

Studying the Z boson resonance

– not with weak loops
– but with the SMATASY/ZFITTER-kit

Tord Riemann, DESY Thanks to: M. Grünewald + S. Riemann



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<https://indico.cern.ch/event/369827/overview>

Present interest in precision approaches to the Z boson

Belle-II

$\sqrt{s} \sim 10 \text{ GeV}$

→ **Belle-II** will measure $10^9 \mu^+ \mu^-$ events

See e.g. T.Ferber [1]

Fcc-ee

$\sqrt{s} \sim M_Z$

→ **Fcc-ee** expects 10^{13} events at the Z resonance

See e.g. A. Freitas [2].

Much work on weak two-loop contributions to the Z resonance has been done by Hollik et al., Czakon et al., Freitas et al.

see [3] and many refs. therein.

Model-independent alternative: How to do?

→ Request by the Fcc-ee physics study group

→ **S-matrix approach** a la SMATASY/ZFITTER

See T.Riemann [4]

Outline

S-matrix approach to the Z line shape

- Developed as a model-independent analysis tool of $e^+e^- \rightarrow (\gamma, Z) \rightarrow f^+f^-$ **around the Z boson resonance**
- Aim: determinations of M_Z and Γ_Z
 - $\rightarrow \sigma_T$: Leike/TR/Rose 1991 [5]
 - $\rightarrow A_{FB,LR,pol}$: TR 1992 [6],
 - \rightarrow SMATASY code: Kirsch/TR 1994 [7]
- First application: LEP/L3 1993 [8], also: Tristan/TOPAZ, VENUS, LEP/OPAL, ...

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- 3 Asymmetries
- 4 Applications
- 5 SMATASY/ZFITTER
- 6 Summary

Introduction

The reaction

$$e^+e^- \rightarrow (\gamma, Z) \rightarrow f^+f^- + (n\gamma) \quad (1)$$

allows to study the Z boson, its mass M_Z , its width Γ_Z , its couplings, and potentially deviations from the Standard Model.

Need correct “model”

See experiences with *constant* and *s-dependent* Z width:

$$\frac{1}{[s - M_Z^2 + iM_Z \Gamma_Z(s)]} \quad \text{versus} \quad \frac{1}{[s - M_Z^2 + iM_Z \Gamma_Z]} \quad (2)$$

To a very good accuracy, it holds: $\Gamma_Z(s) \approx s/M_Z^2 \times \Gamma_Z$

see next slide, \rightarrow Bardin/Leike/Riemann/Schwartz 1988 [9]; also: Berends/Burgers/Hollik/v.Neerven 1988 [10]

Need correct unfolding ..

.. of *Realistic Observables* in order to get *Pseudo Observables*. \rightarrow e.g.: Borrelli/Consoli/Maiani/Sisto 1990 [11], Later: Bardin/Passarino 1999 [12], Bardin/Grünewald/Passarino 1999 [13], Passarino 2003 [14], Passarino 2013 [15] and refs. therein.

Lesson: The model influences numerical results

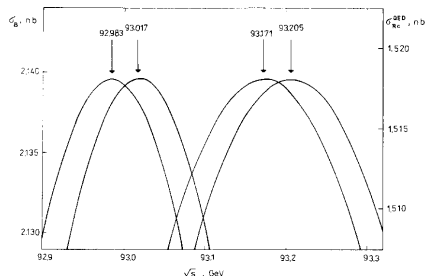


Fig. 1. Total cross sections σ_B , σ_{RC}^{QED} for the reactions $e^+e^- \rightarrow \mu^+\mu^- (\gamma)$ in Born approximation (left scale) and including $O(\alpha)$ QED corrections (right scale). Peaks of σ with energy-dependent width $\Gamma_Z(s)$ are shifted by 34 MeV to the left.

Total cross section for $e^+e^- \rightarrow \mu^+\mu^-$ production at LEP ...

... without (left) and with (right) QED corrections. Both sample data produced with an energy-dependent Z width.

The assumptions on the Z -propagator in the fit formulas influence the location of the peak, but not the “experimental errors”.

Fig.: from [9], license Number: 3557090997554.

	Born		Born + QED	
from fit: \rightarrow	M_Z	Γ_Z	M_Z	Γ_Z
$\Gamma_Z(s)$	93.000 \pm 0.013	2.498 \pm 0.009	93.000 \pm 0.016	2.498 \pm 0.011
Γ_Z	92.966 \pm 0.013	2.498 \pm 0.009	92.966 \pm 0.016	2.498 \pm 0.011

Introduction

Stuart 1991 [16], S-Matrix ansatz for $e^+e^- \rightarrow Z \rightarrow f^+f^-$

$$M = \frac{R}{s - s_0} + F(s), \quad s_0 = M_Z^2 - iM_Z\Gamma_Z \quad (3)$$

Allows to study:

- Mass M_Z and width $\Gamma_Z \rightarrow$ Leike/Riemann/Rose 1991 [5]
- How many independent degrees of freedom? \rightarrow Leike/Riemann/Rose 1991 [5], Kirsch/S.Riemann, L3 [17, 8]
- But also: How to define mass and width of the Z boson at higher orders of perturbation theory? \rightarrow Denner 2014 [18], Freitas 2014 [3] and Fcc-ee [2], Degrassi FCC-ee [19] and refs. therein

Total cross sections

There are immediate questions, from an experimental point:

- What about the photon exchange?
- What about QED corrections, e.g. the $2 \rightarrow 3$ part of the cross sections?
- What about asymmetries, besides σ_{tot} ?

We have to describe

$$e^+ e^- \longrightarrow (\gamma, Z) \longrightarrow f^+ f^- (\gamma) \quad (4)$$

Ansatz in the complex energy plane, for **four helicity matrix elements**:

$$\mathcal{M}^i(s) = \frac{R_\gamma^i}{s} + \frac{R_Z^i}{s - s_Z} + F^i(s), \quad i = 1, \dots, 4. \quad (5)$$

Beware: Eqn. (5) is mathematically not consistent \rightarrow Böhm/Sato 2004 [20]

The poles of \mathcal{M} have complex residua R_Z and R_γ , the latter corresponding to the photon, and the background $F(s)$ is an analytic function without poles:

$$F^i(s) = \sum_{n=0}^{\infty} F_n^i (s - s_0)^n \quad (6)$$

Comment on the photon term (3 Feb 2015)

$$\begin{aligned}
 \frac{R_\gamma^i(s)}{s} &= \frac{\sum_{n=0}^{\infty} R_n^i(s-s_0)^n}{s} & (7) \\
 &= \frac{\sum_{n=0}^{\infty} R_n^i(s-s_0)^n}{s_0 - (s_0 - s)} \\
 &= \sum_{n=0}^{\infty} R_n^i(s-s_0)^n \frac{1}{s_0} \frac{1}{1 - \frac{s_0 - s}{s_0}} \\
 &= \sum_{n=0}^{\infty} R_n^i(s-s_0)^n \frac{1}{s_0} \left[1 + \frac{s_0 - s}{s_0} + \left(\frac{s_0 - s}{s_0} \right)^2 \cdots \right]
 \end{aligned}$$

The term $R_\gamma^i(s)/s$ is part of the the background term $F(s)$.

- *It is useful to sum up a selected part of the photonic background of the Z resonance in order to take explicit notice of physically known pieces of the input expressions.*
- *Compare: It is useful to sum up a selected part of self-energy insertions in the propagators in order to derive the Breit-Wigner resonance form.*

Ansatz for realistic applications

The analysis of the Z line shape will be based here on the cross section

$$\sigma(s) = \sum_{i=1}^4 \sigma^i(s) = \frac{1}{4} \sum_{i=1}^4 s |\mathcal{M}^i(s)|^2, \quad (8)$$

where the sum must be performed over four helicity amplitudes with different residues R_Z^i and functions $F^i(s)$. The result is, with QED corrections folded in:

$$\sigma_T(s) = \frac{4}{3} \pi \alpha^2 \int \frac{ds'}{s} \left[\frac{r^\gamma}{s} + \frac{sR + (s - M_Z^2)J}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \dots \right] \rho_{ini} \left(\frac{s'}{s} \right). \quad (9)$$

The radiation connected with initial-final state interferences can be taken into account by an analogue formula to (9) with a slightly more complicated structure [21, 22]:

$$\sigma_{\text{int}}(s) = \int ds' \sigma(s, s') \rho_{\text{int}}(s'/s). \quad (10)$$

Some details

See [\[23\]](#)

ZFITTER

The Standard Model analysis tool for the Z resonance: ZFITTER (D. Bardin et al.)

- Complete electroweak radiative corrections;
- QED corrections by convolution:
with some $\sigma_0(s')$
beware: for initial-final state interferences with some $\sigma_0(s, s')$;
- **semi-analytical QED** integrations;
- free choice of $\sigma_0(s')$ by **user interfaces**;
- Standard Model interfaces: four weak form factors $\rho, \kappa_e, \kappa_f, \kappa_{ef}$.

ZFITTER is **well-tested**, **flexible**, **accurate** and **fast**.

References: CPC in 2001 [24] and in 2006 [25], also: CPC in 1990 [26]

- The actual Fortran package is v.6.44; version 6.42 is public in CPC program library, with CPC-licence.
- Beware: **Gfitter/GSM (2007-2011) is an illegal clone of ZFITTER**. The code is available at <http://zfitter-gfitter.desy.de/> and <http://fh.desy.de/projekte/gfitter01/Gfitter01.htm>.
For ZFITTER, see also: <http://sanc.jinr.ru/users/zfitter>

catalogue numbers: **ADMJ_v1_0 (2001), ADMJ_v2_0 (2006)**
<http://www.cpc.cs.qub.ac.uk/licence/licence.html>



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Born Asymmetries

On a Sunday in Summer 1992, I had a discussion with Luciano Maiani in the CERN library. He had doubt that, **for asymmetries**, an analogue to the model-independent ansatz for σ_{tot} might be usefully formulated, especially in view of the QED corrections.

I believed one can do that, and I followed the rule “The proof of the pudding is in the eating” [6]. The result:

$$A^{Born}(s) = A_0 + A_1 \left(\frac{s}{M_Z^2} - 1 \right) + A_2 \left(\frac{s}{M_Z^2} - 1 \right)^2 + \dots \quad (11)$$

$$A_{FB} = \frac{\sigma_{FB}}{\sigma_T}, \quad A_{pol} = \frac{\sigma_{pol}}{\sigma_T}. \quad (12)$$

The A_{FB} and A_{pol} are helicity combinations as also σ_T is, i.e. have also S-matrix ansatzes.

The parameters in $A^{Born}(s)$ are in [QED-]Born approximation:

$$A_0 = \frac{R_A}{R_T}, \quad (13)$$

and

$$A_1 = \left[\frac{J_A}{R_A} - \frac{J_T}{R_T} \right] A_0. \quad (14)$$

Some details

See [\[23\]](#)

QED corrections for asymmetries

QED corrections to asymmetries lead to few simple correction factors [6, 7]:

$$A_{LR}^{QED}(s) = A_{0,LR}^{Born} + c_{1,T}(s) A_{1,LR}^{Born} \left(\frac{s}{M_Z^2} - 1 \right) + \dots \quad (15)$$

$$A_{FB}^{QED}(s) = c_{0,FB}(s) A_{0,FB}^{Born} + c_{1,FB}(s) A_{1,FB}^{Born} \left(\frac{s}{M_Z^2} - 1 \right) + \dots \quad (16)$$

The A_0 and A_1 are constant, and the same as in Born approximation.
The QED corrections are contained in the model-independent factor $C(s)$.

$$c_{0,FB}(s) = \frac{C_{FB}^R}{C_T^R}, \quad c_{0,T}(s) = 1 \quad (17)$$

$$c_{1,A}(s) = c_{0,A} \frac{C_T^J}{C_T^R} \quad (18)$$

Sample QED factor

$$C_{T,FB}^R(s) = \int dk \rho_{T,FB}(s'/s) \frac{s'R}{sR} \frac{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}{(s' - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} \quad (19)$$

Sketch of derivation of the expression for $A_{FB}(s)$ with QED corr's

See [23]

QED corrections for asymmetries

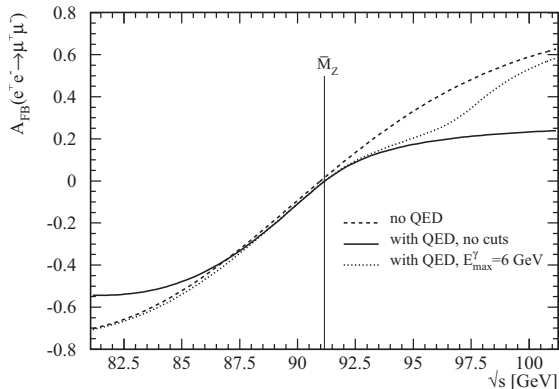


Figure 1 : The forward-backward asymmetry for the process $e^+e^- \rightarrow \mu^+\mu^-$ near the Z boson peak. From Kirsch/Riemann 1994 [7], license Number: 3557090997554.

Applications

In Leike/S.Riemann/Riemann 1992 [27] correlations are discussed.

For the Z **peak position** s_{peak} , one may derive the relation:

$$\Delta\sqrt{s_{peak}} = \Delta M_Z + \frac{1}{4} \frac{\Gamma_Z^2}{M_Z} \Delta \left(\frac{J_T}{R_T} \right) + \dots \quad (20)$$

between an uncertainty in M_Z and an uncertainty in the γZ interference. The latter also influences A_1 .

Similarly, for a **hypothetical heavy gauge boson** Z' , the effects from its virtual exchange transform after a partial fraction decomposition into simple shifts of the γZ interferences [27]:

$$\Delta \left(\frac{J_T}{R_T} \right) = -2 \frac{g'^2}{g^2} \frac{M_{Z'}^2}{M_{Z'}^2 - M_Z^2} \frac{(a_e a'_e + v_e v'_e)(a_f a'_f + v_f v'_f)}{(a_e^2 + v_e^2)(a_f^2 + v_f^2)}, \quad (21)$$

Correlations: L3 at LEP1, 1993

Table 1

Results of the S matrix fit to total cross-sections and forward-backward asymmetries: (a) all parameters except the photon exchange are left free; (b) in addition the γZ interference terms are fixed to the Standard Model expectation.

Parameter	Case (a)	Case (b)
\overline{m}_Z (GeV)	91.152 ± 0.015	91.160 ± 0.010
$\overline{\Gamma}_Z$ (GeV)	2.494 ± 0.012	2.492 ± 0.012
$r_{\text{tot}}^{\text{lep}}$	0.141 ± 0.002	0.140 ± 0.002
$j_{\text{tot}}^{\text{lep}}$	0.032 ± 0.064	fixed to 0.0058
$r_{\text{fb}}^{\text{lep}}$	0.004 ± 0.001	0.004 ± 0.001
$j_{\text{fb}}^{\text{lep}}$	0.674 ± 0.087	0.675 ± 0.087
$r_{\text{tot}}^{\text{had}}$	2.859 ± 0.030	2.855 ± 0.029
$j_{\text{tot}}^{\text{had}}$	0.720 ± 0.700	fixed to 0.219

Table 2

Results of the S matrix fit to total cross-sections, forward-backward asymmetries and τ polarization: (a) all parameters except the photon exchange are left free; (b) in addition the hadronic γZ interference terms for the total cross-section are fixed to the Standard Model expectation.

Parameter	Case (a)	Case (b)
\overline{m}_Z (GeV)	91.155 ± 0.013	91.160 ± 0.010
$\overline{\Gamma}_Z$ (GeV)	2.494 ± 0.012	2.492 ± 0.012
$R_Z^{\text{lep}0}$	0.429 ± 0.012	0.429 ± 0.012
$R_Z^{\text{lep}1}$	-0.370 ± 0.003	-0.370 ± 0.003
$R_Z^{\text{lep}2}$	0.323 ± 0.016	0.323 ± 0.016
$r_{\text{tot}}^{\text{had}}$	2.860 ± 0.030	2.856 ± 0.029
$j_{\text{tot}}^{\text{had}}$	0.620 ± 0.620	fixed to 0.219
m_Z (GeV)	91.189 ± 0.013	91.194 ± 0.010
Γ_Z (GeV)	2.495 ± 0.012	2.493 ± 0.012
$\widehat{g}_n^{\text{lep}}$	-0.037 ± 0.010	-0.037 ± 0.010
$\widehat{g}_a^{\text{lep}}$	-0.4991 ± 0.0019	-0.4988 ± 0.0019
$\sin^2 \widehat{\theta}_W$	0.2317 ± 0.0037	0.2316 ± 0.0037

From: L3-Collaboration, "An S matrix analysis of the Z resonance", Phys. Lett. B315 (1993) 494-502.

See also **LEPEWWG et al., Phys. Rept. 2006**, section 2 [28] and **LEPEWWG et al., Phys. Rept. 2013**, App. A [29]

Correlations: LEP1 + LEP2

A complete analysis of the LEP-2 data in terms of J_{had}^{tot} is lacking.
But see K. Sachs, 2003 [30].

Including more measurements from LEP2 solves this problem, reducing the correlation. The final result of $M_Z = 91\,186.9 \pm 2.3 \text{ MeV}$ ⁸ is in very good agreement with the result of the standard lineshape fit $M_Z = 91\,187.6 \pm 2.1 \text{ MeV}$ ⁹ with only slightly increased error.

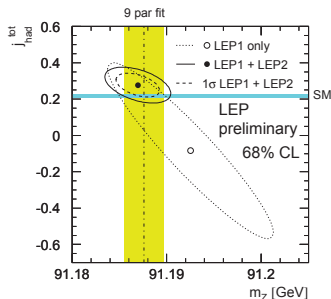


Figure 2: Correlation between the mass of the Z and J_{had}^{tot} . Results are shown for LEP1 data only and for a combined fit to LEP1 and LEP2 data. The yellow band indicates the 1σ error from the 9 parameter fit.

Figure 2 : K. Sachs, “Standard model at LEP II”, talk held at Moriond 2001, fig. 2 [30]

Correlations: LEP1 + TRISTAN

LEP experiments use cross-section and forward-backward asymmetry results from $\sqrt{s} \sim M_Z$ and LEP II. OPAL and L3 have reported preliminary results which are given in Table 1, and are compared to the value obtained by VENUS [9] using data at $\sqrt{s} \sim 60$ GeV and preliminary LEP I S-Matrix results. The results are consistent with each other, and with the SM prediction $j_{had}^{tot} = 0.22$.

Expt	Data	j_{had}^{tot}
L3:	LEP I + LEP II	0.30 ± 0.10
OPAL:	LEP I + LEP II	0.21 ± 0.12
VENUS:	VENUS + LEP I	0.20 ± 0.08

Table 1: Measurements of j_{had}^{tot}

Figure 3 : P. Holt, "Fermion pair production above the Z^0 resonance", talk held at HEP 2001, table 1 [31]

Fortran programs: ZPOLE - ZUSMAT - SMATASY/ZFITTER

ZPOLE – The stand-alone Fortran test package (Leike/Riemann, v.0.5, July 1991) is available on request.

It was used for the numerics of [5].

ZUSMAT ... the S-Matrix interface of older ZFITTER versions.

ZUSMAT was used for analysing the total cross sections, but could not treat asymmetries.

SMATASY/ZFITTER – With interface package **SMATASY** one has the full functionality of **ZFITTER** corrections [24, 25, 32].

The actual Fortran program for the S-matrix Z line shape approach:

M. Grünewald, S. Kirsch, T. Riemann 1994 [7]

SMATASY v.6.42.01 = SMATA642 (2 June 2005)

available at <https://gruenew.web.cern.ch/gruenew/smatasy.html>

ADBG_v1_0 (1995) <http://www.cpc.cs.qub.ac.uk/licence/licence.html>



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Good Scientific Practice – when using software ...

There is practically no literature on that. Software is never mentioned.

The general rules of international academic basic research apply. They include:

- Attribution of a scientific achievement to those who made it.
- Often by a proper citation.
- Other proper attributions are possible.

A proper attribution informs on:

- What was done?
- Who did it?
- How important is it?

In case of ethical or legal problems, questions have to be answered, in this order:

- On facts: What are the initial facts?
- On rules: Are there rules, are they violated?
- On sanctions: in case. Are there sanctions foreseen? By whom?

Reference [33]:

F. Carminati (CERN), D. Perret-Gallix (LAPP), T. Riemann (DESY)

“Summary of the ACAT 2013 round table discussion: Open-source, knowledge sharing and scientific collaboration”

arXiv.1407.0540, DESY-14-079, in proceedings of ACAT 2013

Summary

- The S-matrix approach is absolutely independent of the Standard Model approach.
- The degrees of freedom for σ_{tot} are, at minimum:

M_Z

Γ_Z

R – the residue of the Z resonance, *per scattering channel*

J – the value of the γZ interference, *per scattering channel*

- So we have at least **four degrees of freedom**.
This deserves at least **five data points** as a function of s .

See also: M.Grünewald, S.Kirsch 1993 [34].

- **Asymmetries** may be described as well as σ_{tot} .
- For an exact numerical analysis of data, an **accurate description of QED** corrections is mandatory.
This has been realised by combining SMATASY with ZFITTER.
- With so much more statistics at the Fcc-ee compared to LEP-1 and LEP-2:

The S-matrix approach might gain at the Fcc-ee even more interest as an alternative to the Standard Model approach.

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