

*ATLAS Standard Model Workshop,  
Lessons from Run 1 and Preparation for Run 2  
Annecy, February 2014*

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# Thoughts on QCD for Run 2



Gavin Salam (CERN)

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# Why SM studies?

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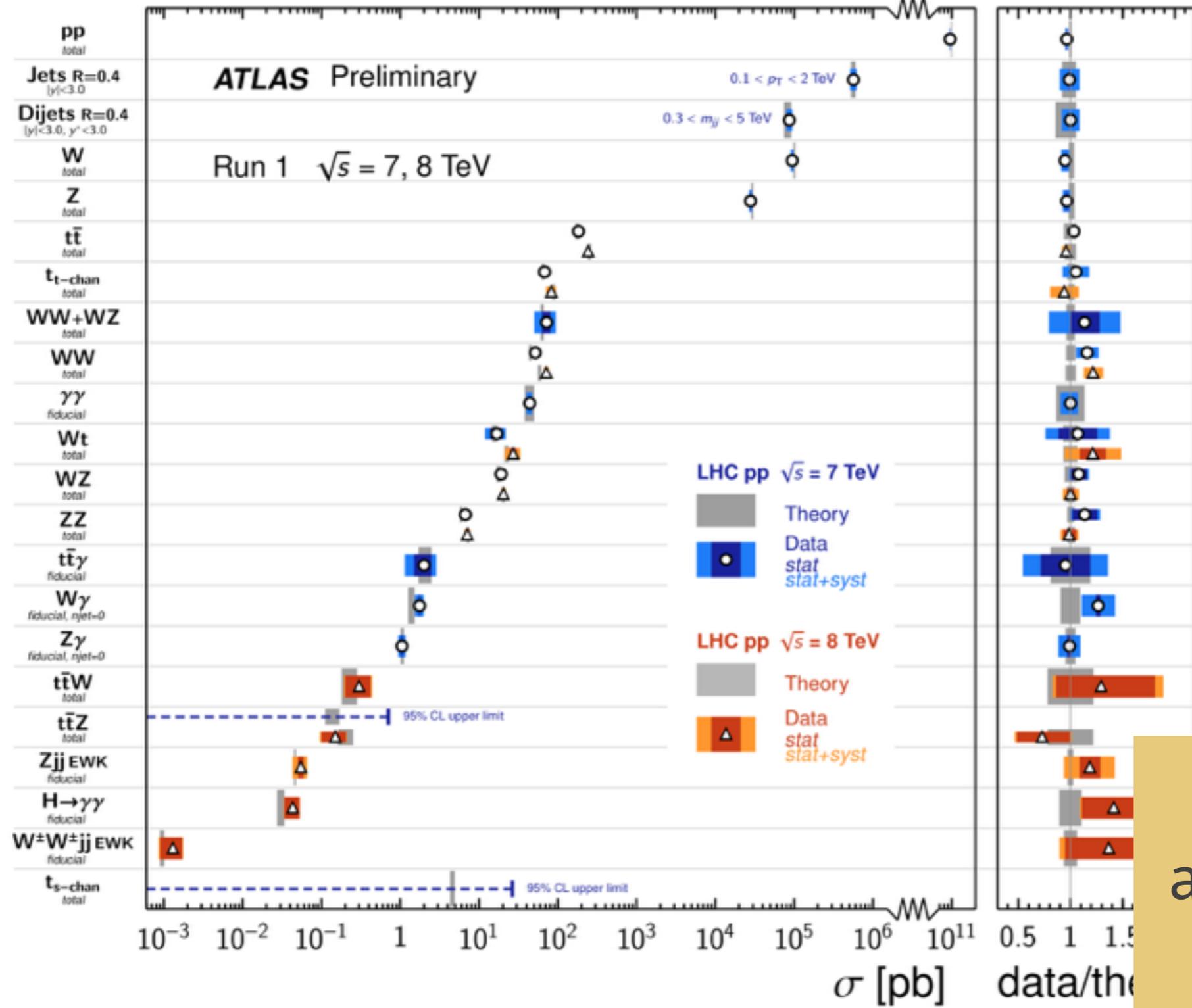
- ❖ Test and extend our understanding of QCD “tools” (MC’s, etc.) – **just how good are the tools?**
- ❖ **Measure fundamental constants** (e.g.  $\alpha_s$ ,  $M_W$ ) and fundamental non-perturbative inputs (**PDFs**) – such measurements can have decade-long staying power.
- ❖ Demonstrate **new physical effects** (cf. what condensed-matter physicists do all the time)

# 20 questions

- [ ] VV tensions, jet vetoes, fiducial cross sections
- [X] Are  $\times 2$  scale variations sufficient? HT/2 v. MINLO?
- [X] Z+j NLO discrepancies; interpretation and prospects for NNLO
- [ ] Practical use of NNLO (not ntuples) and approx NNLO (not threshold); + when will NNLO come
- [ ] Will NNLO V+jets be competitive for  $\alpha_s$ ?
- [ ] Best scale, e.g. with DeltaPhi or DeltaR or jet-rate measurements (MINLO v etc.)
- [ ] jet substructure, with %-level systematics control?
- [ ] fixed-order v. MC: agreement, hadronisation, etc.
- [X] jet-radius dependence: for which measurements, how small in R?
- [ ] EW corrections at  $\sqrt{s}=13\text{TeV}$ , scale 1 TeV; state-of-knowledge
- [ ] jet flavour, e.g. for inclusive X-scts; based on flav-kt?
- [ ] Do TH uncertainties cancel in 7/8/13 TeV x-sct ratios
- [ ] Regions of phase-space to discriminate BFKL/DGLAP
- [ ] Best way to present meas. sensitive to # of quark&gluon-initiated jets
- [ ] Photon isolation: study Frixione isolation?
- [ ] Low-pt photons limited by frag.fns; can they be improved
- [ ] Any new variables for binning x-scts? In order to get smaller uncertainties
- [ ] Any TH reason to avoid small  $\Delta R_{\gamma,\text{jet}}$
- [X] Sym/Asymmetric cuts to avoid divergences in calcs
- [ ] Correlations on scale uncertainties between phase-space regions

# Standard Model Production Cross Section Measurements

Status: July 2014



$\int \mathcal{L} dt$ [fb <sup>-1</sup> ]	Reference
$8 \times 10^{-8}$	ATLAS-CONF-2014-040
4.5	ATLAS-STDM-2013-11
4.5	JHEP 05, 059 (2014)
0.035	PRD 85, 072004 (2012)
0.035	PRD 85, 072004 (2012)
4.6	arXiv:1406.5375 [hep-ex]
20.3	arXiv:1406.5375 [hep-ex]
4.6	arXiv:1406.7844 [hep-ex]
20.3	ATLAS-CONF-2014-007
4.7	ATLAS-CONF-2012-157
4.6	PRD 87, 112001 (2013)
20.3	ATLAS-CONF-2014-033
4.9	JHEP 01, 086 (2013)
2.0	PLB 716, 142-159 (2012)
20.3	ATLAS-CONF-2013-100
4.6	EPJC 72, 2173 (2012)
13.0	ATLAS-CONF-2013-021
4.6	JHEP 03, 128 (2013)
20.3	ATLAS-CONF-2013-020
1.0	ATLAS-CONF-2011-153
4.6	PRD 87, 112003 (2013)
4.6	PRD 87, 112003 (2013)
20.3	ATLAS-CONF-2014-038
4.7	ATLAS-CONF-2012-126
20.3	ATLAS-CONF-2014-033

Beautiful agreement for so many measurements

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# Z $p_T$ ( $>40$ GeV)

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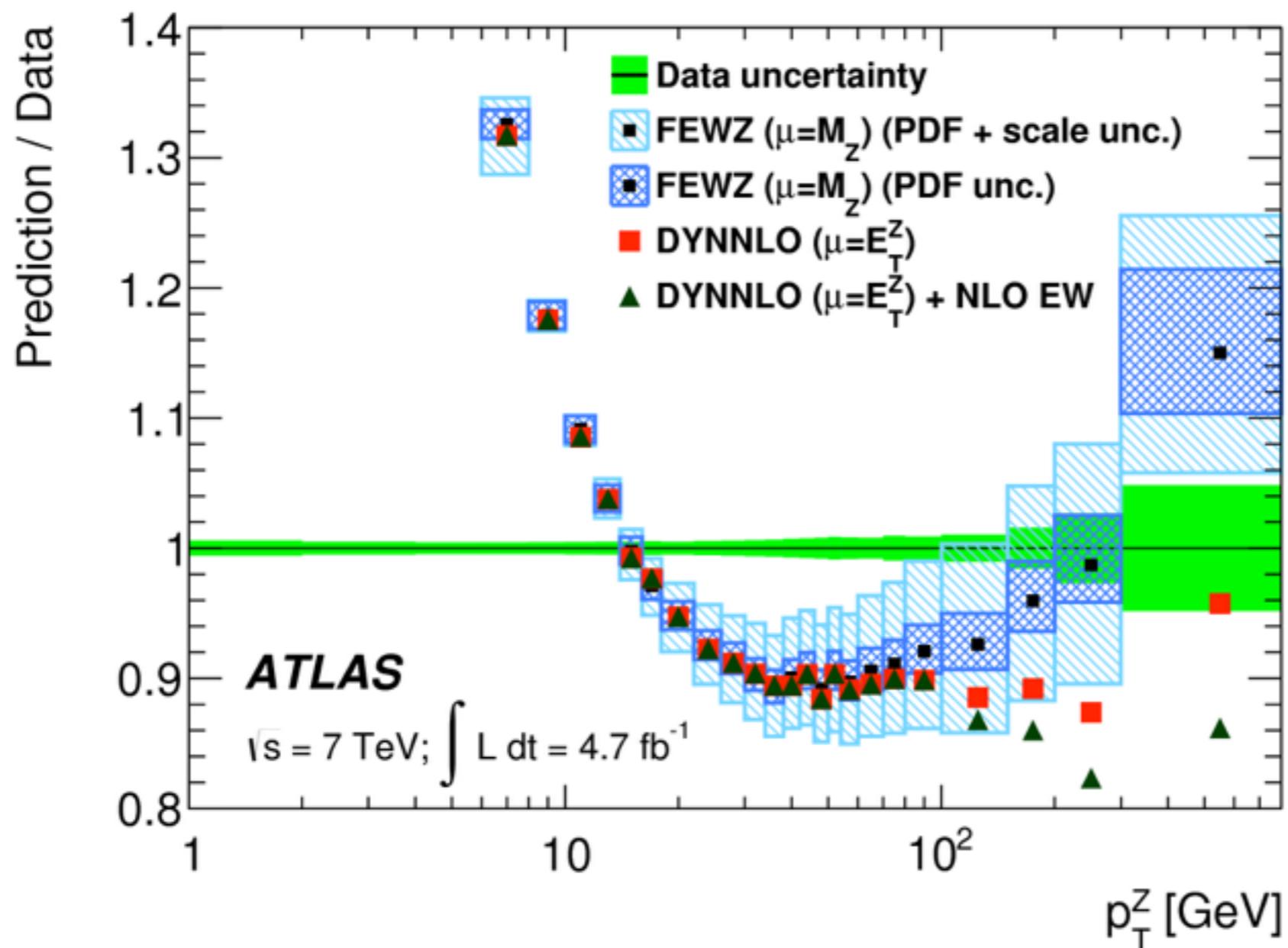
Z+j NLO discrepancies; interpretation and prospects for NNLO

Z  $p_T$  ( $> 40$  GeV) should be a gold-plated observable:

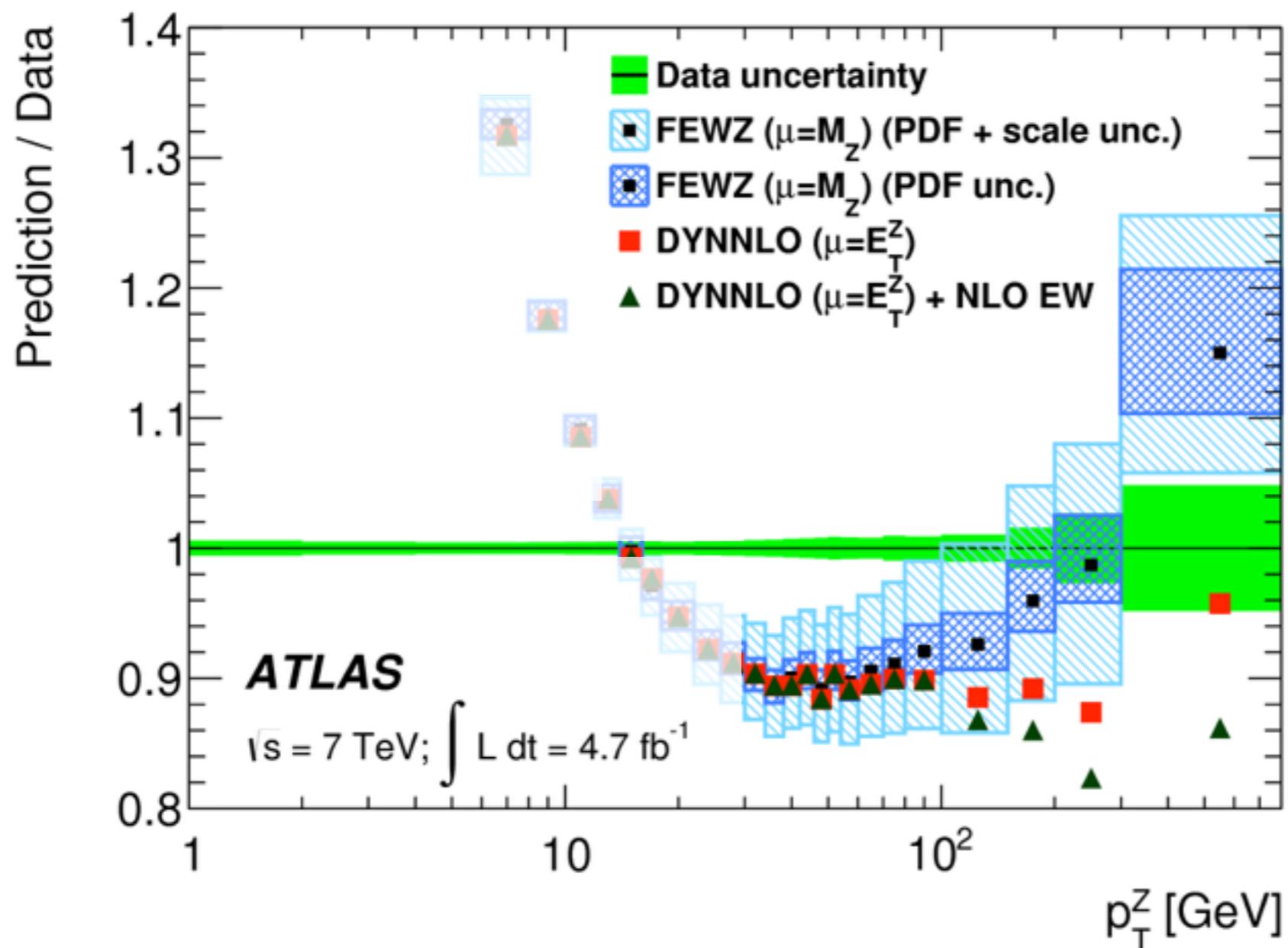
- Experimentally very clean (errors  $< 1\%$ )
- Theoretically clean too: free of large logs  
(NNLO for Z is NLO for Z  $p_T$ ; NNLO still to come)

Important because uniquely sensitive to  **$\alpha_s$  x gluon x quark**

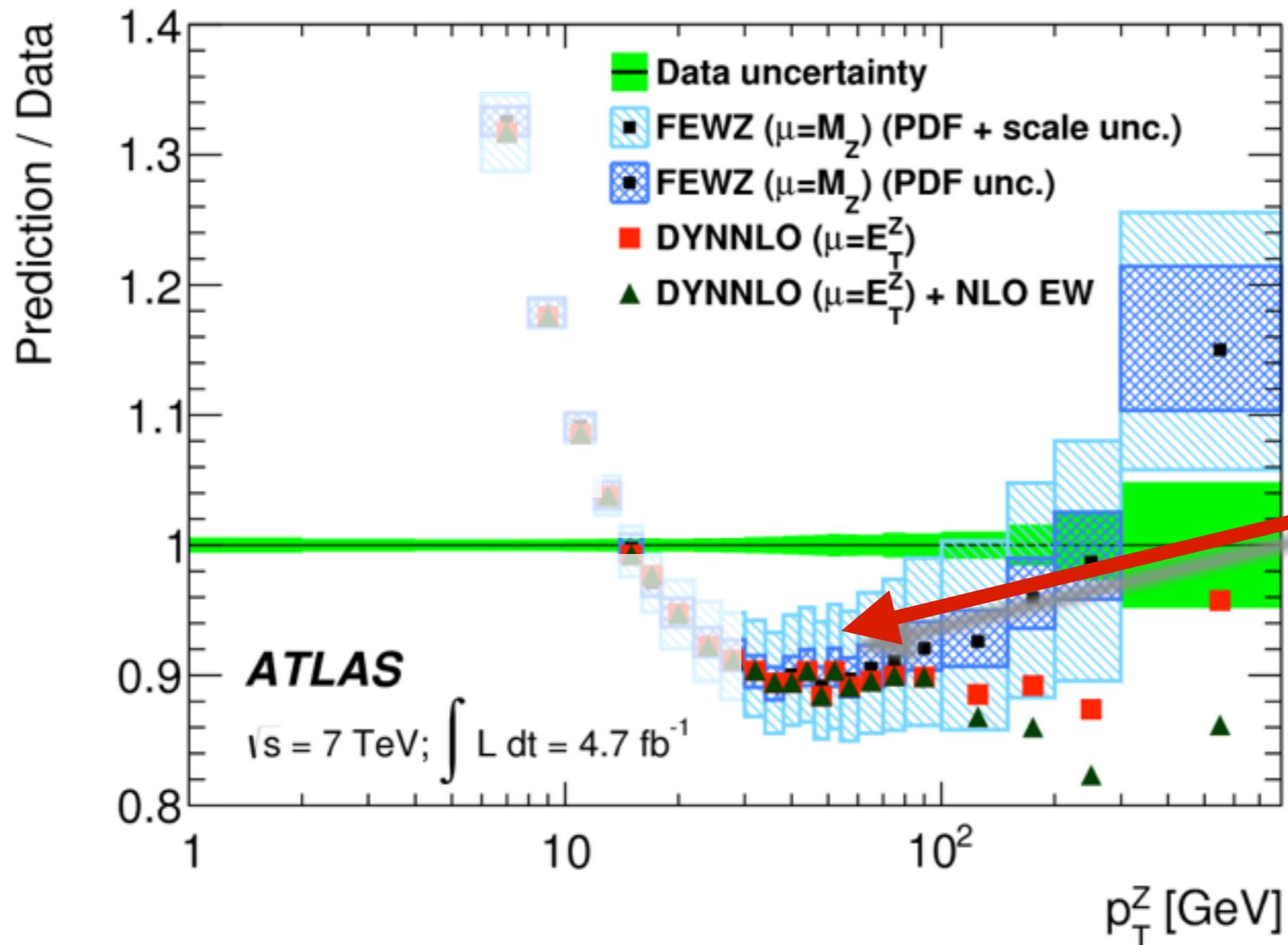
# ATLAS Z pT: NNLO / Data



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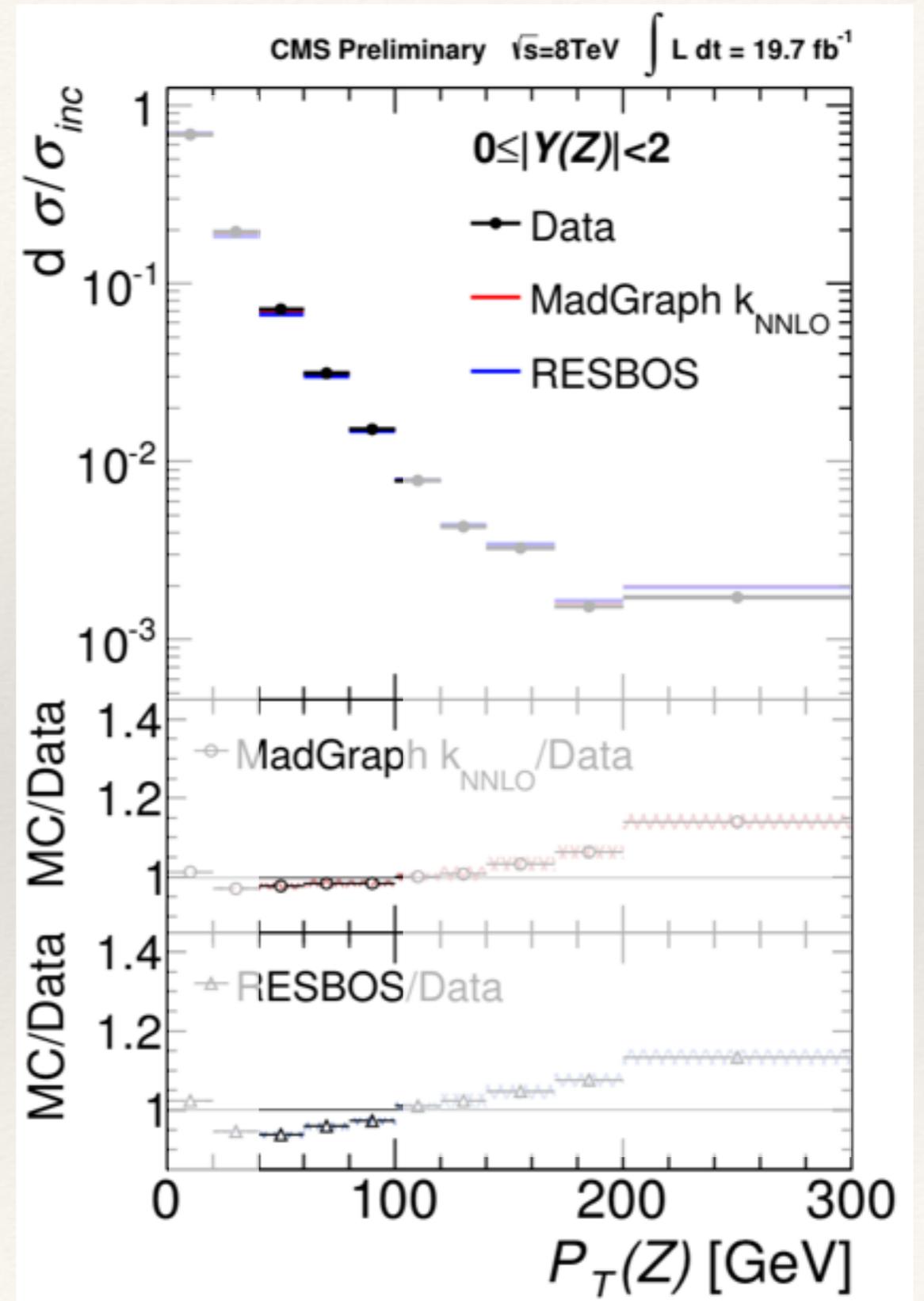
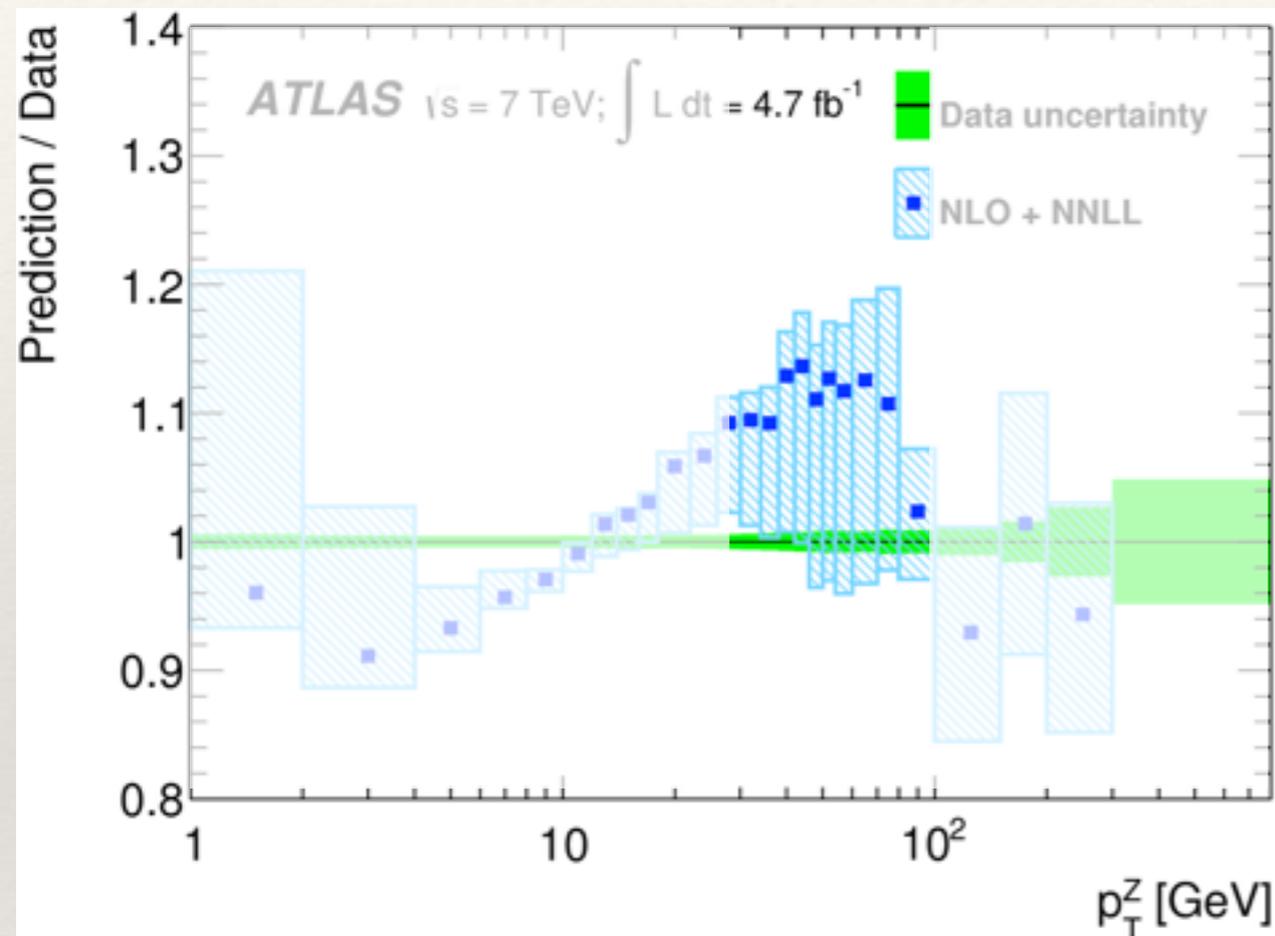
# ATLAS Z pT: NNLO / Data



**2- $\sigma$  discrepancy**  
in region where  
(N)NLO should  
be reliable

CMS doesn't compare to pure fixed-order (sees 5-7% discrepancy with RESBOS)

BDMT prediction looks 20% higher in same region: not clear why.



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# Z $p_T$ mystery needs solving

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The discrepancy feeds into other observables (e.g. jet dist<sup>n</sup> in Z+jet events).

Is theory uncertainty badly underestimated? Will NNLO solve the problem? What's the real scope for resummation to modify distribution for  $p_T > 40$  GeV?

Or are PDFs substantially wrong? (Z  $p_T$  is never an input; while much less precise incl. jets are an input – why?)

Do CMS and ATLAS data agree?

# Scale Variation

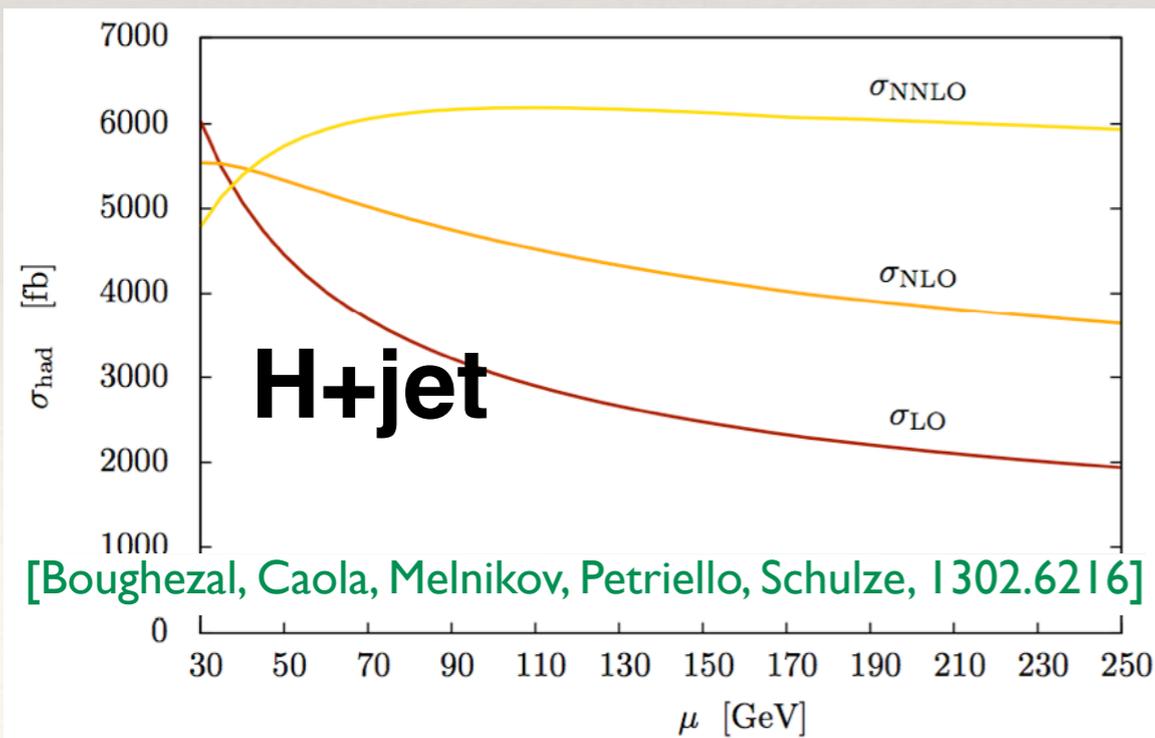
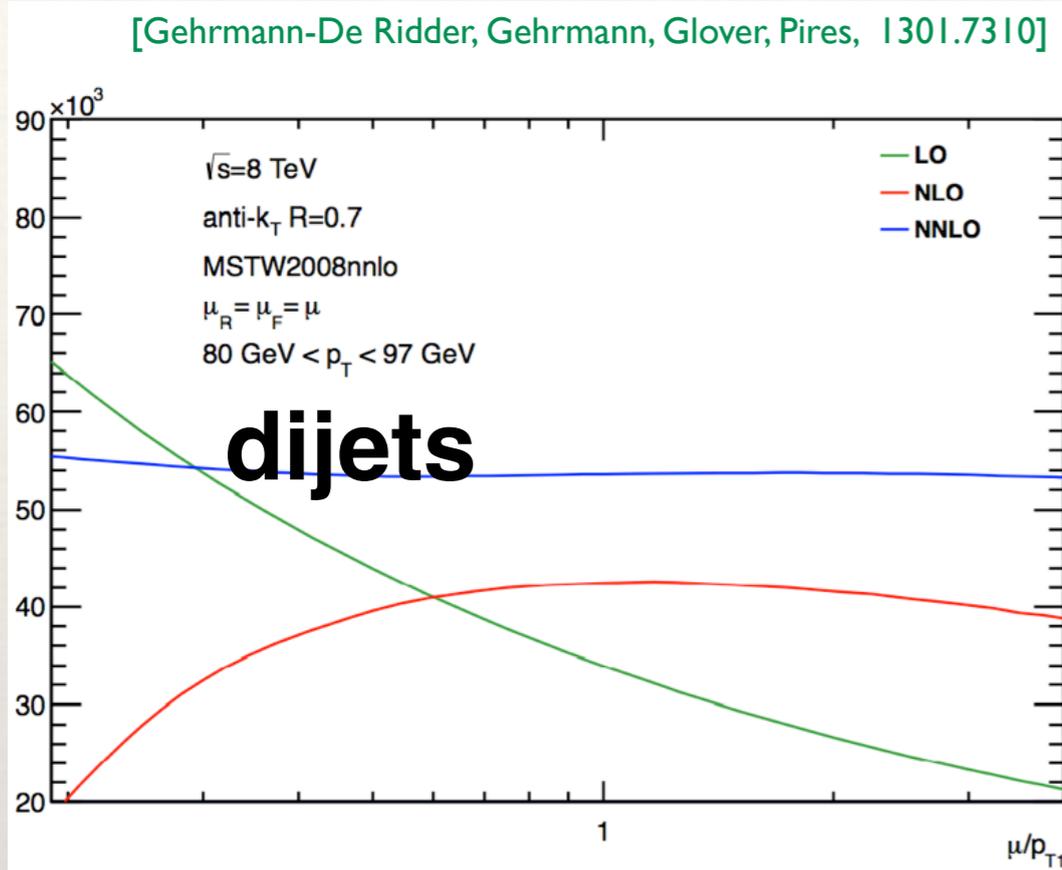
We have a **convention** – choose a central (possibly dynamical) scale, and calculate with  $x_2$  and  $x_{1/2}$  scale variations.

Reflects fact that physical scale choice is genuinely ambiguous and conveniently gives us an uncertainty estimate.

# How reliable is scale variation?

There's no shortage of cases where (sometimes partial) NNLO is at or beyond edge of NLO scale variation

[Czakon, Fiedler & Mitov 1303.6254]



$t\bar{t}$  @ LHC8

LO:  $145^{+49}_{-34}$  pb

NLO:  $213^{+25}_{-27}$  pb

NNLO:  $239^{+9}_{-15}$  pb

top++, MSTW2008NNLO,  $\mu = m_t$

Scale variation gives an uncertainty  
But to what extent is it a measure of *the* uncertainty?

## Toy model:

(1) Take a running coupling where

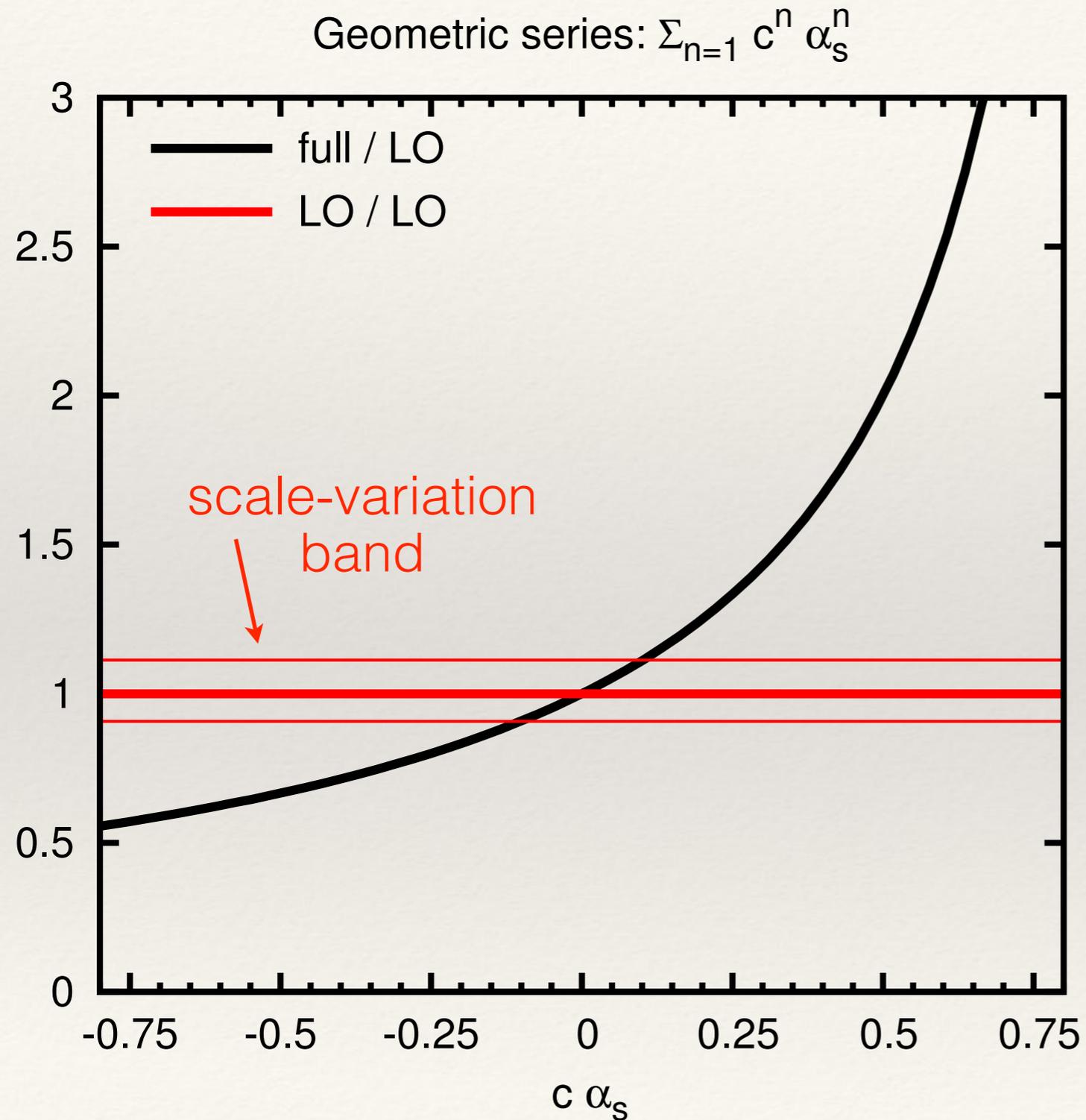
$$\begin{aligned}\beta_0 &= \beta_{0,\text{QCD}} \\ \beta_1 &= \beta_2 = \dots = 0\end{aligned}$$

(2) Consider a simple perturbative series that you can sum to all orders. E.g.

$$\sigma = \frac{c\alpha_s(M)}{1 - c\alpha_s(M)} = c\alpha_s + c^2\alpha_s^2 + c^3\alpha_s^3 + \dots$$

simplest possible series in QCD: corresponds to coupling at one scale expressed in terms of coupling at another (reference) scale  $M$

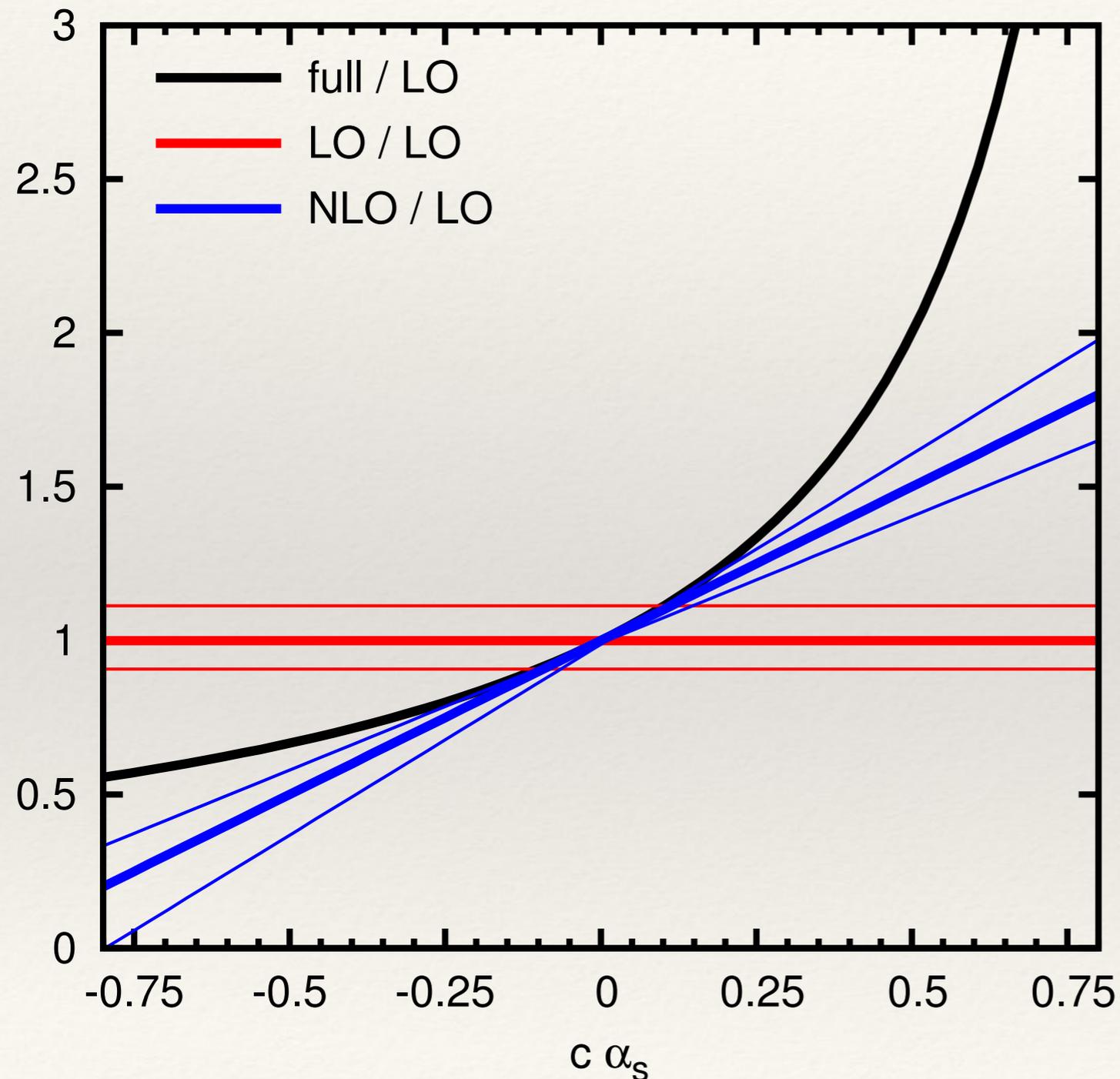
Now examine truncations of series,  
as a function of  $c$  for  $\alpha_s = 0.12$



LO: scale variation  
mostly useless.

Now examine truncations of series,  
as a function of  $c$  for  $\alpha_s = 0.12$

Geometric series:  $\sum_{n=1} c^n \alpha_s^n$

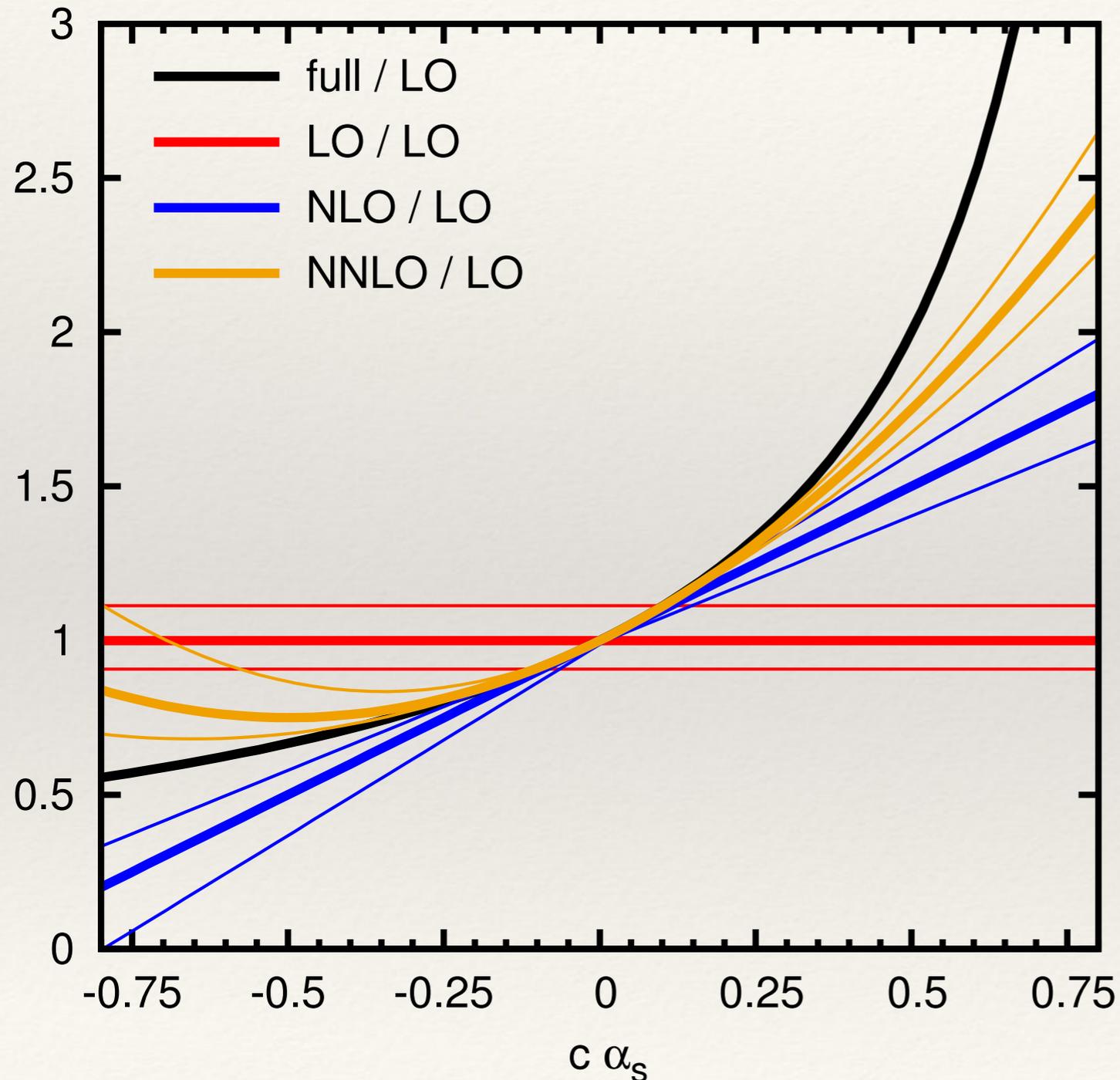


LO: scale variation  
mostly useless.

NLO: its usefulness  
extends further, but at  
some point breaks  
down.

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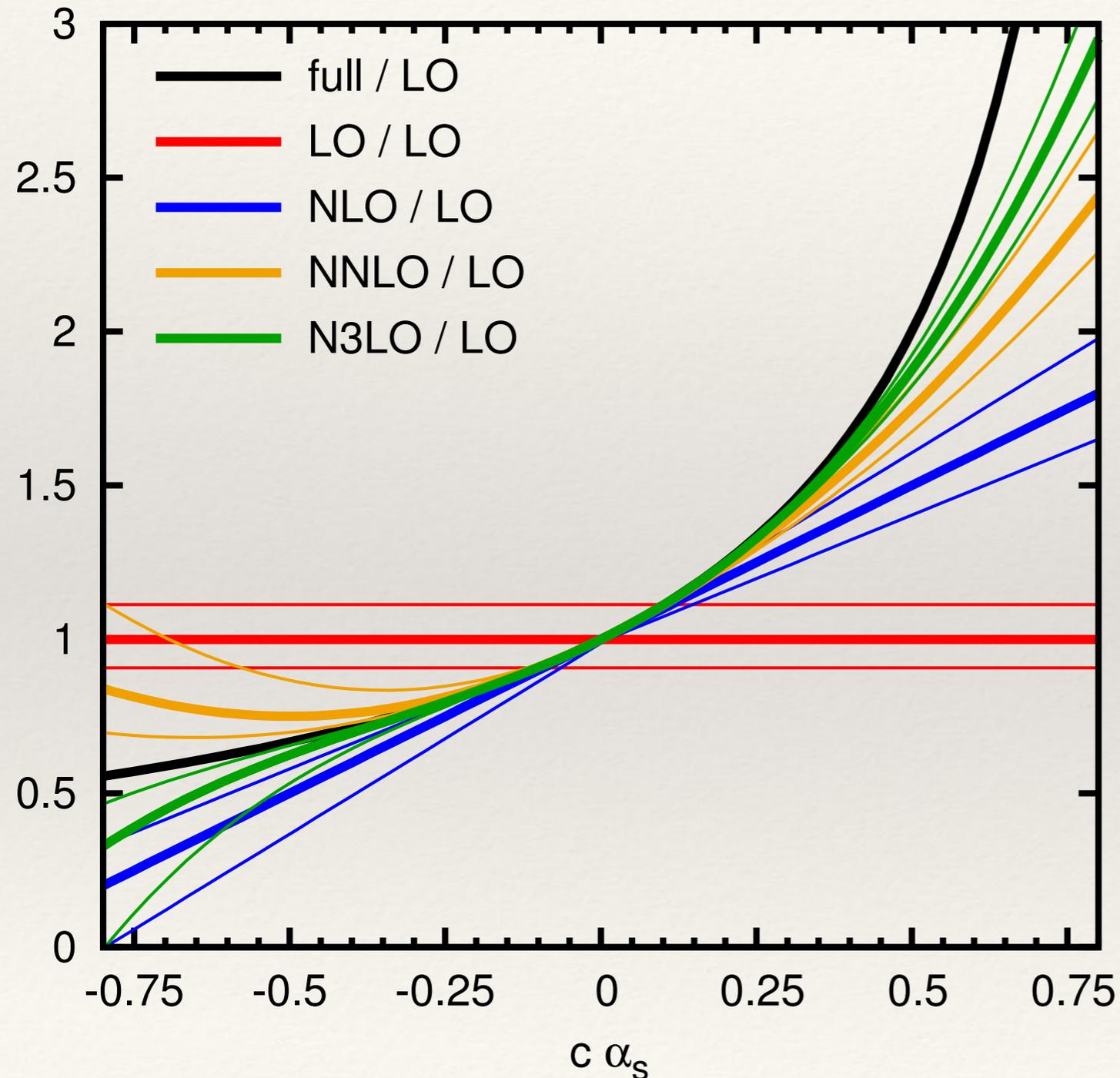
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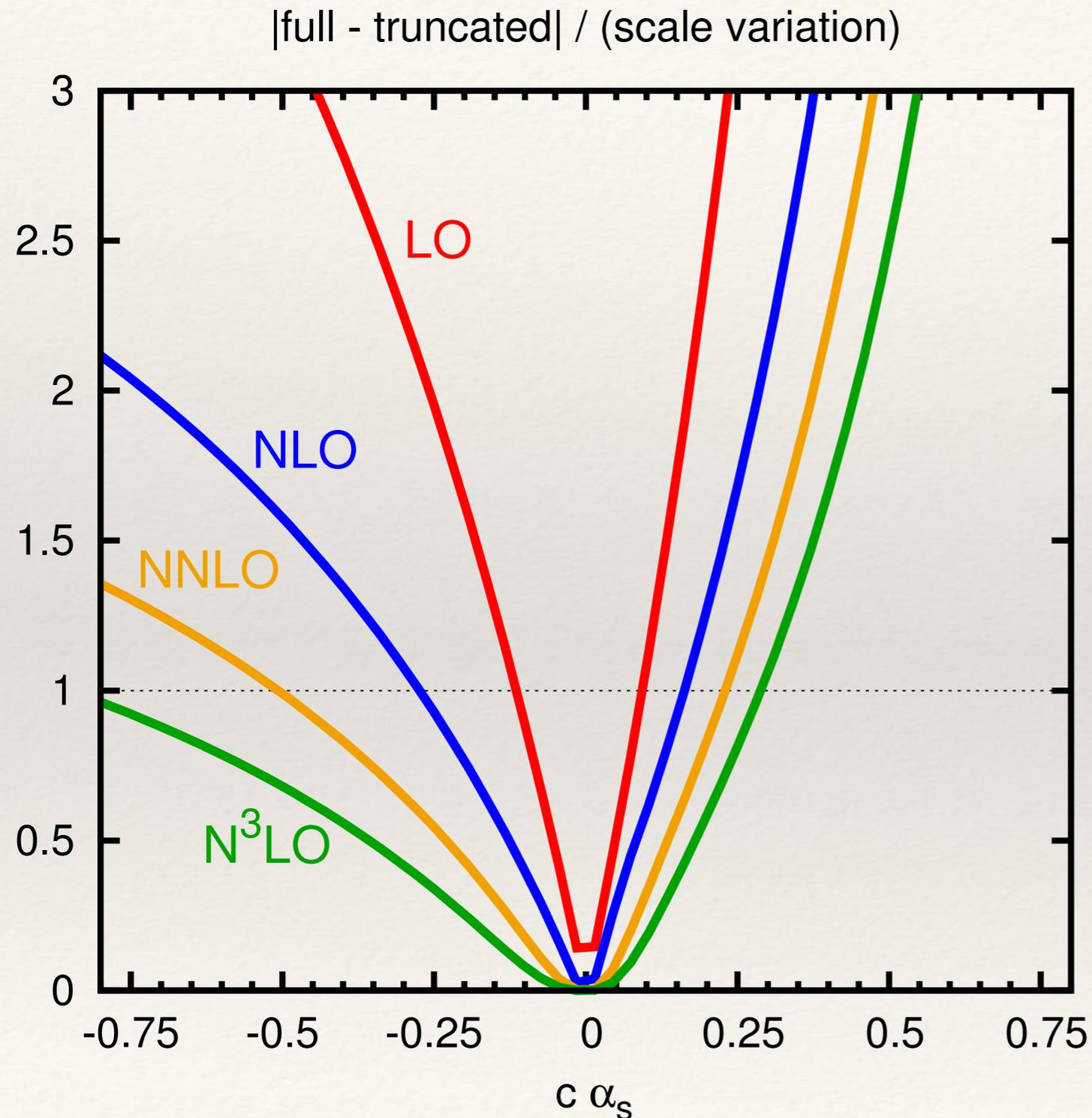


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Now examine truncations of series,  
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LO: scale variation  
mostly useless.

NLO: it's usefulness  
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NNLO: ditto  
NNNLO: ditto

$$\sigma = \sum_{n=1} (c \alpha_s)^n$$

$$\text{Higgs } (\mu = m_H) \\ \mathcal{N} \times (\alpha_s^2 + 11\alpha_s^3 + 62\alpha_s^4)$$

Normalised to LO, what's missing from N<sup>p</sup>LO is:

$$\sim c^{p+1} \alpha_s^{p+1}$$

Scale var<sup>n</sup> ( $c \gg 1$ ) gives:

$$\sim (p + 1) \cdot c^p \alpha_s^{p+1} \quad *$$

Ratio scale uncertainty/  
true missing higher  
orders:

$$\sim \frac{p + 1}{c}$$

For poorly converging series ( $c \gg 1$ ), scale variation **parametrically** underestimates the uncertainty.

At higher orders  
( $\equiv$  for larger  $p$ )  
scale variation works further, but for large enough  $c$  inevitably breaks down

\* coefficient is  $\frac{23}{6\pi} \ln 2 \simeq 0.85$

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# Other approaches?

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**Cacciari-Houdeau** tries to get a prob. distribution for uncertainty; but shares limitations with scale variation (cannot detect large geometric growth of series)

**David-Passarino**: attempts “series acceleration” – does detect coefficient growth, though arguably fairly complicated

**MINLO**: tries to solve a different problem, i.e. scale & Sudakov setting in multi-scale problems; uncertainty is somehow a separate problem

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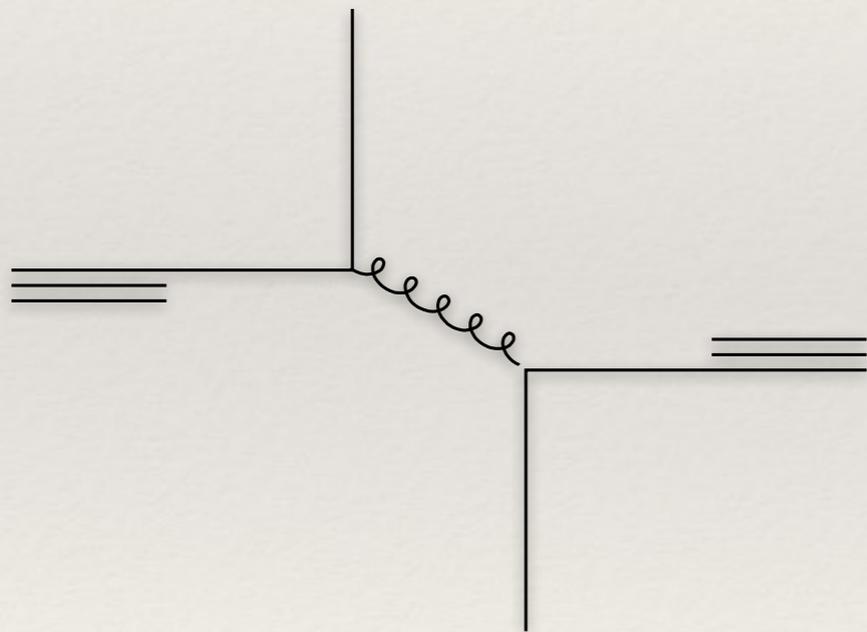
# Summary on scale variation

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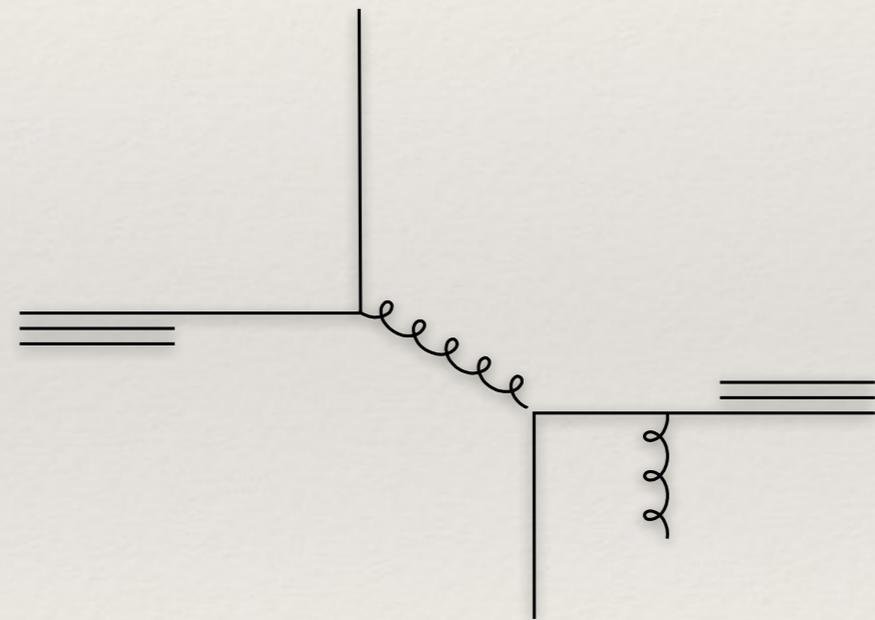
- ❖ Choice of scale is a genuine ambiguity
- ❖ But size of scale variation knows little about physics, only about coefficients of the series
- ❖ Scale variation doesn't correctly handle case when coefficients grow large.

Can one do better? Possibly, e.g. by supplementing scale variation uncertainties with information on growth of coefficients (à la David–Passarino, maybe with simplifications)

# (A) Symmetric cuts



LO dijet/digamma/etc. configurations  
are symmetric



But symmetry broken even by  
tiny amount of ISR

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# Mathematical illustration

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Hard  $2 \rightarrow 2$  cross section:  $\sigma_0(p_t) \sim \frac{1}{p_t^n} \quad [n \sim 5 \gg 1]$

NLO with symmetric cuts:  
( $p_{t1}, p_{t2} > p_t$ )  $\sigma_{\text{sym-cut}}(p_t) \sim \sigma_0(p_t)(1 - \alpha_s \ln^2 n)$

[Due to ISR, which imbalances the event]

Large double-log is considered dangerous, so symmetric cuts are widely deprecated. [Frixione-Ridolfi '97]

# Asymmetric cuts

$$p_{t1} > p_t$$

$$p_{t2} > (1-\epsilon) p_t$$

$\epsilon=0$ : equivalent to symmetric cuts

$$\sigma_{\text{sym-cut}}(p_t) \sim \sigma_0(p_t)(1 - \alpha_s \ln^2 n)$$

Take  $\epsilon \sim 0.5$ , i.e. cut mostly on hardest jet

$$\sigma_{p_{t1}\text{-cut}}(p_t) \sim \sigma_0(p_t)(1 + \alpha_s \ln^2 n)$$

[schematically...]

Take  $\epsilon \sim 0.1-0.2$ , i.e. standard asym. cut

$$\sigma_\epsilon(p_t) \sim \sigma_0(p_t)(1 + \alpha_s f(n, \epsilon))$$

**Asymmetric cuts just make a bad problem more complex**

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# A possible solution?

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$$\frac{1}{2}(p_{t1} + p_{t2}) \equiv \frac{1}{2}H_{T2} > p_t, \quad p_{t2} > (1 - \epsilon)p_{t1}$$

ISR of momentum  $p_{t,ISR}$  leaves  $H_{T2}$  almost unchanged

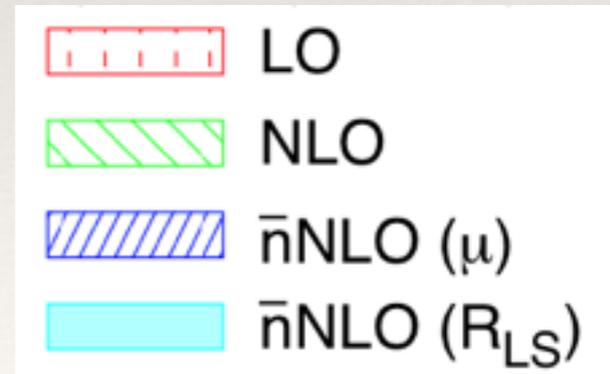
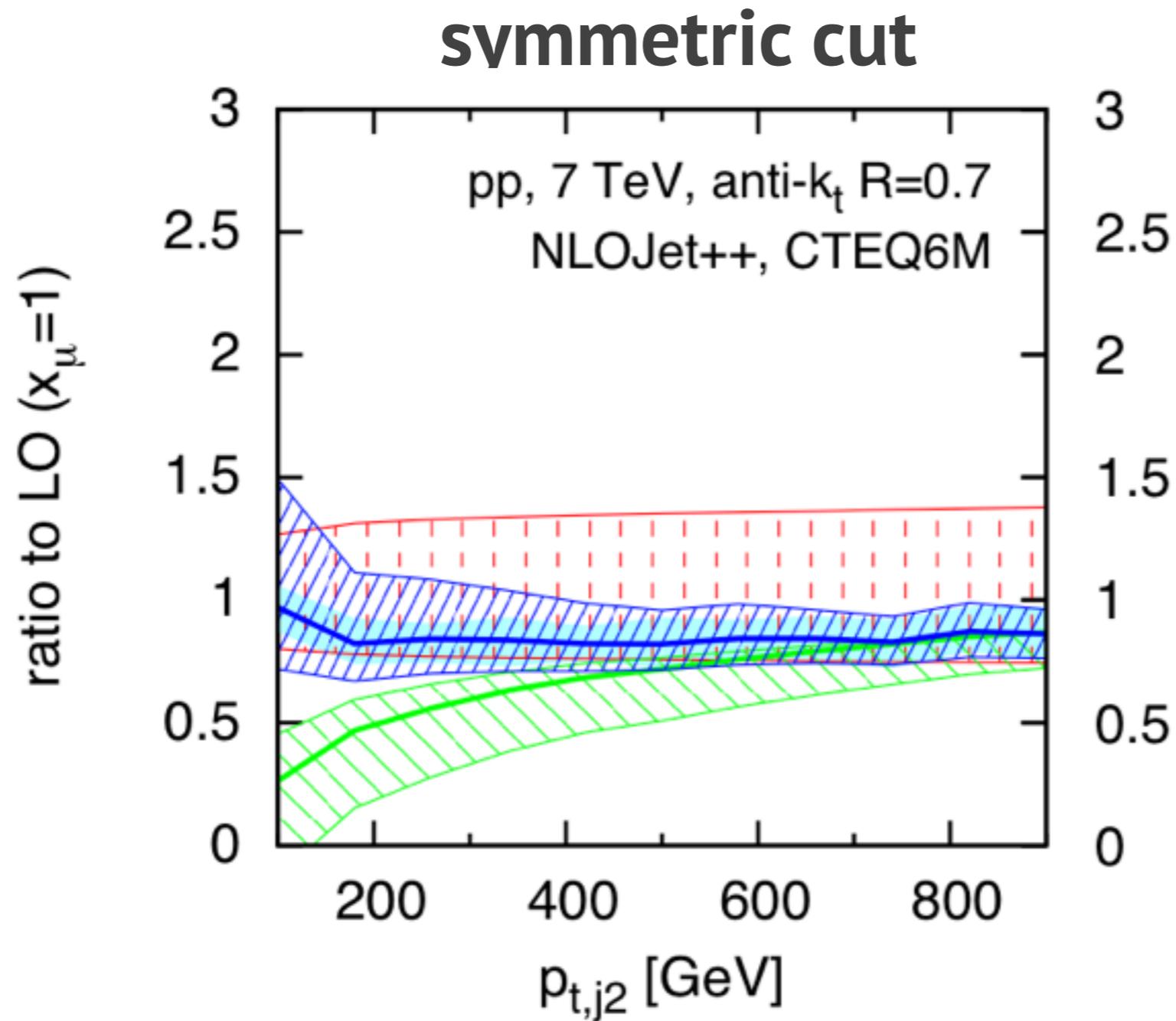
$$\frac{1}{2}H_{T2} \rightarrow \frac{1}{2}H_{T2} + \mathcal{O}\left(\frac{p_{t,ISR}^2}{H_T}\right)$$

because impact of ISR mostly balances in jets 1 and 2.

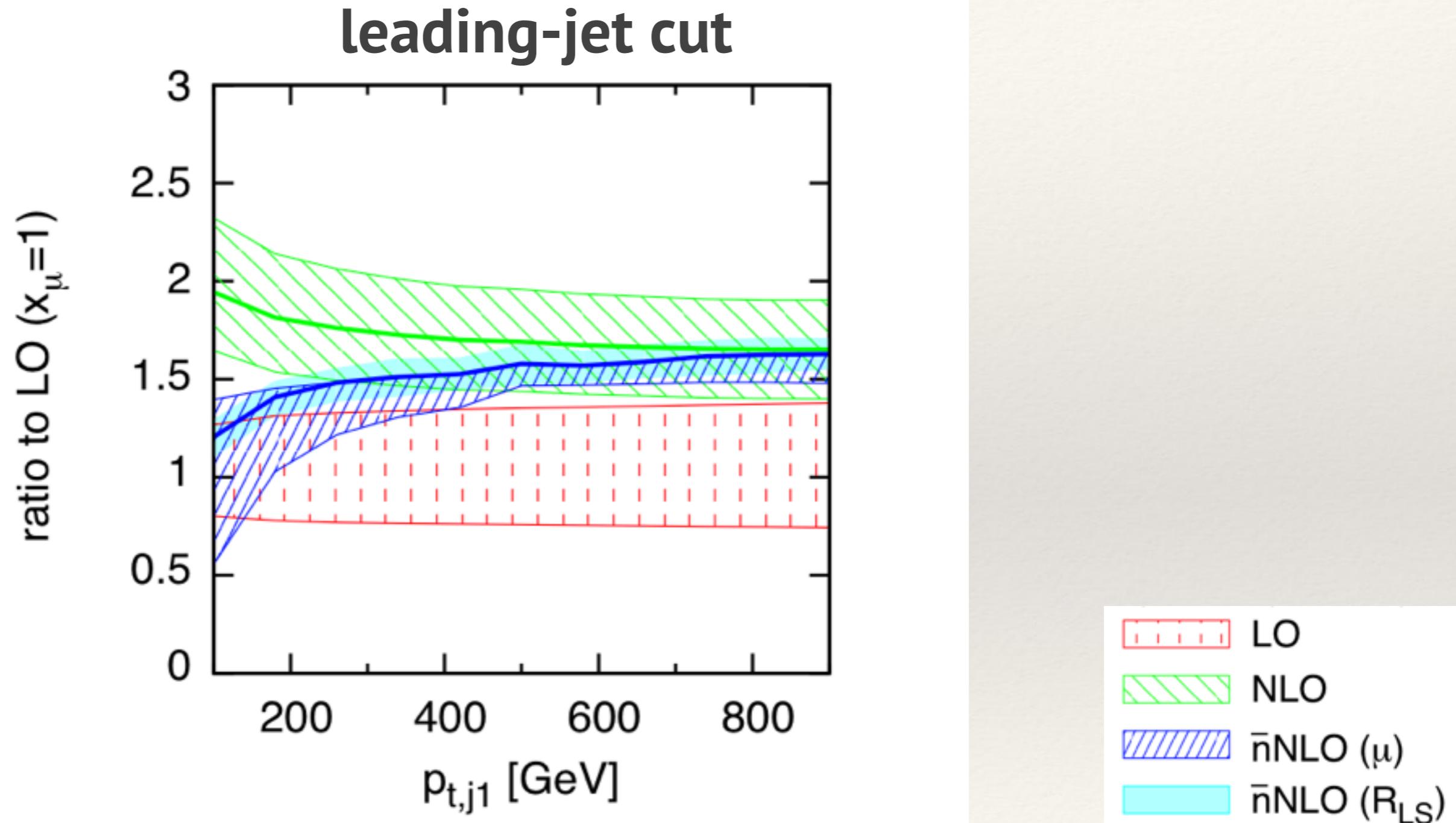
If  $\epsilon$  is moderate, e.g.  $\epsilon = 0.5$ , perturbation theory should

be well behaved (no large logs of  $n$  or of  $\epsilon$ )

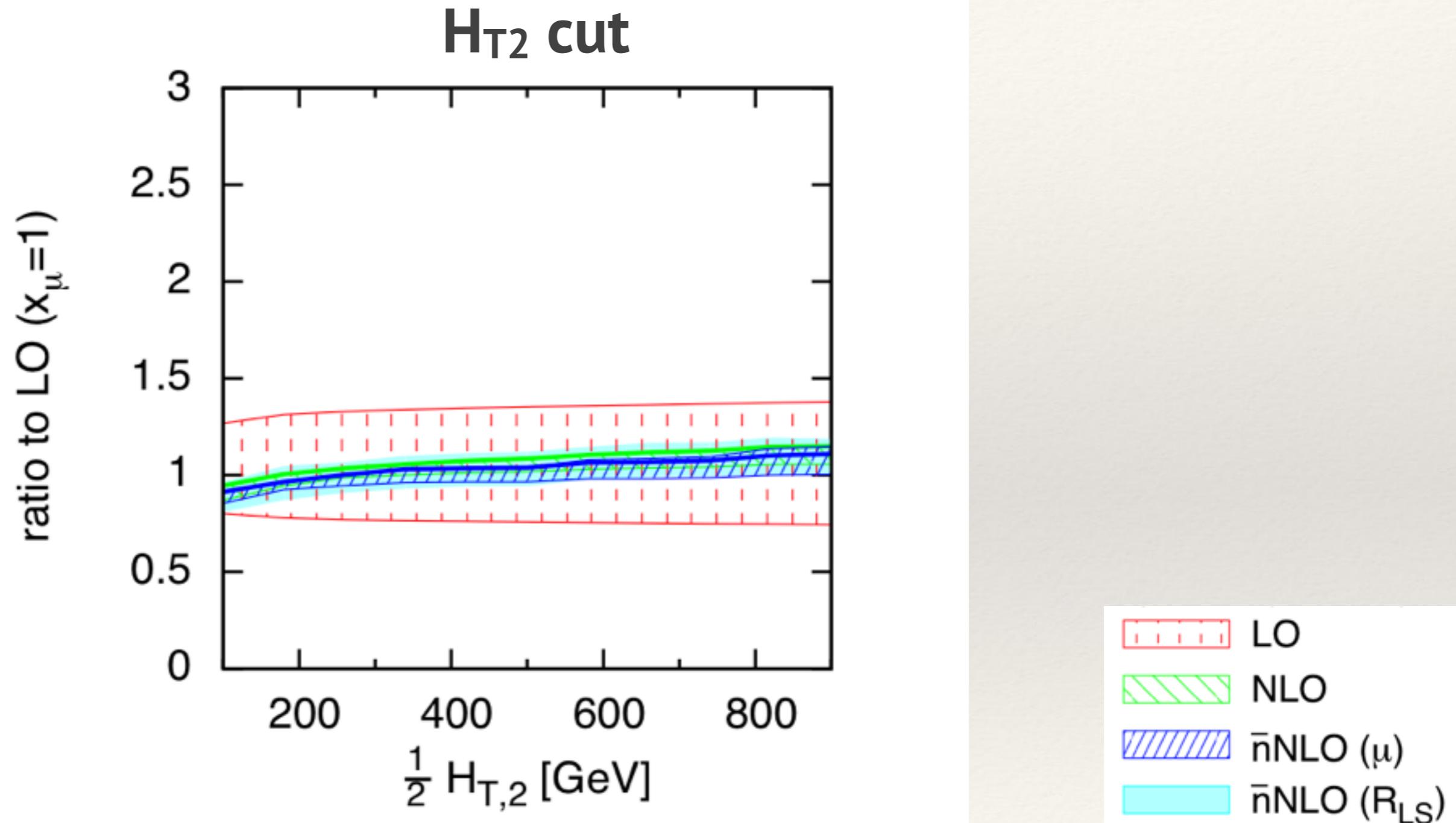
# Performance?



# Performance?



# Performance?



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# Asymmetric cuts summary

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- ❖ Asymmetric cuts fine tune cancellation between two different physical effects (interplay of ISR with steeply falling cross section, and logs of  $\epsilon$ )
- ❖ That's dangerous.
- ❖ **Self-balancing cuts**, e.g. on  $H_{T2}$ , appear to do better based both on analytic considerations and (n)NLO convergence (also from some simple Pythia studies with ISR turned on/off)

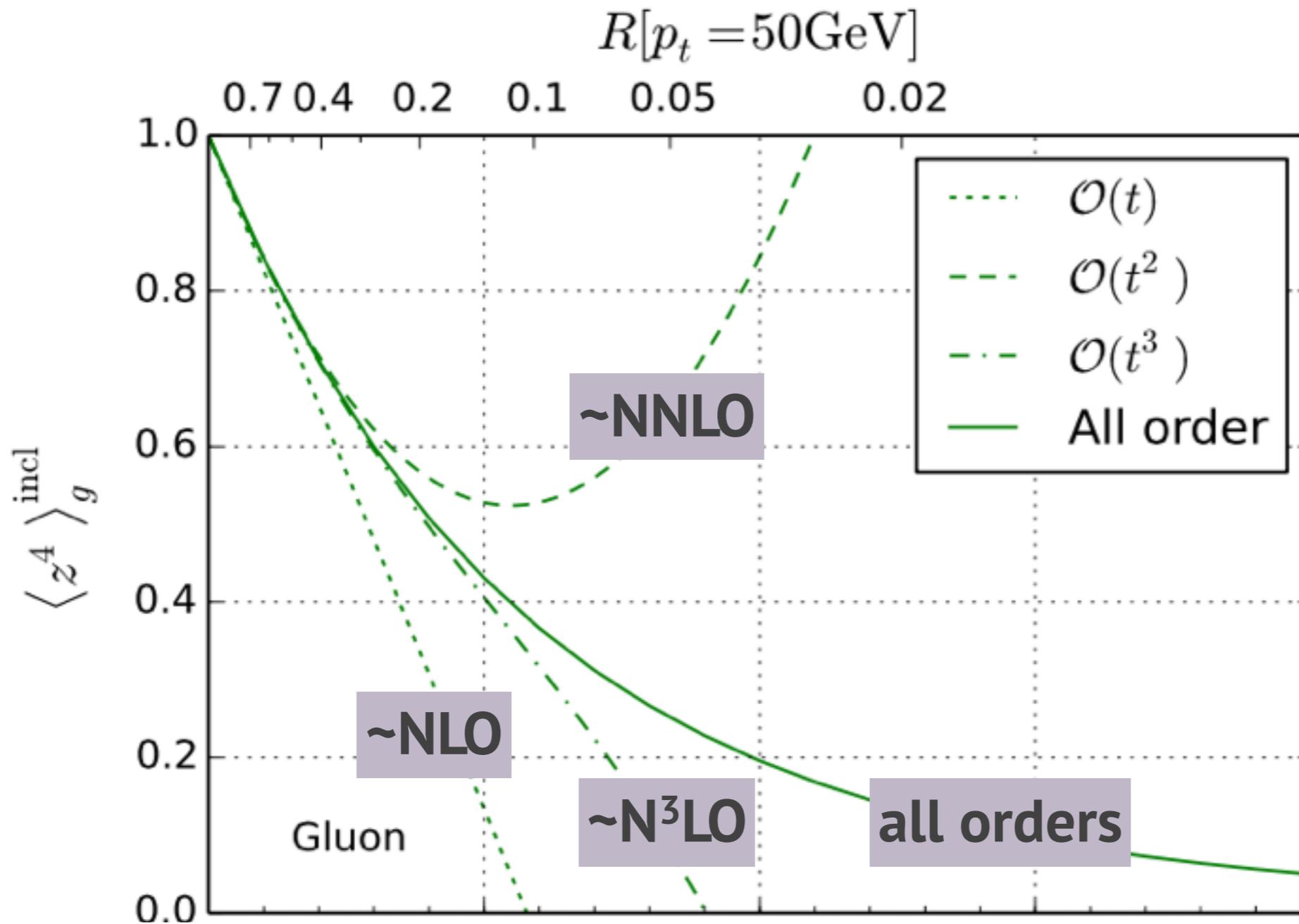
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# Small jet radii — how small?

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- ❖ Heavy-ion physics regularly uses small  $R$  (e.g. 0.2)
- ❖ Small  $R$  can be useful in  $pp$  to reject pileup jets (e.g. for VBF, where tracking not accessible for forward jets)
- ❖ I think it's interesting to explore  $R$  as small as experimentally possible
- ❖ New calculations give us insight into what happens perturbatively at small  $R$ . [Dasgupta, Dreyer, GPS Soyez '14]

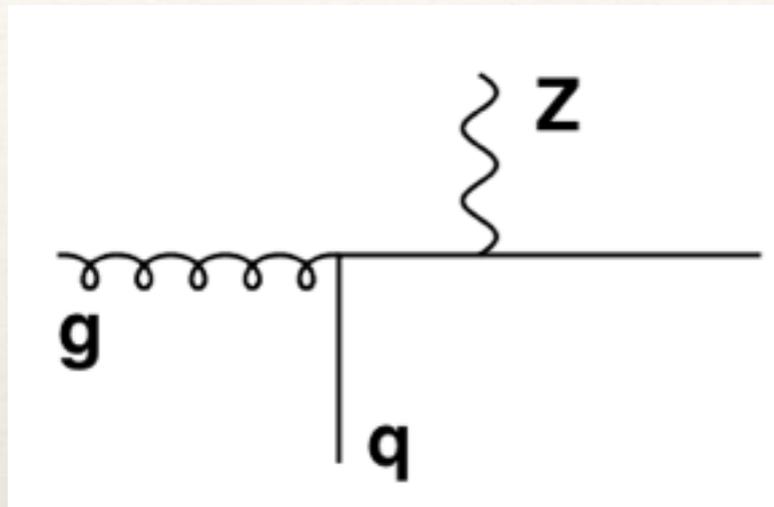
# Small-R corrections for gluon-induced part of inclusive spectrum



Suggests that NLO *may* be insufficient for  $R=0.4$  (cf. caveats)

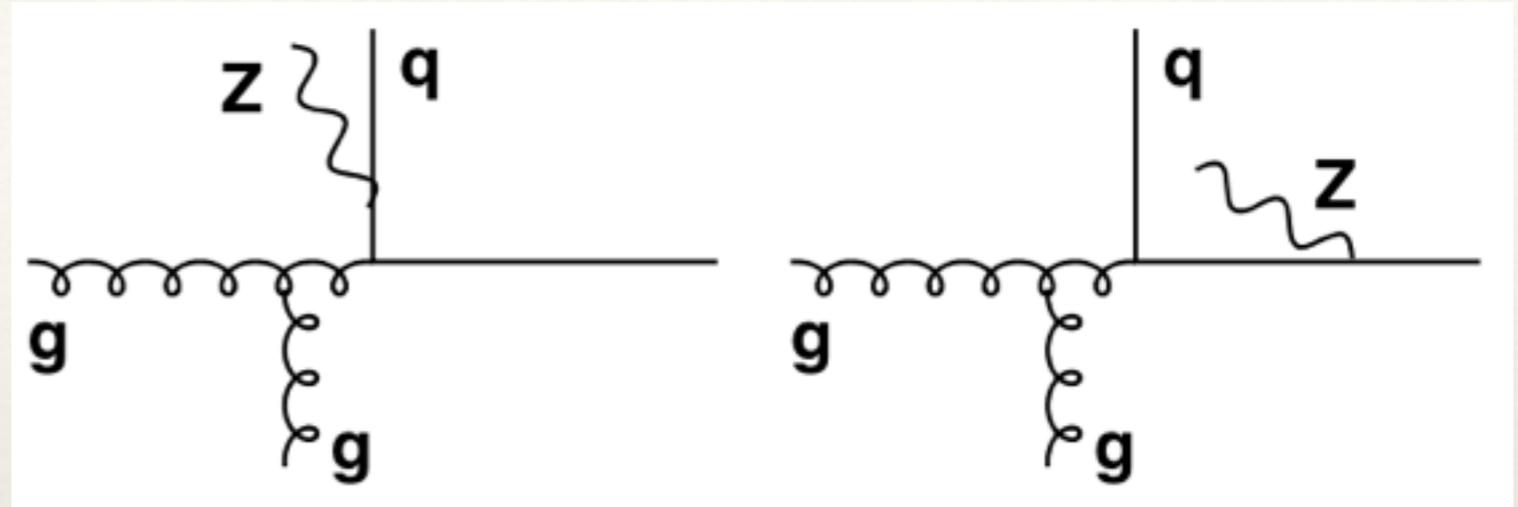
But all-order small-R effects can be resummed

# New physical effects



## LO Z+jet topology

Dominates if bin in  
Z  $p_T$



## new NLO Z+jet topologies

Dominate if bin in leading jet  $p_T$   
Z effectively a light degree of freedom  
(restoring EW symmetry)

It would be fun to demonstrate  
dominance of these topologies  
at high  $p_T$

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# Conclusions

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- ❖  $Z$   $p_T$  distribution is important; deserves dedicated effort to resolve TH-EXP discrepancies (e.g. ATLAS-CMS, understanding different theory calculations, etc.)
- ❖ Scale variation is guaranteed to fail for some observables; arguably needs to be supplemented with other info from perturbative series
- ❖ Making series look better always helps: asymmetric cuts is one area that needs revisiting
- ❖ Exploring new phase space (small- $R$ , “light” EW bosons) is interesting in its own right

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