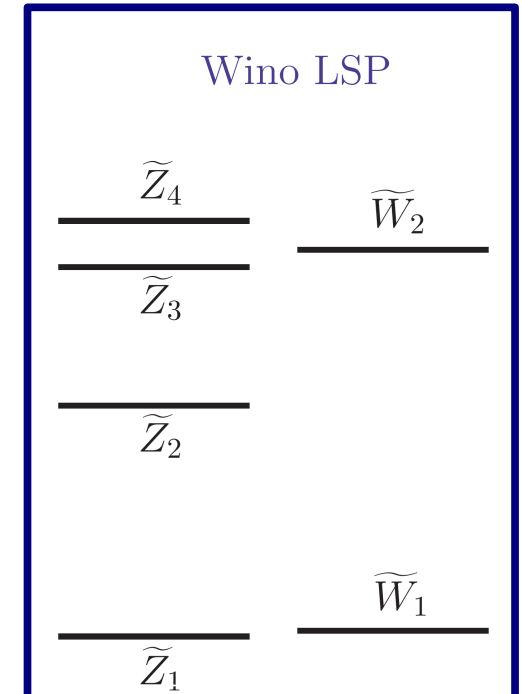
EW-ino spectrum



canonical case
mSUGRA/CMSSM



Natural SUSY



AMSB

- Higgsino mass splitting is small enough to see decay products, but too large to see disappearing charge track

Mono-jets and mono-photons

arXiv: 1401.1162

- Higgsinos have compressed spectrum with mass gap 10-30 GeV - only soft visible energy ==> higgsinos mostly appear as MET.

- $pp \rightarrow \tilde{Z}_{1,2}\tilde{Z}_{1,2}/\tilde{Z}_{1,2}\tilde{W}_1/\tilde{W}_1\tilde{W}_1 + (j \text{ or } \gamma)$

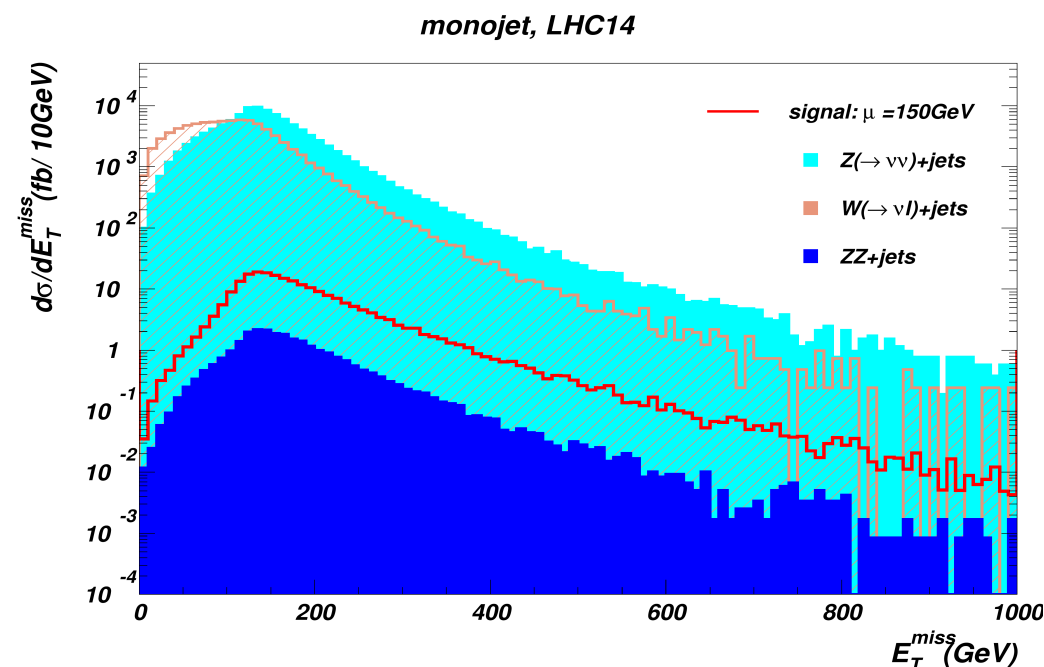
- Contact interaction, used in ATLAS/CMS search, not applicable:

mediator mass $M_Z \ll \sqrt{\hat{s}} \sim p_T(\text{jet}) + E_T^{\text{miss}}$

leading to extra $1/s$ suppression for ME

- Signal has same shape as BG and $S/B \approx 1\%$.

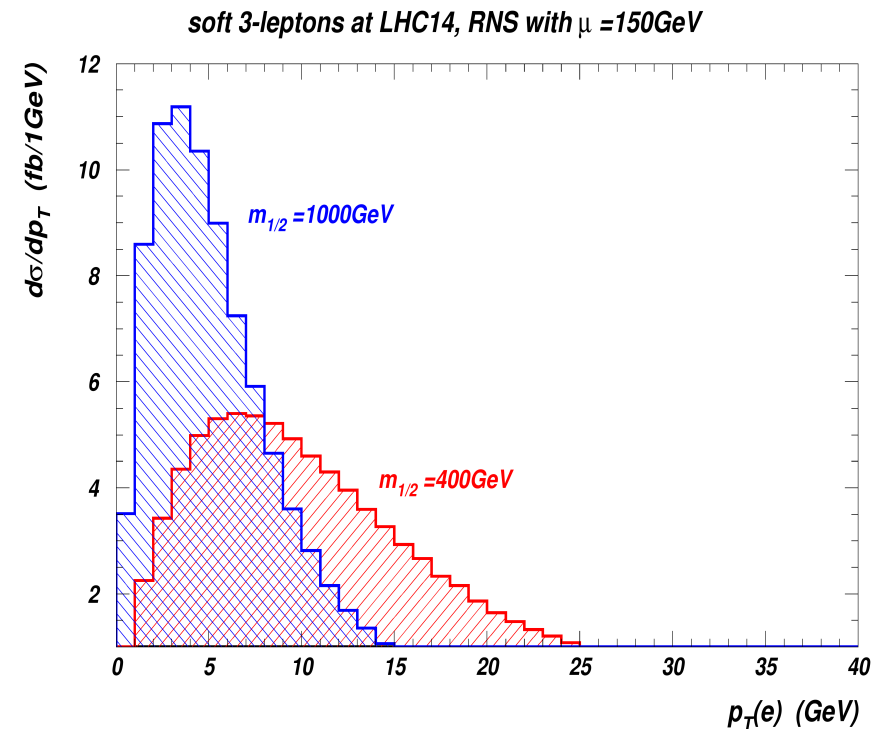
Detection is very challenging!



Soft trileptons

arXiv: 1302.5816

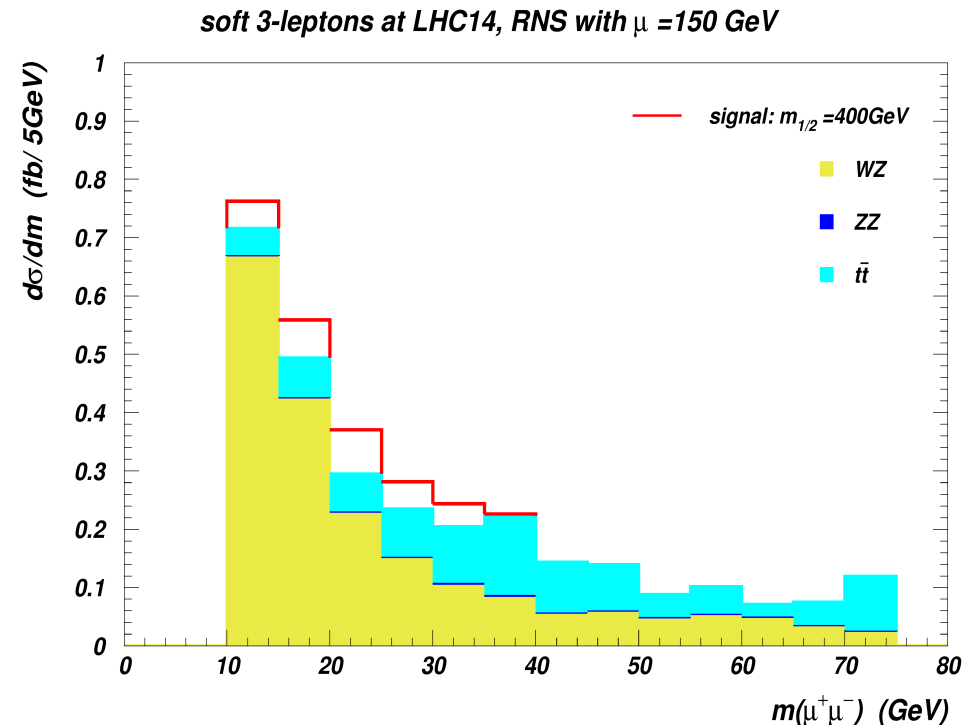
- Mass gap between higgsinos is $< 30 \text{ GeV} \Rightarrow$ very soft leptons for higgsino pair production. But ATLAS/CMS can detect soft leptons $p_T(\mu) > 5 \text{ GeV}$, $p_T(e) > 10 \text{ GeV}$
- $pp \longrightarrow \widetilde{W}_1 \widetilde{Z}_2 \longrightarrow (e\nu_e \widetilde{Z}_1) + (\mu^+ \mu^- \widetilde{Z}_1)$
- For $m_{1/2} < 400 - 500 \text{ GeV}$ and $\mu = 150 \text{ GeV}$ most e pass trigger threshold



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- A shape analysis may allow to claim a signal with high luminosity - confirmation channel



Mono-jets plus soft dilepton

arXiv: 1409.7058

- Soft leptons can provide additional handles, improving monojet detection prospects [Giudice, T.Han, Wang and L.T.Wang PRD'10]
- $pp \rightarrow \tilde{Z}_1 \tilde{Z}_2 / \tilde{Z}_2 \tilde{Z}_2 / \tilde{Z}_{1,2} \tilde{W}_1 / \tilde{W}_1^+ \tilde{W}_1^- + j$
 $\tilde{W}_1 \rightarrow \ell \nu \tilde{Z}_1, \tilde{Z}_2 \rightarrow \ell^+ \ell^- \tilde{Z}_1$
- Production cross section $\sim (10-100)$ fb times lepton c BF.
- Signals:
 - $1l+j+\text{MET}$ from $\tilde{W}_1 \tilde{Z}_1 j$ - obscured by large Wj and $t\bar{t}$
 - $3l+j+\text{MET}$ from $\tilde{W}_1 \tilde{Z}_2 j$ - severely rate limited
 - **$2l+j+\text{MET}$ is the most promising**
- Recent works: Schwaller and Zurita (2014), Z.Han et al (2014)

Backgrounds

arXiv: 1409.7058

- $(Z/\gamma^* \rightarrow \ell\bar{\ell}) + j$ removed by MET cut
- $(Z/\gamma^* \rightarrow \tau\bar{\tau}) + j$, W^+W^-j , $(Z \rightarrow \nu\bar{\nu}) + (Z/\gamma^* \rightarrow \ell\bar{\ell}/\tau\bar{\tau}) + j$
- ttbar and single-top, where leptons come from W and b
- $(Z \rightarrow \nu\bar{\nu}) + b\bar{b} + j$, where leptons come from b quarks
- $(W \rightarrow \ell/\tau + \nu) + (Z/\gamma^* \rightarrow \ell\bar{\ell}/\tau\bar{\tau}) + j \rightarrow \ell\ell\ell'j + E_T^{\text{miss}}$

	$\tau\bar{\tau}j$	$t\bar{t}$	WWj	$Z\gamma^*j$	$W\gamma^*j$	tW	tq	Signal
before cuts	20660	953600	460.53	5.59	171.83	62330	68080	398.7
1-event level	0.011	0.0954		0.000016		0.0312	0.068	
$N_b = 0$	20395	386450	455.0	5.54	169.7	37270	41550	393.2
$N_j = 1$	8020.4	33470	290.9	4.24	100.2	6945	7752	236.8
$p_T(j_1) > 100$ GeV, $ \eta(j_1) < 2.5$	2961	15590	162.7	1.95	58.8	2351	3646	133.7
$E_T^{\text{miss}} > 100$ GeV	935.0	12950	82.28	1.54	19.16	1700	2257	120.6
$N_\ell \geq 2$	171.3	880.8	52.26	0.56	15.0	72.79	19.75	2.20

Suppressing Z+j

arXiv: 1409.7058

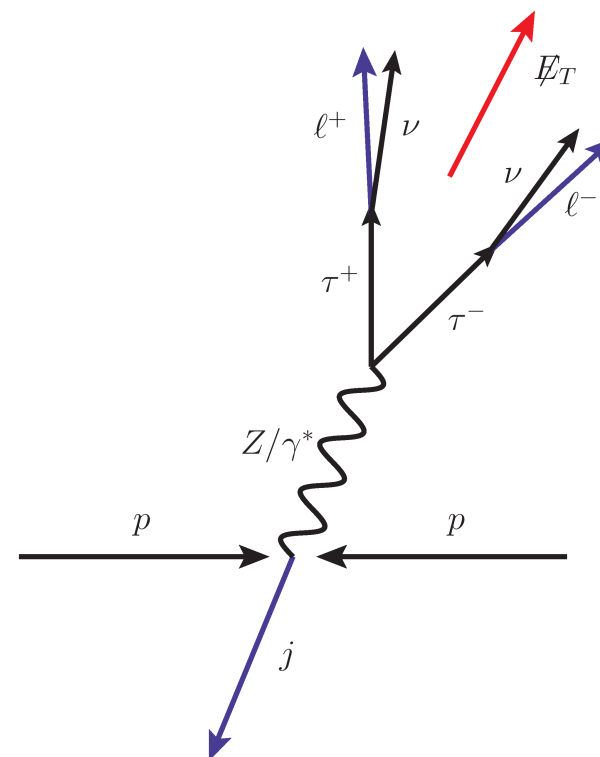
- $(Z/\gamma^* \rightarrow \tau\bar{\tau}) + j$ allows reconstruction of Z mass in the collinear approximation [Ellis, Hinchliffe, Soldate, van der Bij NPB'88]

- Z is boosted and taus decay products are collinear

$$-\vec{p}_T(j) = (1 + \xi_1)\vec{p}_T(\ell_1) + (1 + \xi_2)\vec{p}_T(\ell_2)$$

- Solve for $\xi_{1,2}$

$$m_{\tau\tau}^2 = (1 + \xi_1)(1 + \xi_2)m_{\ell\ell}^2$$



Suppressing Zj

arXiv: 1409.7058

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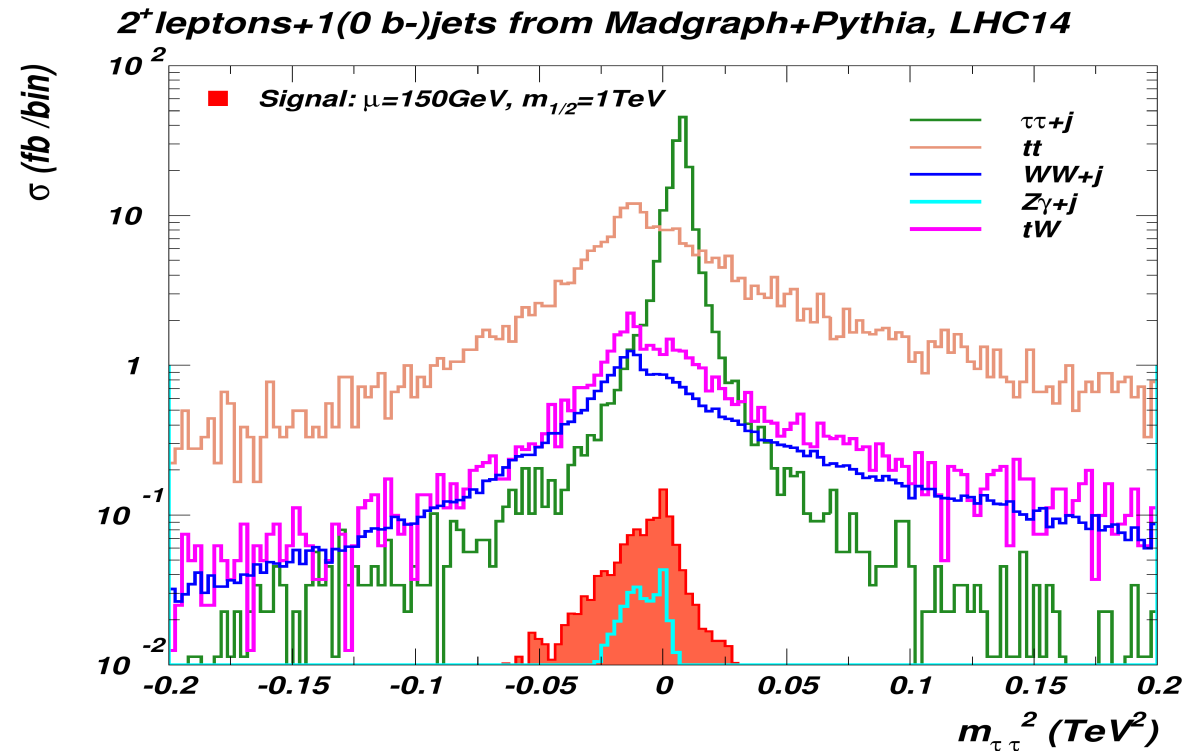
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- Solve for $\xi_{1,2}$

$$m_{\tau\tau}^2 = (1 + \xi_1)(1 + \xi_2)m_{\ell\ell}^2$$

- Require $m_{\tau\tau}^2 < 0$



Enhancing the Signal

- Require OS/SF leptons, since most of the signal come from

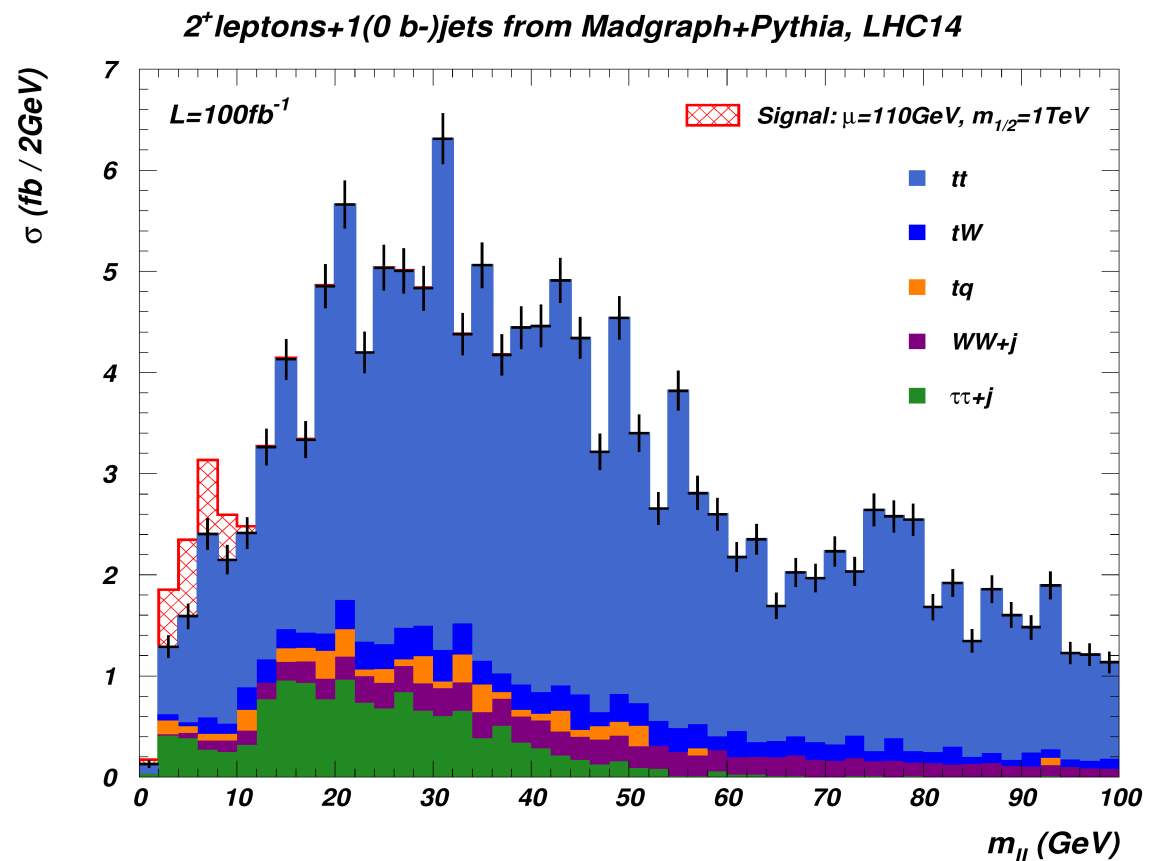
$$\tilde{Z}_2 \rightarrow \ell^+ \ell^- \tilde{Z}_1$$

- Dilepton mass is bounded by $\tilde{Z}_2 - \tilde{Z}_1$ mass gap.

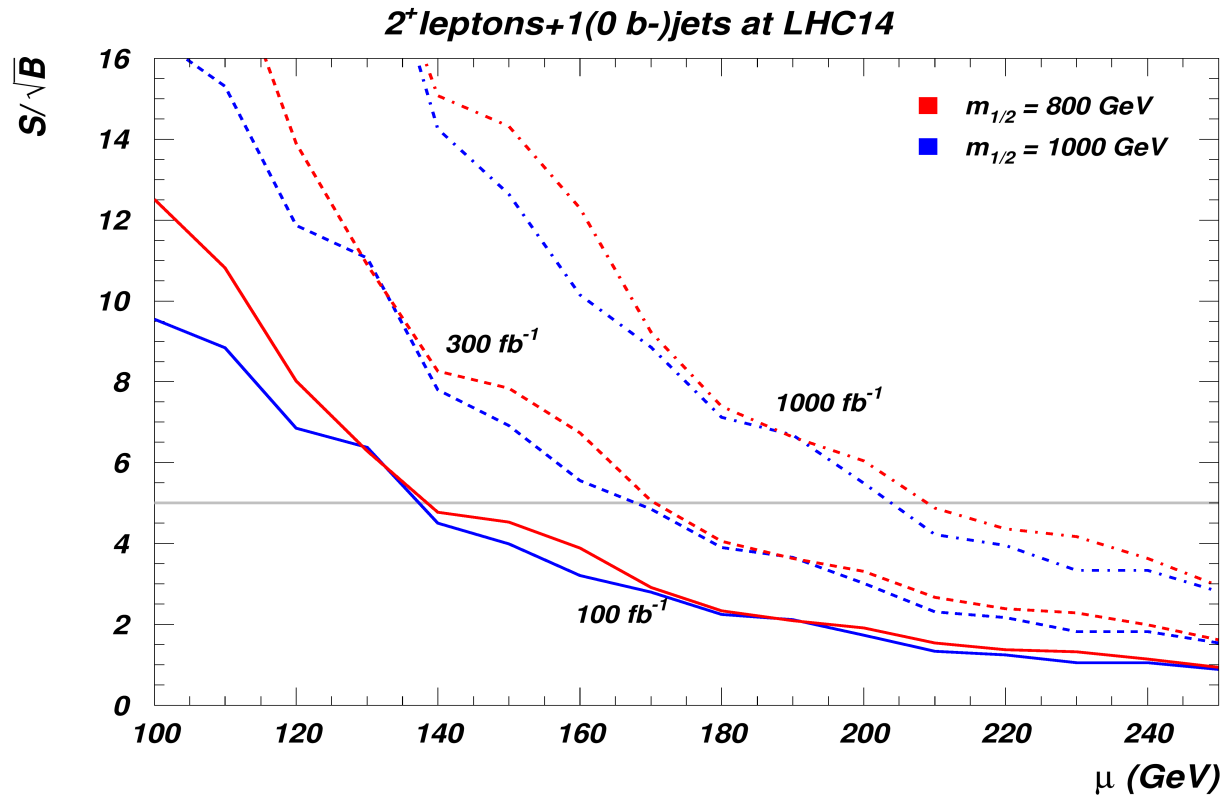
Require

$$0 \leq m_{\ell\ell} \leq m^{\text{cut}}$$

to maximize significance



Reach via cut-and-count



Note: requires control of systematic errors on BG normalization at better than 5% to claim detection.

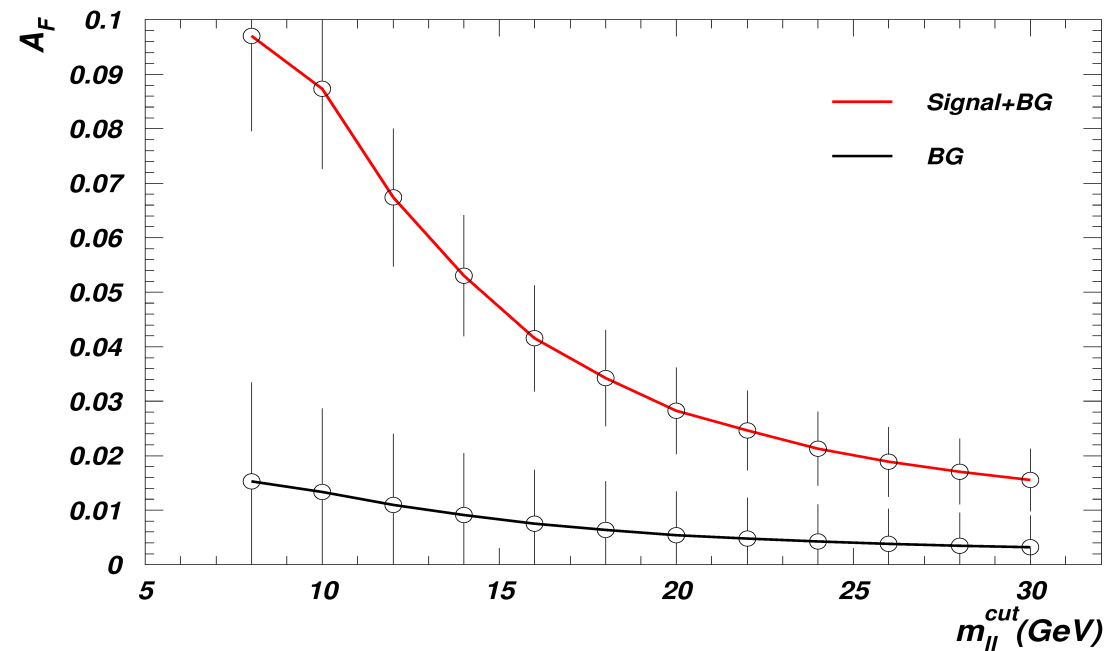
Lepton flavor asymmetry

- Signal mostly has SF dileptons, while BG has equal number of SF and OF \Rightarrow use dilepton flavor asymmetry

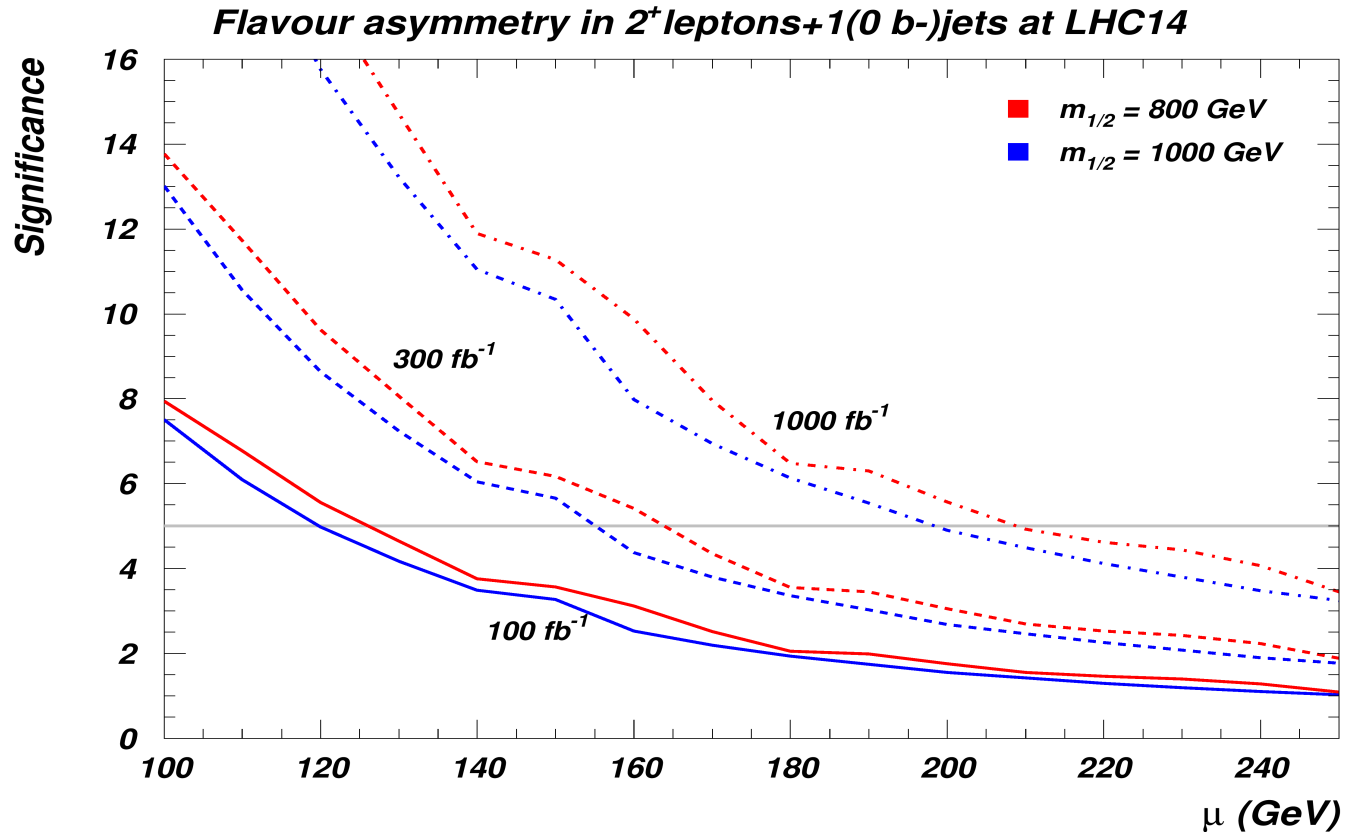
$$\mathcal{A}_F = \frac{N(SF) - N(OF)}{N(SF) + N(OF)}$$

- Insensitive to BG normalization!

$2^+ \text{ leptons} + 1(0 \text{ } b\text{-})\text{jets at LHC14 for } 300 \text{ fb}^{-1}$



Reach via lepton flavor asymmetry



- Smaller by 10-15 GeV reach due to $\sim\sqrt{2}$ smaller significance, but not sensitive to BG systematics.

Signatures at ILC

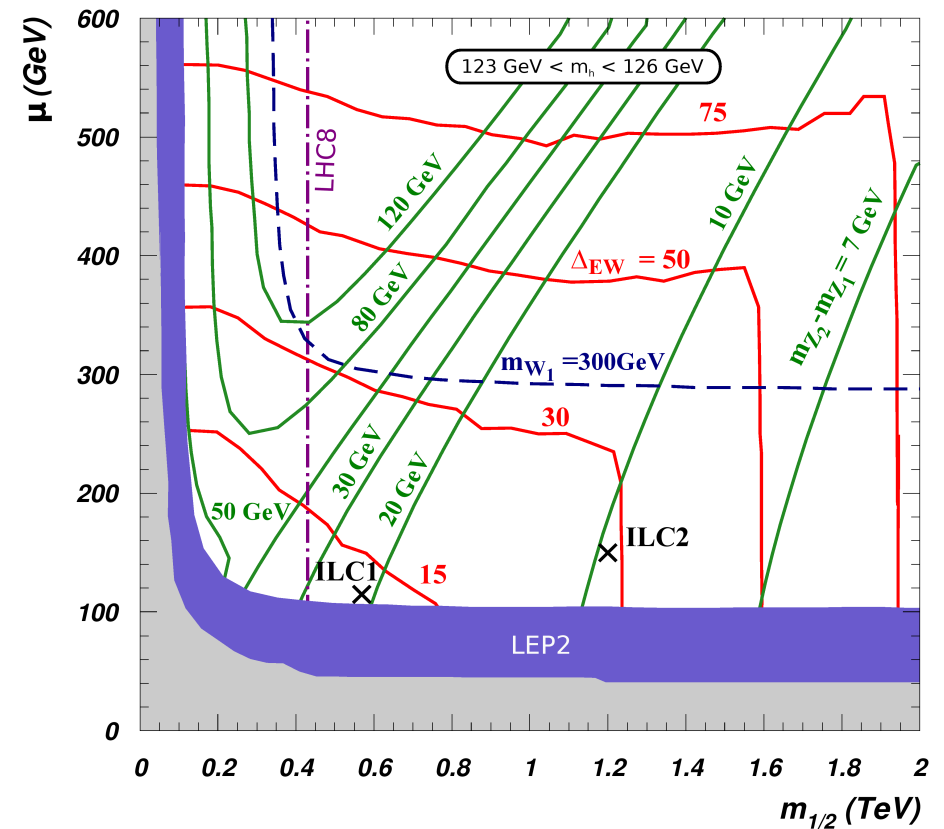
arXiv: 1404.7510

- Detection of light higgsinos is very challenging at LHC due to compressed spectrum. ILC with ~ 600 GeV will be able to produce charginos covering region with

$$\Delta_{EW} < 30$$

- Higgsino production ~ 100 - 1000 fb is 5-10 larger than Zh
 \Rightarrow Higgsino factory!
- Beam polarization allows to establish higgsino nature of inos
- Consider two cases:
 ILC1 with mass gap ~ 20 GeV,
 ILC2 with near minimal mass gap ~ 10 GeV

NUHM2: $m_0=5$ TeV, $\tan\beta=15$, $A_0=-1.6m_0$, $m_A=1$ TeV, $m_t=173.2$ GeV



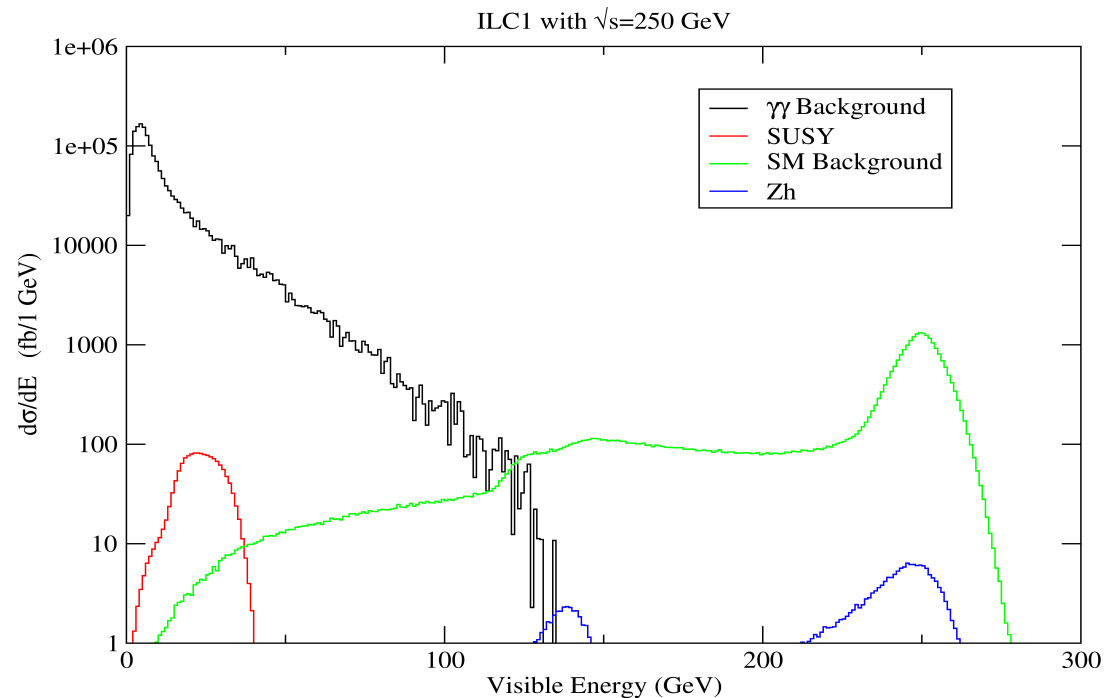
Benchmark ILC1 at 250GeV

arXiv: 1404.7510

- $m_{\tilde{Z}_1} = 102.7 \text{ GeV}$, $m_{\tilde{Z}_2} = 124 \text{ GeV}$, $m_{\tilde{W}_1} = 117.3 \text{ GeV}$
- 2->2 SM background mostly has larger visible energy E_{vis}
- 2->4 "two-photon" BG are back-to-back with essentially no MET
- $e^+e^- \rightarrow \tilde{W}_1\tilde{W}_1 \rightarrow q\bar{q}\tilde{Z}_1 + l\nu\tilde{Z}_1$

leads to 2j+1l+MET with
6.43 fb after cuts;
BG $\approx 0.018 \text{ fb}$.

- $e^+e^- \rightarrow \tilde{Z}_1\tilde{Z}_2 \rightarrow \tilde{Z}_1 + l\bar{l}\tilde{Z}_1$
leads to acoplanar dileptons
with 19.55 fb after cuts and
BG = 0.44 fb. Need polarized
beam to reduce WW BG.



ILC1 mass measurements

arXiv: 1404.7510

- For chargino pairs, $E(jj)$ upper endpoint is sensitive to \tilde{W}_1 and \tilde{Z}_1 masses. Fit shapes of theory samples to "data" \rightarrow masses to 2-3%

$$m_{\tilde{W}_1} = 117.8 \pm 2.8 \text{ GeV } (1\sigma)$$

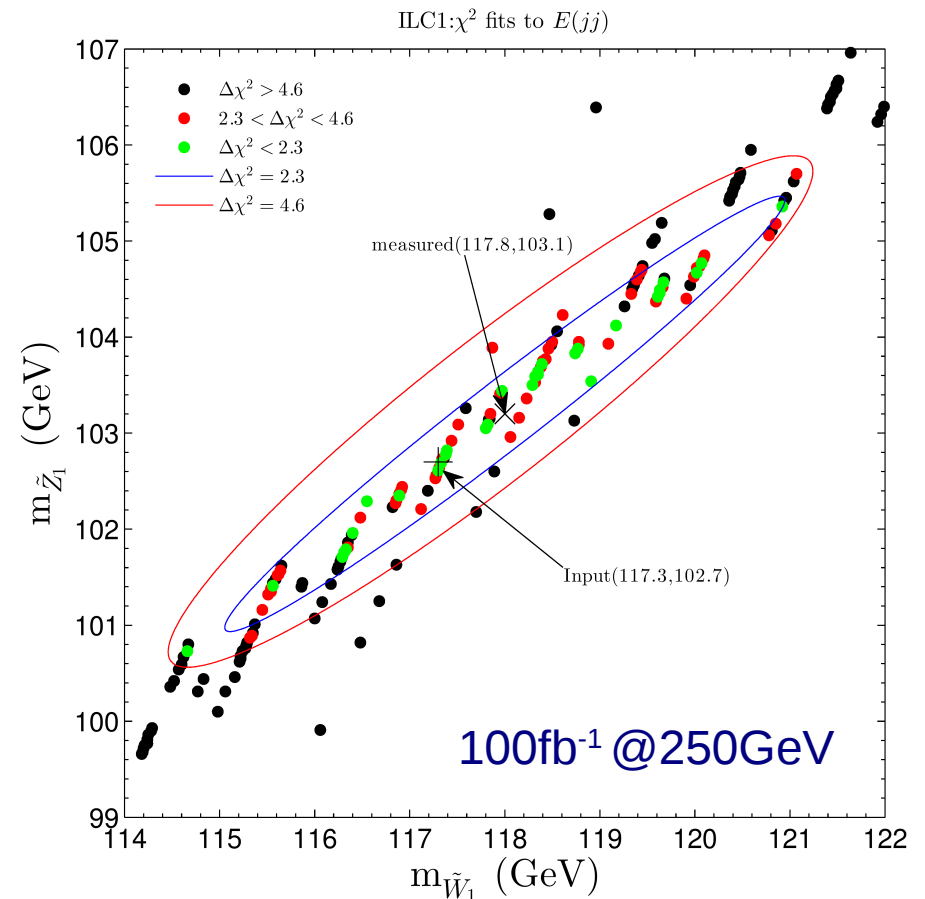
$$m_{\tilde{Z}_1} = 103.1 \pm 2.2 \text{ GeV } (1\sigma)$$

- For neutralino pairs, $m(l+l-)$ shape fit gives mass gap to 1%

$$m_{\tilde{Z}_2} - m_{\tilde{Z}_1} = 21.0 \pm 0.2 \text{ GeV } (1\sigma)$$

- Energy of daughter leptons depends on \tilde{Z}_2 boost \Rightarrow fit to $E(l+l-)$ shape gives

$$m_{\tilde{Z}_2} = 123.7 \pm 0.2 \text{ GeV } (1\sigma)$$



Benchmark ILC2 at 340GeV

arXiv: 1404.7510

- $m_{\tilde{Z}_1} = 148 \text{ GeV}$, $m_{\tilde{Z}_2} = 157.8 \text{ GeV}$, $m_{\tilde{W}_1} = 158.3 \text{ GeV}$
- $e^+e^- \rightarrow \tilde{W}_1\tilde{W}_1 \rightarrow q\bar{q}\tilde{Z}_1 + l\nu\tilde{Z}_1$
 - $2j+1l$ -- signal resolution is difficult due to smaller mass gap
 - $1j+1l+\text{MET}$ is 7.1 fb with BG = 2.8 fb: discovery with a few fb⁻¹

- $e^+e^- \rightarrow \tilde{Z}_1\tilde{Z}_2 \rightarrow \tilde{Z}_1 + l\bar{l}\tilde{Z}_1$

signal of 2.6 fb with BG = 0.15 fb after cuts

- $\tilde{Z}_2 - \tilde{Z}_1$ mass gap measured at 2% from $m(l^+l^-)$
- \tilde{Z}_2 mass measured from $E(l^+l^-)$ with sub-GeV precision

Conclusions

- Small μ is necessary (but not sufficient) condition for naturalness leading to light higgsino-like chargino and neutralinos with compressed spectrum
- Light Higgsinos are difficult to detect at LHC due to low visible energy release from their decays.
- Mono-jets and mono-photons are inefficient probes for higgsinos.
- Soft trileptons from compressed higgsinos can be used as confirmatory channel.
- Combining soft leptons with monojet in $2l+j+MET$ channel allows LHC to achieve 5σ for higgsinos < 170 GeV with $S/B \sim 8.5\%$ with 300 /fb, and extend over 200 GeV for 1000 /fb. Systematics can be controlled with use of lepton flavor asymmetry.
- ILC is necessary to fully probe light higgsinos. ILC can discover and measure higgsino masses at 2-3% even for small mass gap ~ 10 GeV.

EW Fine-tuning

arXiv: 1207.3343,
1404.1386

- Minimization condition for higgs scalar potential (1-loop)

$$\frac{M_Z^2}{2} = \frac{(m_{H_d}^2 + \Sigma_d^d) - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2$$

- Naturalness requires all terms on RHS comparable to LHS
⇒ fine-tuning parameter

$$\Delta_{EW} \equiv \max \left(\frac{m_{H_u}^2}{M_Z^2/2} \frac{\tan^2 \beta}{\tan^2 \beta - 1}, \frac{\Sigma_u^u}{M_Z^2/2} \frac{\tan^2 \beta}{\tan^2 \beta - 1}, \frac{\mu^2}{M_Z^2/2}, \dots \right)$$

- $\Delta_{EW}=100$ corresponds to 1% EWFT, $\Delta_{EW}=30$ is ~3% EWFT
- Limited value of $\Delta_{EW} \Rightarrow$ upper limit on $m_{H_u}^2$ and μ^2
- Δ_{EW} is not a measure but a bound on FT, it measures minimal FT present in given spectrum