

Top FCNC decays: review of theoretical expectations

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March, 2015

On the tree level only charged current top decays are allowed in the Standard Model

$$t \rightarrow W^+ b \quad \text{dominant, BR} = 91\%$$

$$t \rightarrow W^+ s/d \quad \text{CKM suppressed}$$

FCNC top decays are only possible on loop level.

Four two-particle final states can be considered:

$$t \rightarrow q\gamma, qZ, qg, qH \quad q = u, c$$

Current experimental limits are (RPP2014):

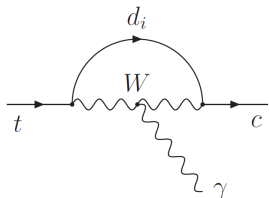
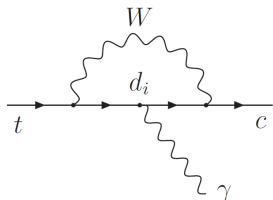
$$BR(t \rightarrow \gamma q) < 5.9 \cdot 10^{-3} \quad 95\% \text{ CL}$$

$$BR(t \rightarrow Z q) < 2.1 \cdot 10^{-3}$$

Leading order diagrams for FCNC decay $t \rightarrow c \gamma$

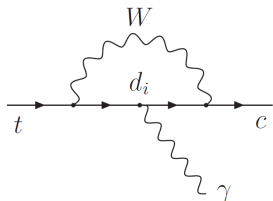
When neglecting down quark masses, the decay amplitude is suppressed (GIM):

$$\mathcal{M} \sim \sum_{d_i} V_{td_i}^* V_{cd_i} = 0$$



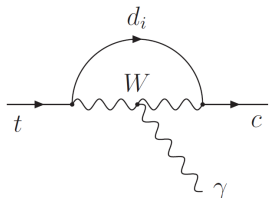
Leading order diagrams for FCNC decay $t \rightarrow c \gamma$

However, taking into account quark masses, GIM cancelation is not perfect



$$\mathcal{M} \sim \sum_{d_i} V_{td_i}^* V_{cd_i} \mathcal{F}(x_{d_i})$$

$$x_{d_i} = \frac{m_{d_i}^2}{M_W^2}$$



Leading order diagrams for FCNC decay $t \rightarrow c \gamma$

Assuming $m_d = m_s \ll m_b$ the leading contribution is:

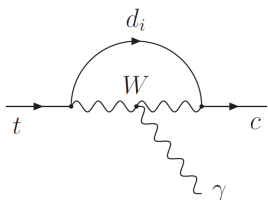
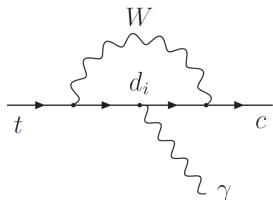
$$\mathcal{M} \sim V_{tb}^* V_{cb} [\mathcal{F}(x_b) - \mathcal{F}(0)]$$

Resulting decay width:

$$\Gamma(t \rightarrow c \gamma) \sim |V_{bc}|^2 \alpha_{em}^3 m_t \left(\frac{m_b}{M_W} \right)^4$$

Double suppression due to

- CKM: $|V_{bc}| \sim 0.04$
- GIM: $\frac{m_b}{M_W} \sim 0.04$



Standard Model expectations for the FCNC top decays (Snowmass 2013):

$$BR(t \rightarrow c \gamma) \sim 5 \cdot 10^{-14}$$

$$BR(t \rightarrow c Z) \sim 1 \cdot 10^{-14}$$

$$BR(t \rightarrow c g) \sim 5 \cdot 10^{-12}$$

$$BR(t \rightarrow c H) \sim 3 \cdot 10^{-15}$$

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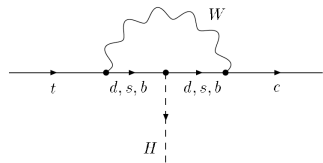
$$BR(t \rightarrow c g) \sim 5 \cdot 10^{-12}$$

$$BR(t \rightarrow c H) \sim 3 \cdot 10^{-15}$$

Same suppression mechanism in all channels (CKM+GIM).

Only for $t \rightarrow c H$ channel, GIM mechanism is not applicable (in one of the diagrams) due to Higgs coupling proportional to mass.

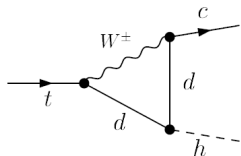
But the contribution of this diagram is still suppressed by $\frac{m_b}{M_W}$ (Higgs coupling)



Two Higgs Doublet Model

Probably the simplest possible extension of the SM.

Decay channel $t \rightarrow c h$ is affected by modified Higgs couplings:



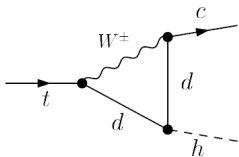
hdd :

$$g = g_{SM} \times (\sin(\beta - \alpha) - \tan \beta \cdot \cos(\beta - \alpha))$$

possible enhancement at large $\tan \beta$

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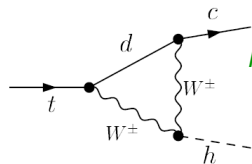
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hWW :

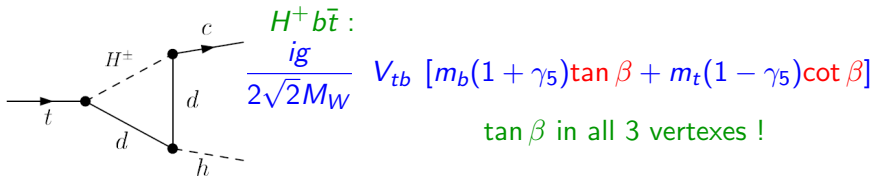
$$g = g_{SM} \times \sin(\beta - \alpha)$$

no enhancement possible

Two Higgs Doublet Model

New contributions to $t \rightarrow c h$ (as well as to $t \rightarrow c \gamma$, $c Z$, $c g$) from diagrams with H^\pm in the loop (instead of W^\pm).

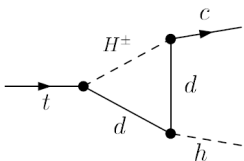
In case of 2HDM(II) (as an example):



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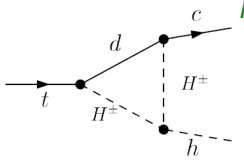


The diagram shows a top quark (t) entering from the left, splitting into a bottom quark (b) and a charm quark (c). The bottom quark and charm quark meet at a vertex where a H^\pm Higgs boson is exchanged. The H^\pm boson then splits into a bottom quark (b) and a top quark (t), which then meet at another vertex where a charm quark (c) and a Higgs boson (h) are produced.

$H^+ b \bar{t}$:

$$\frac{ig}{2\sqrt{2}M_W} V_{tb} [m_b(1 + \gamma_5)\tan\beta + m_t(1 - \gamma_5)\cot\beta]$$

$\tan\beta$ in all 3 vertexes !



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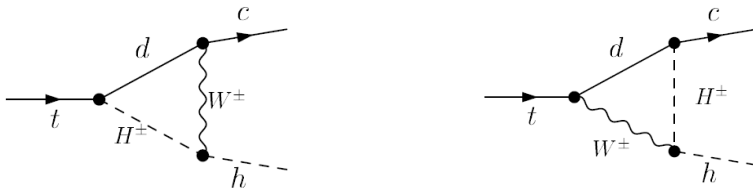
$H^\pm H^\pm h$:

$$\frac{-ig}{M_W} \left[\frac{M_h^2}{\sin 2\beta} (\cos^3\beta \cos\alpha - \sin^3\beta \sin\alpha) - M_{H^\pm}^2 \sin(\alpha - \beta) \right]$$

enhancement possible for both large and small $\tan\beta$

Two Higgs Doublet Model

One also has to consider diagrams with both H^\pm and W^\pm :



In the “standard” 2HDM scenarios, loop contributions can be enhanced significantly. However, FCNC remain suppressed at the tree level due to assumed flavour diagonal Higgs couplings.

However, one can also consider “non standard” scenarios, eg. “Top 2HDM”, where one of Higgs doublets couple to top quark only, where tree level FCNC couplings are possible!...

Expected maximal branching ratios for different scenarios

Significant differences between different papers - overall limit ranges given

Model	$BR(t \rightarrow c h)$	$BR(t \rightarrow c \gamma)$	$BR(t \rightarrow c g)$	$BR(t \rightarrow c Z)$
SM	$3 \cdot 10^{-15}$	$5 \cdot 10^{-14}$	$5 \cdot 10^{-12}$	10^{-14}
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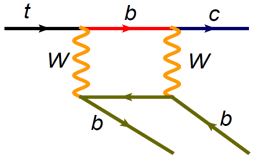
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Little Higgs	10^{-5}	$1.3 \cdot 10^{-7}$	$1.4 \cdot 10^{-2}$	$2.6 \cdot 10^{-5}$
Quark Singlet	$4.1 \cdot 10^{-5}$	$7.5 \cdot 10^{-9}$	$1.5 \cdot 10^{-7}$	$1.1 \cdot 10^{-4}$
Randal-Sundrum	10^{-4}	10^{-9}	10^{-10}	10^{-3}

Remark

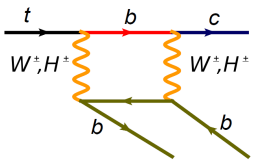
For both $t \rightarrow c h$ and $t \rightarrow c Z$ channels, non-resonant (box) diagrams have to be considered for each final state, in addition to the resonant channels.



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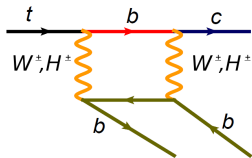


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For hadronic channels, one also has to consider dominant $t \rightarrow b W^+$ decay.

For the $cb\bar{b}$ final state all contributions should be considered together:

$$t \rightarrow c h \rightarrow c b\bar{b}$$

$$t \rightarrow c Z \rightarrow c b\bar{b}$$

$$t \rightarrow b W^+ \rightarrow b c\bar{b}$$

$$t \rightarrow c b\bar{b} \quad (\text{non-resonant})$$

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Decay $t \rightarrow c h$ in 2HDM is an interesting scenario, allowing for very large enhancement of signal with minimal extension of the model and no new final states - well suited for the case study.

Selected references for theoretical BR estimates in different models:

- *Snowmass 2013: Top quark working group report*, arXiv:1311.2028
- B.Yang, N.Liu¹, J.Han, arXiv:1308.4852
- I.Bauma, G.Eilamb, S.Bar-Shalomb, arXiv:0802.2622
- Hou Hong-Sheng, arXiv:hep-ph/0703067
- J.J.Cao, G.Eilam, M.Frank, et al., arXiv:hep-ph/0702264
- J.L.Diaz-Cruz, C.Pagliarone, arXiv:hep-ph/0612120
- J.A.Aguilar-Saavedra, arXiv:hep-ph/0409342
- S.Bejar, J.Guasch, J.Sola, arXiv:hep-ph/0101294
- Barbara Mele, arXiv:hep-ph/0003064