

# Unveiling the nature of Dark Matter in the LOFAR sky

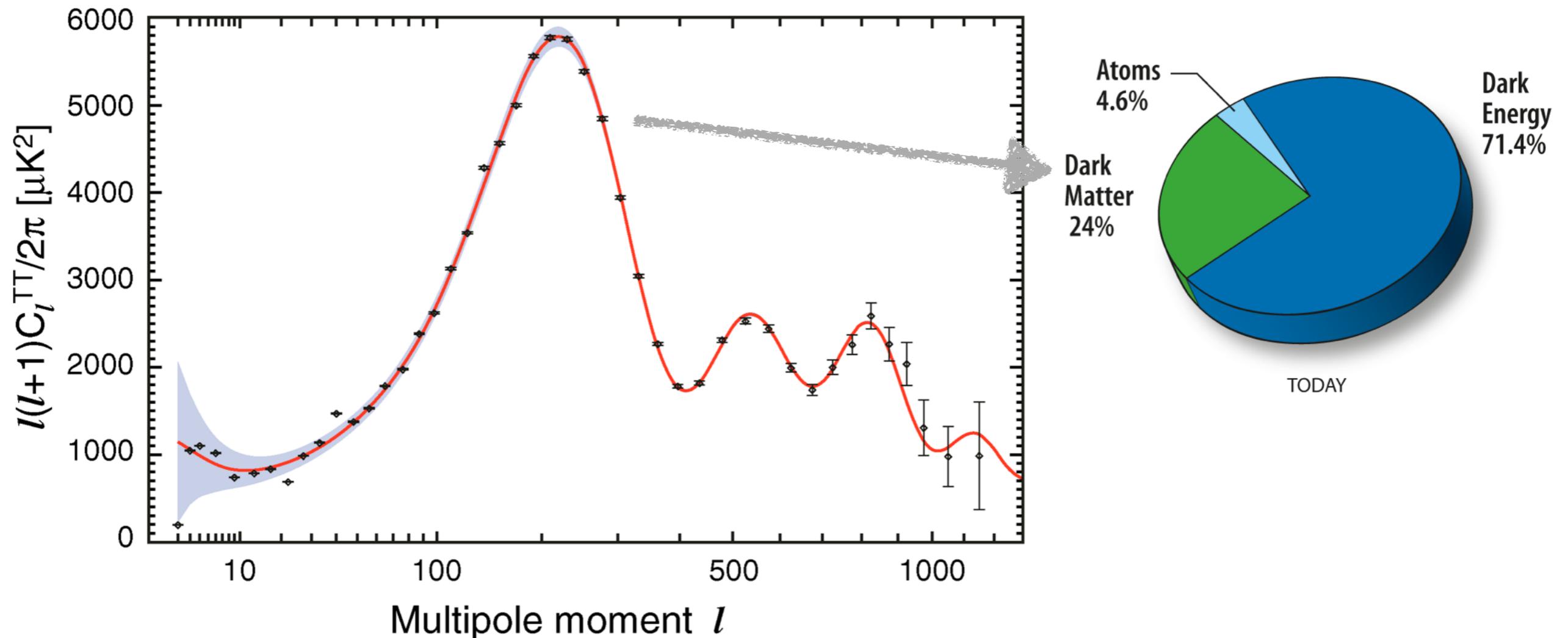
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Carmelo Evoli ( Universität Hamburg )



Dark-MALT | Munich | 18th of February 2015

# Dark Matter on cosmological scales



- Dark Matter (DM) needed on all scales.
- DM model requires electrically neutral particles moving (almost) non-relativistically when structures formed.
- Therefore, Standard Model is not a complete and fundamental theory.

# Dark Matter on galactic scales

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## 21-CM LINE STUDIES OF SPIRAL GALAXIES. II. THE DISTRIBUTION AND KINEMATICS OF NEUTRAL HYDROGEN IN SPIRAL GALAXIES OF VARIOUS MORPHOLOGICAL TYPES

A. BOSMA<sup>a)</sup>

Kapteyn Astronomical Institute, University of Groningen, Groningen, The Netherlands

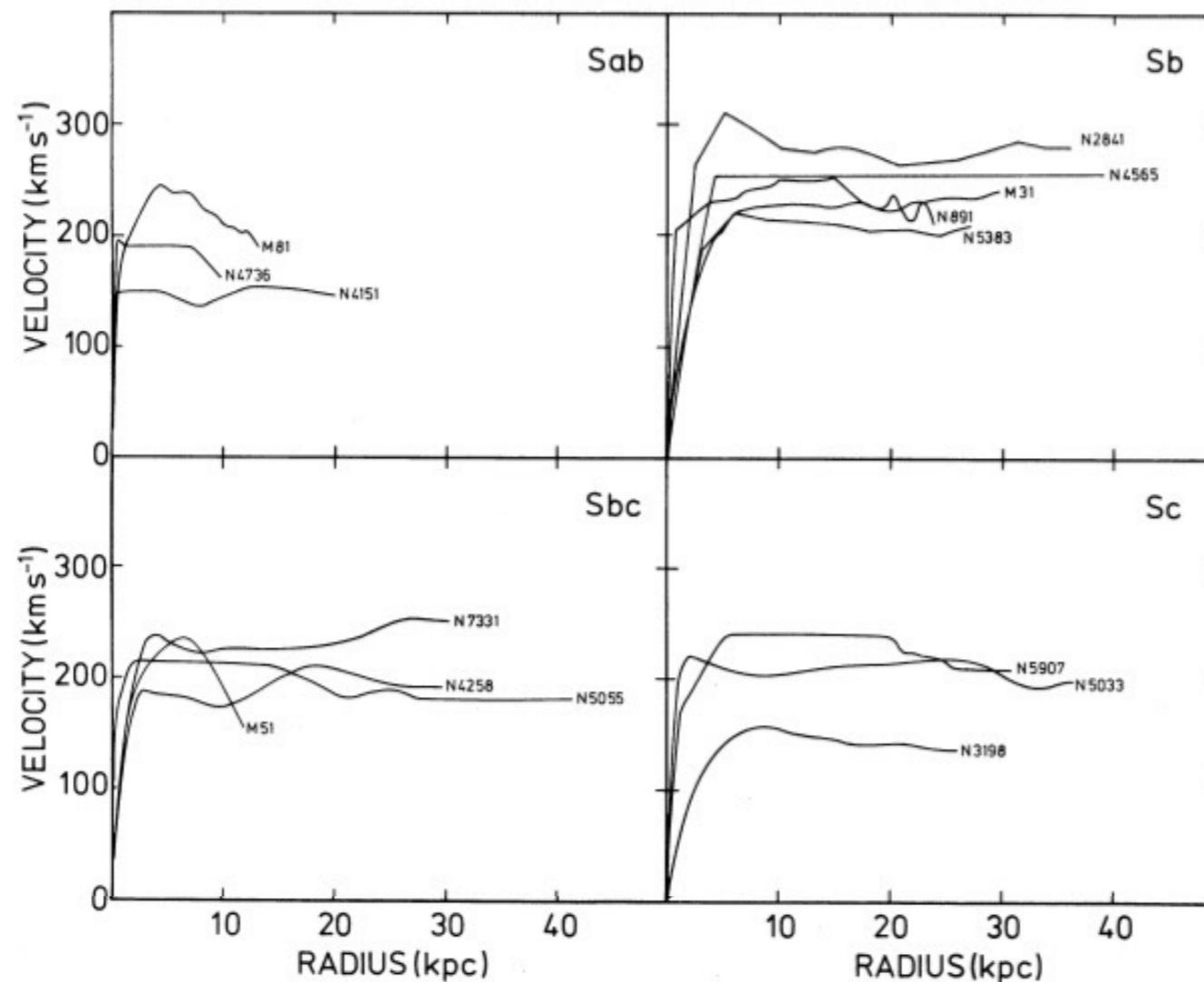


FIG. 3. ROTATION CURVES OF GALAXIES OF VARIOUS HUBBLE TYPES.

# The WIMP “miracle”

E. Kolb and M. Turner, “The Early Universe” (1994)

Let  $\chi$  be a stable particle carrying a non-zero charge under the SM gauge group:



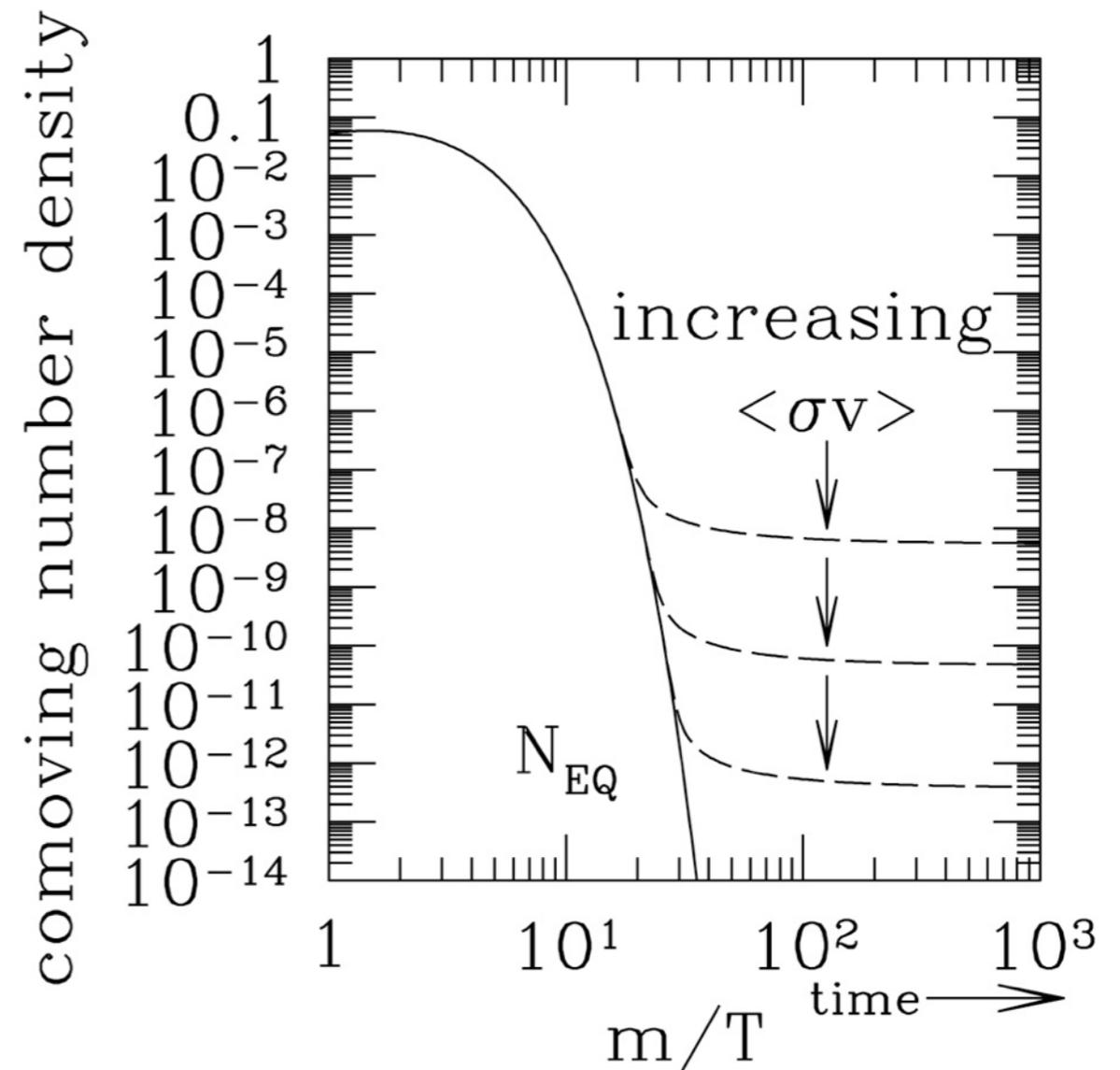
The evolution of the number density is described by the Boltzmann equation:

$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle\sigma v\rangle_T [(n_\chi)^2 - (n_\chi^{\text{eq}})^2] \quad (2)$$

At freeze-out, when  $\Gamma \ll H$ , the number density per comoving volume becomes constant. For a species which is not-relativistic at freeze-out:

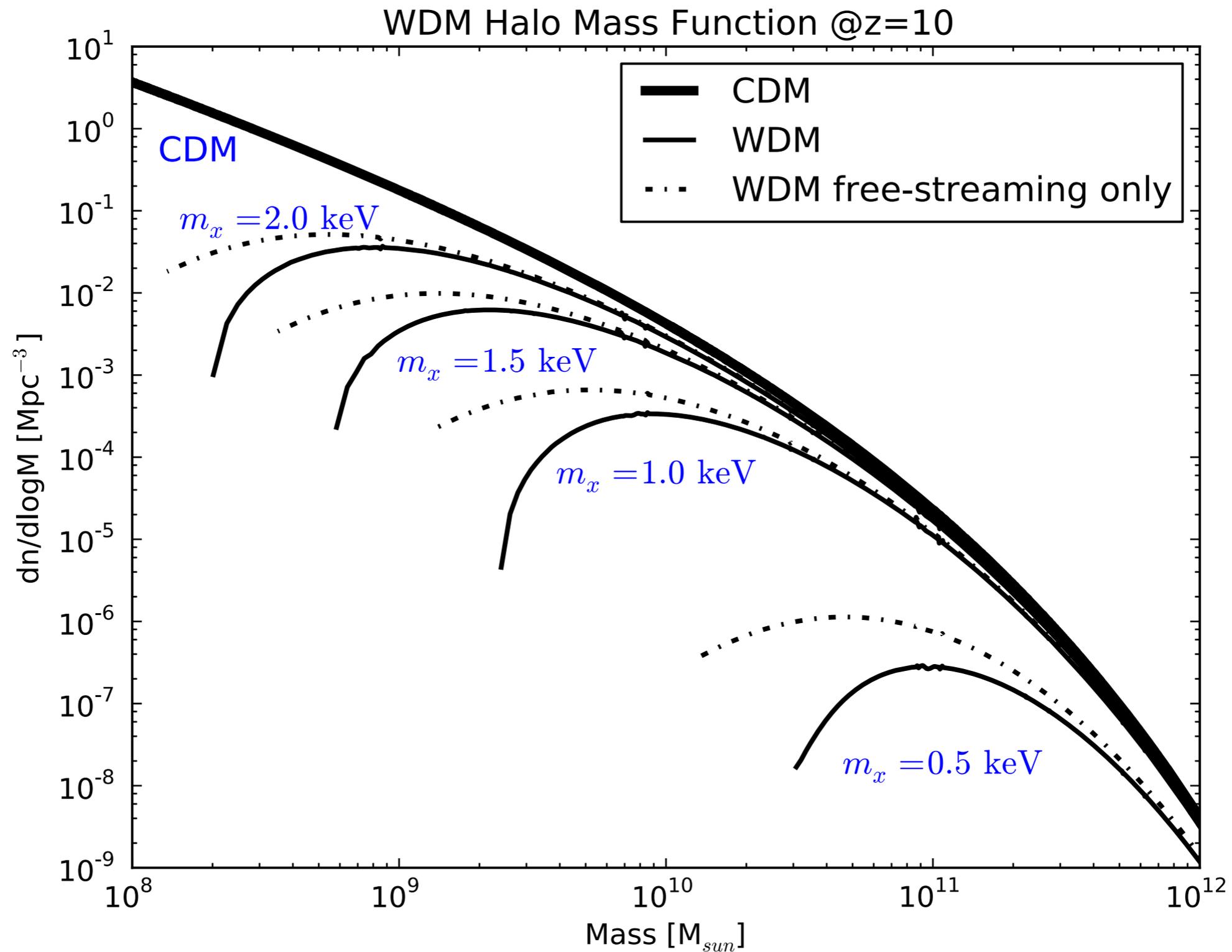
$$\Omega h^2 \sim \frac{3 \times 10^{-27} \text{ cm}^3/\text{s}}{\langle\sigma v\rangle_{T=T_f}} \quad (3)$$

The WIMP recipe to embed a DM candidate in a SM extension foresees an extra particle  $\chi$  that is stable, massive and weakly-interacting!

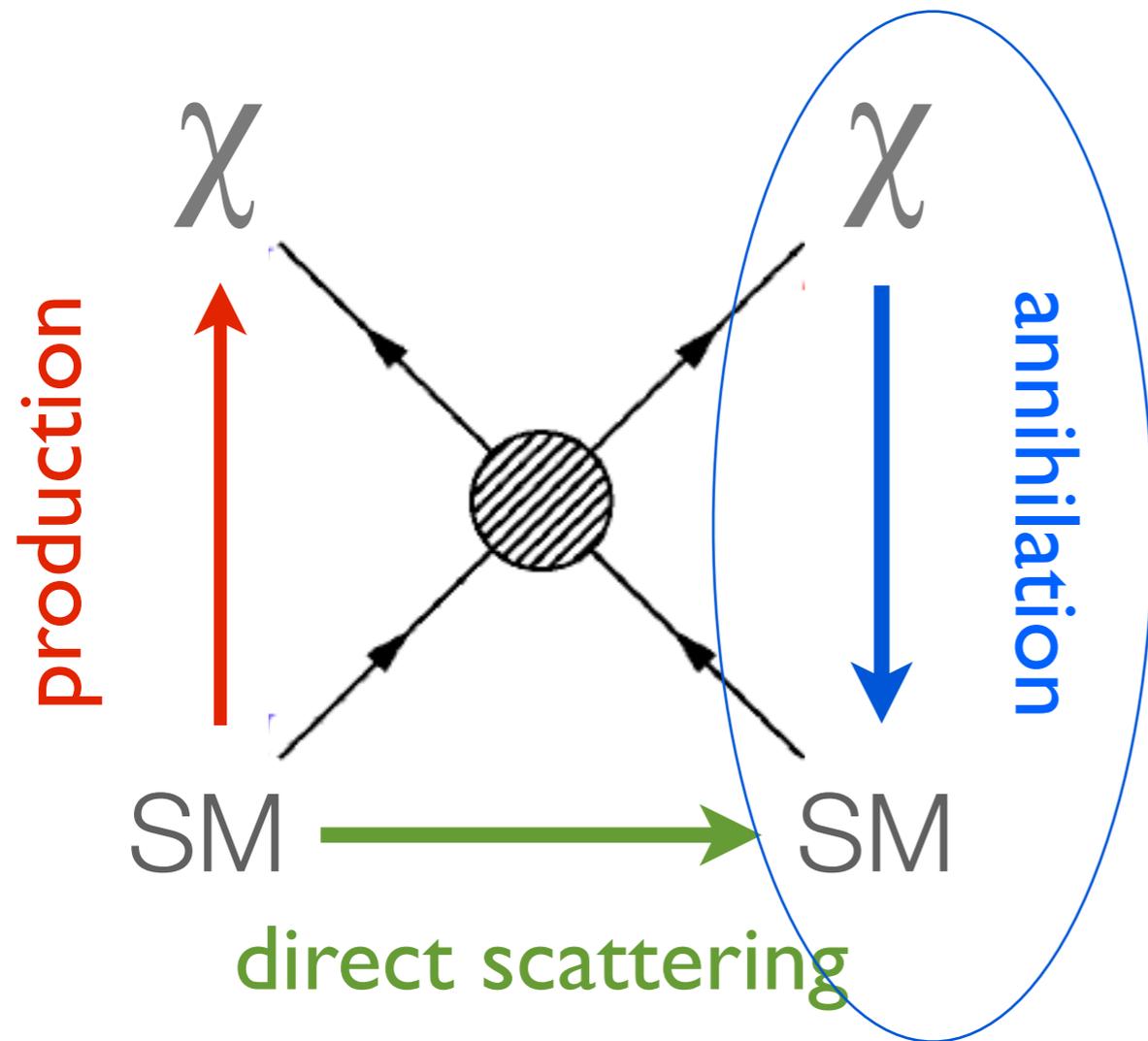


# How Cold Is Cold Dark Matter?

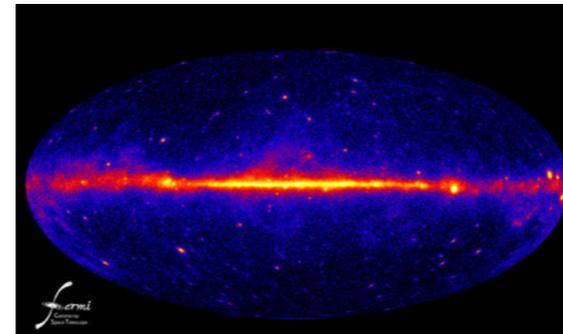
F. Pacucci, A. Mesinger and Z. Haiman, MNRAS, 2013



# WIMPs detection strategies



$$\Gamma_{ann} \propto n_{\chi}^2 \sigma v$$



Galactic Center



Dwarf galaxies

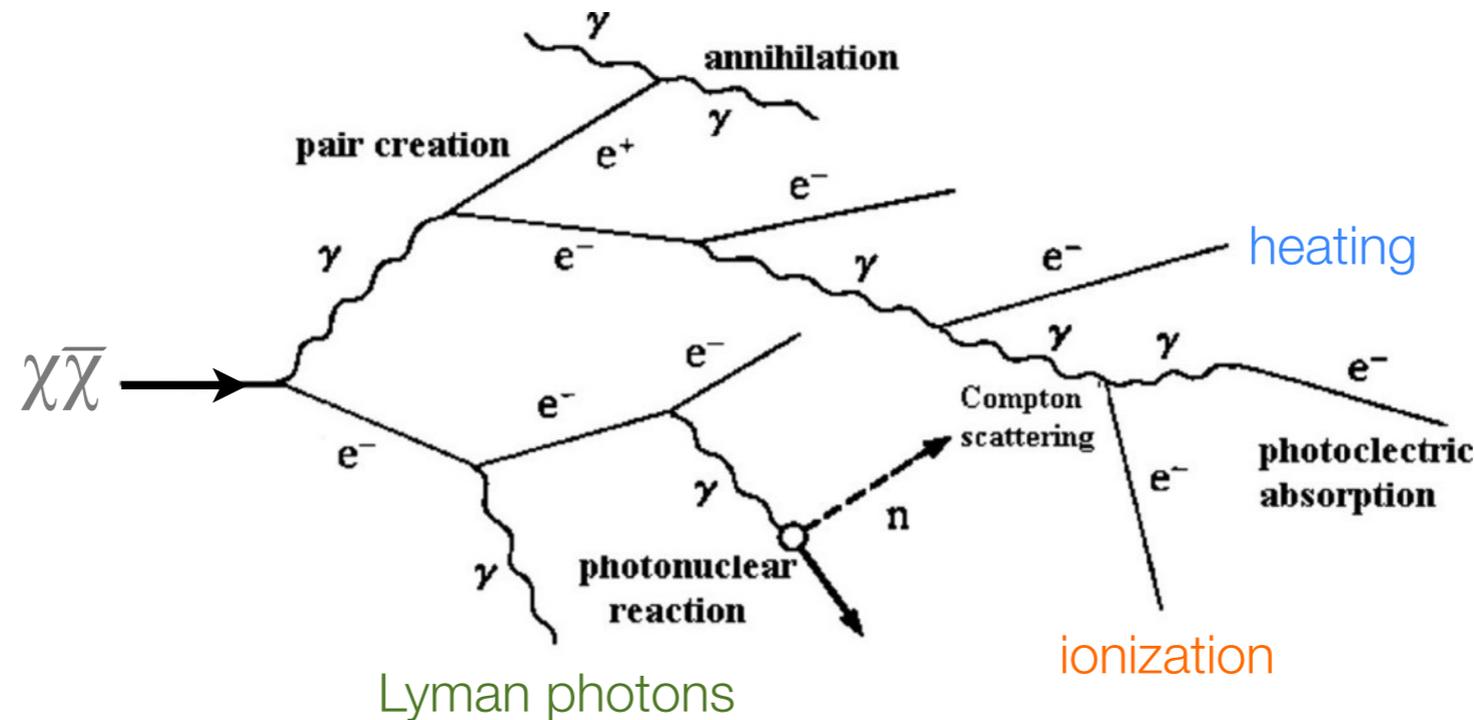


Early Universe

All these approaches are complementary!

# DM during Dark Ages

CE, M.Valdés, A.Ferrara, N. Yoshida, MNRAS, 2012

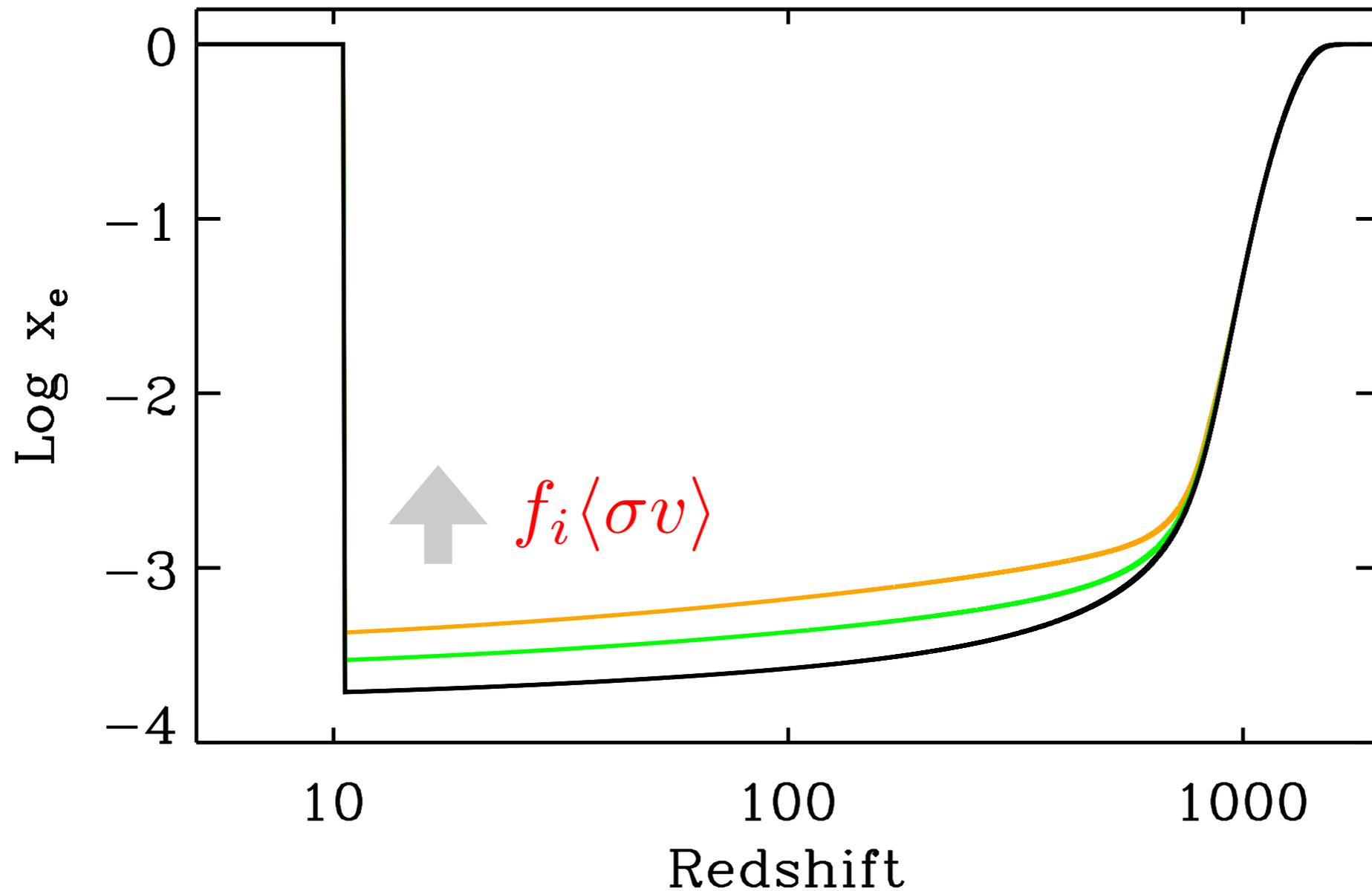


- WIMP annihilation produce gamma-rays, cosmic-rays, anti-matter: release energy which, in principle, **ionize/heat** up baryonic gas.
- Energy released enormous: **1 GeV** energy is sufficient to ionize  **$\sim 10^8$  hydrogen atoms**, only a tiny fraction of DM decay/annihilate is enough to affect ionization history.
- By looking for departures from the “standard recombination” scenario, we can place limits on energy injection at  **$z=100-1000$**  and translate these limits to exclusions in WIMP parameter space (e.g. the cross-section / mass plane, etc.).

# Baryon history in the Universe

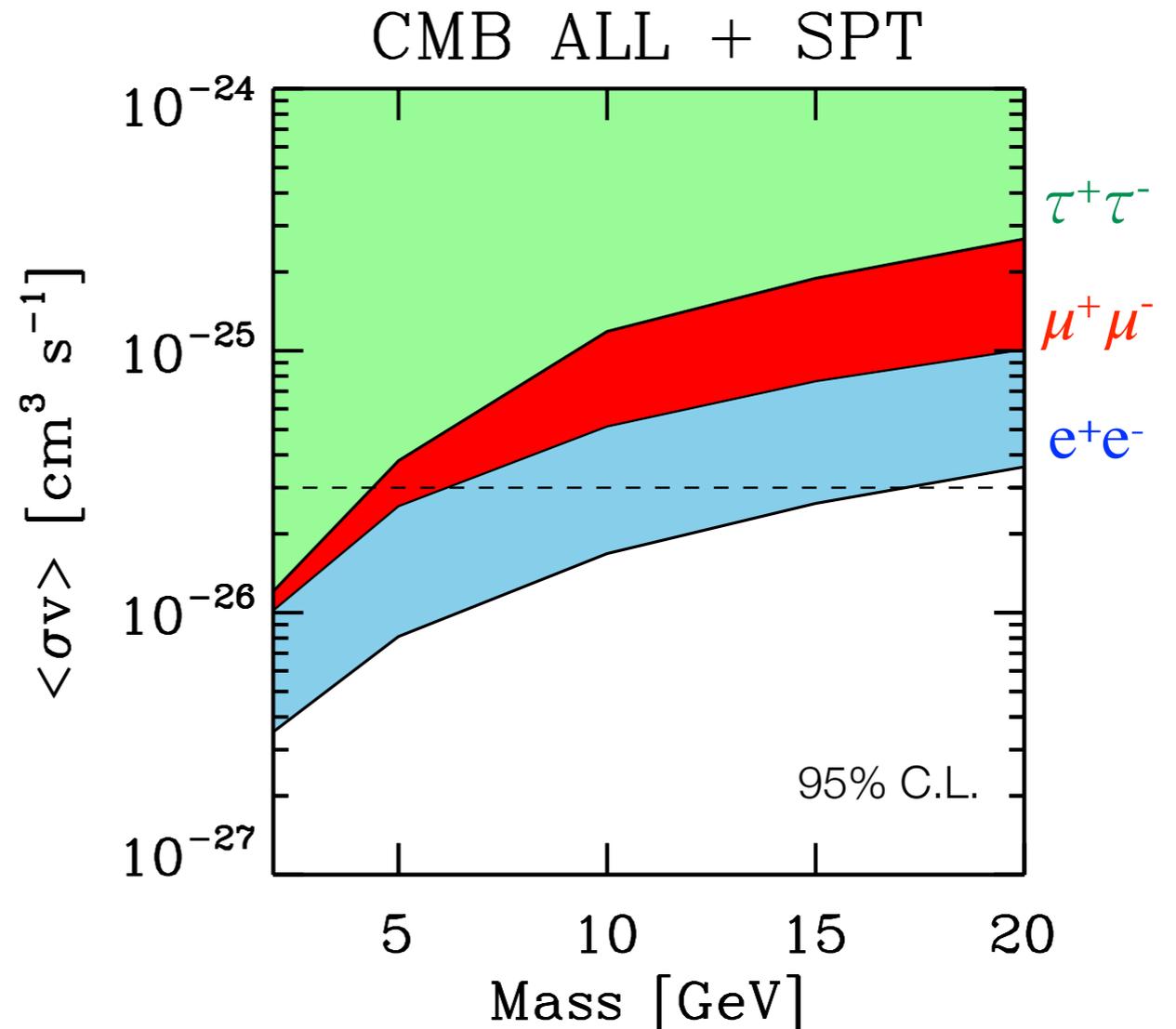
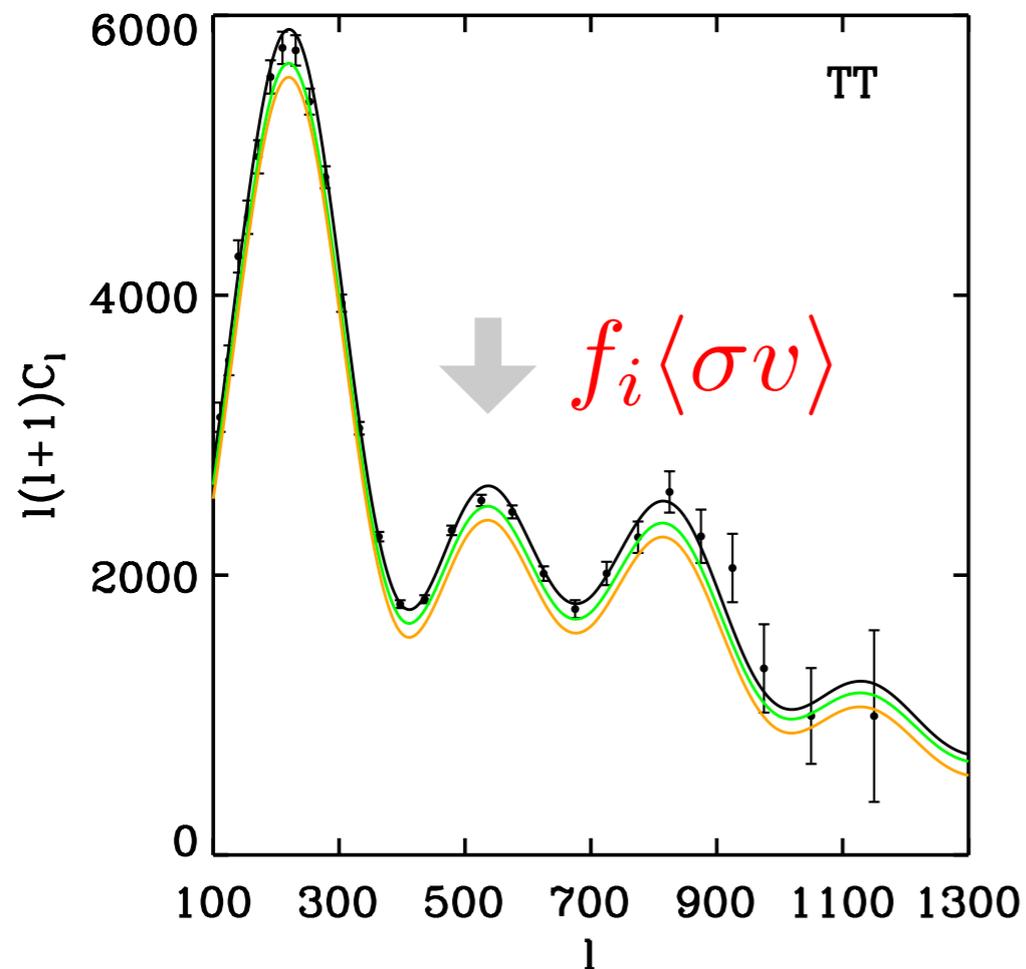
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$$\frac{dE_\chi}{dt} = 2M_\chi f_i \langle \sigma v \rangle n_\chi^2$$



# CMB constraints

CE, S.Pandolfi and A.Ferrara, MNRAS, 2013



- more **free electrons** survive at low redshift after recombination compared to the standard case. Increase the width of the LSS, and consequently the width of the visibility function.
- these results are quite robust since we understand recombination and the **CMB** quite well and astrophysical backgrounds are less relevant with respect to constraints based on, e.g., annihilations in late-time halos.

# A Schematic Outline of the Cosmic History

“Dark Ages”

Time since the Big Bang (years)

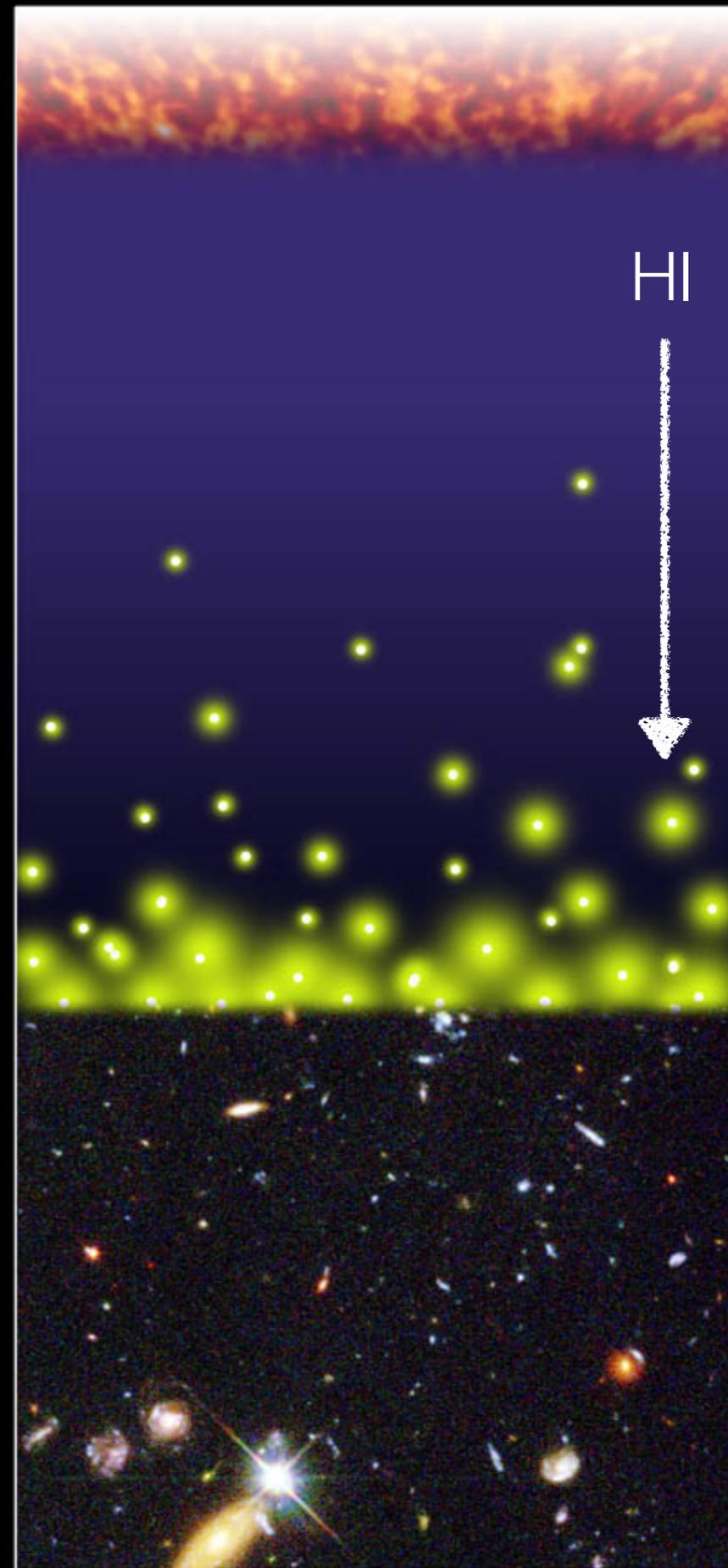
~ 300 thousand

~ 500 million

~ 1 billion

~ 9 billion

~ 13 billion



← The Big Bang

The Universe filled with ionized gas

← The Universe becomes neutral and opaque

The Dark Ages start

Galaxies and Quasars begin to form  
The Reionization starts

The Cosmic Renaissance  
The Dark Ages end

← Reionization complete, the Universe becomes transparent again

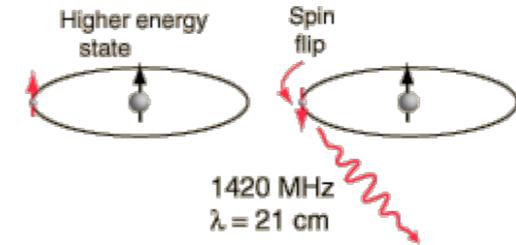
Galaxies evolve

The Solar System forms

Today: Astronomers figure it all out!

# 21 cm physics in a nutshell

J.Pritchard and A.Loeb, arXiv:1109.6012



- H I emission (predicted by van de Hulst in 1945, student of Oort):

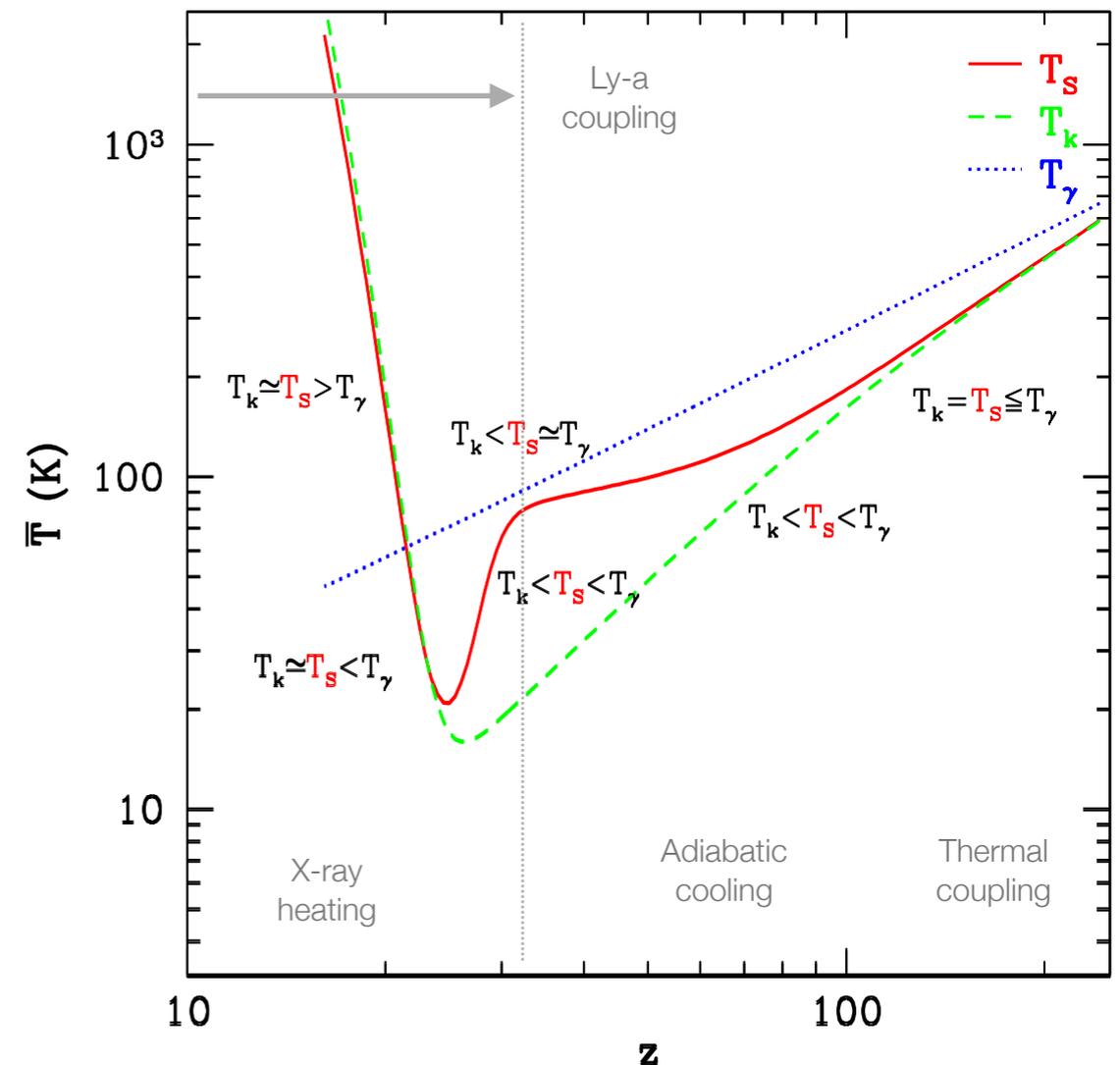
$$\frac{n_1}{n_0} = 3 \exp\left(-\frac{T_{21\text{cm}}}{T_S}\right)$$

- 21 cm brightness temperature of an IGM gas parcel at a redshift  $z$ , relative to the CMB:

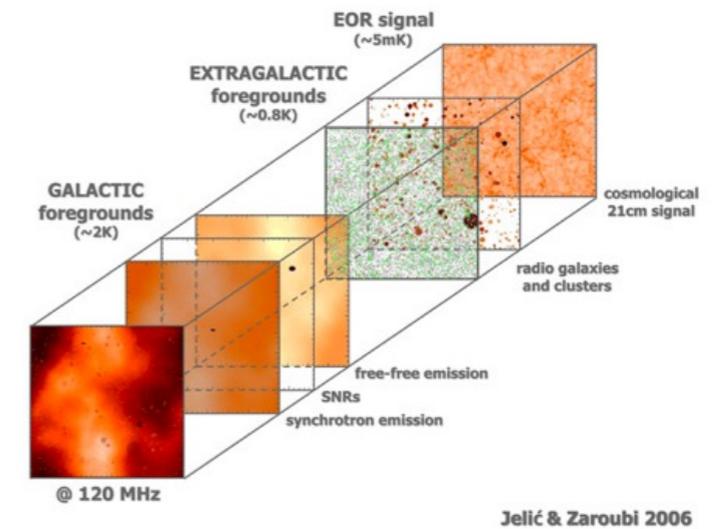
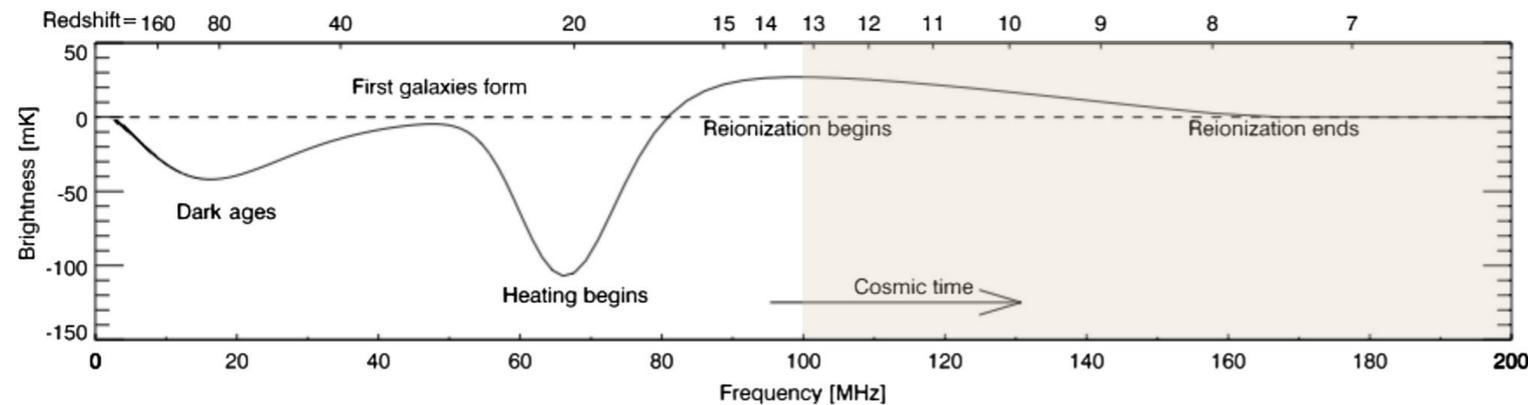
$$\delta T_b \sim 25 x_{\text{HI}} (1 + \delta) \left(\frac{1+z}{10}\right)^{1/2} \left[\frac{T_S - T_{\text{CMB}}}{T_S}\right] \text{ mK}$$

- In *emission* if  $T_S > T_{\text{CMB}}$ ; in *absorption* if  $T_S < T_{\text{CMB}}$
- In the absence of coupling mechanisms:  $T_S = T_{\text{CMB}}$ .  
But! Two mechanisms couple  $T_S$  to  $T_K$ : *collisions* and *Ly $\alpha$  pumping* (also known as Wouthuysen-Field Effect):

$$T_S = \frac{T_{\text{CMB}} + y_\alpha T_k + y_c T_k}{1 + y_\alpha + y_c}$$



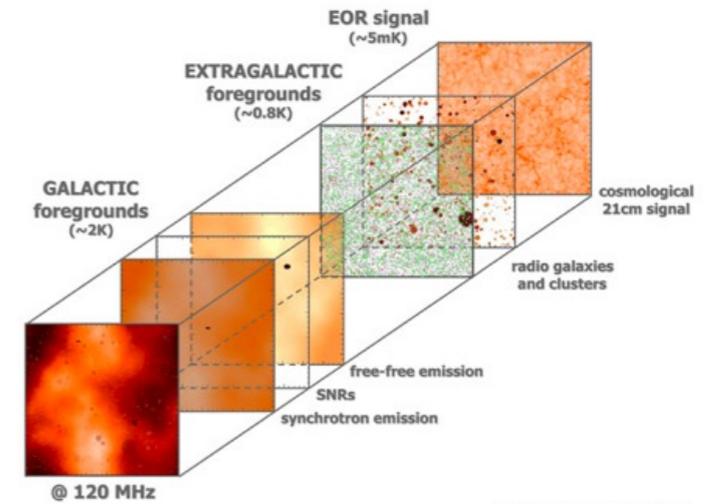
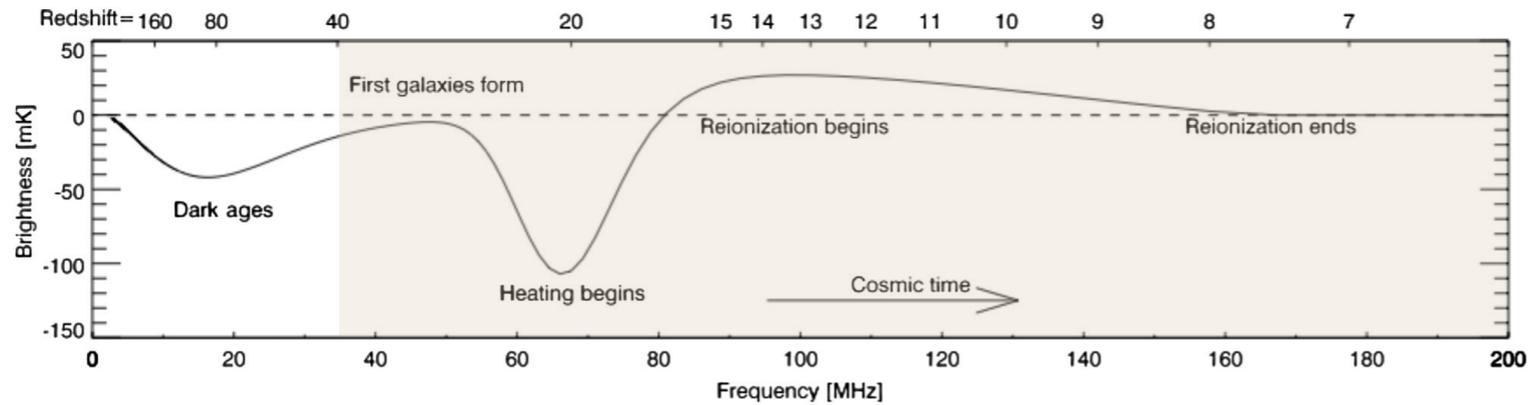
# The LOFAR EoR experiment



frequency coverage: 120-240 Mhz  
 site: Netherlands  
 angular resolution: 5"-3"  
 effective area: 93,479 m<sup>2</sup>  
 already in data-taking

**Fig. 19** *Left hand panel:* An artists impression of the layout of the LOFAR telescope over Western Europe. For the EoR, only the central part of the telescope is relevant. (courtesy of Peter Prijs). *Right hand panel:* The very central area of LOFAR. This circular area is known as the superterp and is the heart of the LOFAR core. The high-band array stations (covered in black blastic sheets) are clearly seen in this picture. In between one can also see the Low-Band Array antennas (Zaroubi, 2013)

# Square Kilometer Array (SKA)



(artistic view)

frequency coverage: 40-30,000 Mhz  
 site: Australia + South Africa  
 angular resolution:  $<0.1''$   
 effective area: 1,000,000 m<sup>2</sup>  
 first light in 2020!

# DM annihilation effects on 21 cm

S.R.Furlanetto, S.P.Oh & E.Pierpaoli, PRD, 2006 ; A.Natarajan & D.J.Schwarz, PRD, 2009

$$T_S = \frac{T_{\text{CMB}} + y_\alpha T_k + y_c T_k}{1 + y_\alpha + y_c},$$

$$(1+z) \frac{dT_k}{dz} = 2T_k + \frac{l_\gamma x_e}{H(z)(1+f_{\text{He}}+x_e)} (T_k - T_{\text{CMB}}) - \frac{2\chi_h \dot{E}_x}{3k_b H(z)(1+f_{\text{He}}+x_e)},$$

$$J_\alpha(z) = \frac{\mathcal{N}_H^2 hc}{4\pi H(z)} \left[ x_e x_p \alpha_{2P}^{\text{eff}} + x_e x_{\text{HI}} \gamma_{\text{eH}} + \frac{\chi_\alpha \dot{E}_x(z)}{\mathcal{N}_H h \nu_\alpha} \right],$$

heating fraction

Lya fraction

# M.E.D.E.A. code

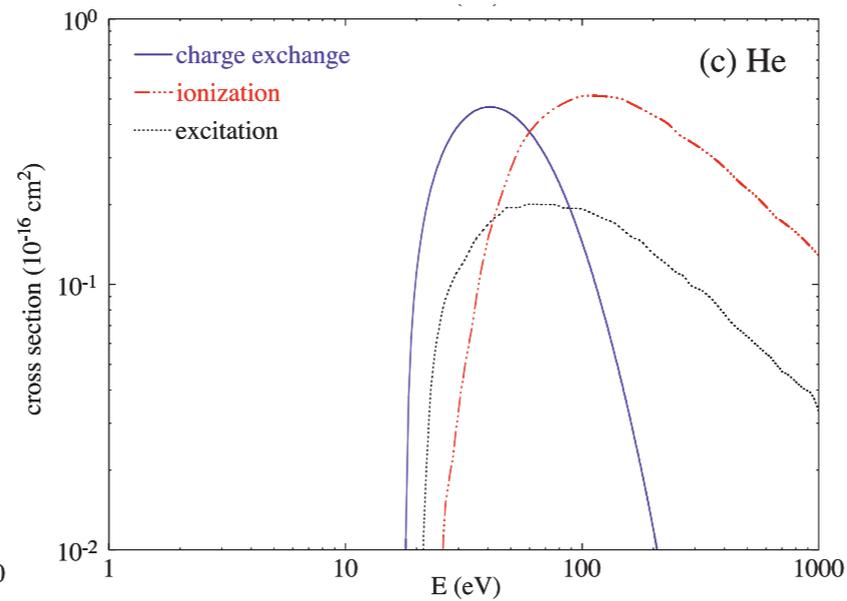
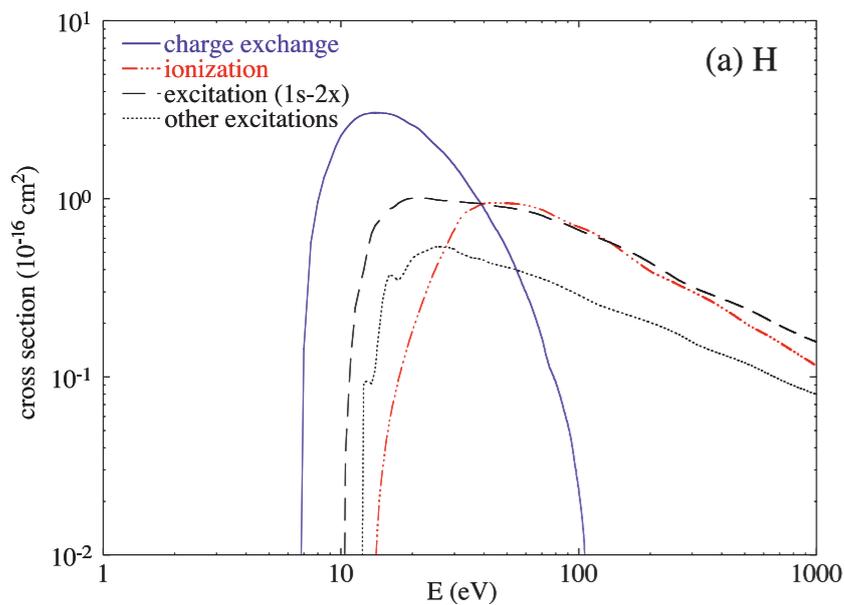
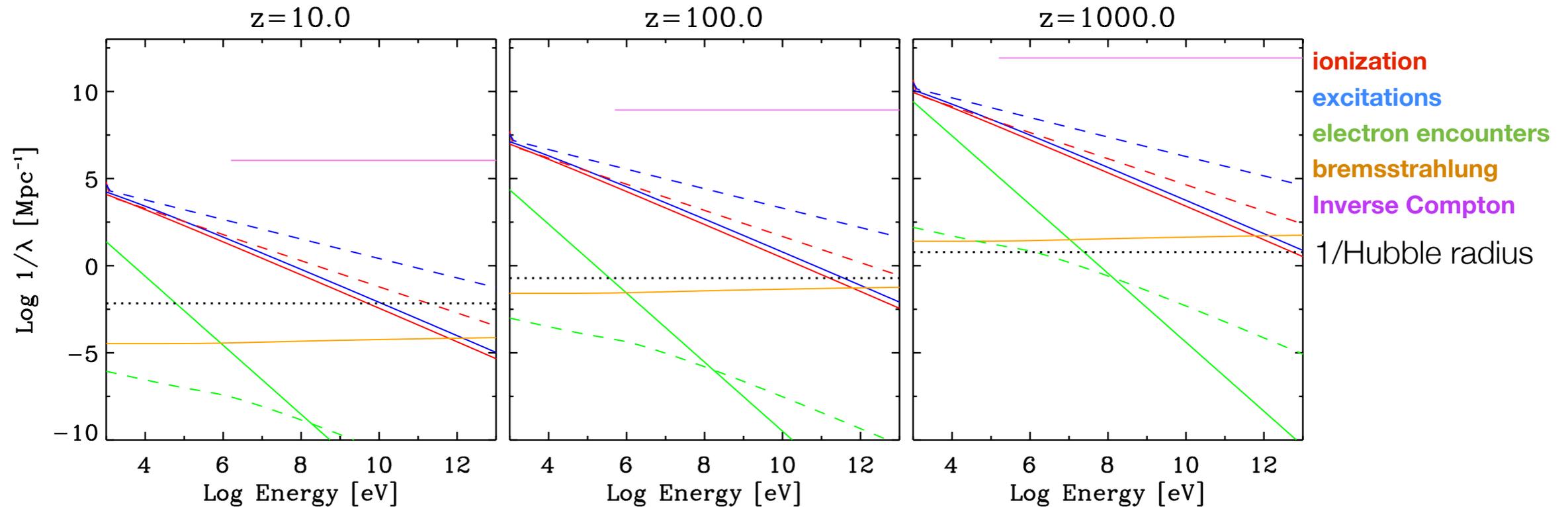
M.Valdés, CE, A.Ferrara, MNRAS, 2011

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- Several astrophysical and cosmological sources produce high energy particles through different acceleration processes:  
**how is this energy transferred to the surrounding environment?**
- **Monte Carlo method:** repeated random sampling of the relevant physical quantities and processes ( i.e. cross-sections and interaction probabilities ) to follow the evolution of a relativistic electron.
- An high energy electron in the IGM produces a chain of reactions: start with one particle, end up with many → the code calculates for every particle the probability of the main interaction channels and then selects one by a random number generator, until the energy is deposited into the IGM.
- **Montecarlo Energy DEposition Analysis.**

# Lepton interactions

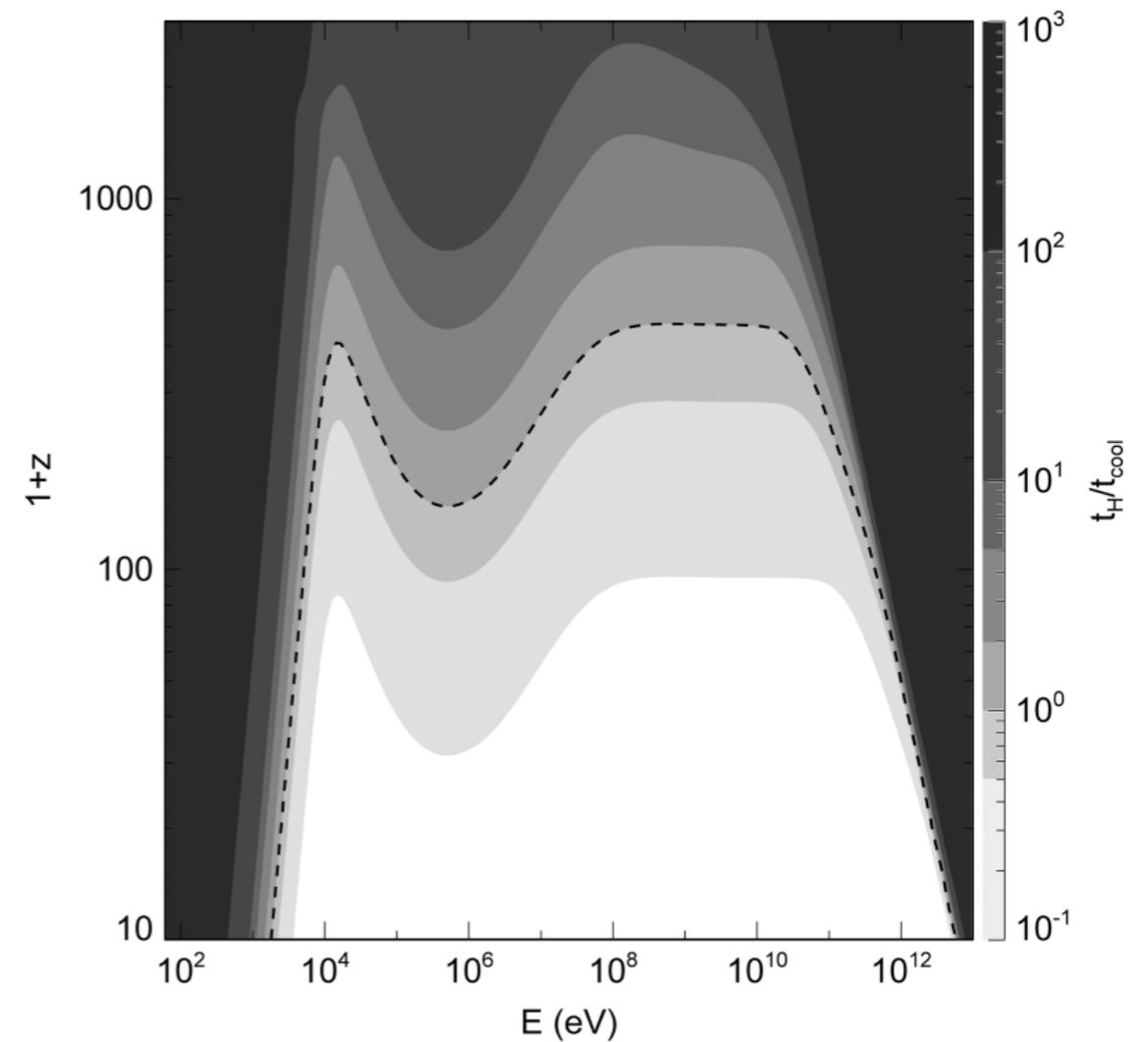
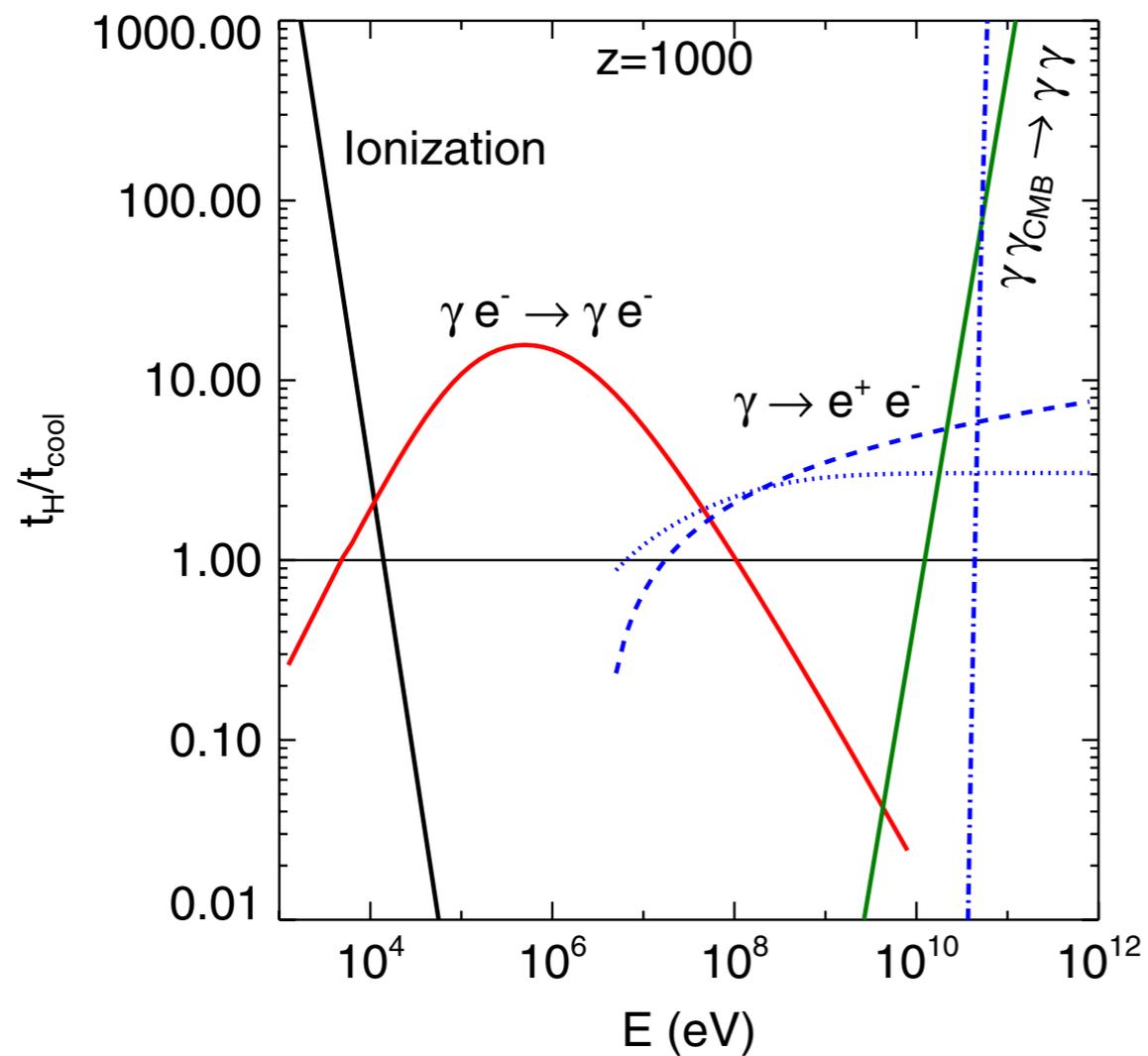
N.Guessoum et al., A&A, 2005 ; CE, M.Valdés, A.Ferrara and N.Yoshida, MNRAS, 2012



- IC dominant down to a O(MeV) energy
- Positronium formation in positron collision with H and He.

# Photon interactions

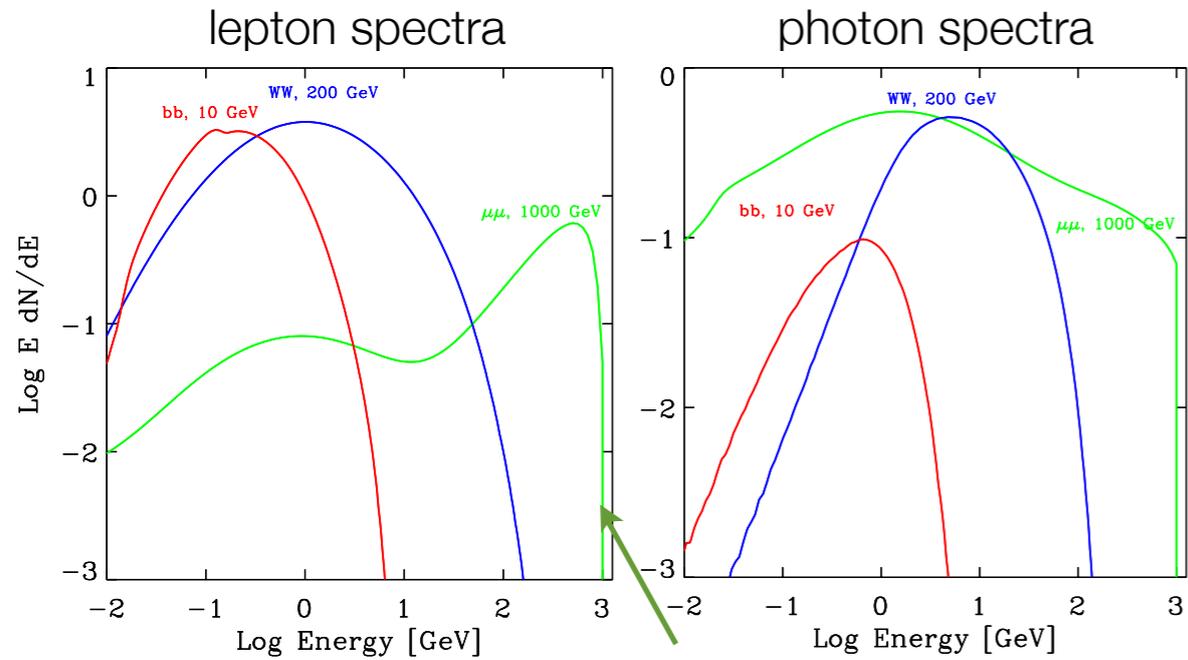
A.A.Zdziarski and R.Svensson, APJ, 1989 ; T.Slatyer et al., PRD, 2009



- There is a  $z$ -dependent transparency window.
- In the  $z \sim 100 - 1000$  region energy deposition is almost “local” (as for electrons).

# DM energy depositions

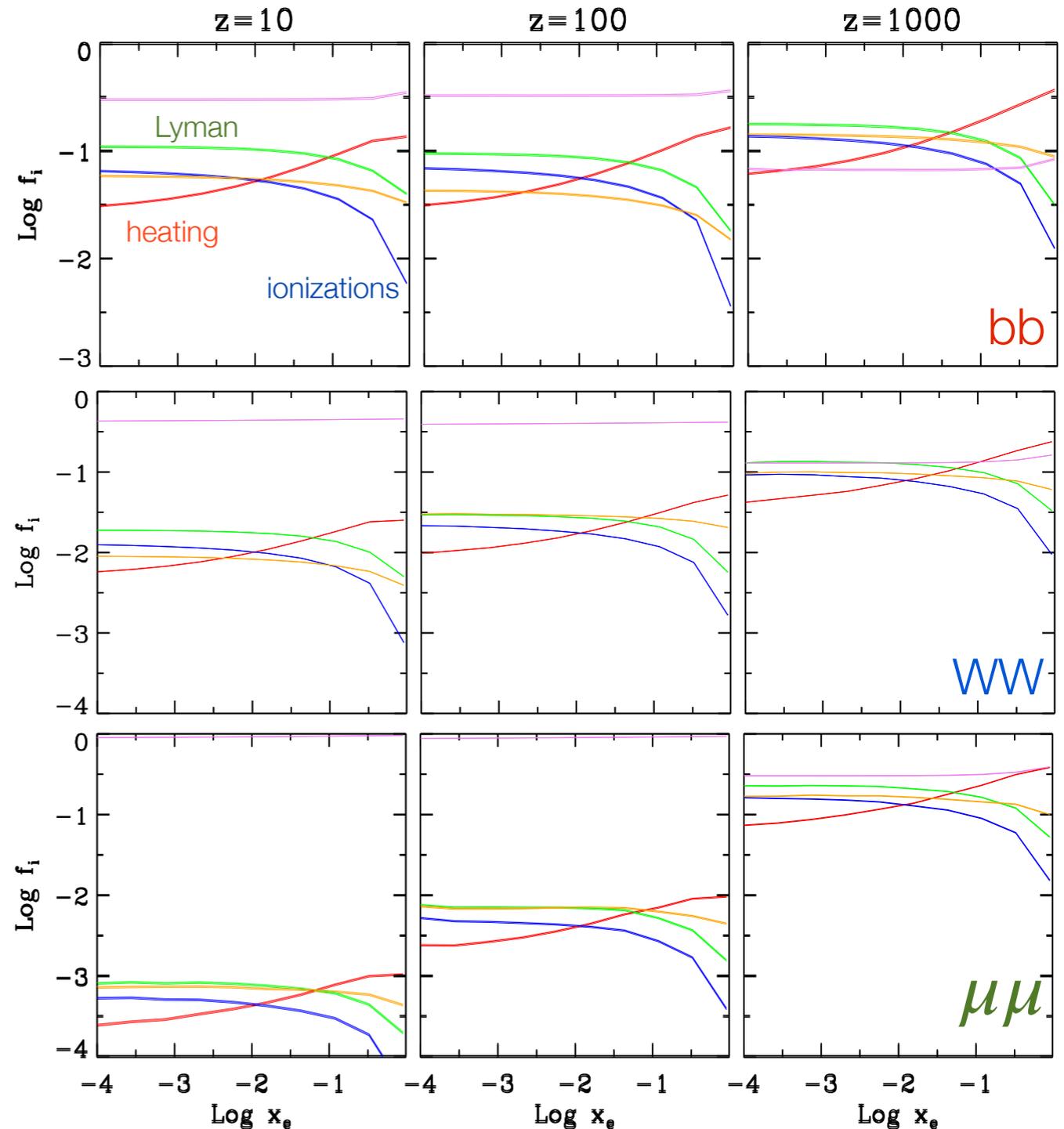
CE, M.Valdés, A.Ferrara and N.Yoshida, MNRAS, 2012



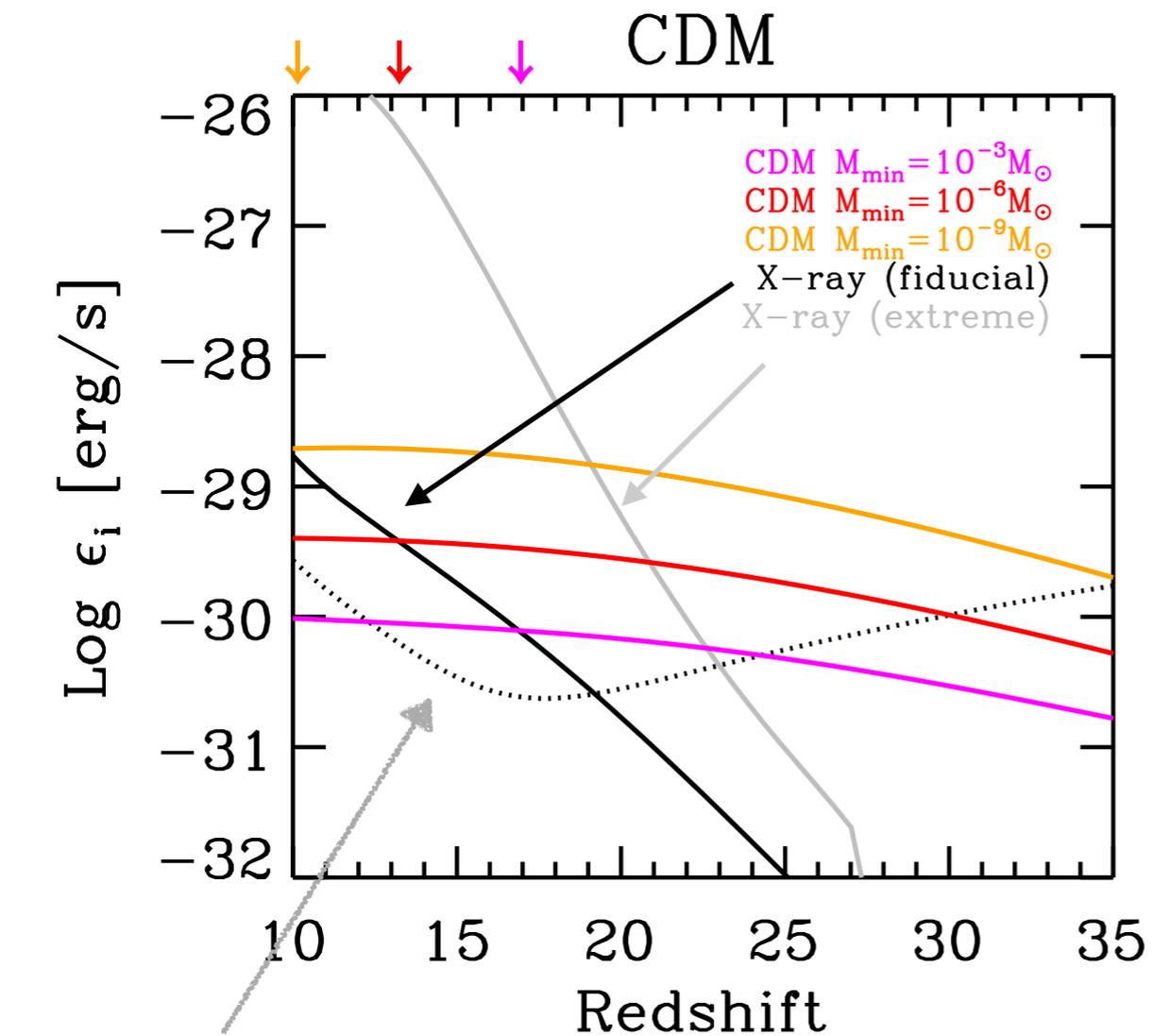
Ciafaloni et al., JCAP, 2011

DarkSUSY v.5.0.5

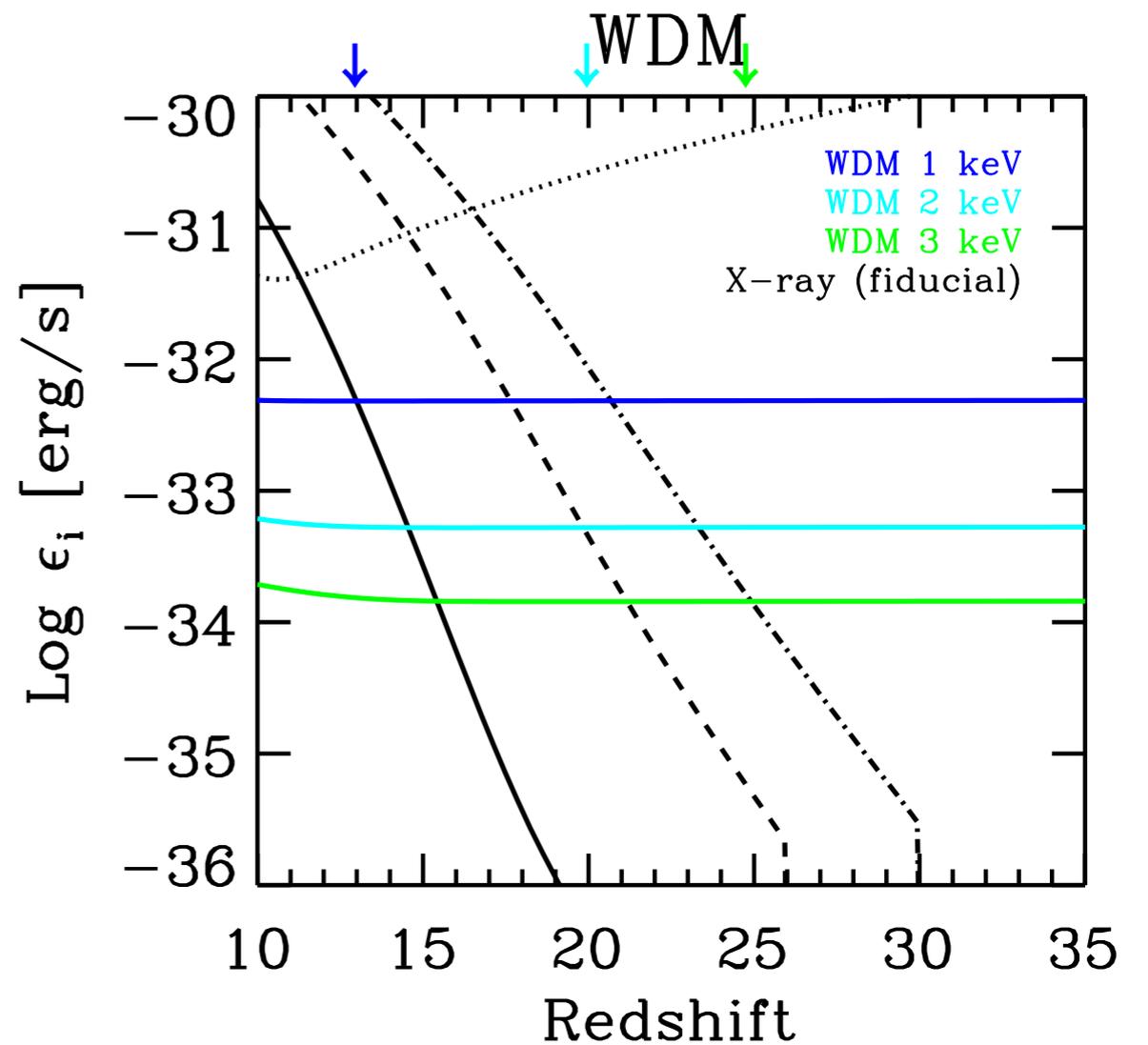
- heating grows with  $x_e$
- ionizations and Lyman photons grows with  $x_H$
- “high”-energy photons dominate energy depositions above 1 GeV
- “low”-energy photons dominate energy depositions below 10 MeV



# Energy deposition rates



CDM  $\chi\chi \rightarrow e^+e^-$

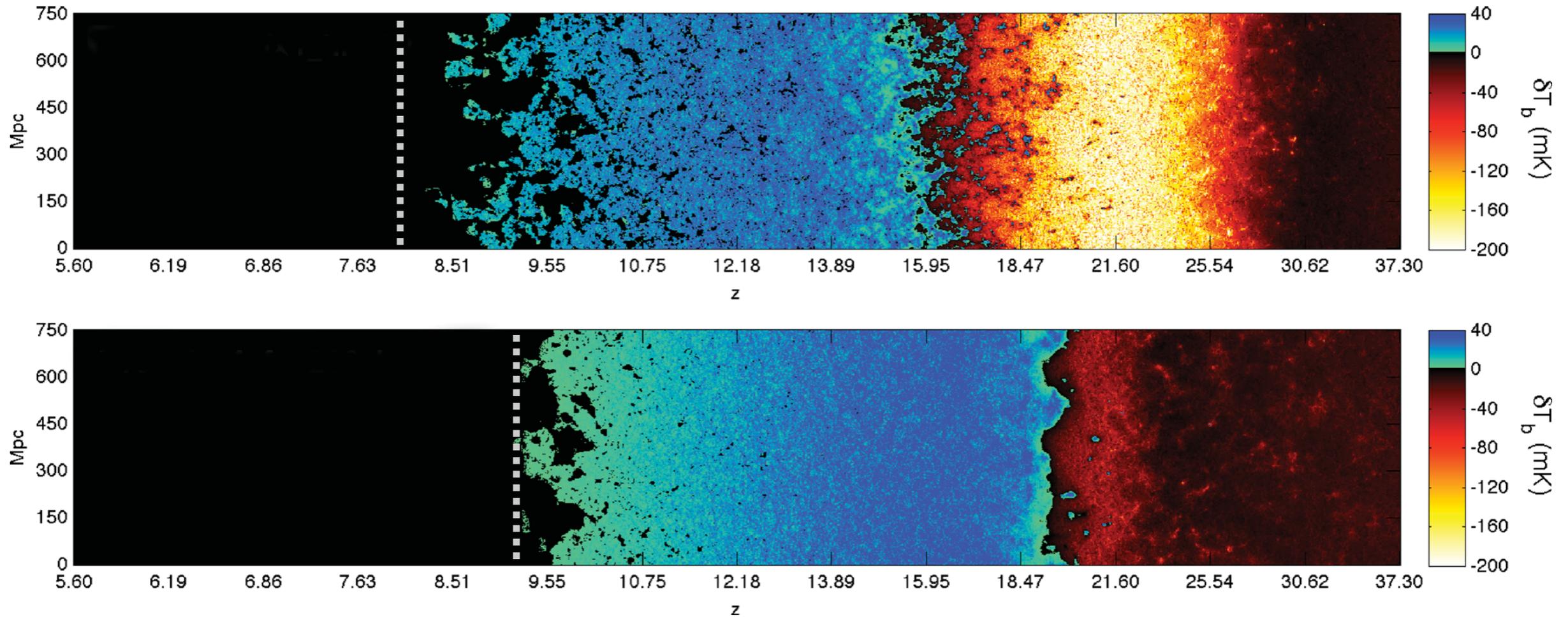


WDM  $\nu_s \rightarrow \nu_a + \gamma$

adiabatic cooling

# Reionization models

reionization



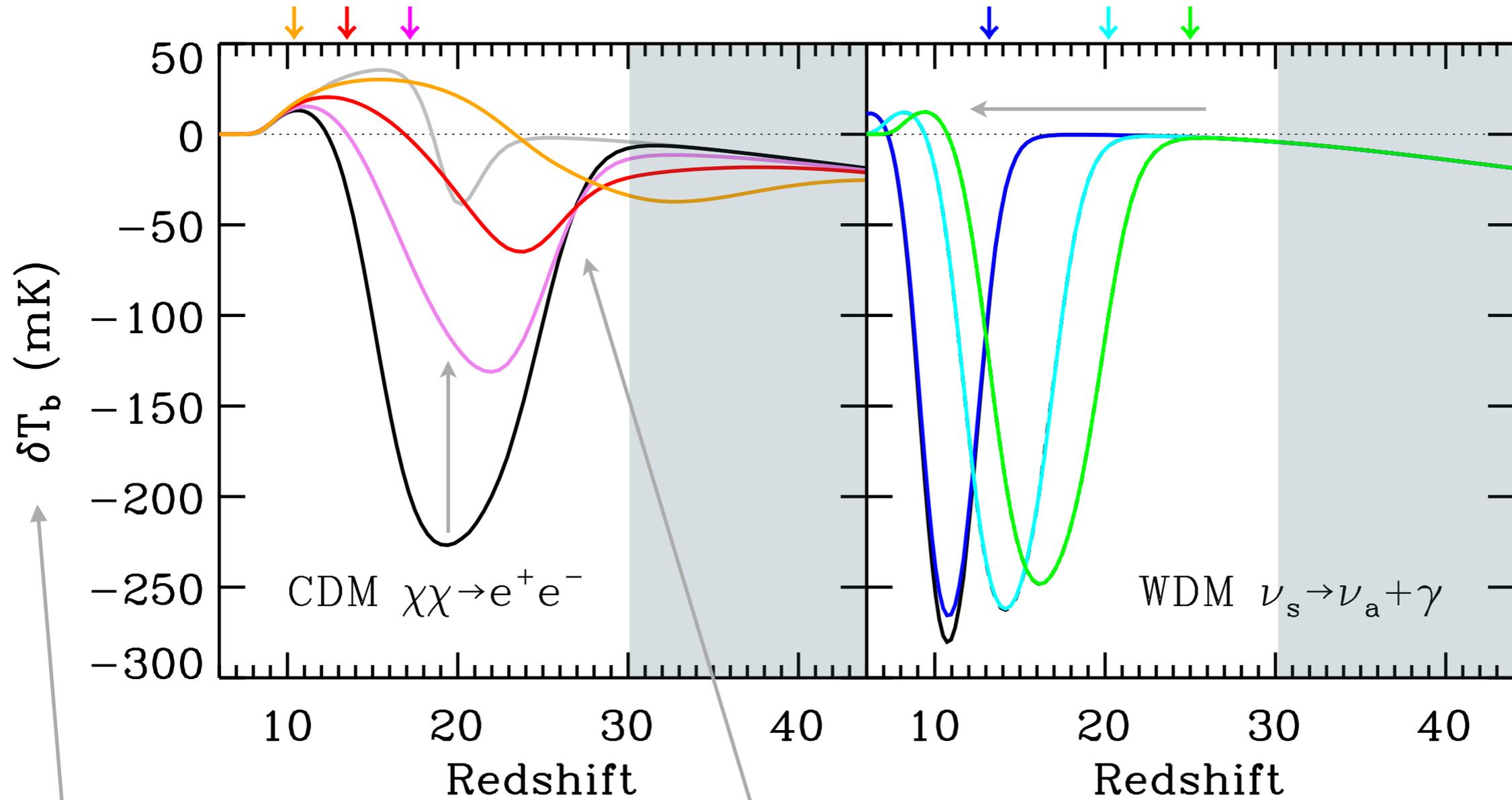
*fiducial (0.2 X-ray per stellar baryon, UV dominated)*

VS

*extreme (4000 X-ray per stellar baryon, X-ray dominated)*

# The 21 cm global signal

M.Valdés, CE, A.Mesinger, A.Ferrara & N.Yoshida, MNRAS, 2013

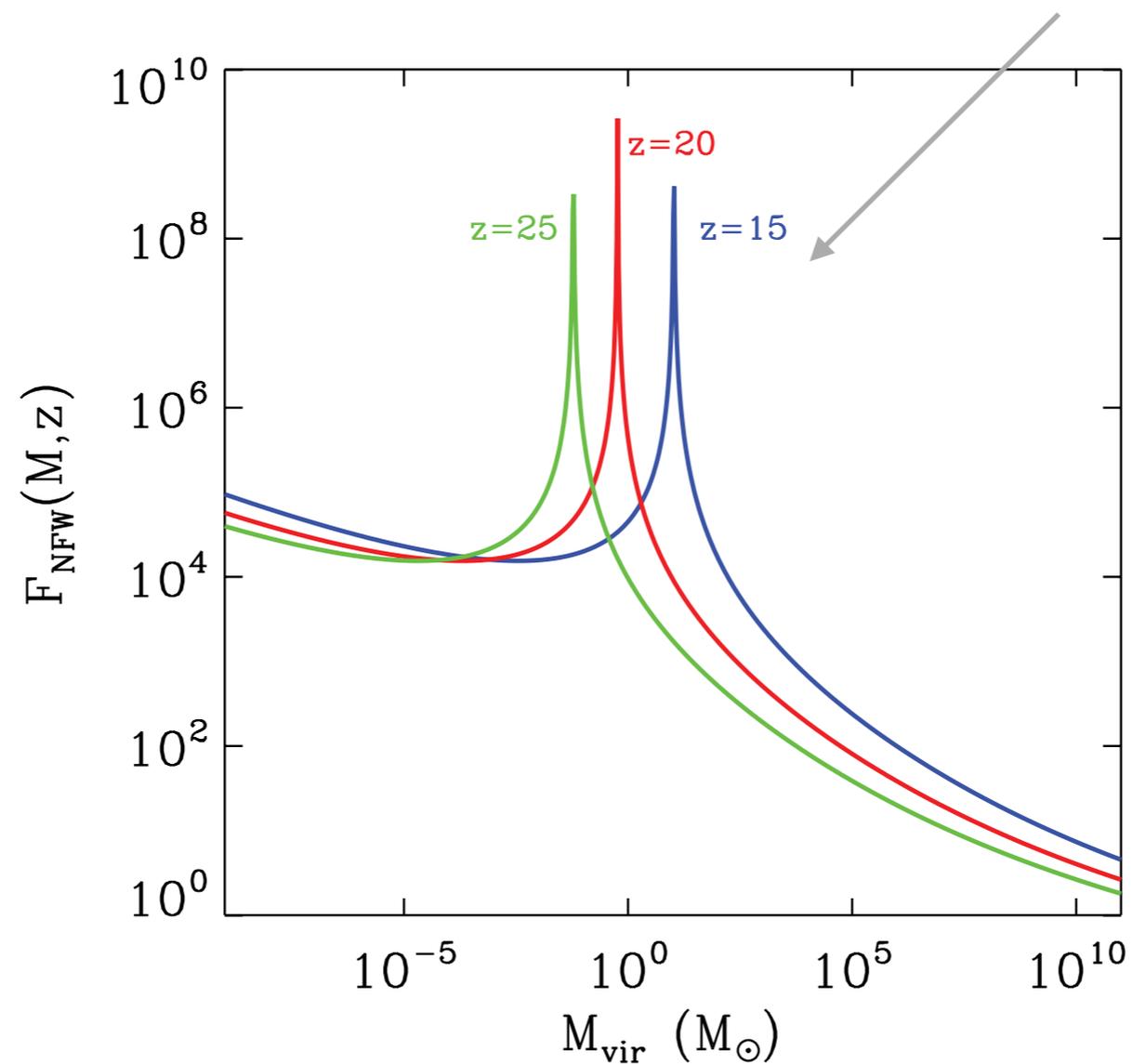


differential brightness  
Temperature

DM Ly- $\alpha$  couples  
 $T_s \rightarrow T_k$  earlier

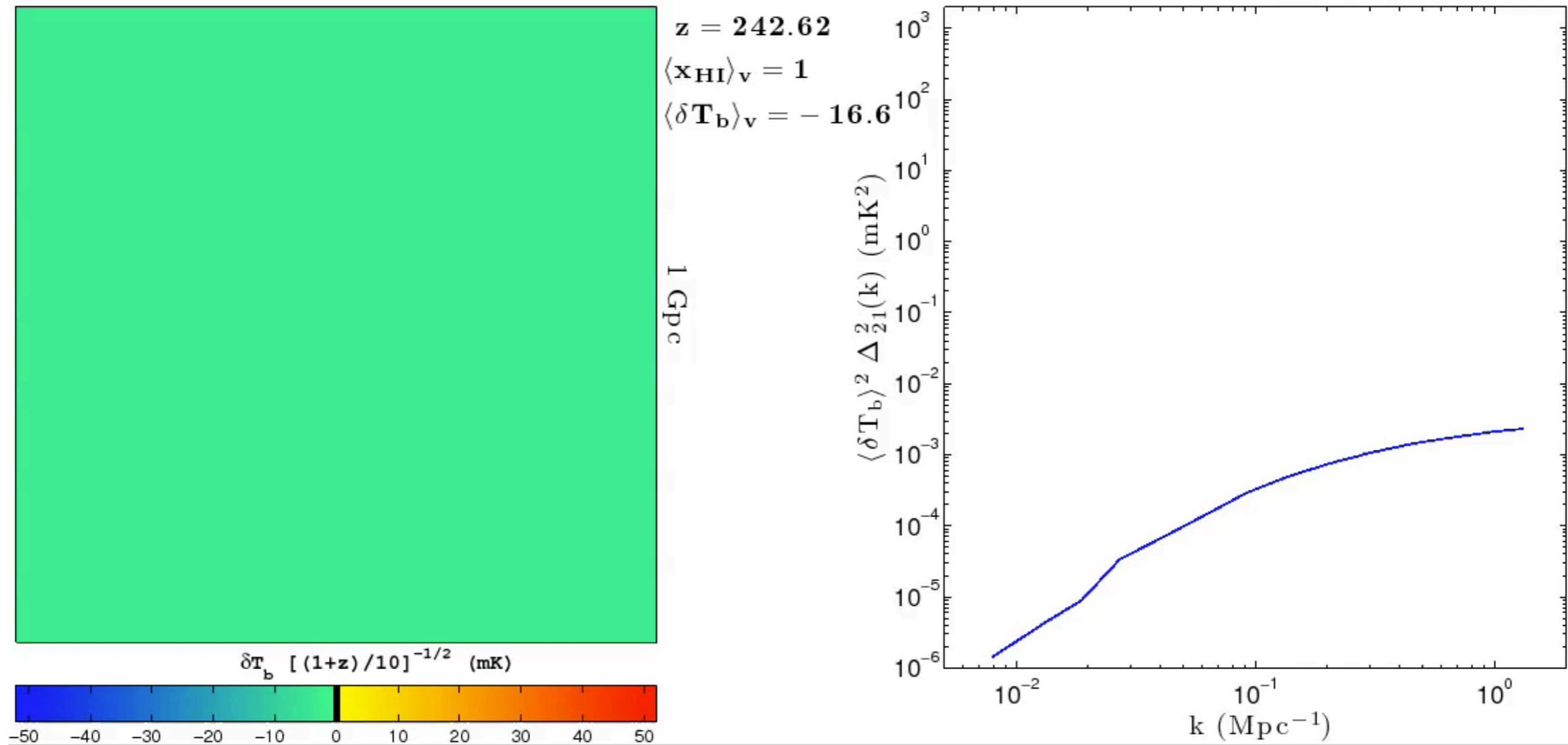
# DM heating is uniform

$$A^{\text{struct}}(z) = \frac{\langle \sigma v \rangle}{2 m_\chi^2} \int dM \frac{dn}{dM}(z, M) (1+z)^3 \int dr 4\pi r^2 \rho_i^2(r, M(z))$$



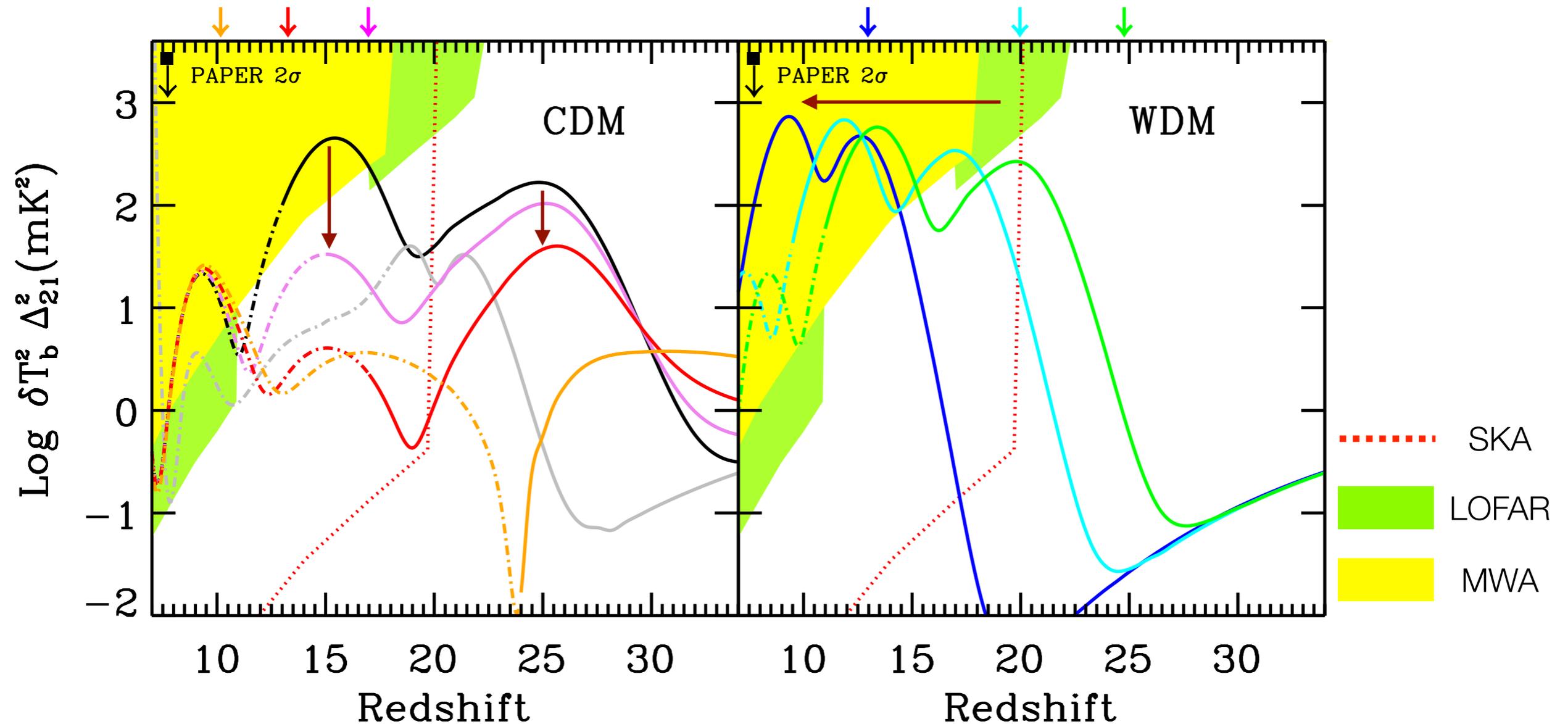
# Astrophysical heating sources are biased

21cmFAST code



# Power-spectrum

CE, A.Mesinger and A.Ferrara, *to be finished soon*



- *Robust* DM identification:
- when DM: the “heating” peak is lower than the “Ly-alpha” peak.
- when DM: the “heating” peak is in emission.

# Conclusions

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- Energy injection by DM annihilation at  $z \sim 100-1000$  can be properly constrained after accounting for degeneracies with cosmological parameters.
- We have developed a new model, which includes lepton and photon interactions with the IGM, allowing us to compute the energy partition into heating, excitations and ionizations as a function of the **electron initial energy**, the **ionization fraction** and the **redshift**.
- Our results can be applied to calculate DM contribution to the 21 cm cosmological evolution. Observations of the 21 cm signal from the *Dark Ages* could eventually detect a **robust** signal of exotic processes like DM decays or annihilations.

