

Towards NLL threshold resummation for $2 \rightarrow 3$ type processes

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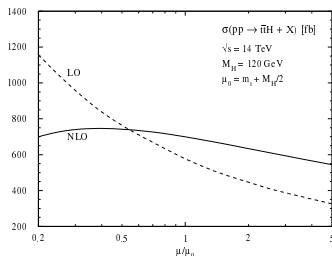
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Resummation for more final state particles

- General formulation available for $2 \rightarrow n$ [Bonciani et al. , '03][Mert-Aybat et al. , '06]
- Recent advances in automatization of resummation in $2 \rightarrow 5$ (massless) [Gerwick et al. , '14]
- Relation to event shapes and N-jettiness in SCET
- Practical applications of threshold resummation (in traditional, conjugate space formulation) to production of more than two particles in the final state?

Importance of $pp \rightarrow t\bar{t}H$

- A Higgs boson found with a mass of 125 GeV
- Precision study needed to determine if it is SM Higgs
- Direct way to access Yukawa coupling
- QCD Corrections up to NLO [*Beenakker et al. , '02*] [*Dawson et al. , '02*]
- Matched to parton showers by: aMC@NLO [*Frederix et al. , '11*], PowHel [*Garzelli et al. , '11*], Sherpa [*Hoeche et al., '12*], POWHEG-BOX [*Hartanto et al. , '14*]
- Electroweak correction [*Frixione et al. , '14*][*Zhang, '14*]



Currently used results

\sqrt{S}	m_H [GeV]	σ [pb]	Scale%	(PDF+ α_s) %
8.0 TeV	125.0	0.1293	+3.8 -9.3	+8.1 -8.1
13.0 TeV	125.0	0.5085	+5.7 -9.3	+8.8 -8.8
13.5 TeV	125.0	0.5608	+5.8 -9.3	+8.9 -8.9
14.0 TeV	125.0	0.6113	+5.9 -9.3	+8.9 -8.9

[LHCHSWG, '15]

Why resummation for $t\bar{t}H$?

Gains

- NNLO corrections out of reach
- Resummation can help reduce scale uncertainty
- Good process to start:
 - Simple color structure
 - Massive particles \rightarrow no final state collinear divergences

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Pitfalls

- $2 \rightarrow 3$ phase space suppressed near threshold ($\sigma \propto \beta^4$)
- Resummation not an approximation for NNLO

Definition of Threshold

Threshold variable $\hat{\tau} = \frac{M^2}{\hat{s}}$ (Other definitions also possible)

$$M = m_H + 2m_t$$

$\sqrt{\hat{s}}$: the partonic center of mass energy

$$\beta^2 = 1 - \hat{\tau} = 1 - \frac{M^2}{\hat{s}} \sim \frac{\text{maximum energy of the emitted gluons}}{\text{total available energy}}$$

At NLO: [Beenakker et al. , '02]

$$\sigma_{\text{LO}} \times \frac{2\alpha_s}{\pi} \left\{ (T_1^2 + T_2^2) \left[2 \log^2 \beta - 3 \log \beta - 2 \log \beta \log \left(\frac{\mu_F}{2m_t + m_H} \right) \right] - (T_3 + T_4)^2 \log \beta \right\}$$

Coulomb enhanced terms not treated in the resummation

Mellin Transform

Mellin transform is used with respect to τ (needed for factorization of phase space):

$$\begin{aligned}\tilde{\sigma}_{pp \rightarrow t\bar{t}H}(N) &\equiv \int_0^1 d\tau \tau^{N-1} \sigma_{pp \rightarrow t\bar{t}H}(\tau, \mu_R, \mu_F) \\ &= \sum_{i,j} \tilde{f}_{i/p}(N+1, \mu_F) \tilde{f}_{j/p}(N+1, \mu_F) \tilde{\sigma}_{ij \rightarrow t\bar{t}H}(N, \mu_R, \mu_F)\end{aligned}$$

- $\tilde{f}_{i/p}(N+1, \mu_F)$: Mellin transform with respect to x
- $\tilde{\sigma}_{ij \rightarrow t\bar{t}H}(N, \mu_R, \mu_F)$: Mellin transform with respect to $\hat{\tau}$

$\log^n(1 - \hat{\tau}) \Rightarrow \log^n N$ and threshold $\hat{\tau} \rightarrow 1 \sim N \rightarrow \infty$

First application for 2 \rightarrow 3 in Mellin space

Orders of Resummation

Large logarithms $\log N \equiv L$ for $N \rightarrow \infty$

Perturbation needs to be reordered in α_s and L :

$$\tilde{\sigma} \sim \tilde{\sigma}_{LO} \times \mathcal{C}(\alpha_s) \exp [Lg_1(\alpha_s L) + g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \dots]$$

With orders of precision:

\Downarrow LL \Downarrow $\alpha_s^n \log^{n+1}(N)$	\Downarrow NLL \Downarrow $\alpha_s^n \log^n(N)$	\Downarrow NNLL \Downarrow $\alpha_s^{n+1} \log^n(N)$
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Exponential functions are universal for initial state emission

[Kodaira, Trentadue, '82][Sterman, '87][Catani, d'Emilio, Trentadue, '88][Catani, Trentadue, '89]

Soft anomalous dimension

Color space

- Using s-channel color basis
- Same as for example $pp \rightarrow t\bar{t}$

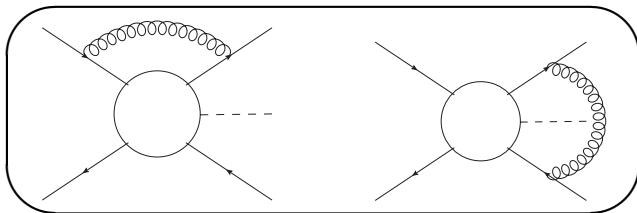
Soft anomalous dimension

Calculated at the hand of UV divergences of eikonal integrals

$$\Gamma_{ij \rightarrow klB,IJ}^{(1)} = -C_{IJ}^{ab} \text{Res} \{ \omega^{ab} \} - \Gamma_{ij \rightarrow C}^{(1)} \delta_{IJ}$$

- I and J : color indices
- a and b : colored particle indices
- C_{IJ}^{ab} : color factor of the exchange
- ω^{ab} : UV divergent terms of eikonal integrals

Eikonal integrals



Using feynman rules approximated for soft gluon emission (eikonal)

Soft anomalous dimension matrix

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,11}^{(1)} = -\frac{\alpha_s}{2\pi} 2C_F (L\beta_{34} + 1)$$

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,12}^{(1)} = \frac{\alpha_s}{2\pi} \frac{2C_F}{N_C} \Omega_3 = \frac{C_F}{2N_C} \Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,21}^{(1)}$$

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,22}^{(1)} = \frac{\alpha_s}{2\pi} [(N_C - 2C_F)(L\beta_{34} + 1) + N_C \Lambda_3 + (8C_F - 3N_C) \Omega_3]$$

$$\Omega_3 = (T_{13} + T_{24} - T_{14} - T_{23}) / 2$$

$$\Lambda_3 = (T_{13} + T_{24} + T_{14} + T_{23}) / 2$$

$$T_{ij} = \log \left(\frac{m_j^2 - t_{ij}}{m_j \sqrt{s}} \right) + \frac{i\pi - 1}{2}$$

with $t_{ij} = (p_i - p_j)^2$ and $\beta_{34}^2 = 1 - (m_3 + m_4)^2 / s_{34}$

Soft anomalous dimension matrix

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,11}^{(1)} = -\frac{\alpha_s}{2\pi} 2C_F (L\beta_{34} + 1) \quad \text{Agrees with } q\bar{q} \rightarrow Q\bar{Q}$$

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,12}^{(1)} = \frac{\alpha_s}{2\pi} \frac{2C_F}{N_C} \Omega_3 = \frac{C_F}{2N_C} \Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,21}^{(1)} \quad \text{for } p_5 \rightarrow 0 \text{ and } m_5 \rightarrow 0$$

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For 2 → 2:

$$\begin{array}{cc} \Downarrow & \Downarrow \\ t_{13} = t_{24} & s_{34} = \hat{s} \\ t_{23} = t_{14} & \beta_{34} = \beta \end{array}$$

Soft wide-angle emission

Threshold limit for soft anomalous dimension:

$$2\mathcal{R}e\Gamma_{q\bar{q} \rightarrow Q\bar{Q}B,IJ,\text{thr.}}^{(1)} = \frac{\alpha_s}{\pi} \text{diag}(0, -N_C)$$

$$2\mathcal{R}e\Gamma_{gg \rightarrow Q\bar{Q}B,IJ,\text{thr.}}^{(1)} = \frac{\alpha_s}{\pi} \text{diag}(0, -N_C, -N_C)$$

results in soft wide-angle contribution:

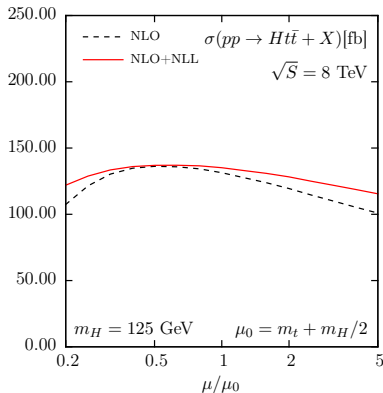
$$\Delta_I^{\text{NLL}} = \exp[h_2(\alpha_s L, -C_I)]$$

With C_I the quadratic Casimir invariant

Results

[Kulesza, Motyka, Stebel, VT, in preparation]

PDFs used: MMHT2014NLO



NLO+NLL is NLL matched to NLO

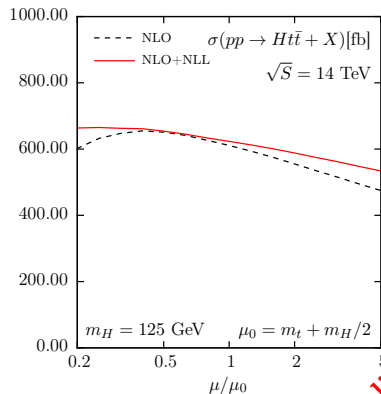
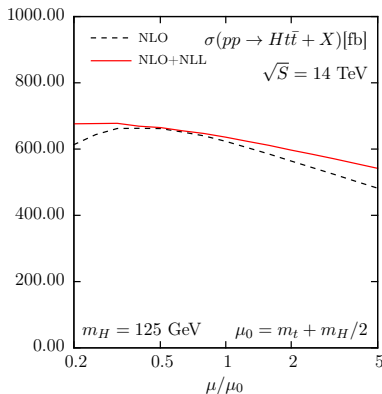
preliminary

Results

[Kulesza, Motyka, Stebel, VT, in preparation]

PDFs used: MSTW2008NLO

PDFs used: MMHT2014NLO



NLO+NLL is NLL matched to NLO

preliminary

Results

[Kulesza, Motyka, Stebel, VT, in preparation]

PDFs used: MMHT2014NLO

\sqrt{S} [TeV]	NLO [fb]	NLO+NLL [fb]	Pdf uncertainty	K -factor
8	131.4 ^{+3.6%} _{-9.0%}	135.1 ^{+1.3%} _{-5.0%}	+3.0% -2.7%	1.028
13	506.3 ^{+5.7%} _{-9.3%}	517.0 ^{+4.0%} _{-5.9%}	+2.3% -2.3%	1.021
14	610.9 ^{+6.5%} _{-9.1%}	623.4 ^{+4.9%} _{-5.6%}	+2.3% -2.2%	1.020

preliminary

Resummation for differential observables

$$pp \rightarrow Q\bar{Q}B + X \quad \Rightarrow Q^2$$

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$$v_5 = \frac{(p_1 + p_2 - p_5)^2 - s_{34}}{\hat{s}} = \frac{\hat{s} + \tilde{t}_{15} + \tilde{t}_{25} + m_5^2 - s_{34}}{\hat{s}}$$

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$$pp \rightarrow Q\bar{Q} + X[B] \quad \Rightarrow P_{T,34}, y_{34}, Q_{34}^2$$

$$s_5 = \frac{(p_1 + p_2 - p_3 - p_4)^2 - m_5^2}{\hat{s}} = \frac{\hat{s} + \tilde{t}_{13} + \tilde{t}_{14} + \tilde{t}_{23} + \tilde{t}_{24} + s_{34} - m_5^2}{\hat{s}}$$

Threshold resummation in differential notation

$$\begin{aligned}
 W &= (1 - x_1) c_1 + (1 - x_2) c_2 + \omega \\
 &= \omega_1 c_1 + \omega_2 c_2 + \omega_s
 \end{aligned}$$

$$\begin{aligned}
 &\int dW e^{-NW} \frac{d\sigma_{ij \rightarrow c+X}}{d\hat{\Pi}_3} \\
 &= \int dW e^{-NW} H_{ij}(\hat{\Pi}_3) \int d\omega_1 d\omega_2 d\omega_s \delta(W - \omega_1 c_1 - \omega_2 c_2 - \omega_s) \\
 &\psi_{i/i}(\omega_1, p_i, n) \psi_{j/j}(\omega_2, p_i, n) S(\omega_s Q^2 / \mu^2, v_i, n) \\
 &= H_{ij}(\hat{\Pi}_3) \tilde{\psi}_{i/i}(Nc_1, p_i, n) \tilde{\psi}_{j/j}(Nc_2, p_i, n) \tilde{S}\left(\frac{Q^2}{\mu^2 N}, v_i, n\right)
 \end{aligned}$$

- $\omega_{1/2}$ Collinear weights
- ω_s Soft weight

Summary

Conclusions

- Improvement in scale uncertainty (factor 2)
- Small corrections (2%)

Outlook

- Different threshold definitions
- Differential distributions

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Thank you for your attention