

# Sensitivity of neutrino telescopes to solar dark matter annihilations into light quarks, electrons and muons

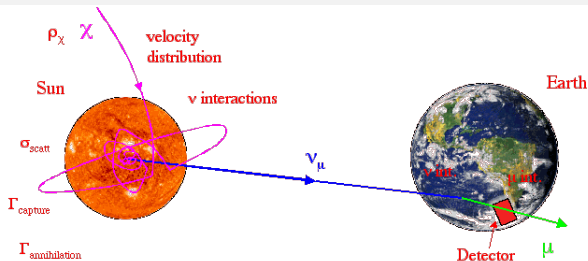
Sebastian Wild

Technische Universität München



MIAPP mini-workshop on dark matter and neutrinos, February 20, 2015

# Neutrinos from dark matter annihilations in the Sun



[Figure from J. Edsjo]

- If dark matter is a WIMP with non-zero  $\sigma_{\chi N}$ , it can be **captured** in the Sun
- After accumulating in the center of the Sun, dark matter starts to **annihilate**  
↪ among all stable SM particles in the final state, only neutrinos can escape the Sun
- Experimental signature: excess of  $\nu$ 's correlated with the direction of the Sun  
↪ Smoking-gun signature for particle dark matter!
- If capture and annihilation is in equilibrium, then the neutrino flux is  $\propto \sigma_{\chi N}$   
↪ results can be compared to direct detection experiments

# Stopping of particles in the Sun

Center of the Sun:  $\rho_{\text{matter}} \simeq 160 \text{ g/cm}^3$

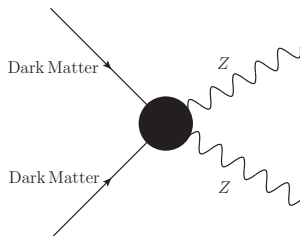
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by interactions with the solar medium!

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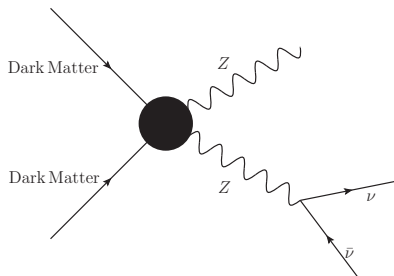


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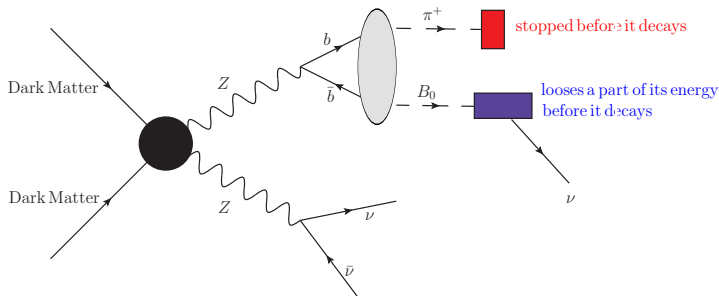
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•  $\tau_{B\text{-mesons}} \simeq \tau_{\text{interaction}}$

$\Rightarrow B$  mesons lose part of energy before decaying

•  $\tau_{\pi^\pm} \gg \tau_{\text{interaction}}$

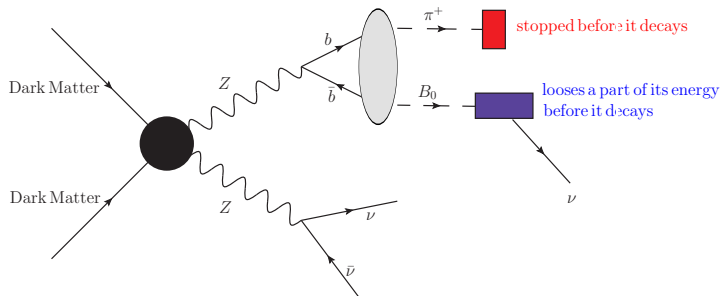
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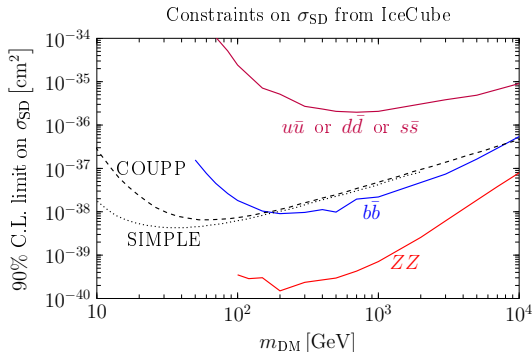
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- $\tau_{B\text{-mesons}} \simeq \tau_{\text{interaction}}$   $\Rightarrow B$  mesons lose part of energy before decaying
- $\tau_{\pi^\pm} \gg \tau_{\text{interaction}}$   $\Rightarrow \pi^+$  decay at rest,  $\pi^-$  are captured by nuclei
- $\tau_{\mu^\pm} \gg \tau_{\text{interaction}}$   $\Rightarrow \mu^\pm$  decay at rest

# Constraints on $\sigma_{SD}$ : standard annihilation channels



**“Naive”, first-order** classification of annihilation channels:

- $\nu\bar{\nu}$ ,  $ZZ$ ,  $W^+W^-$ ,  $\tau^+\tau^-$ : strong neutrino signal  $\Rightarrow$  strong constraints
- $b\bar{b}$ : softer neutrino spectrum, but still competitive constraints on  $\sigma_{SD}^p$
- $u\bar{u}$ ,  $d\bar{d}$ ,  $s\bar{s}$ : mainly produce pions  $\Rightarrow$  considerably weaker constraints
- $e^+e^-$ ,  $\mu^+\mu^-$ : no high-energy neutrinos at all  $\Rightarrow$  no signal (?)



# Constraining annihilations into $q\bar{q}, e^+e^-, \mu^+\mu^-$

Two (rather) recent ideas to constrain annihilations into  $q\bar{q}, e^+e^-, \mu^+\mu^-$ :

- (1) **Higher-order effects** can lead to the production of weak gauge bosons  
     $\hookrightarrow$  high energy  $\nu$ 's from  $Z, W^\pm$  decays

Ibarra, Totzauer, SW [1311.1418, 1402.4375]  
see also Baratella et. al. [1312.6408]

- (2) Search for the **MeV neutrinos** produced from decays of stopped  $\pi^+$  and  $\mu^\pm$

Bernal, Martin-Albo, Palomares-Ruiz [1208.0834]  
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$\hookrightarrow$  covered in the talk by Sergio Palomares-Ruiz!

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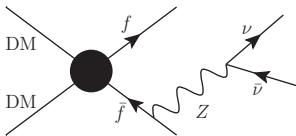
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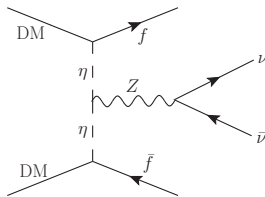
# Higher-order annihilations in the Sun

- Higher-order effects **generically** lead to the emission of  $Z$ ,  $W^\pm$  bosons in these annihilation channels! Examples are:

weak final state radiation



virtual internal bremsstrahlung

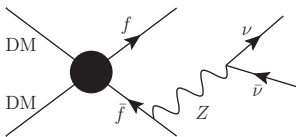


Generation of a  $\nu$  flux in previously unconstrained annihilation channels

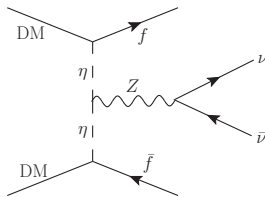
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Generation of a  $\nu$  flux in previously unconstrained annihilation channels

- A generic DM model falls into one of **two classes**:
  - dominant annihilation channel is  $DM DM \rightarrow f\bar{f}$
  - $DM DM \rightarrow f\bar{f}$  is helicity/velocity suppressed
- We derive separate limits on  $\sigma_{SD}$  for both scenarios

## Case a): weak final state radiation in DM DM $\rightarrow f\bar{f}$

### a) Models with dominant annihilation DM DM $\rightarrow f\bar{f}$

$\hookrightarrow$  relevant if no helicity suppression is present

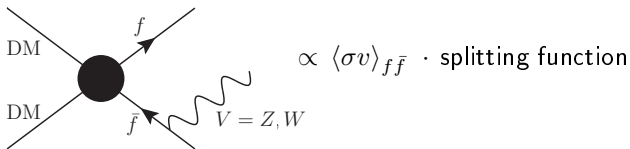
$\hookrightarrow$  relevant e.g. for Dirac DM or for Majorana DM annihilating into  $f_L\bar{f}_R$

- Most important source of neutrinos for  $m_{\text{DM}} \gtrsim 100$  GeV:

production of  $Z$  and  $W$  bosons via **weak final state radiation** (FSR)

$\hookrightarrow$  calculation is completely model-independent!

(We use PYTHIA 8.176 for that)

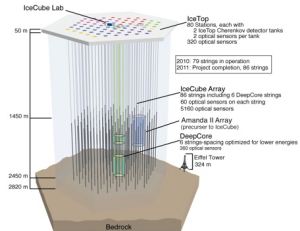
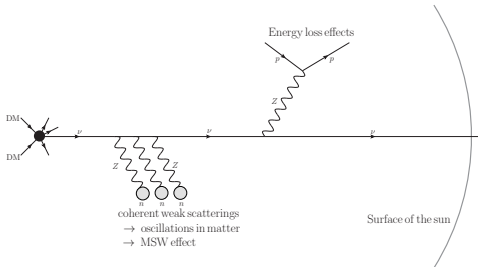


- Particularly interesting example: FSR splitting  $l_L^- \rightarrow W^- \nu_l$

$\hookrightarrow$  high-energetic  $\nu$  from the final state leg:  $E_{\nu_l} \simeq m_\chi$

# Propagation and detection of high-energy neutrinos

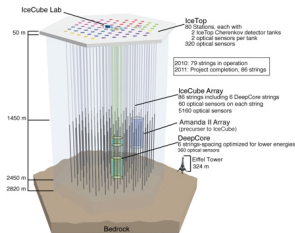
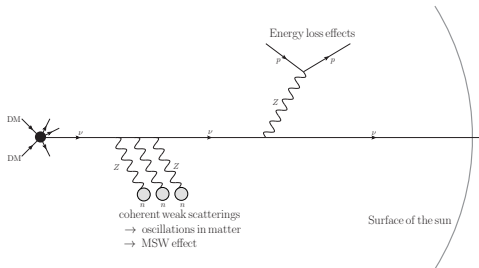
Once produced by DM annihilations, the neutrinos propagate out of the Sun:



- We take into account matter oscillation and absorption effects by using the event-by-event Monte Carlo *WimpSim* [Blennow, Edsjö, Ohlsson 0709.3898]

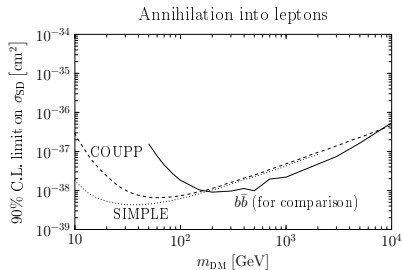
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- We derive constraints from the **IceCube 79-string data**, using a likelihood ratio test based on the number of events within  $\theta_{\text{Sun}} < \theta_{\text{max}}$   
↪ we (indirectly) include the spectral information of the  $\nu$  signal by optimizing  $\theta_{\text{max}}$  separately for every dark matter mass and annihilation channel

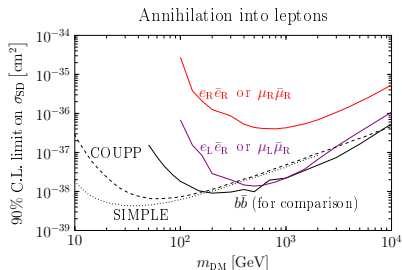
# Case a): weak final state radiation in DM $DM \rightarrow f\bar{f}$



Ibarra, Totzauer, SW [1402.4375]



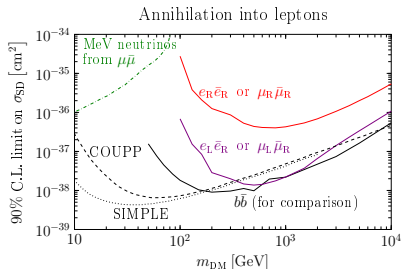
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↪ tighter constraints for annihilation into left-handed fermions

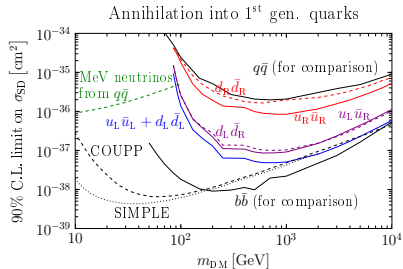
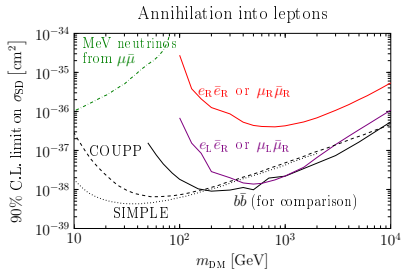
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- These results are complementary to the recent idea of using MeV neutrinos for these annihilation channels

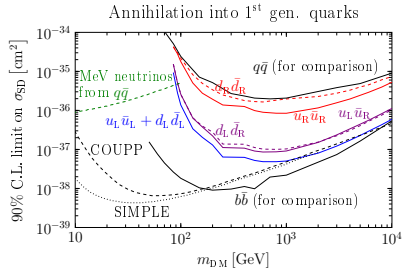
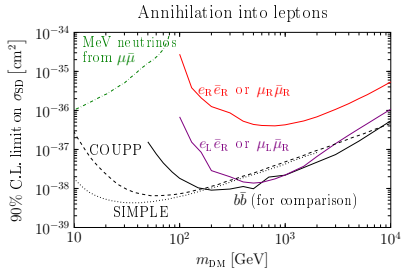
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In several cases, the resulting constraints are

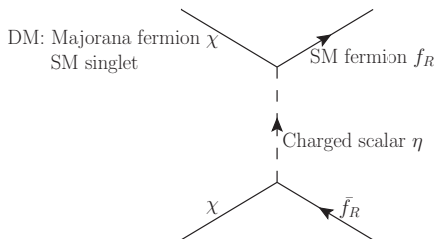
- significantly stronger** when taking into account weak FSR
- comparable to the limits from annihilation into e.g.  $b\bar{b}$
- stronger than the best direct detection limits on  $\sigma_{SD}$

## Case b): DM DM $\rightarrow f\bar{f}$ is helicity suppressed

### b) Models where DM DM $\rightarrow f\bar{f}$ is helicity/velocity suppressed

- $\hookrightarrow$  relevant e.g. for Majorana DM annihilating into  $f_L\bar{f}_L$  or  $f_R\bar{f}_R$
- $\hookrightarrow$  in this case, the  $\nu$ -flux is not model-independent anymore

#### • **Simplified model** for Dark Matter with helicity suppression:



$$\langle\sigma v\rangle_{f\bar{f}} \propto (\dots) \cdot (m_f/m_\chi)^2 + (\dots) \cdot v_{\text{rel}}^2$$

- Dark Matter particle: Majorana fermion  $\chi$ , SM singlet
- Next-to-lightest new particle: Charged scalar  $\eta$
- $\mathcal{L}_{\text{int}} = y \bar{\chi} f_R^\dagger \eta + \text{h.c.}$

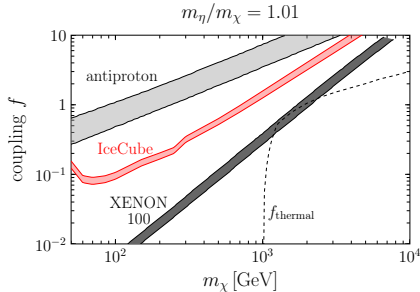
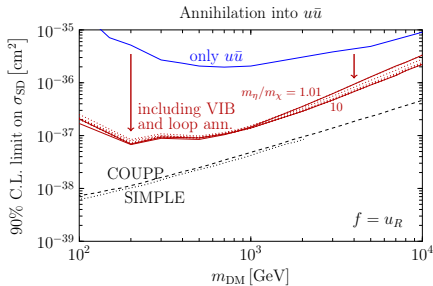
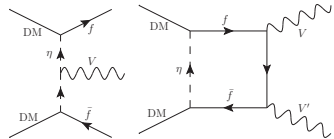
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One example: dark matter coupling to  $u_R$

$\hookrightarrow$  dominant annihilation channels:

virtual internal bremsstrahlung  $\chi\chi \rightarrow u\bar{u}V$

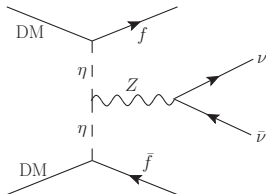
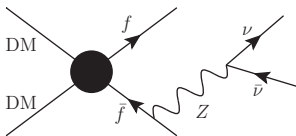
one loop annihilations  $\chi\chi \rightarrow VV$



Ibarra, Totzauer, SW [1311.1418]

Taking into account VIB and loop annihilations, the limits improve by  $\simeq 1$  order of magnitude with respect to pure  $u\bar{u}$  annihilation

# Conclusions



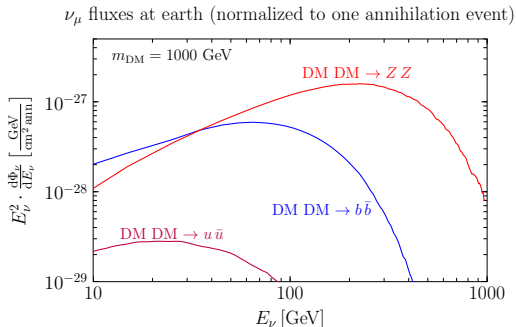
If dark matter annihilates dominantly into light quarks,  $e^-$  or  $\mu^-$ , no high-energetic neutrinos are produced at tree-level

- ↪ **Higher-order effects** generically lead to production of high-energetic  $\nu$ 's:  
weak final state radiation, virtual internal bremsstrahlung,  
one-loop annihilation into gauge bosons
- ↪ The resulting **limits on  $\sigma_{SD}$  from IceCube** are quite strong for some cases,  
and sometimes even better than the best direct detection limits on  $\sigma_{SD}$
- ↪ this hold for both helicity suppressed and non-suppressed scenarios

Backup slides



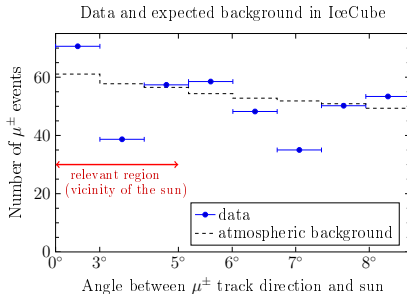
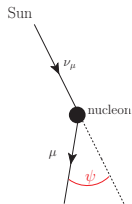
# Neutrino fluxes at earth: some examples



- $\text{DM DM} \rightarrow ZZ$ : hard spectrum due to prompt  $\nu$ 's from  $Z \rightarrow \nu\bar{\nu}$   
 $\hookrightarrow$  favorable for detection as  $A_{\text{eff}} \propto E_\nu^2$
- $\text{DM DM} \rightarrow b\bar{b}$ : softer spectrum, as  $B$  hadrons and mesons lose part of their energy before decaying
- $\text{DM DM} \rightarrow u\bar{u}$ : almost no high-energy  $\nu$ 's, as all pions and kaons are stopped in the Sun before decaying

# Constraints from IceCube 79-string data

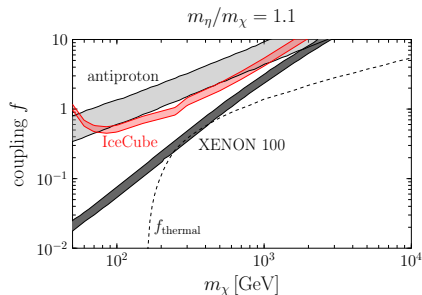
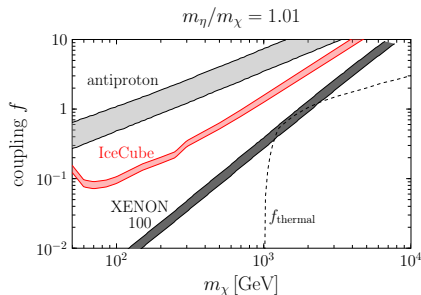
[IceCube collaboration 2013]



Assuming equilibrium, we use the IceCube 79-string data to derive limits on  $\sigma_{\text{SD}}^P$ :

- We use a likelihood ratio test to constrain the number of signal events within a given cone angle around the Sun
- We (indirectly) include the spectral information of the  $\nu$  signal by **optimizing the cone angle** around the Sun separately for every dark matter mass and annihilation channel

# Limits on the Yukawa coupling $f$ - comparing approaches



Ibarra, Totzauer, SW '13

- For a compressed spectrum, **IceCube constraints** on the Yukawa coupling  $f$  are **competitive**, in particular with PAMELA  $\bar{p}/p$   
↪ for this, taking into account the  $2 \rightarrow 3$  channels is crucial!
- XENON 100 constraints are still the most stringent one in this scenario
- More details in the paper: [Ibarra, Totzauer, SW \[1311.1418\]](#)
  - Scenario of coupling to  $b_R$
  - Thorough **analysis of the equilibration condition**, including its interplay with XENON 100 and collider limits