



Département
de Physique
—
École Normale
Supérieure



Anomalies, Chern-Simons Terms, and Black Hole Entropy

Tatsuo Azeyanagi (ENS)

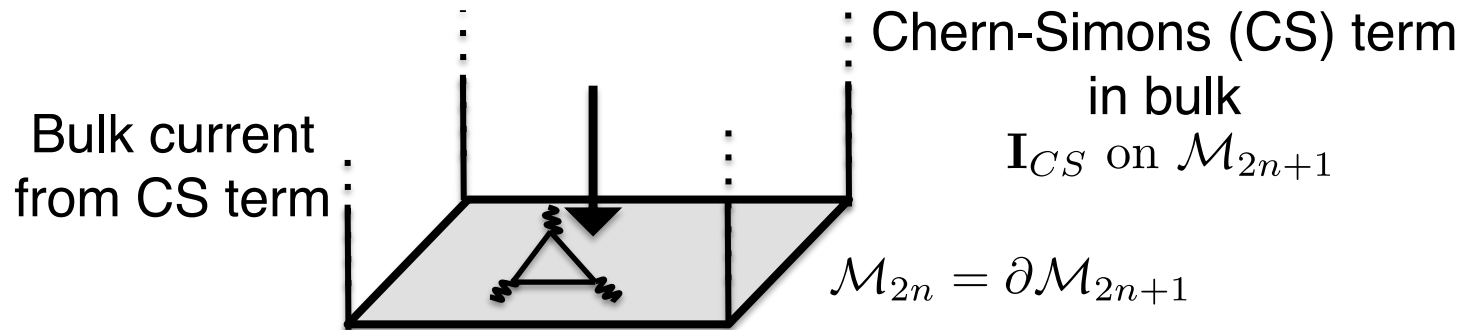
Based on arXiv:1311.2940, 1407.6364, 1505.02816

Collaborators: R. Loganayagam (IAS), G.S. Ng (McGill), M.J. Rodriguez (Harvard)

Anomalies at Zero Temperature

Beautiful mathematical structure associated with anomalies
(we consider **Global Anomalies in Even-dim. QFT**)

Anomaly Inflow Mechanism [Callan-Harvey]



Anomalies are classified by **Anomaly Polynomials**

$$P_{anom}(\mathbf{F} = d\mathbf{A}, \mathbf{R} = d\mathbf{\Gamma} + \mathbf{\Gamma} \wedge \mathbf{\Gamma}) \equiv d\mathbf{I}_{CS}$$

example

CS term

Anomaly
Polynomial

2d gravitational anomaly

$$\text{tr} \left(\mathbf{\Gamma} \wedge \mathbf{R} - \frac{1}{3} \mathbf{\Gamma} \wedge \mathbf{\Gamma} \wedge \mathbf{\Gamma} \right)$$

$$\text{tr}(\mathbf{R} \wedge \mathbf{R})$$

4d mixed anomaly

$$\mathbf{A} \wedge \text{tr}(\mathbf{R} \wedge \mathbf{R})$$

$$\mathbf{F} \wedge \text{tr}(\mathbf{R} \wedge \mathbf{R})$$

Anomalies at Finite Temperature

Anomaly-Induced Transport

[Son-Surowka, Bhattacharyya et.al.
Erdmenger et.al., Torabian-Yee, ...]

In hydrodynamic limit, anomalies generate new transports

Replacement Rule

[Loganayagam, Loganayagam-Surowka
Jensen-Logayanagam-Yarom, ...]

Anomaly-induced transport coefficients are determined by anomaly polynomial

$$\text{Stress-Energy Tensor} \quad (T_{\alpha\beta})_{anom} = -n\mathfrak{F}[\mu, T](u_\alpha V_\beta + u_\beta V_\alpha) + \dots$$

$$\text{U(1) current} \quad (J_\alpha)_{anom} = -\frac{\partial\mathfrak{F}}{\partial\mu}V_\alpha + \dots$$

$$\text{Entropy current} \quad (J_\alpha^S)_{anom} = -\frac{\partial\mathfrak{F}}{\partial T}V_\alpha + \dots$$

$$V^\mu = \epsilon^{\mu\nu\rho_1\cdots\rho_{2n-2}}u_\nu(\partial_{\rho_1}u_{\rho_2})\cdots(\partial_{\rho_{2n-3}}u_{\rho_{2n-2}})$$

with

$$\mathfrak{F}[T, \mu] = \mathbf{P}_{anom}(\mathbf{F} \rightarrow \mu, tr(\mathbf{R}^{2k}) \rightarrow 2(2\pi T)^{2k})$$

Anomaly polynomials classify anomaly-induced transports

Gauge/Gravity Duality

Setup on CFT Side

Fluid with non-trivial anomaly-induced transports

→ U(1) charged rotating (conformal) fluid in $2n$ -dim.

Setup on Gravity Side

Theory

- $(2n+1)$ -d Einstein-Maxwell-**Chern-Simons** theory with negative cosmological const.
- CS Terms: U(1), gravitational, mixed
 - Same as anomaly inflow mechanism

Configuration

U(1) charged rotating black hole (BH) on AdS_{2n+1}

Cf. some quantities in 5d case

[Landsteiner-Megias-Melgar-Pena-Benitez, Karzheev-Yee,...]

Why Gravity Side ?

1. Proof of replacement rule in QFT is formal
2. Interesting properties of CS terms :

- Higher derivative terms
- **Non-covariant** under diffeo and gauge trans.

$$\delta_{\chi} \mathbf{I}_{CS} = \mathcal{L}_{\xi} \mathbf{I}_{CS} + d(\dots) \quad \chi = \{\xi : \text{diffeo}, \Lambda : \text{gauge}\}$$

while EoM is **covariant**

3. Basics are not well-understood :
 - **Boundary stress tensor for general CS terms?**
 - **BH entropy with general CS term?**
 - **Charged rotating AdS BH in higher dim.?**

Our Work : Anomaly Polynomial Plays Crucial Roles!

Implication to BH Microstates

“Traditional” Microstate Counting for BH Entropy

Map to CFT₂ entropy counting problem → Cardy Formula

(example) BTZ BH, (near) extremal BHs

BH in higher-dim. AdS spacetime

- Dual higher-dim. CFT does not have neither infinite dimensional symmetry nor modular invariance
→ **Hard to compute entropy fully from higher dim. CFT**

cf. supersymmetric cases in 4d and 6d [Komargodski et.al.]

If Chern-Simons BH entropy = Replacement Rule ...

With replacement rule, we can compute anomaly-induced part of entropy for **higher-dim finite temp. BH from CFT!**

“BH Entropy is Noether Charge”

Wald Entropy Formula

[Wald, Lee-Wald, Iyer-Wald],
[Kang-Jacobson-Myers]

• BH entropy formula for general **diff. invariant** Lagrangian

$$\left(\delta_\chi L_{cov} = \mathcal{L}_\xi L_{cov} \quad (\text{example}) \quad R_{abcd} R^{abcd} \right)$$

$$S_{\text{Wald}} = 2\pi \int_H \epsilon_{ab} \epsilon_{cd} \frac{\delta L_{cov}}{\delta R_{abcd}}$$

Differential Noether Charge

Killing vector

$$\xi = \partial_t + \Omega_H \partial_\phi$$



Conserved charge

$$d\oint \mathbf{Q}_{\text{Noether}} = 0$$

$\oint(\dots)$: linear combination of $\delta(\text{something})$

$$\rightarrow \frac{\int_\infty \oint \mathbf{Q}_{\text{Noether}}}{\delta M + \Omega_H \delta J} = \frac{\int_H \oint \mathbf{Q}_{\text{Noether}}}{T_H \delta S}$$

A Sketch of Derivation

Some Steps of Noether Procedure

Point 1 : Variation of Lagrangian

$$\delta\mathbf{L}(\phi) = \not\delta\mathbf{E} + d(\not\delta\Theta) \quad \not\delta(\dots) : \text{linear combination of } \delta(\text{something})$$

Point 2 : Pre-symplectic current Ω

$$d(\Omega(\delta_1\phi, \delta_2\phi)) = \delta_1(\not\delta_2\mathbf{E}) - \delta_2(\not\delta_1\mathbf{E})$$

⋮

Point 3 : Differential Noether charge

$$d(\not\delta\mathbf{Q}_{Noether}) = \Omega(\delta\phi, \delta_\chi\phi) + (\text{on-shell vanishing terms})$$

Key Point of Wald Formalism

[Wald, Lee-Wald, Iyer-Wald]

$$d(\Omega(\delta_1\phi, \delta_2\phi)) = \delta_1(\not\delta_2\mathbf{E}) - \delta_2(\not\delta_1\mathbf{E}) = d(-\delta_1(\not\delta_2\Theta) + \delta_2(\not\delta_1\Theta))$$

“Lagrangian-Based Prescription”

Extension to CS Term

Generalization to CS Term

- Minimum modification of Wald formalism to take into account

$$\delta_{\chi} \mathbf{I}_{CS} = \mathcal{L}_{\xi} \mathbf{I}_{CS} + \underline{d(\dots)} \quad \text{[Tachikawa]}$$

- In 5d and higher, an appropriate coordinate & gauge need to be taken to get desirable results ... ??????????

[Bonora et. al.]

Origin of Non-Covariance

$$\delta \Theta = \delta \Gamma^b_a \left(\underline{\frac{\partial \mathbf{I}_{CS}}{\partial \mathbf{R}^a_b}} \right) + \delta \mathbf{A} \left(\underline{\frac{\partial \mathbf{I}_{CS}}{\partial \mathbf{F}}} \right) + \dots \longrightarrow \begin{matrix} \text{Non-covariant} \\ \Omega \quad \delta \mathbf{Q}_{Noether} \end{matrix}$$

Chern-Simons BH Entropy

Manifestly Covariant Formalism

CS contribution to EoM \sim derivatives of anomaly polynomials

$$\rightarrow d(\Omega(\delta_1\phi, \delta_2\phi)) = \delta_1(\not{\delta}_2\mathbf{E}) - \delta_2(\not{\delta}_1\mathbf{E})$$

$$\rightarrow \Omega \not{\delta} \mathbf{Q}_{Noether} \sim \text{derivatives of anomaly polynomial}$$

\rightarrow Covariant !

“EoM-Based Prescription”

BH Entropy Formula for CS Term

For general anomaly polynomial $\mathbf{P}_{anom}(\mathbf{F}, \mathbf{R}) = \mathbf{F}^l \wedge \prod_i \text{tr} \mathbf{R}^{2k_i}$

$$S_{CS} = \int_H \sum_{k=1}^{\infty} 8\pi k \Gamma_N (d\Gamma_N)^{2k-2} \frac{\partial \mathbf{P}_{anom}}{\partial \text{tr} \mathbf{R}^{2k}} \quad \Gamma_N = \frac{1}{2} \epsilon_a{}^b \Gamma^b{}_a$$

(Generalized) Tachikawa's Proposal Proved Covariantly!

Gravity Dual of Anomalous Fluid

- Fluid/gravity solution for charged rotation AdS BH in Einstein-Maxwell theory (**up to 2nd order**)
- Back reaction due to CS term

$$ds^2 = -2u_\mu dx^\mu dr + r^2[-f(r, m, q)u_\mu u_\nu + P_{\mu\nu}]dx^\mu dx^\nu + \dots \\ + g_V(r, m, q)(u_\mu V_\nu + u_\nu V_\mu)dx^\mu dx^\nu + \dots$$

- Evaluation of diff. Noether charge at horizon & boundary
→ **Replacement rule reproduced!**

...but anomaly-induced transports are higher-order:

$$V^\mu = \epsilon^{\mu\nu\rho_1\cdots\rho_{2n-2}} u_\nu (\partial_{\rho_1} u_{\rho_2}) \cdots (\partial_{\rho_{2n-3}} u_{\rho_{2n-2}})$$

Is 2nd order enough?

- **“Non-renormalization” proved for any odd dim. by using CS diff. Noether charge \sim deriv. of anomaly polynomial !**

Summary

Anomaly Polynomials Play Crucial Roles!

CS Terms are Nice Class of Higher-Derivative Terms!

- Manifestly covariant diff. Noether charge for CS term
→ BH entropy formula for CS terms
- Systematic study of holography for anomalies in the hydrodynamics limit.
→ Replacement rule reproduced