

SUSY Higgs-portal and X-ray lines

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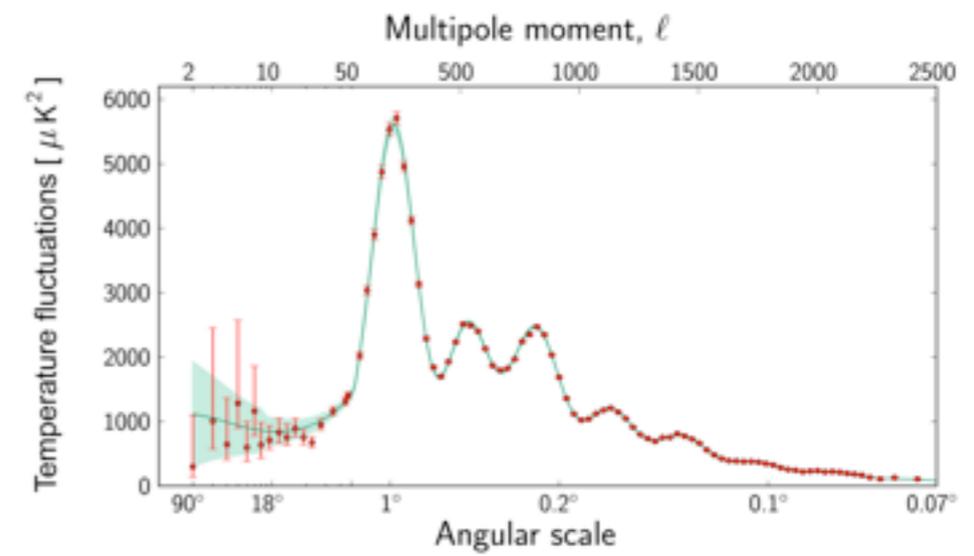
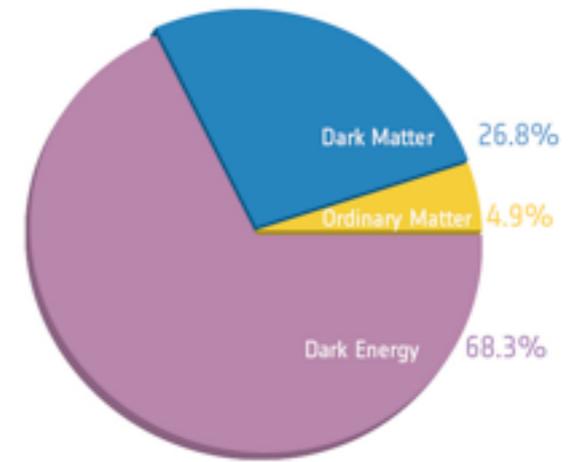
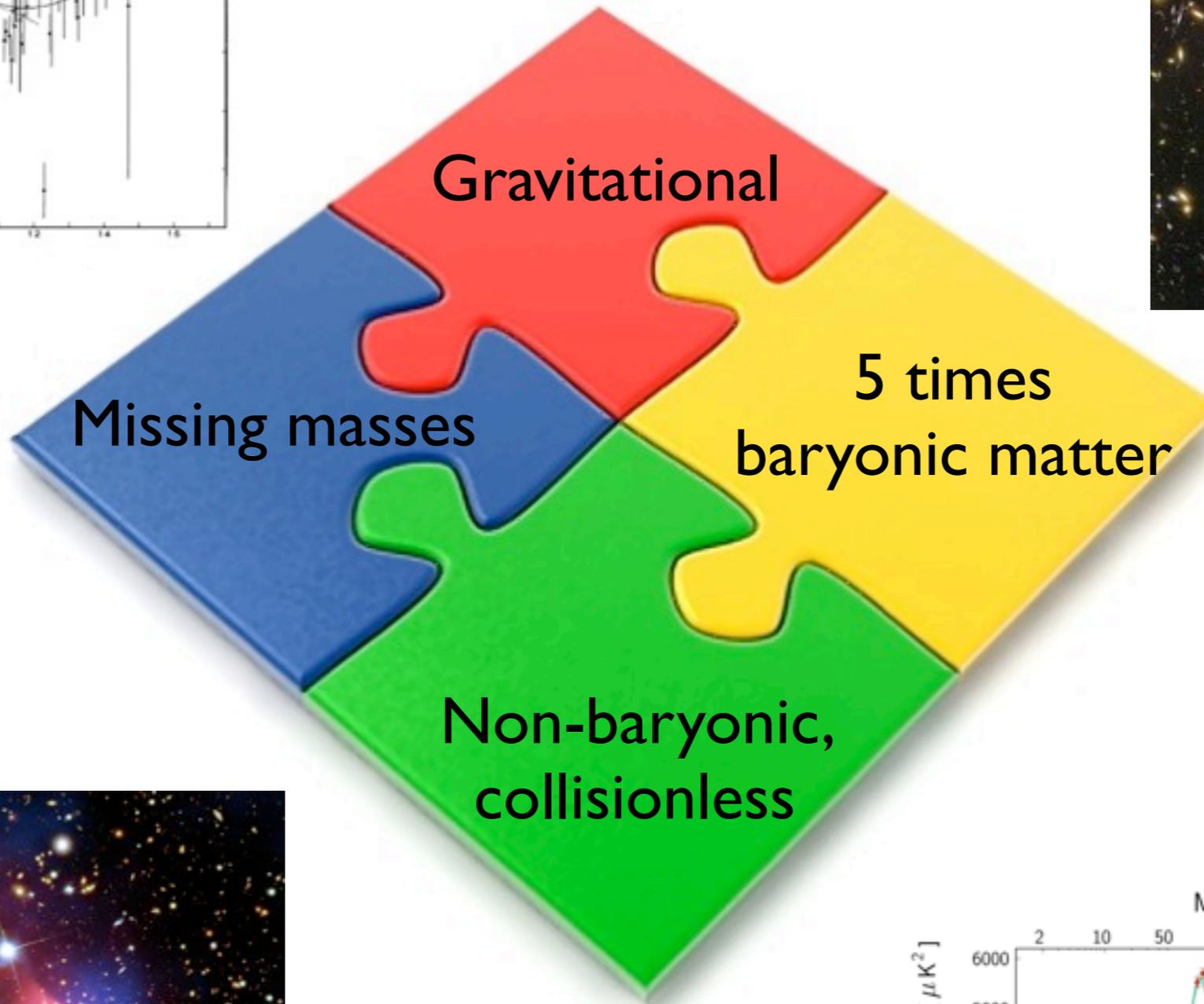
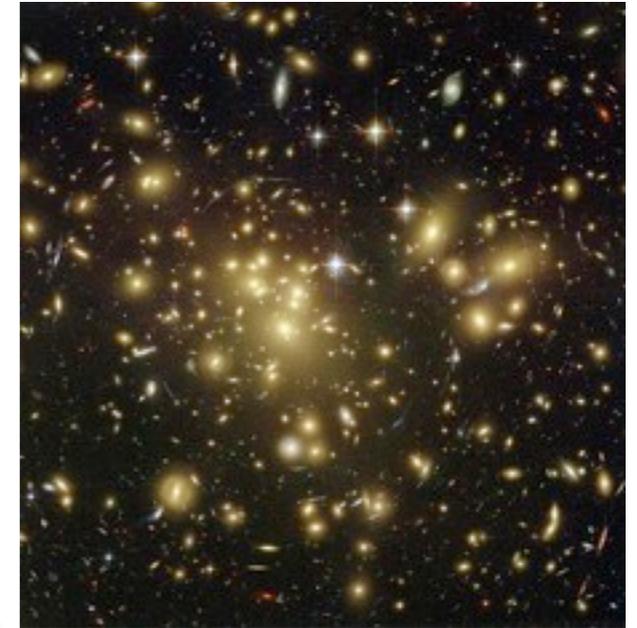
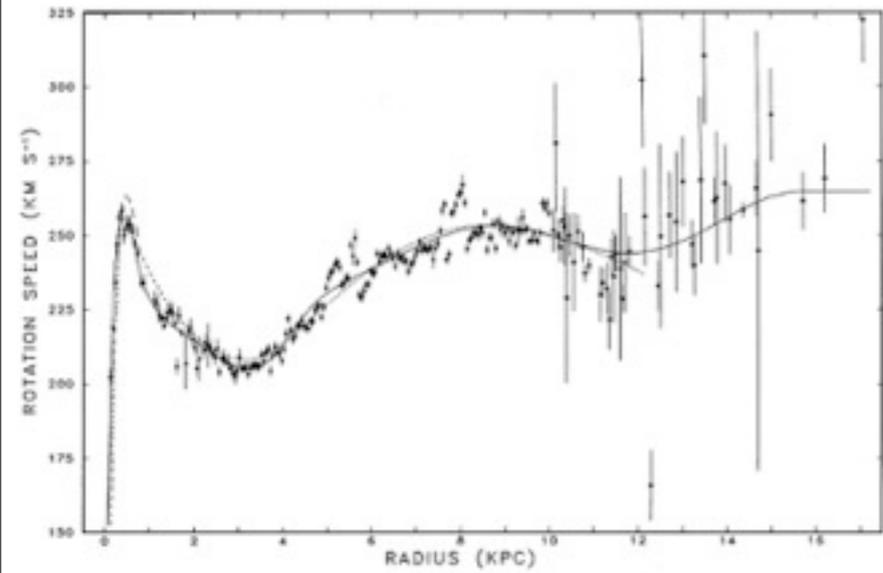
Ref. HML, C.B. Park, M. Park, 1501.05479 [hep-ph].

MIAPP Workshop on Dark Matter
Feb 23, 2015

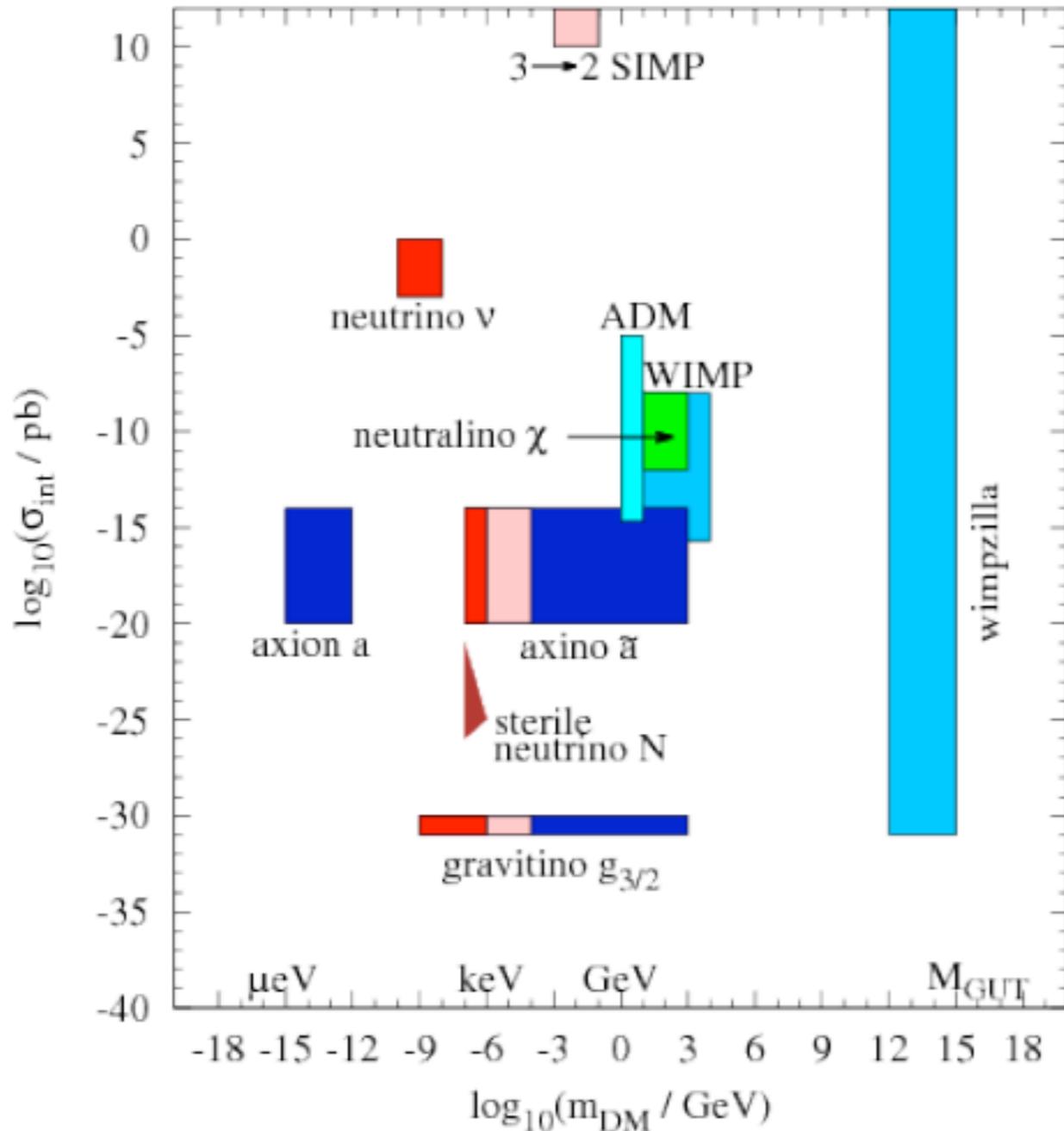
Outline

- Introduction
- Magnetic dark matter and X-ray lines
- Dark matter in SUSY Higgs-portal
- Conclusions

Dark Matter



Professor DM



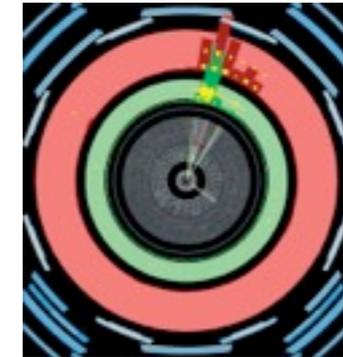
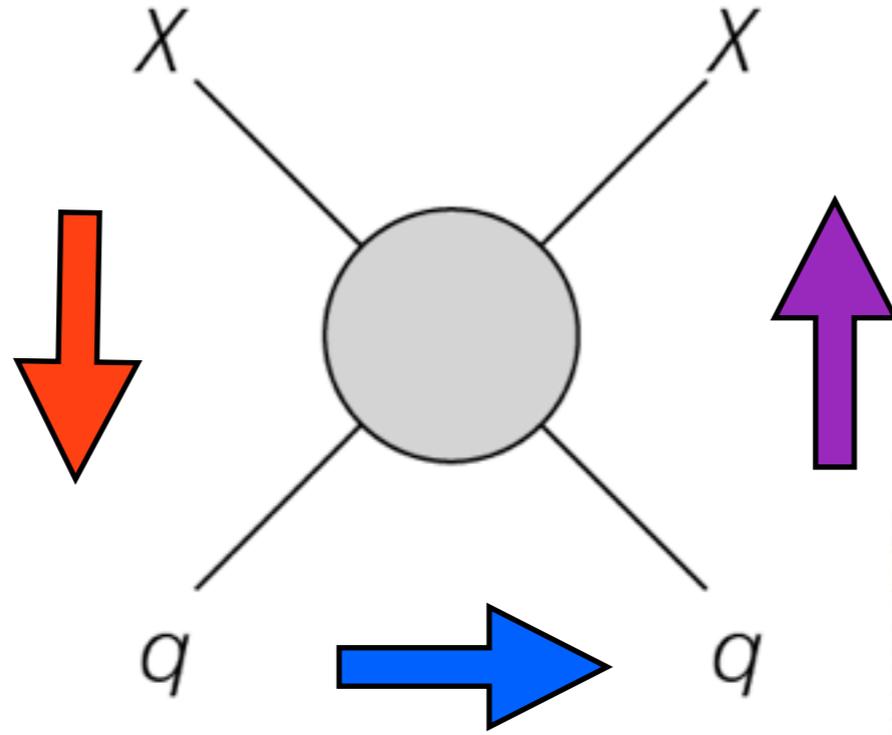
[Choi, Kim, Roszkowski, 1407.0017]

Are we simplifying it too much?

Testing dark matter

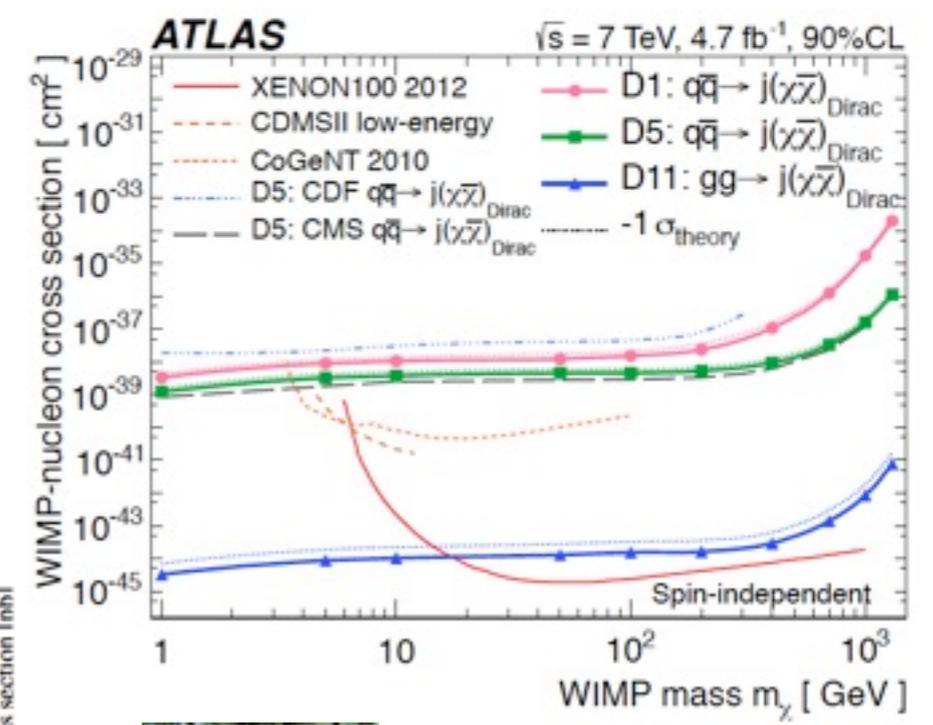
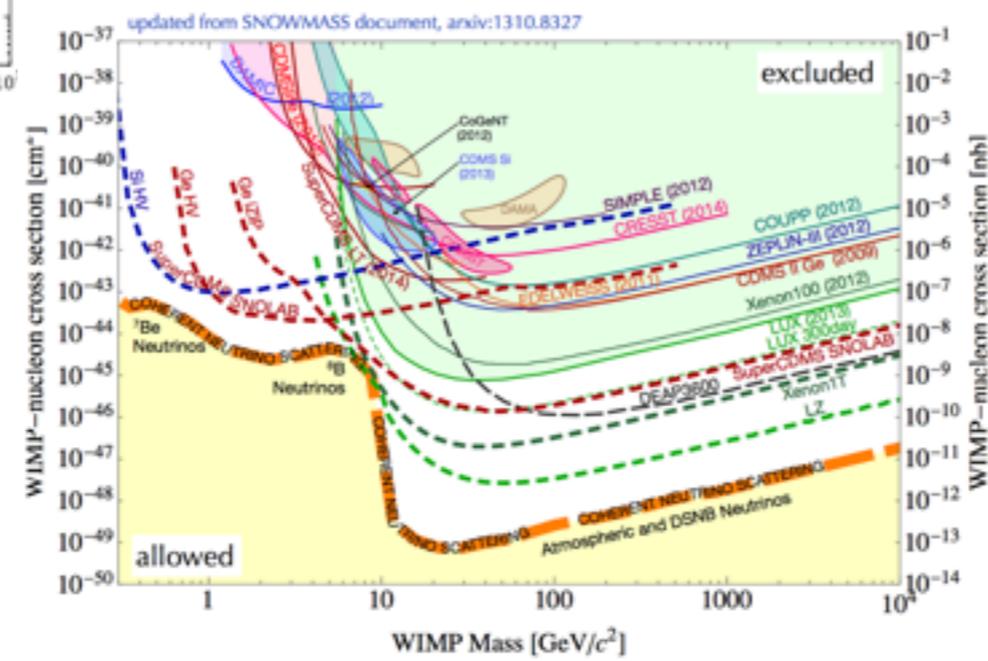
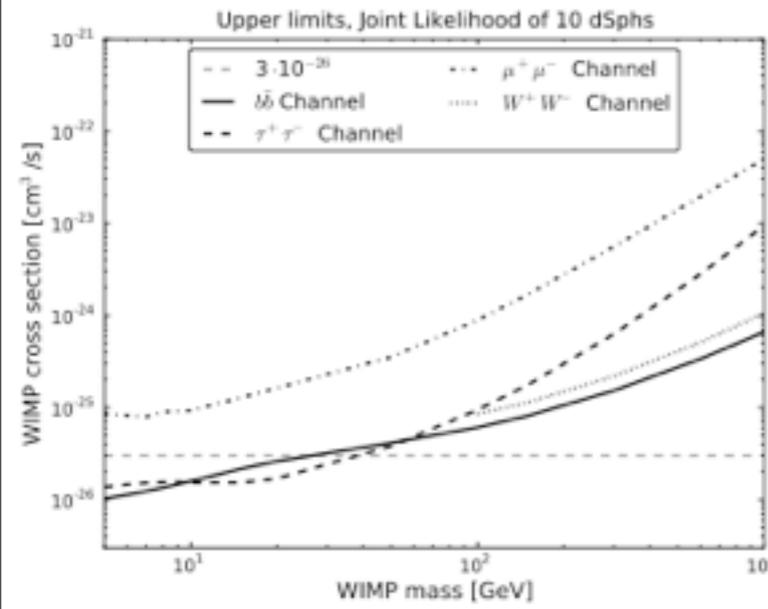


Indirect detection



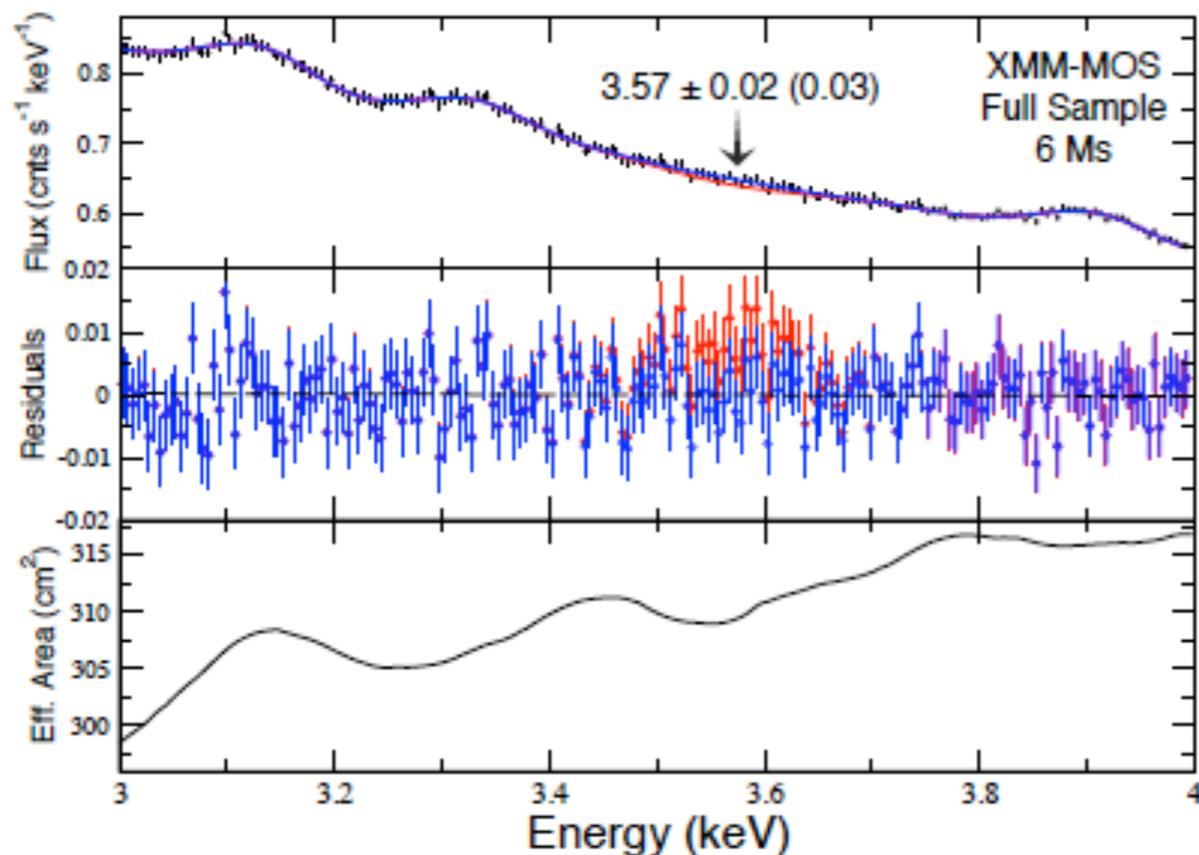
Particle colliders

Direct detection

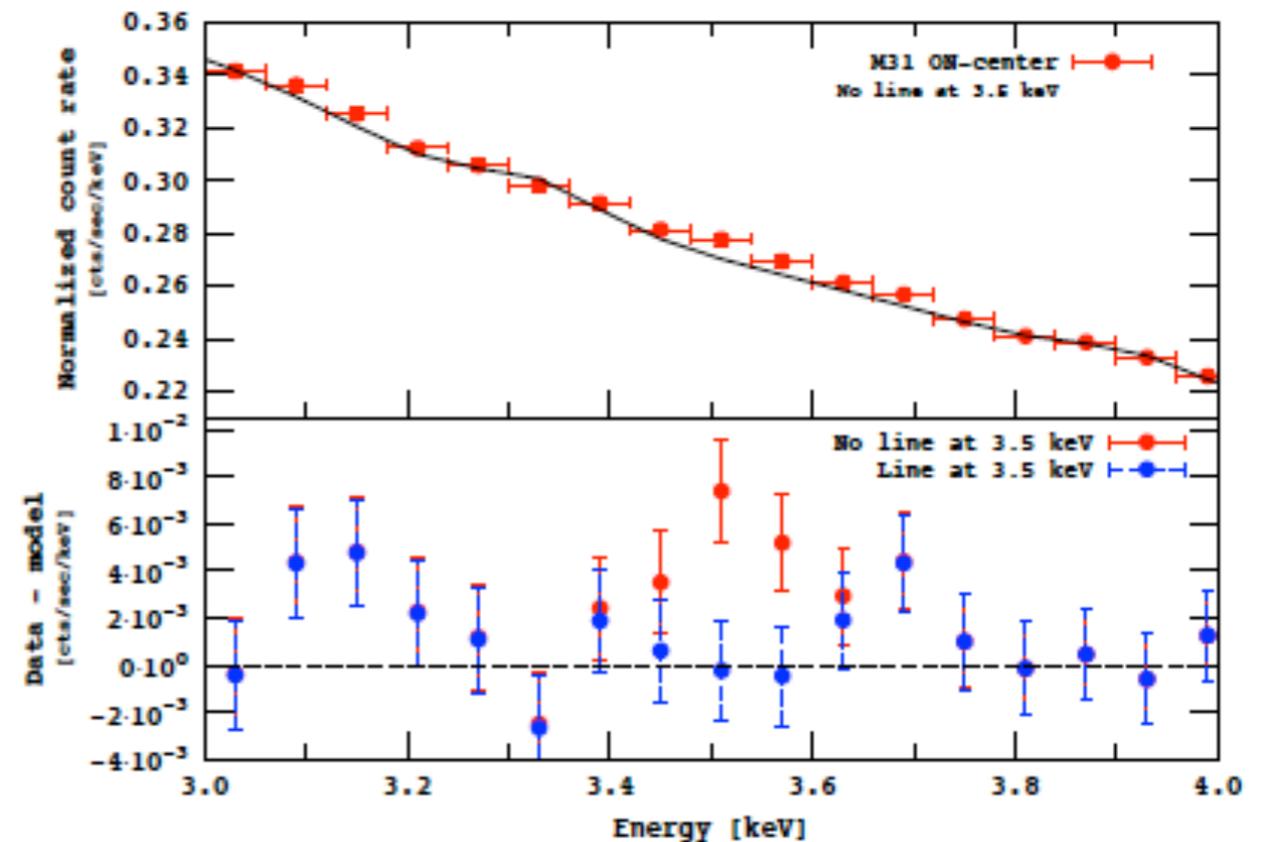


X-ray line from galaxies

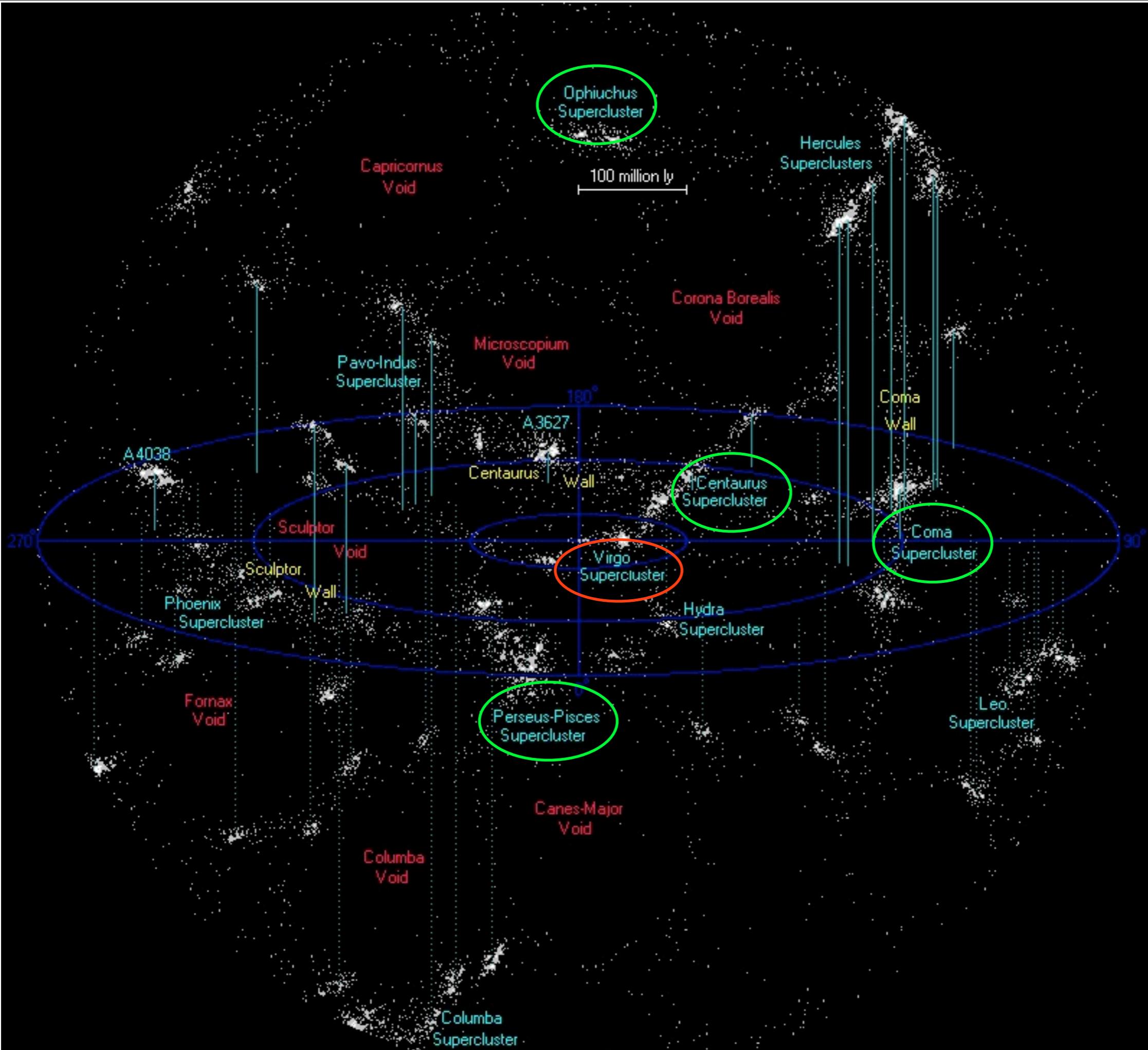
- Monochromatic photon lines probe well indirect signals of dark matter due to directional information.
- X-ray spectra (XMM-Newton) of various galaxy clusters and Andromeda show an unidentified line signal at 3.55 keV at $\sim 4\sigma$.



[Bulbul et al, 1402.2301]



[Boyarsky et al, 1402.4119;
1408.2503 (also Milky way)]



Signal samples

- Bulbul et al XMM-Newton



{ Perseus cluster
Nearby clusters: Coma, Centaurus, Ophiuchus
All stacked 73 galaxy clusters

➔ $E = 3.55 - 3.57 \pm 0.03 \text{ keV}$ at $\gtrsim 3\sigma$.

Chandra

excess from Perseus,
no line from Virgo cluster.

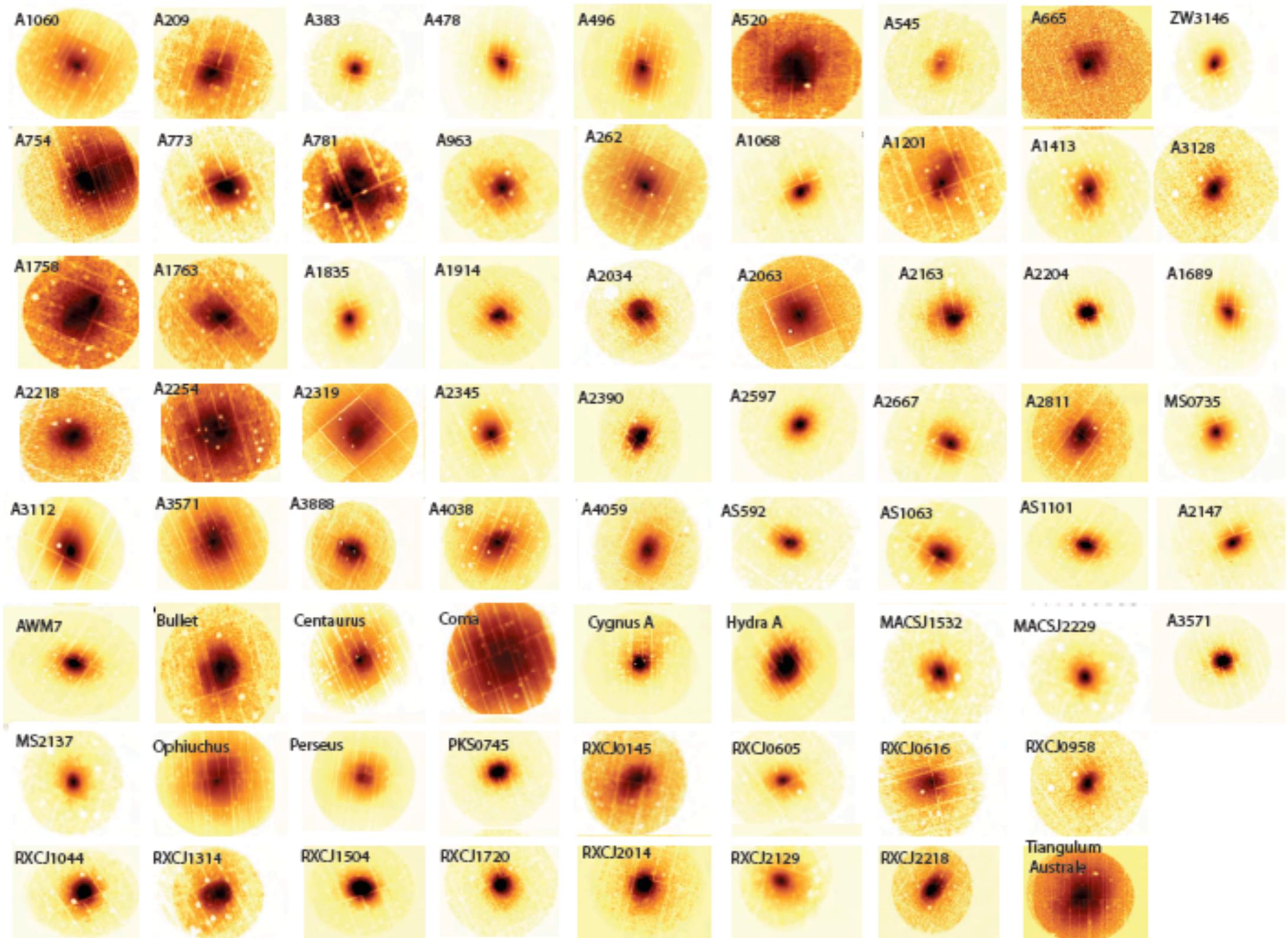


- Boyarsky et al XMM-Newton

{ Perseus cluster
Andromeda galaxy (M31)

➔ $E \sim 3.5 \text{ keV}$

Sample Selection



Taken from Randall Smith's talk (April 2014)

73 galaxy clusters for z=0.01-0.35

Perseus + Andromeda

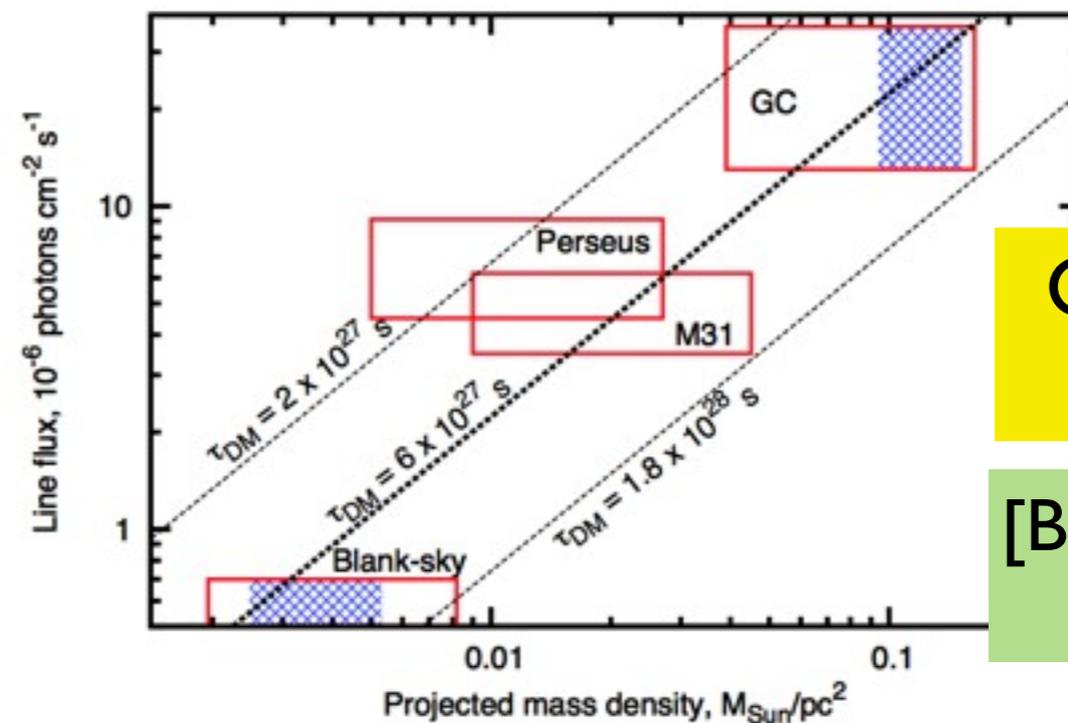
Sample	(1) Inst.	(2) Energy (keV)	(3) Flux (10^{-6} phots $\text{cm}^{-2} \text{s}^{-1}$)
Full Sample	XMM	MOS	3.57 ± 0.02 (0.03)
		PN	3.51 ± 0.03 (0.06)
		PN	3.57^*
Coma + Centaurus + Ophiuchus	XMM	MOS	3.57^*
		PN	3.57^*
Perseus (without the core)	XMM	MOS	3.57^*
		PN	3.57^*
Perseus (with the core)	XMM	MOS	3.57^*
		PN	3.57^*
All Other Clusters	XMM	MOS	3.57^*
		PN	3.57^*
Perseus	Chandra	ACIS-S	3.56 ± 0.02 (0.03)
		ACIS-I	3.56^*
Virgo?	Chandra	ACIS-I	3.56^*

Dataset	Exposure [ksec]	$\chi^2/\text{d.o.f.}$	Line position [keV]	Flux 10^{-6} cts/sec/cm ²	$\Delta\chi^2$
M31 ON-CENTER	978.9	97.8/74	3.53 ± 0.025	$4.9^{+1.6}_{-1.3}$	13.0
M31 OFF-CENTER	1472.8	107.8/75	3.53 ± 0.03	< 1.8 (2 σ)	...
PERSEUS CLUSTER (MOS)	528.5	72.7/68	$3.50^{+0.044}_{-0.036}$	$7.0^{+2.6}_{-2.6}$	9.1
PERSEUS CLUSTER (PN)	215.5	62.6/62	3.46 ± 0.04	$9.2^{+3.1}_{-3.1}$	8.0
PERSEUS (MOS) + M31 ON-CENTER	1507.4	191.5/142	$3.518^{+0.019}_{-0.022}$	$8.6^{+2.2}_{-2.3}$ (Perseus) $4.6^{+1.4}_{-1.4}$ (M31)	25.9 (3 dof)
BLANK-SKY	15700.2	33.1/33	3.53 ± 0.03	< 0.7 (2 σ)	...

$$F_{DM} \approx 2.0 \times 10^{-6} \frac{\text{cts}}{\text{cm}^2 \cdot \text{sec}} \left(\frac{\Omega_{fov}}{500 \text{ arcmin}^2} \right) \times$$

$$\left(\frac{S_{DM}}{500 M_{\odot}/\text{pc}^2} \right) \frac{10^{29} \text{ s}}{\tau_{DM}} \left(\frac{\text{keV}}{m_{DM}} \right) : \text{DM flux}$$

$$S_{DM} = \int_{l.o.s.} \rho_{DM}(r) dr : \text{DM column density}$$



Consistent signals!

[Boyarsky et al, 1408.2503]

Line backgrounds

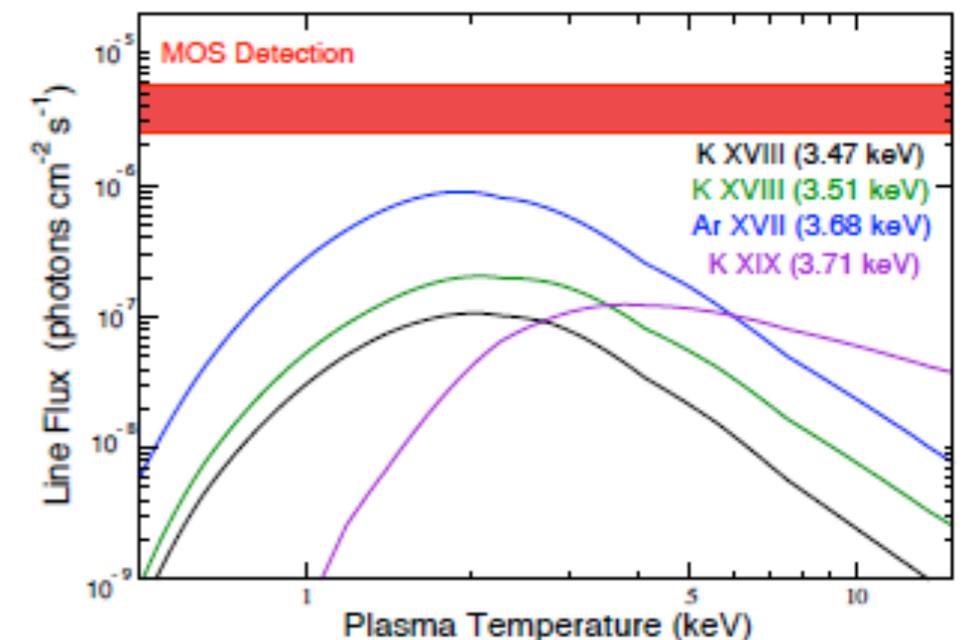
- K-shell excitation of ions leads to X-ray plasma emissions near 3.5 keV (K XVIII: 3.476 and 3.515 keV).
- Expected flux for thermal emission lines is smaller by factor 10-20.

Table 3

Estimated maximum fluxes of K XVIII at the rest energies 3.47 keV, 3.51 keV, Ar XVII at the rest energies 3.68 keV, and K XIX at the rest energy 3.71 keV lines obtained from AtomDB in the units of photons $\text{cm}^{-2} \text{s}^{-1}$. Estimates were performed based on best-fit fluxes obtained from the fluxes of S XVI, Ca XIX, and Ca XX lines in the *line-free* *apec* model. The fits were allowed to go a factor 3 above these estimates. The maximum flux for the Ar XVII DR at 3.62 keV line was initially set to 1% of the Ar XVII line at 3.12 keV in the spectral fits.

Sample	Inst.	Flux K XVIII (3.47 keV) (10^{-7})	Flux K XVIII (3.51 keV) (10^{-7})	Flux Ar XVII (3.62 keV) (10^{-7})	Flux Ar XVII (3.68 keV) (10^{-6})	Flux K XIX (3.71 keV) (10^{-6})
Full Sample	MOS	1.3 ± 0.7	3.5 ± 1.8	0.12	1.0 ± 0.5	1.2 ± 0.6
	PN	0.9 ± 0.4	1.8 ± 0.9	0.14	0.7 ± 0.3	0.3 ± 0.1
Coma + Centaurus + Ophiuchus	MOS	2.7 ± 2.1	8.2 ± 6.3	7.0	2.5 ± 1.9	5.2 ± 4.1
	PN	3.3 ± 2.3	6.8 ± 4.7	1.4	2.5 ± 1.8	0.8 ± 0.6
Perseus	MOS	18.5 ± 9.9	45.7 ± 24.4	6.4	15.1 ± 8.1	11.6 ± 6.2
	PN	13.8 ± 6.8	36.0 ± 17.8	1.99	10.8 ± 5.4	9.15 ± 4.5
All Other Clusters	MOS	0.5 ± 0.2	1.3 ± 0.5	0.10	0.4 ± 0.1	0.29 ± 0.1
	PN	1.3 ± 0.5	2.6 ± 0.9	0.90	1.1 ± 0.4	1.2 ± 0.4

[Bulbul et al, 1402.2301]



Debates on thermal emission

- Jeltima & Profumo I: Pottasium & Chlorine lines?
[1408.1699]

Measured emission in G.C. (XMM-Newton)

↓ **X** $\epsilon(T) = \epsilon(T_{\text{peak}}) \frac{N(T)}{N(T_{\text{peak}})}$

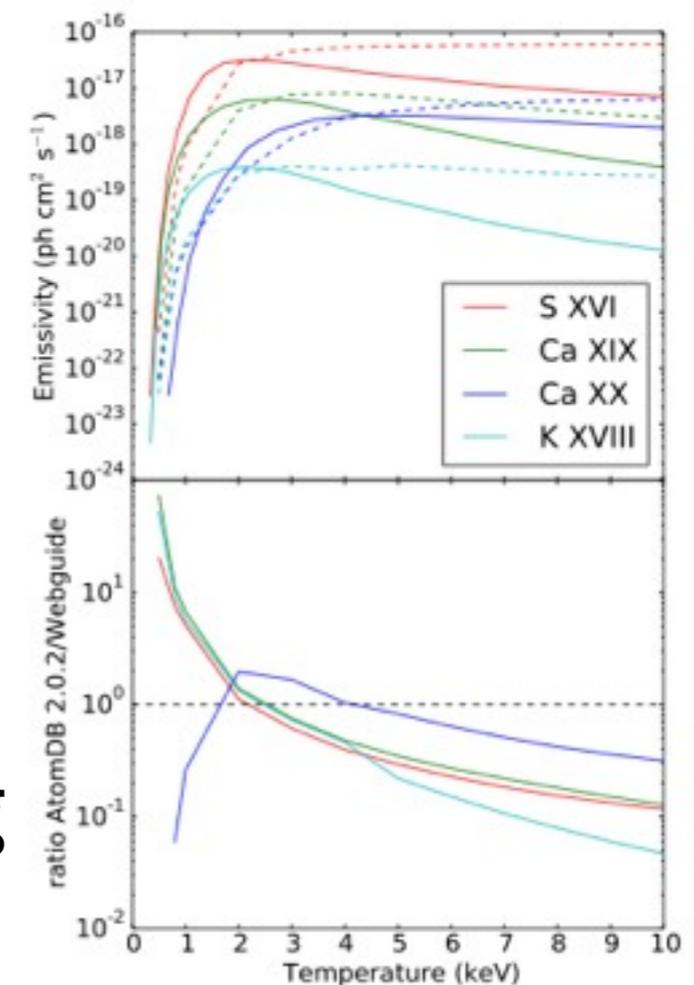
Multi-temperature emission calculated for Perseus & Andromeda (0.5-8 keV)

$\epsilon(T)$: emissivities

$N(T)$: number of ions

- Bulbul et al: incorrect spectral modeling away from peaks ? [1409.4143]

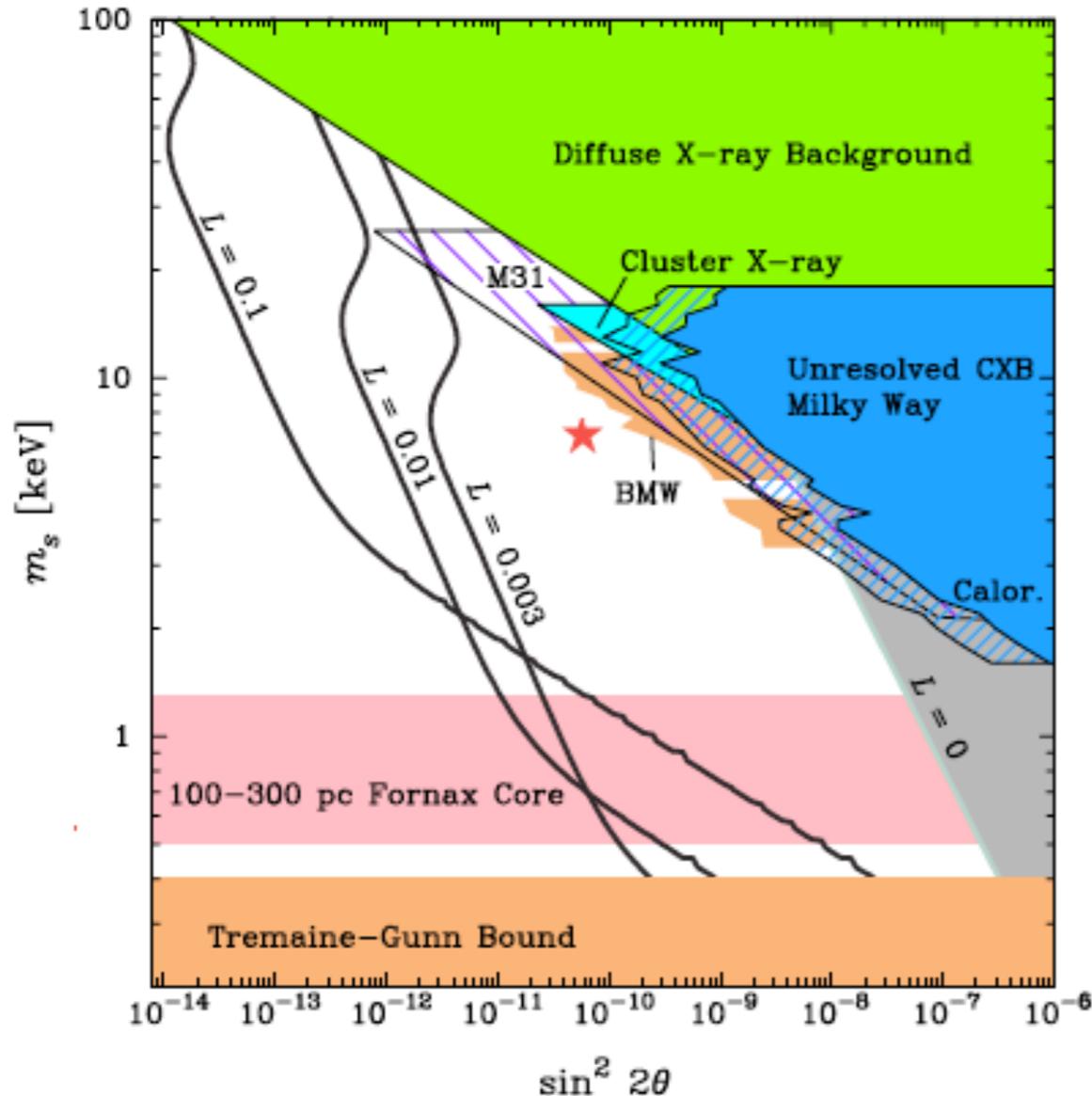
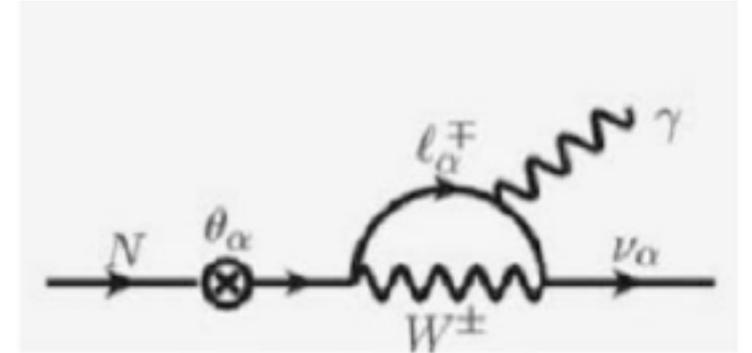
- Jeltima & Profumo II: line ratios for similar peak energies are only relevant and temperature-independent.



Sterile neutrino

- Minimal extension for neutrino masses needs three right-handed neutrinos.

$$\tau_{\text{DM}} = \frac{1024\pi^4}{9\alpha G_F^2 \sin^2(2\theta) m_{\text{DM}}^5} 7.2 \times 10^{29} \text{ sec} \left[\frac{10^{-8}}{\sin^2(2\theta)} \right] \left[\frac{1 \text{ keV}}{m_{\text{DM}}} \right]^5$$



Neutrino masses:

$$m_\nu \sim \frac{\lambda_N v^2}{M_N} \sim 0.05 \text{ eV},$$

X-ray line:

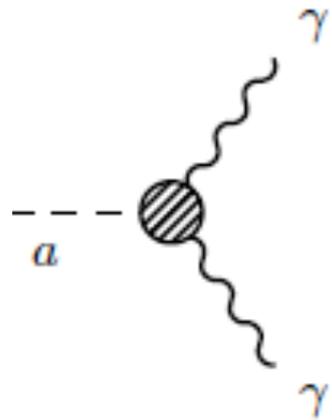
$$\theta^2 \sim \frac{\lambda_N^2 v^2}{M_N^2} \sim 10^{-11},$$

$$m_s \simeq M_N = 7.1 \text{ keV}.$$

Flavor structure is needed for suppressing a mixing angle smaller than $\theta^2 \sim 10^{-5}$.

Axion-like scalar

- Axion-like scalar decay:

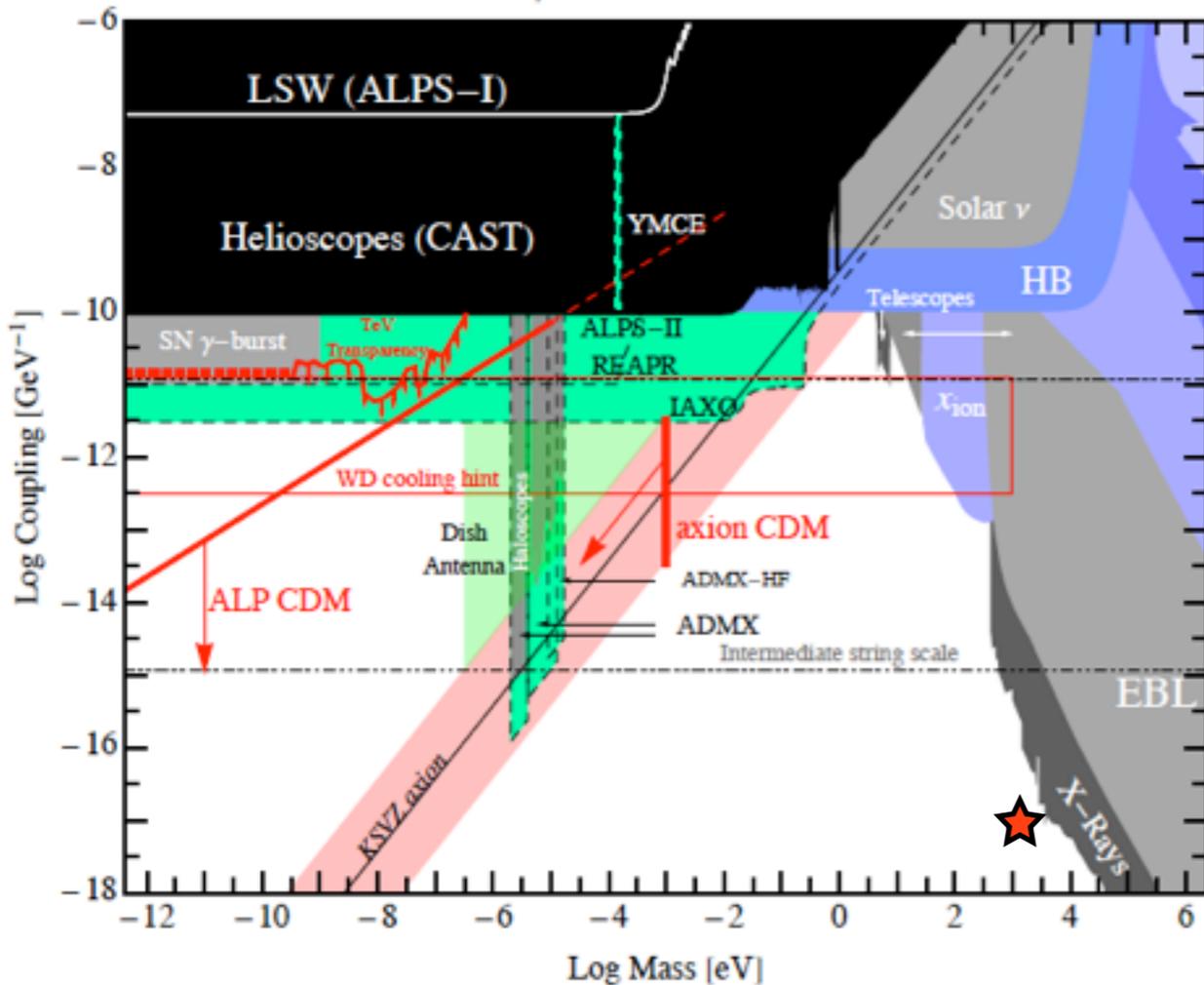


$$\Gamma(a \rightarrow \gamma\gamma) = \frac{\alpha m_a^3}{64\pi f_a^2}$$

[Higaki et al; HML, Park, Park; Jaeckel et al (2014)]

$$m_a = 7.1 \text{ keV},$$

$$f_a = 4 \times 10^{14} \text{ GeV}.$$



- Cold axion (misalignment)

$$Y_a \simeq 6.9 \times 10^{-5} \left(\frac{7 \text{ keV}}{m_a} \right) = \frac{1}{8} \frac{T_R}{m_a} \theta_{\text{osc}}^2 \left(\frac{f_a}{M_{\text{Pl}}} \right)^2$$

$$\Rightarrow T_R \lesssim 1 \text{ GeV} \left(\frac{0.4}{\theta_{\text{osc}}} \right)^2$$

- Hot axion (inflaton decay)

$$Y_a \simeq 6.9 \times 10^{-5} \left(\frac{7 \text{ keV}}{m_a} \right) = \frac{3}{4} \frac{T_R}{m_I} \frac{\text{Br}_{I \rightarrow aa}}{\text{Br}_{I \rightarrow \text{SM}}}$$

Structure formation: $\frac{T_R}{m_I} \gtrsim 0.2$

$$\Rightarrow \text{Br}_{I \rightarrow aa} \lesssim 4 \times 10^{-4}$$

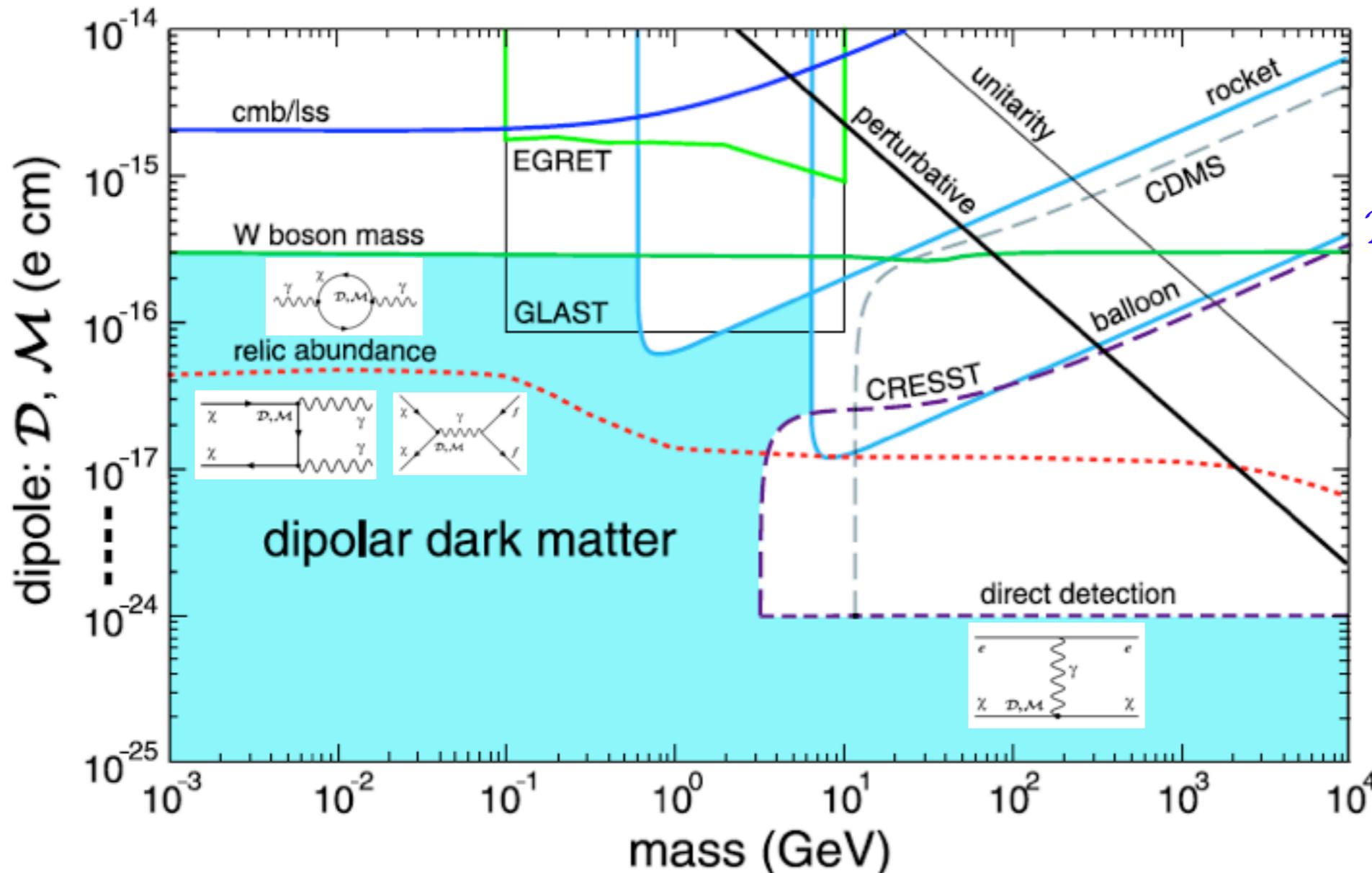
cf. CMB (isocurvature mode) $\Rightarrow H_I \lesssim 4.6 \times 10^9 \text{ GeV} \left(\frac{\theta_{\text{osc}}}{0.4} \right)$

Magnetic dark matter (MDM) and X-ray lines

Dipolar dark matter

- Dark matter is usually assumed neutral but it can have a small E/M dipole moment.

$$\mathcal{L}_{\gamma\chi} = -\frac{i}{2}\bar{\chi}\sigma_{\mu\nu}(\mathcal{M} + \gamma_5\mathcal{D})\chi F^{\mu\nu}, \quad \mathcal{D} \lesssim em_{\chi}^{-1} \simeq 2 \times 10^{-14} (m_p/m_{\chi}) e \text{ cm.}$$



$$\mathcal{D} \sim \mathcal{M} \lesssim 3 \times 10^{-16} e \text{ cm} \\ \sim 1.6 \times 10^{-5} \mu_B.$$

[Sigurdson et al,
astro-ph/0406355]

X-ray line from MDM

- Majorana fermion mDM vanishes. We consider a magnetic transition moment of Majorana fermion dark matter to another.

e.g. sterile-active neutrino transition, inelastic dark matter

$$\mathcal{L} = \frac{m_{\chi_2}}{\Lambda^2} \bar{\chi}_2 i\sigma^{\mu\nu} \chi_1 F_{\mu\nu}$$

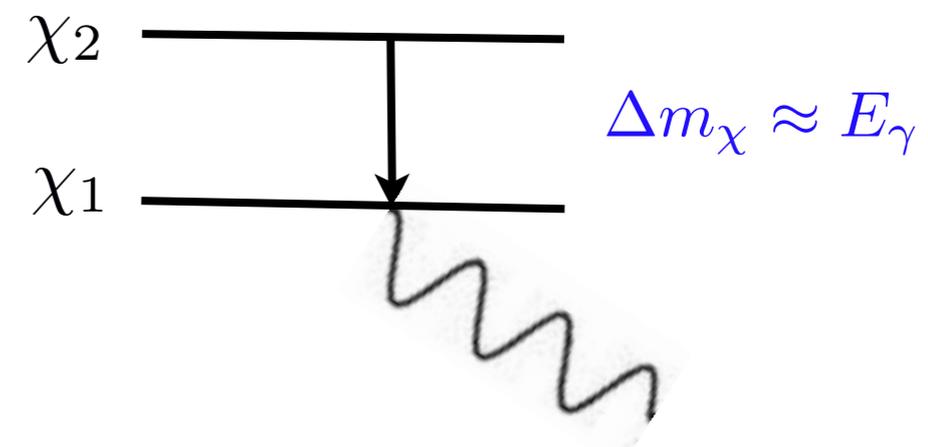
$$\Rightarrow \Gamma(\chi_2 \rightarrow \chi_1 + \gamma) = \frac{4m_{\chi_2}^2}{\pi\Lambda^4} E_\gamma^3, \quad E_\gamma = \frac{1}{2}m_{\chi_2} \left(1 - \frac{m_{\chi_1}^2}{m_{\chi_2}^2}\right)$$

X-ray line @ $E_\gamma = 3.55 \text{ keV}$

$$\Gamma_{\chi_2} = (0.36 - 3.3) \times 10^{-52} \text{ GeV} \left(\frac{m_{\chi_2}}{7.1 \text{ keV}}\right)$$

or $\tau_{\chi_2} = (0.20 - 1.8) \times 10^{28} \text{ s} \left(\frac{7.1 \text{ keV}}{m_{\chi_2}}\right)$

$$\Rightarrow m_{\chi_2} \simeq 10 \text{ GeV} : \Lambda = (5.0 - 8.7) \times 10^7 \text{ GeV.}$$



$$\mathcal{D} \sim 10^{-28} e \text{ cm}$$

Fermion-portal model

Two singlet Majorana fermions, distinguished by Z_2 , have different Yukawa couplings to SM RH fermion and charged scalar.

	SM RH fermion	Singlet fermions	Charged scalar	Singlet scalar	
	ψ_R	χ_{1L}	χ_{2L}	ϕ	S
$U(1)_Y$	q_{ψ_R}	0	0	q_{ψ_R}	0
$U(1)_X$	0	+1	-1	+1	-2
Z_2	+	+	-	-	+

$$-\mathcal{L}_{\text{DM}} = (\lambda \bar{\psi} P_L \chi_2 \phi + m_\psi \bar{\psi}_R \psi_L + \text{c.c.}) + m_\phi^2 |\phi|^2 + \frac{1}{2} y_1 S \bar{\chi}_1 P_L \chi_1 + \frac{1}{2} y_2 S^* \bar{\chi}_2 P_L \chi_2 + \frac{\kappa}{\Lambda_{UV}} S \bar{\psi} P_L \chi_1 \phi + \text{c.c.}$$

Explicit breaking of Z_2

$$\langle S \rangle \neq 0 \longrightarrow \frac{1}{2} m_{\chi_1} \bar{\chi}_1 \chi_1 + \frac{1}{2} m_{\chi_2} \bar{\chi}_2 \chi_2 + \epsilon \bar{\psi} P_L \chi_1 \phi, \quad \epsilon = \frac{\kappa \langle S \rangle}{\Lambda_{UV}} \ll 1.$$

$U(1)_X$ SSB

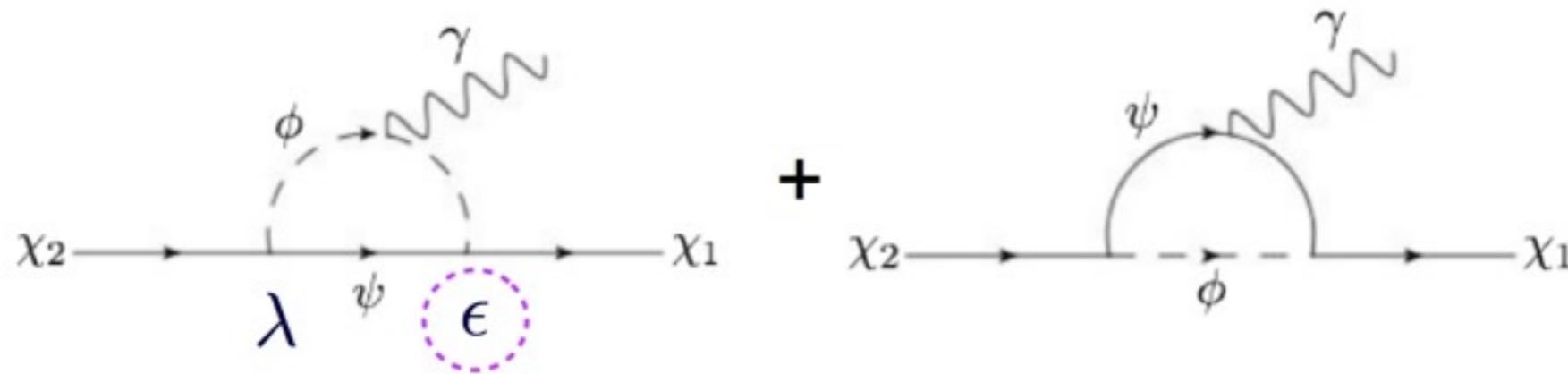
Similar Yukawa couplings \longrightarrow Almost degenerate masses

$$y_1 \approx y_2$$

$$m_{\chi_1} \approx m_{\chi_2}$$

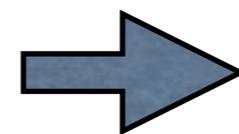
MDM + Thermal DM

- The same fermion-portal coupling and charged scalar are responsible for DM transition and annihilation.

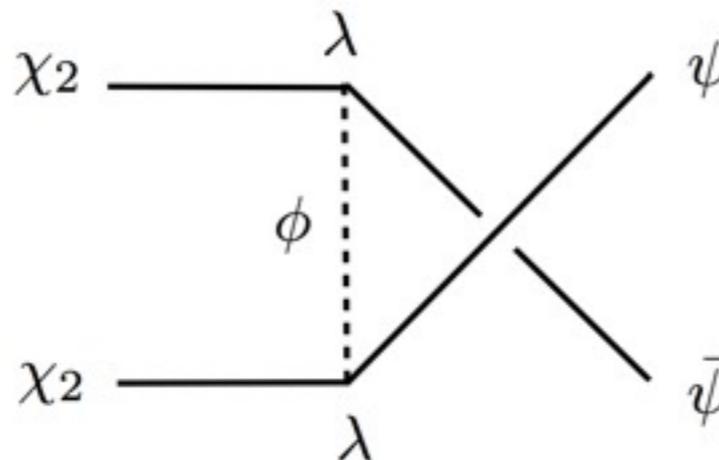
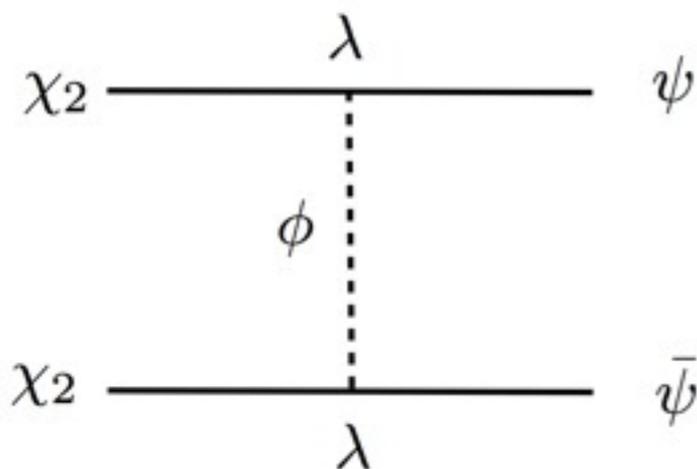


$$\Lambda_{UV} = \sqrt{\frac{64\pi^2}{e|q_\phi|} \frac{m_\phi}{\sqrt{|\text{Im}(\epsilon^* \lambda)|}}}$$

+ fermion lines flipped



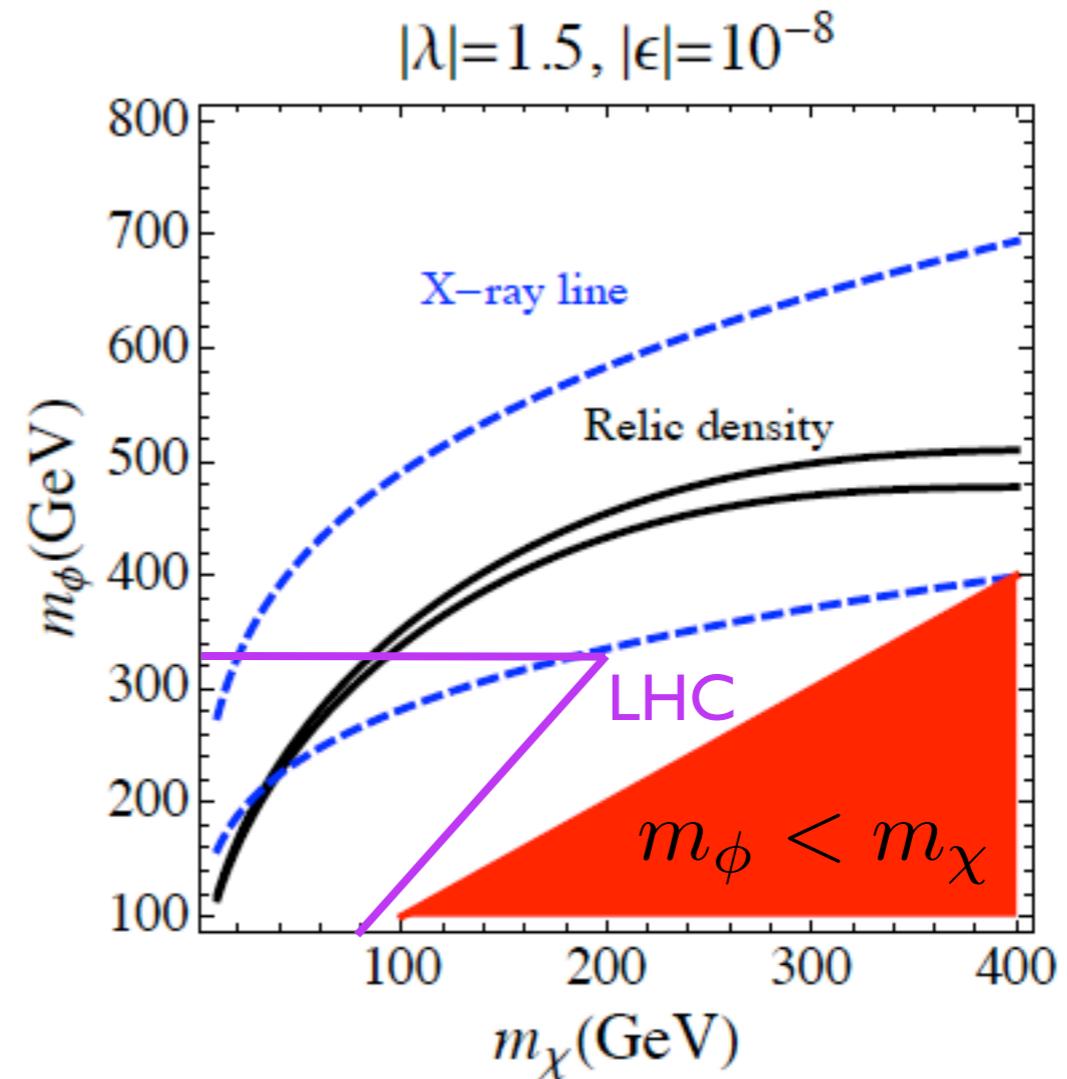
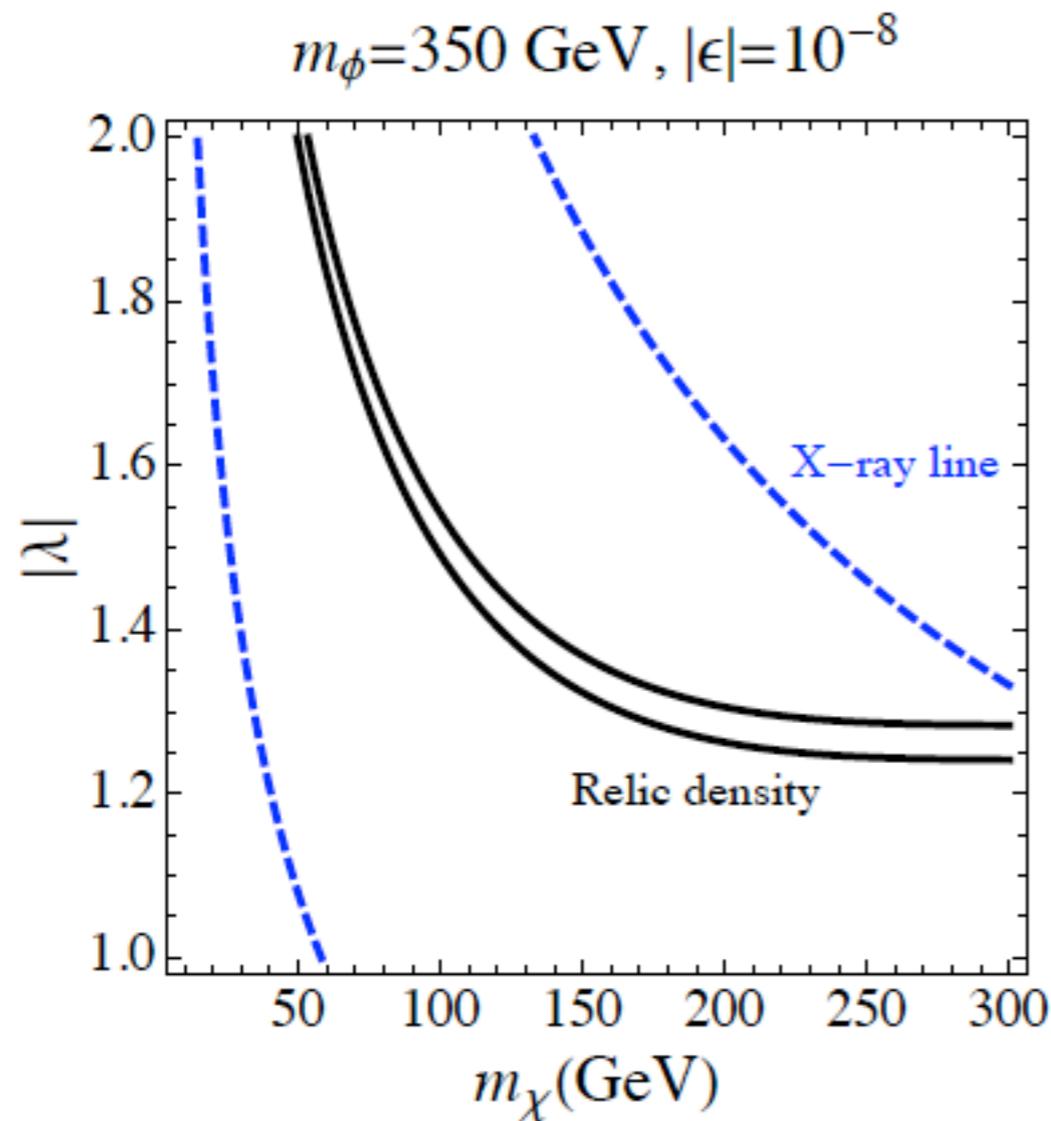
$$\lambda = \mathcal{O}(1), \quad \epsilon = \frac{\alpha \langle S \rangle}{\Lambda_{UV}} \sim 10^{-9} - 10^{-7}$$



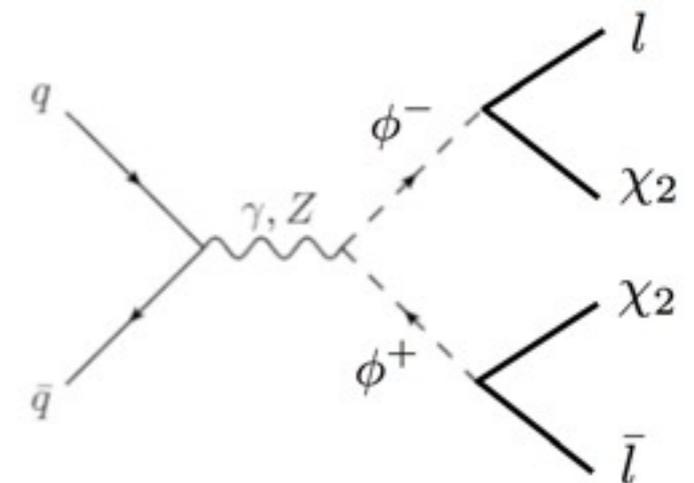
p-wave, but relevant at freeze-out.

$$\langle \sigma v \rangle = \frac{|\lambda|^4}{16\pi} \frac{m_{\chi_2}^2 (m_{\chi_2}^4 + m_\phi^4)}{(m_{\chi_2}^2 + m_\phi^2)^4} \cdot \frac{T}{m_{\chi_2}}$$

Relic density+X-ray line+LHC



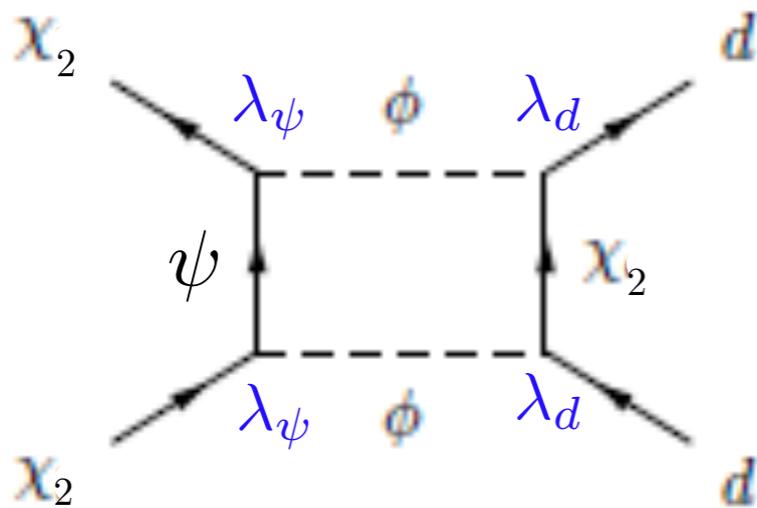
- A wide parameter space exists.
- Slepton-like charged scalar masses below 325 GeV excluded by LHC.



Bounds on DM

- **For baryophilic dark matter**, the WIMP-nucleon spin-independent cross section can be constrained by LUX.

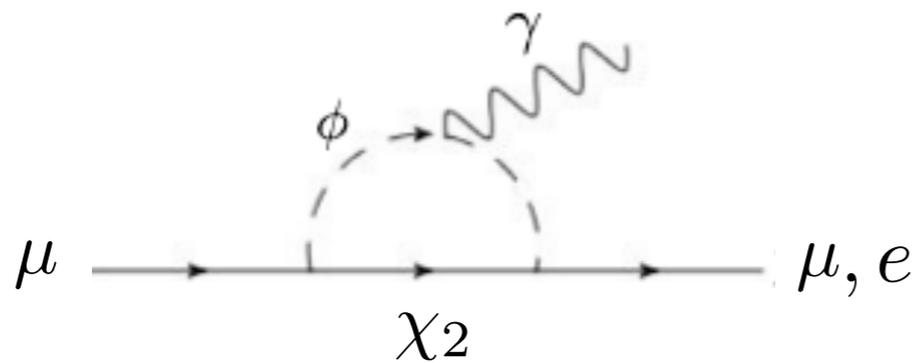
[Batell et al (2014); Nelson et al (2014)]



$$\mathcal{L}_{\text{box}} \sim \frac{\lambda_d^2 \lambda_\psi^2 m_\chi m_d}{32\pi^2 m_\phi^2 (m_\chi^2 - m_\psi^2)} \bar{\chi} \chi \bar{d} d.$$

$\lambda_\psi = 1.5, m_\phi = 400 \text{ GeV}, m_\chi = 200 \text{ GeV}$
 $\Rightarrow \lambda_d \lesssim 4.4$

- **Leptophilic dark matter** is bounded by muon $g-2$ and lepton flavor violation.



$$|\Delta a_\mu| < 3.45 \times 10^{-12} \quad (\text{BNL821})$$

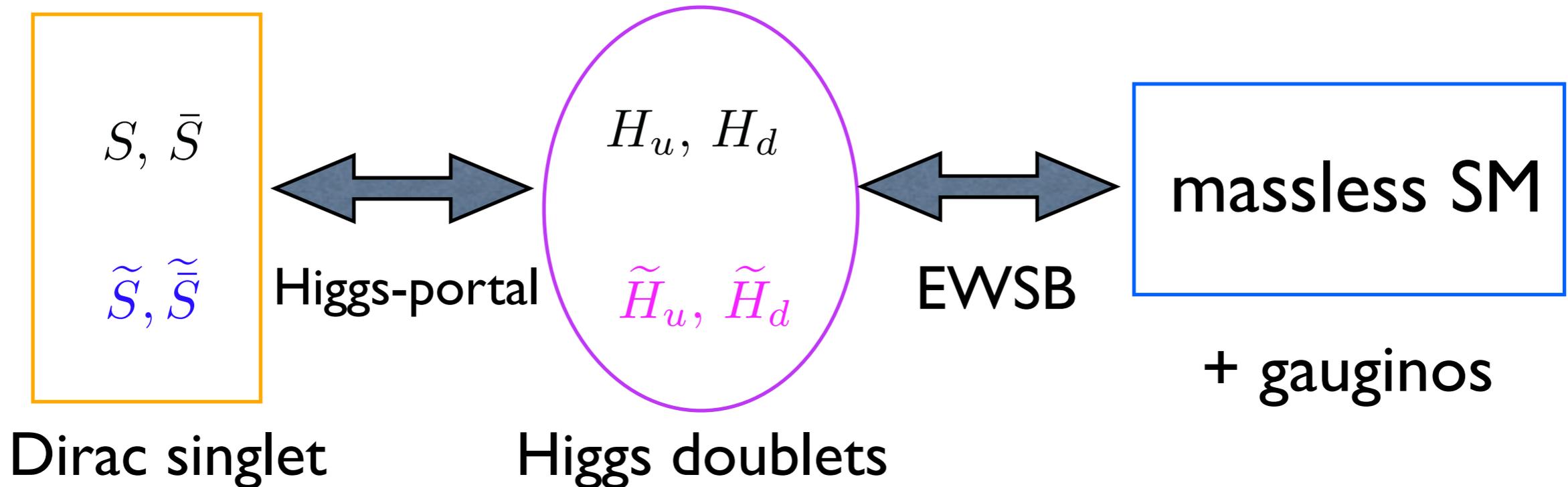
$$\Rightarrow |\lambda_\mu| < 0.27 (m_\phi / 500 \text{ GeV}).$$

$$\text{BR}(\mu \rightarrow e\gamma) < 5.7 \times 10^{-13} \quad (\text{MEG})$$

$$\Rightarrow |\lambda_e| < 0.03 (m_\phi / 500 \text{ GeV}).$$

Dark matter in SUSY Higgs-portal

SUSY Higgs-portal



- SUSY and chiral (PQ) symmetry leads to light Higgsino & singlet fermions, which interact with each other by SUSY Higgs-portal.
- R symmetry can make gauginos light too, while scalar superpartners are decoupled.

→ “Split” SUSY with extra singlets

Dirac singlino

- Superpotential with approximate PQ:

$$W = \lambda_S S H_u H_d + \lambda_{\bar{S}} \bar{S} H_u H_d$$

PQ symmetric

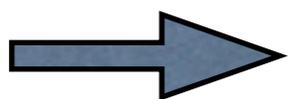
$$+ \mu_H H_u H_d + M_S \bar{S} S + \frac{1}{2} M_1 S^2 + \frac{1}{2} M_2 \bar{S}^2$$

PQ violating

e.g. $W \supset \frac{1}{M_P} P^2 H_u H_d + \frac{1}{M_P} Q^2 \bar{S} S + \dots$

- We assume a singlet flavor symmetry (e.g. U(2)) that singlino gets Dirac mass only.

$$M_1 = M_2 = 0, \quad \text{small } \lambda_S, \lambda_{\bar{S}}$$



Dirac singlino gets a small mass splitting only after EWSB.

Singlino mass splitting

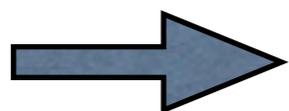
- Neutralino mass matrix for $(\tilde{B}, \tilde{W}^0, \tilde{H}_d^0, \tilde{H}_u^0, \tilde{S}, \tilde{\bar{S}})$

$$M_{\tilde{\chi}^0} = \begin{pmatrix} M_1 & 0 & -\frac{1}{2}g'v_d & \frac{1}{2}g'v_u & 0 & 0 \\ 0 & M_2 & \frac{1}{2}gv_d & -\frac{1}{2}gv_u & 0 & 0 \\ -\frac{1}{2}g'v_d & \frac{1}{2}gv_d & 0 & -\mu_{\text{eff}} & -\frac{1}{\sqrt{2}}\lambda_S v_u & -\frac{1}{\sqrt{2}}\lambda_S v_u \\ \frac{1}{2}g'v_u & -\frac{1}{2}gv_u & -\mu_{\text{eff}} & 0 & -\frac{1}{\sqrt{2}}\lambda_S v_d & -\frac{1}{\sqrt{2}}\lambda_S v_d \\ 0 & 0 & -\frac{1}{\sqrt{2}}\lambda_S v_u & -\frac{1}{\sqrt{2}}\lambda_S v_d & 0 & M_S \\ 0 & 0 & -\frac{1}{\sqrt{2}}\lambda_{\bar{S}} v_u & -\frac{1}{\sqrt{2}}\lambda_{\bar{S}} v_d & M_S & 0 \end{pmatrix}$$

- Singlino mass splitting: $M_{1,2} \gg \mu_{\text{eff}}, M_S, \lambda_S, \lambda_{\bar{S}} \ll 1$

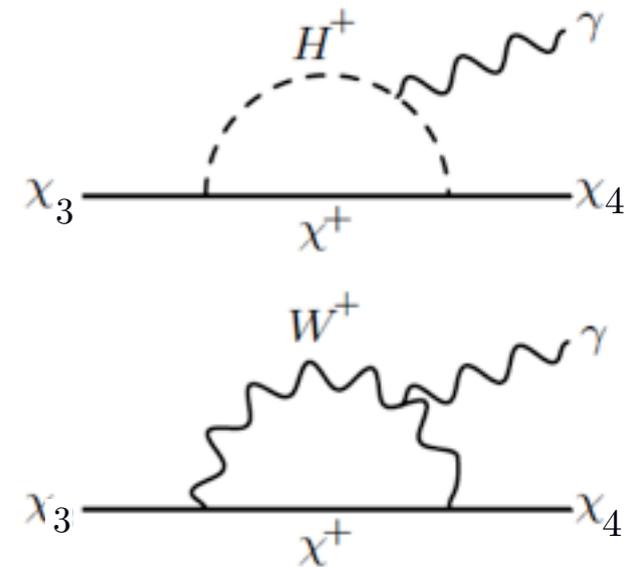
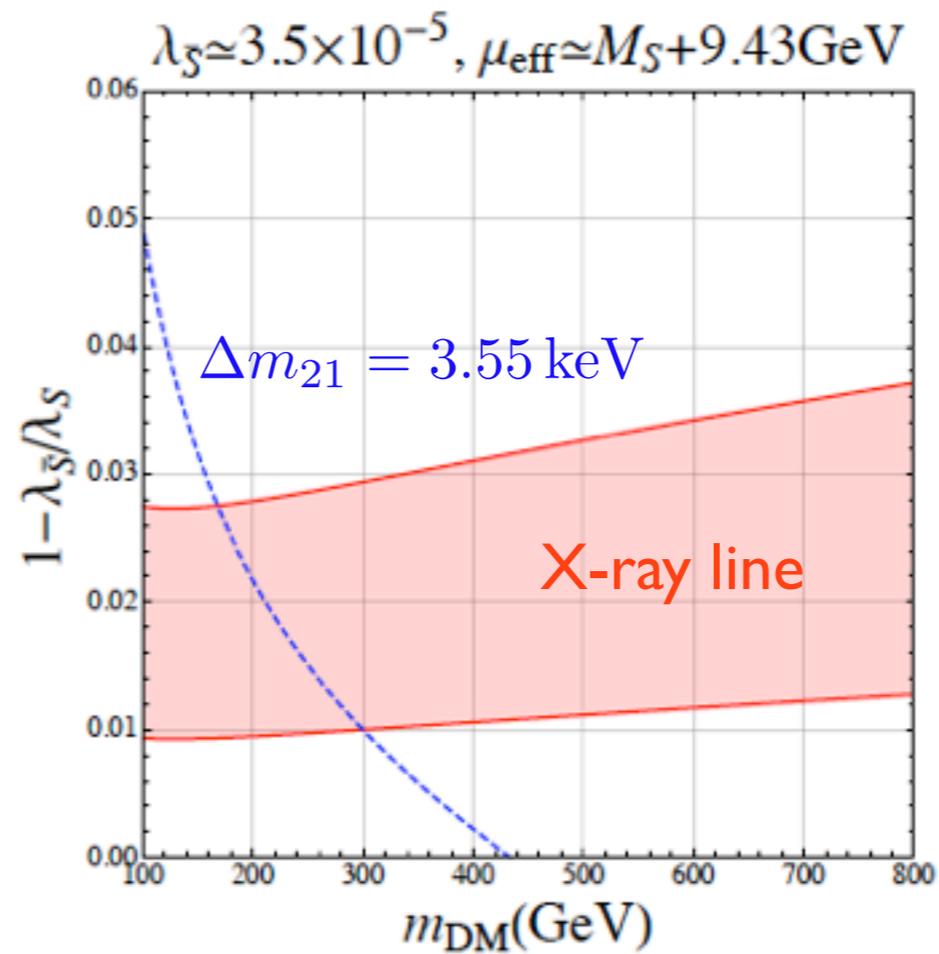
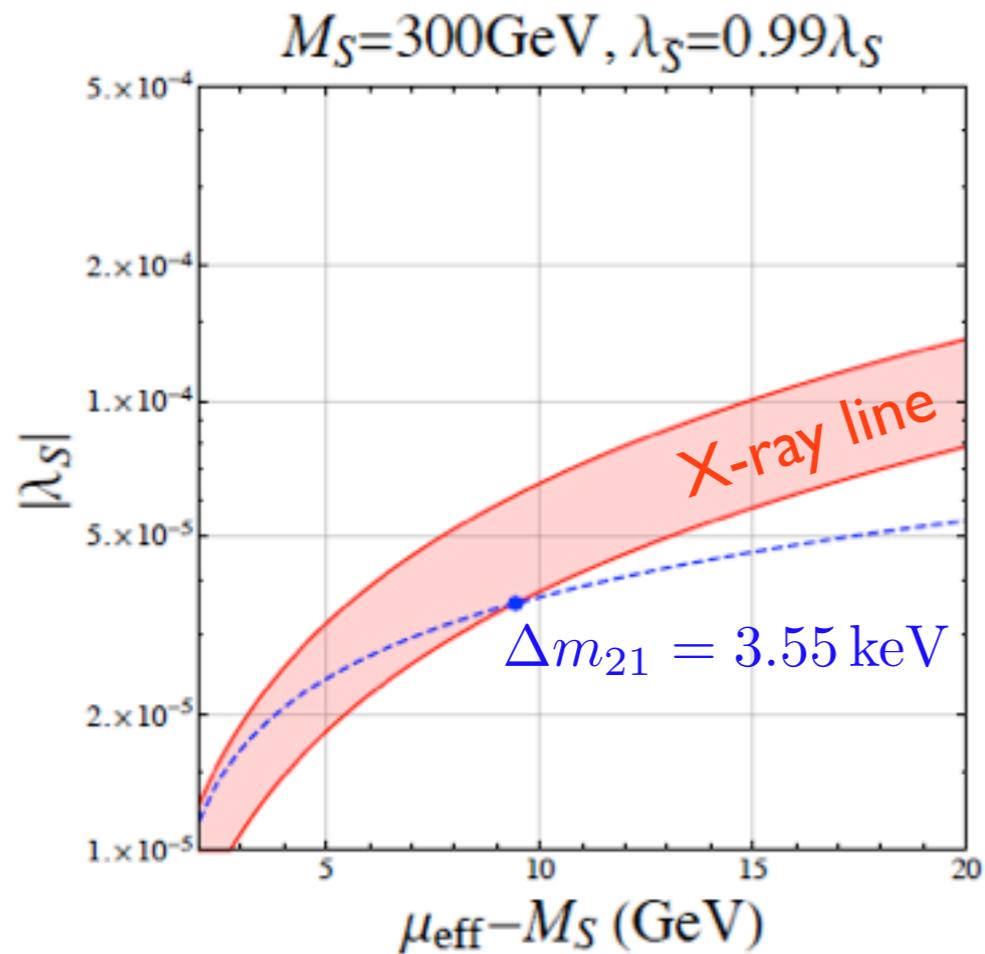
$$\Delta m_{34} \equiv m_{\tilde{\chi}_3^0} - m_{\tilde{\chi}_4^0} \approx \frac{1}{2} \frac{v^2}{\mu_{\text{eff}}^2 - M_S^2} \left((\lambda_+^2 - \lambda_-^2) M_S - (\lambda_+^2 + \lambda_-^2) \mu_{\text{eff}} \sin(2\beta) \right) + \frac{1}{8} (\lambda_+^2 + \lambda_-^2) \frac{v^4 \cos^2(2\beta) \mu_{\text{eff}}^2}{(\mu_{\text{eff}}^2 - M_S^2)^2} \left(\frac{g'^2}{M_1} + \frac{g^2}{M_2} \right),$$

“keV-scale” mass splitting for weak-scale $\mu_{\text{eff}} > M_S$.



$$\lambda_S, \lambda_{\bar{S}} \sim 10^{-5}, \quad M_{1,2} \gtrsim 1 \text{ TeV}.$$

X-ray line from singlinos



$m_{H^\pm} = 1 \text{ TeV}, M_1 = 0.5M_2 = 3 \text{ TeV}$ and $\tan \beta = 10$.

- “Charged Higgs, W-boson”-chargino loops induce magnetic transition moment for singlinos.

- X-ray line requires $\lambda_{\bar{S}}/\lambda_S \simeq 0.97 - 0.99$ for $\lambda_S \gtrsim 10^{-5}$.

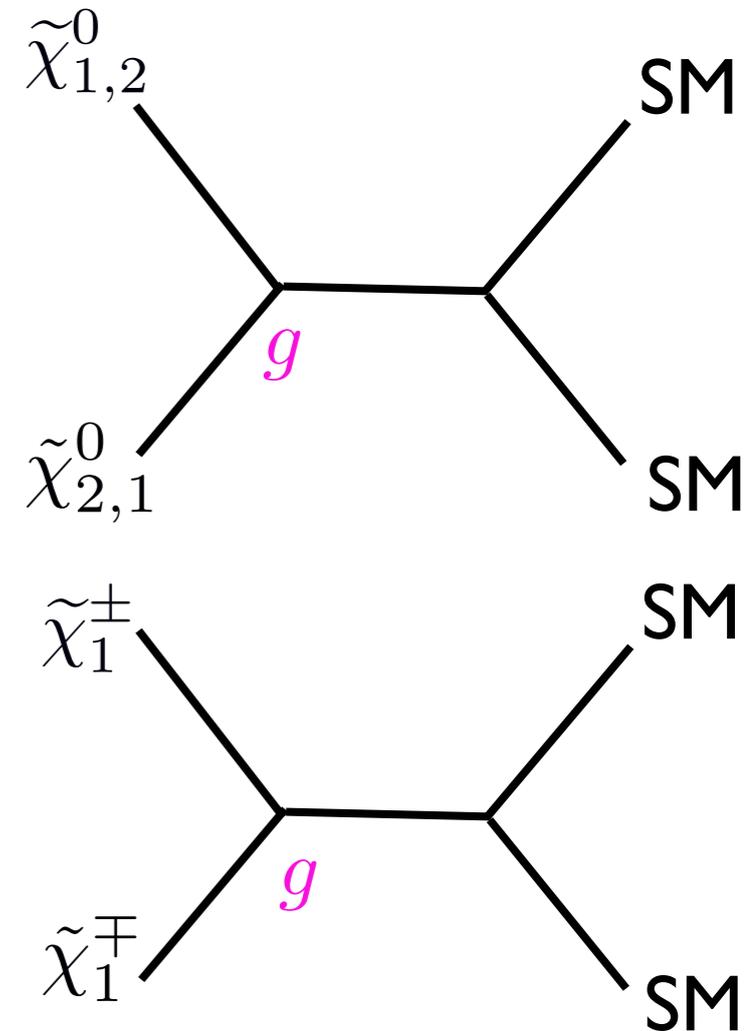
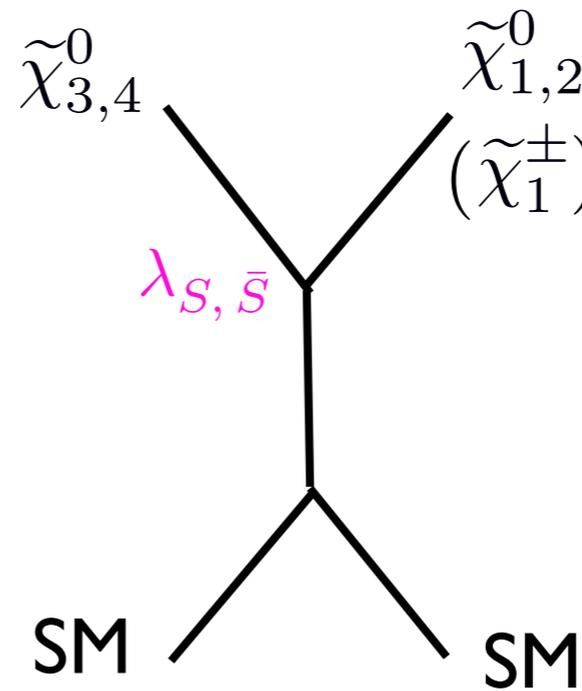
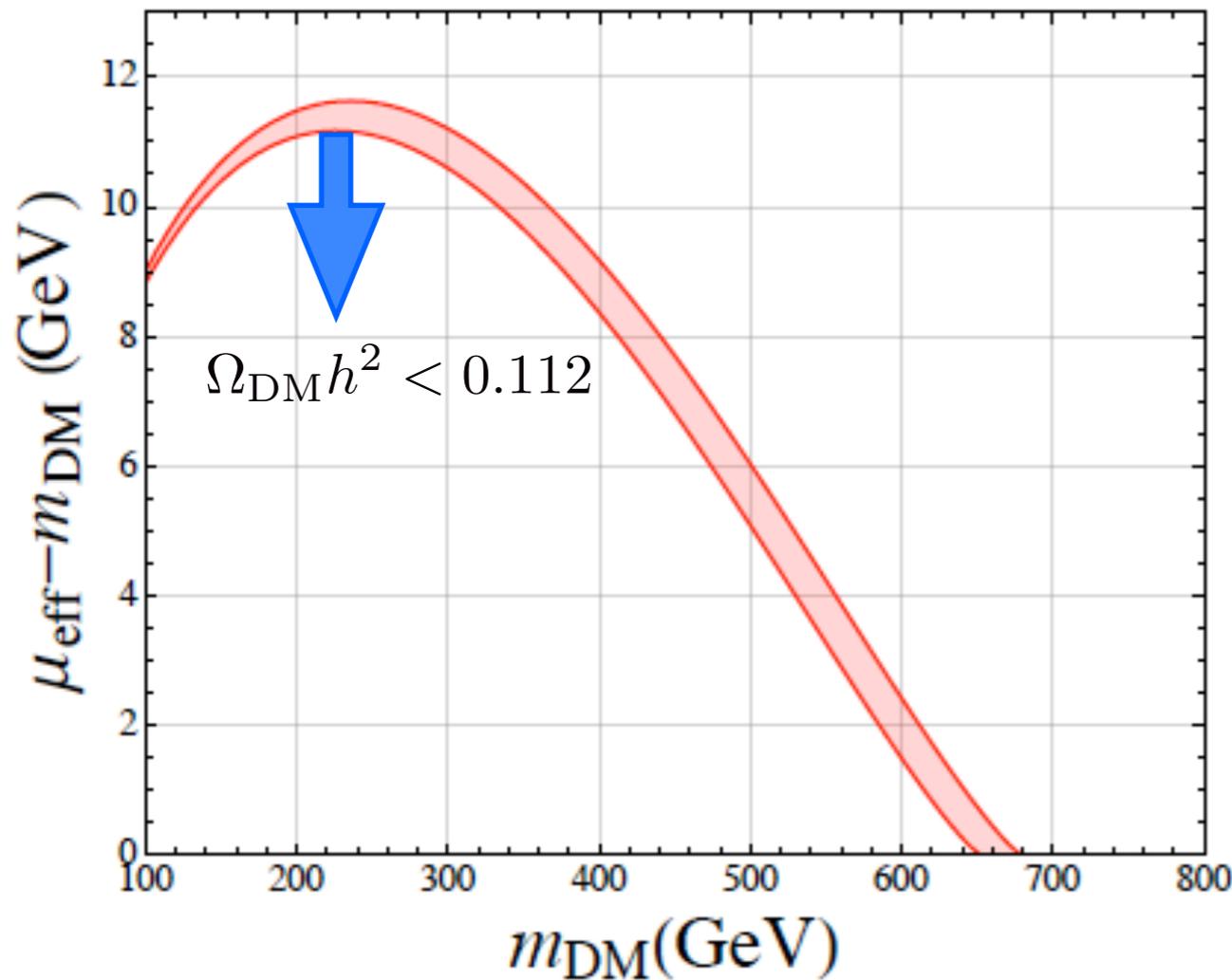
DM relic density

- Singlino DM keeps in thermal equilibrium until freeze-out, due to kinetic scattering with the SM.

$$\frac{\Gamma_{\text{kin}}}{\Gamma_{\text{th}}} \sim \frac{n_{\text{SM}} \langle \sigma v \rangle_{\text{kin}}}{n_{\text{DM}} \langle \sigma v \rangle_{\text{th}}} \sim |\lambda_{S, \bar{S}}|^2 \cdot 10^{10} \gtrsim 1.$$

$$\Rightarrow |\lambda_{S, \bar{S}}| \gtrsim 10^{-5}.$$

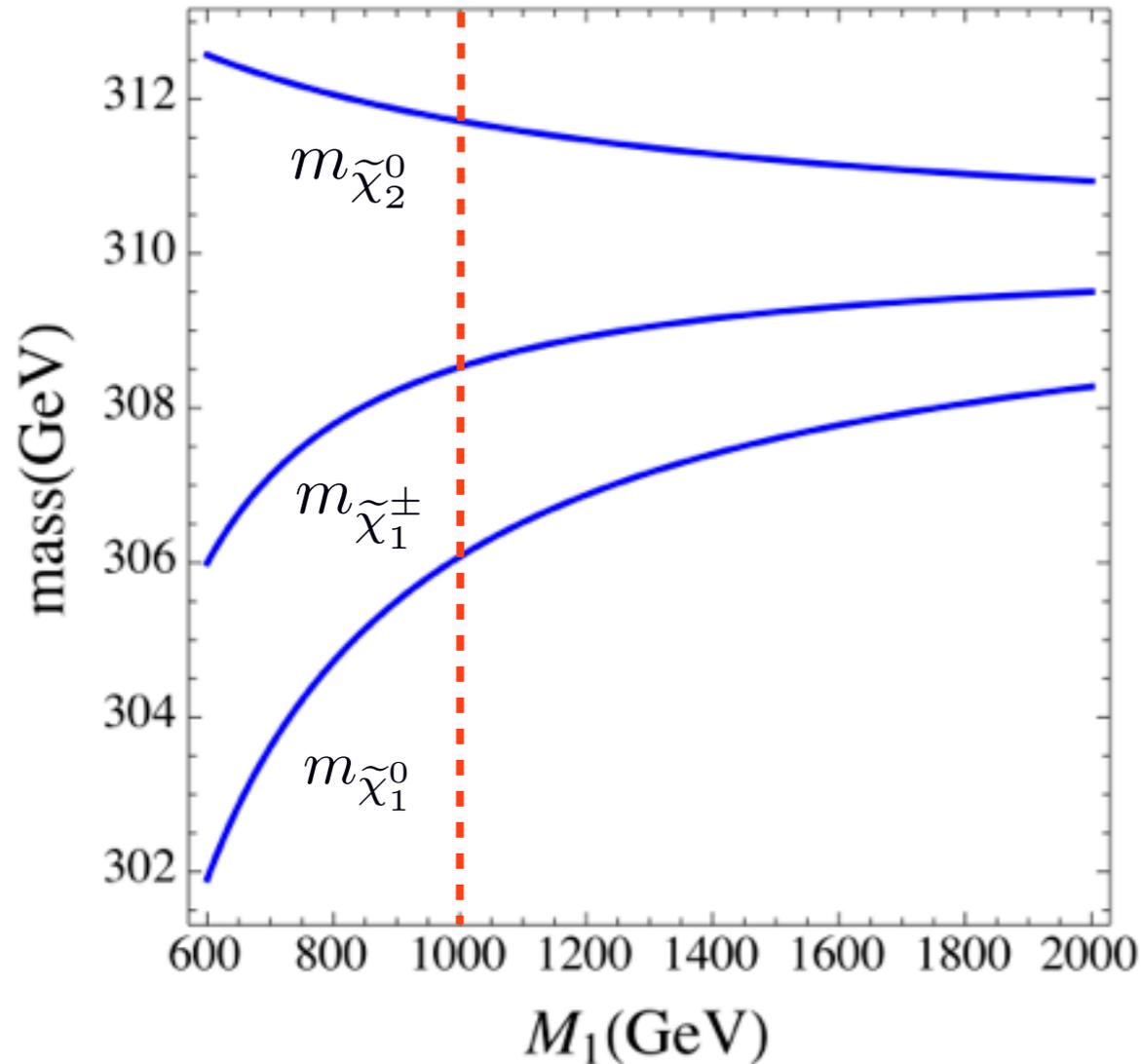
$\Omega_{\text{DM}} h^2$ (Planck 3σ)



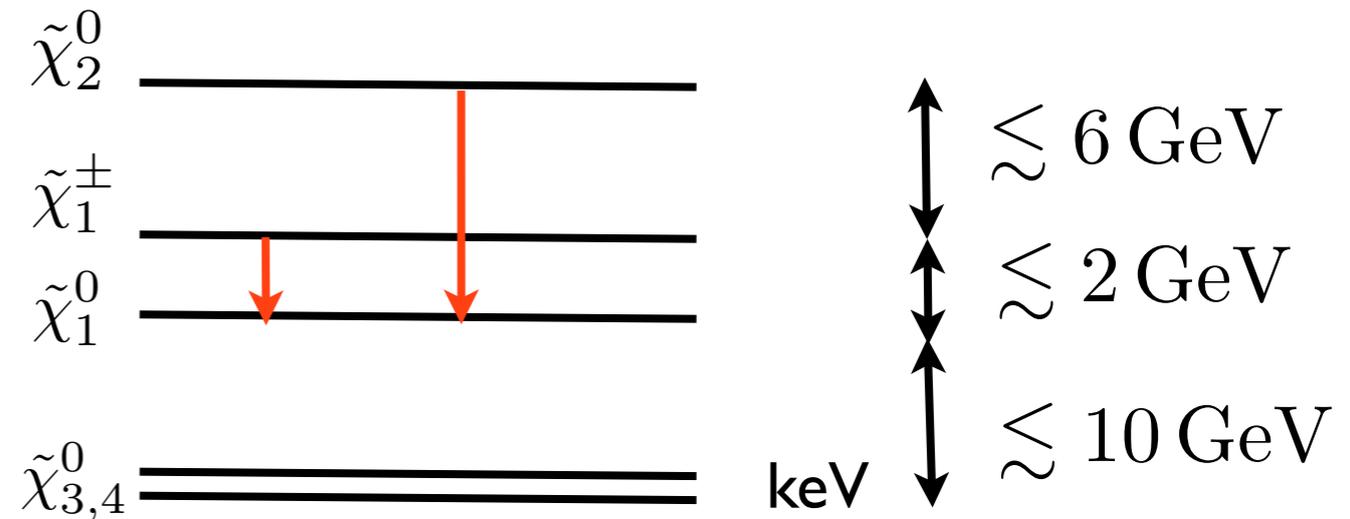
- The relic density conditions are consistent with the X-ray line.

Degenerate Higgsinos at LHC

$M_2=2M_1, M_S=300\text{GeV}, \mu_{\text{eff}}=310\text{GeV}$



For gauginos $> \text{TeV}$,



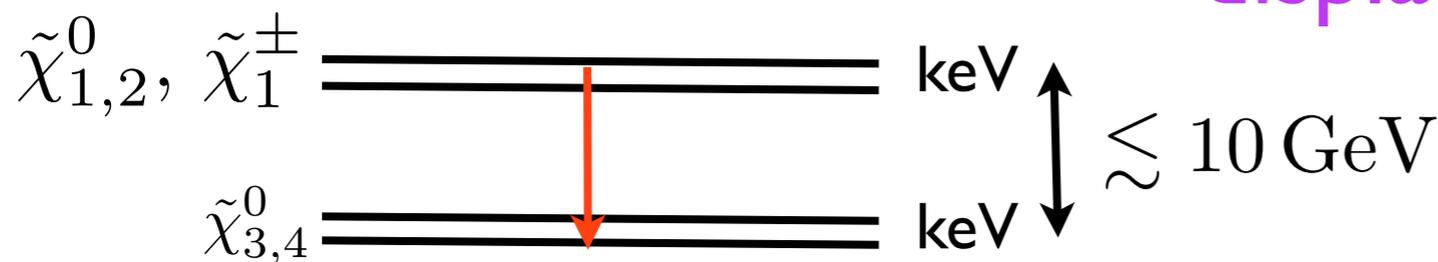
➔ 3-body cascade decays (mostly prompt)

$$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 Z^*(h^*), \quad \tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 W^*.$$

For decoupled gauginos,

Higgsinos decay to singlinos:

“displaced” vertices + MET.



$$\begin{aligned} \tilde{\chi}_{1,2}^0 &\rightarrow \tilde{\chi}_{3,4}^0 Z^*(h^*), \\ \tilde{\chi}_1^\pm &\rightarrow \tilde{\chi}_{3,4}^0 W^*. \end{aligned}$$

Conclusions

- **Dipolar dark matter** can explain the X-ray line and is thermally produced due to an accompanied charged particle.
- **Fermion-portal coupling of order one and charged scalar** can be constrained by $g-2$, LFV, LHC and direct detection.
- In SUSY Higgs-portal model, **Dirac singlino with a naturally small mass splitting** can explain the X-ray line. But, almost degenerate Higgsinos/singlinos would be challenging for the LHC Run II searches.



Dream Improvisation
by Wassily Kandinsky

Dark sector might be dynamic like our world!